

Mean Field Approach to BCS Theory

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1 BCS Hamiltonian

$$H_I = \frac{-V}{\Omega} \sum_{\vec{k}, \vec{p}} c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger c_{-\vec{p}\uparrow} c_{-\vec{p}\downarrow}$$

$$H_K = \sum_{\vec{k}, \sigma} \xi(\vec{k}) c_{\vec{k}\sigma}^\dagger c_{\vec{k}\sigma}$$

$$\xi(\vec{k}) = \frac{\hbar^2 k^2}{2m} - \mu$$

Order Parameter:

$$\Delta = \frac{V}{\Omega} \sum_{\vec{k}} \langle c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger \rangle$$

2 Mean Field Decoupling

Mean Field Decouple on dagger and non-dagger terms to get:

$$c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger = \langle c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger \rangle + \left(c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger - \langle c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger \rangle \right)$$

$$c_{-\vec{p}\uparrow} c_{-\vec{p}\downarrow} = \langle c_{-\vec{p}\uparrow} c_{-\vec{p}\downarrow} \rangle + \left(c_{-\vec{p}\uparrow} c_{-\vec{p}\downarrow} - \langle c_{-\vec{p}\uparrow} c_{-\vec{p}\downarrow} \rangle \right)$$

Define for simplicity of notation, for $\vec{k} \neq \vec{p}$:

$$a_{\vec{k}}^\dagger = c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger$$

$$a_{\vec{p}} = c_{-\vec{p}\uparrow} c_{-\vec{p}\downarrow}$$

Rewrite the interaction Hamiltonian as,

$$H_I = \frac{-V}{\Omega} \sum_{\vec{k}, \vec{p}} \left[\langle a_{\vec{k}}^\dagger \rangle + \left(a_{\vec{k}}^\dagger - \langle a_{\vec{k}}^\dagger \rangle \right) \right] \left[\langle a_{\vec{p}} \rangle + \left(a_{\vec{p}} - \langle a_{\vec{p}} \rangle \right) \right]$$

Multiplying out and ignoring the double fluctuation term for mean field approximation,

$$H_I = \frac{-V}{\Omega} \left(\sum_{\vec{k}, \vec{p}} \langle a_{\vec{k}}^\dagger \rangle \langle a_{\vec{p}} \rangle + \langle a_{\vec{p}} \rangle a_{\vec{k}}^\dagger - \langle a_{\vec{p}} \rangle \langle a_{\vec{k}}^\dagger \rangle + \langle a_{\vec{k}}^\dagger \rangle a_{\vec{p}} - \langle a_{\vec{k}}^\dagger \rangle \langle a_{\vec{p}} \rangle \right)$$

Using the order parameter notation, we get,

$$H_I = \sum_{\vec{k}} -\Delta^\dagger c_{-\vec{k}\downarrow} c_{\vec{k}\uparrow} - \Delta c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\downarrow}^\dagger - \frac{\Omega}{V} \Delta^\dagger \Delta$$

Adding the site hopping Hamiltonian,

$$H_k = \sum_{\vec{k}} \xi(\vec{k}) c_{\vec{k}\uparrow}^\dagger c_{(\vec{k})\uparrow} + \xi(\vec{k}) c_{\vec{k}\downarrow}^\dagger c_{\vec{k}\downarrow}$$

Thus for the full Hamiltonian,

$$H = H_k + H_I = \sum_{\vec{k}} -\Delta^\dagger c_{-\vec{k}\downarrow} c_{\vec{k}\uparrow} - \Delta c_{\vec{k}\uparrow}^\dagger c_{-\vec{k}\uparrow}^\dagger - \frac{\Omega}{V} \Delta^\dagger \Delta + \xi(\vec{k}) c_{\vec{k}\uparrow}^\dagger c_{(\vec{k})\uparrow} + \xi(\vec{k}) c_{\vec{k}\downarrow}^\dagger c_{\vec{k}\downarrow}$$

Using the Nambu Spinor defined as,

$$\Psi = \begin{bmatrix} c_{\vec{k}\uparrow} \\ c_{-\vec{k}\downarrow}^\dagger \end{bmatrix} \Psi^\dagger = \begin{bmatrix} c_{\vec{k}\uparrow}^\dagger & c_{-\vec{k}\downarrow} \end{bmatrix} \quad (1)$$

The full Hamiltonian can be written as

$$H + \xi(\vec{k}) = \sum_{\vec{k}} \Psi_{\vec{k}}^\dagger \begin{bmatrix} \xi(\vec{k}) & -\Delta \\ -\Delta^\dagger & -\xi(\vec{k}) \end{bmatrix} \Psi_{\vec{k}} - \frac{\Omega}{V} \Delta^\dagger \Delta$$

3 Brogliubuv Transformation

4 Diagonalization