

## ASTR 507 Problem Set 2

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### Problem 1

The heat capacity at constant pressure is defined as

$$C_p = \left( \frac{dU}{dT} \right)_P$$

And the heat capacity at constant volume is:

$$C_v = \left( \frac{dU}{dT} \right)_V$$

Starting with the first law of thermodynamics,  $dU = dQ - PdV$ , we note that the heat capacity at constant volume can also be written

$$C_v = \frac{dQ}{dT}$$

Setting this fact aside for now, we can differentiate the internal energy of the fluid to find the heat capacity at constant volume:

$$\left( \frac{dU}{dT} \right)_V = \frac{d}{dT} \left( \frac{f}{2} NkT \right) = \frac{f}{2} Nk$$

where  $f$  is the number of degrees of freedom in the energy of the molecule in question. Now for  $C_p$ , we have, from the first law,

$$C_p = \frac{dQ}{dT} - \frac{PdV}{dT}$$

The first term is just  $C_v$ . The second term can be found from the ideal gas law:  $PdV = d(NkT) = Nk$  and therefore

$$C_p = C_v - Nk$$

Combining these relationships, we arrive at  $C_v - C_p = Nk$ , and

$$\frac{C_p}{C_v} = \frac{C_v - Nk}{C_v} = 1 - \frac{Nk}{C_v} = 1 - \frac{2}{f} = \gamma$$

### Problem 2

We can find the speed of the fastest electron by solving the equation  $nVB(v) = 1$  to find the lowest velocity above which we would expect to find only a single electron. Doing this numerically for the Boltzmann distribution with the given parameters ( $m = m_e$ ,  $T = 10^6$  K,  $n = 0.01 \text{ cm}^{-3}$ , and  $V = 40 \text{ pc}^3$ ), I find  $v_{\text{max}} \sim 6 \times 10^7 \text{ m/s}$ . This is only about a factor of 5 less than the velocity of a cosmic ray, but that difference is still significant. The probability of finding a particle in this distribution with velocity  $v \sim c$  is essentially zero. Of course, the Maxwellian velocity distribution is also non-relativistic. A properly relativistic velocity distribution should drop off even more quickly at velocities near  $c$ . Otherwise a sufficiently high temperature could drive electrons to faster than light speed. Thus cosmic rays cannot have a thermal origin.

## Problem 3

### Part A

The mean free path at the escape density is:

$$l = 1/n_{\text{esc}}\sigma$$

Setting this equation to  $H = kT/(mg)$ , we find:

$$n_{\text{esc}} = \frac{mg}{kT\sigma}$$

### Part B

Integrate:

$$\phi(v)dv = \frac{f(v)}{4\pi}dv \int_0^{2\pi} d\phi \int_0^{\pi/2} v \cos \theta \sin \theta d\theta$$

The integral over  $d\phi$  yields a factor of  $2\pi$ , and the second integral is evaluated as

$$\int_0^{\pi/2} v \cos \theta \sin \theta d\theta = \frac{v}{2} \sin^2 \theta \Big|_0^{\pi/2} = \frac{v}{2}$$

So the complete integral evaluates to

$$\phi(v)dv = \frac{f(v)}{4}v dv$$

### Part C

To simplify things, let's consider the integral

$$\phi = \frac{A}{4} \int_{v_{\text{esc}}}^{\infty} v^3 e^{-bv^2}$$

At the end of the evaluation we can replace  $A$  with the proper set of constants for the Maxwellian velocity distribution, and  $b$  with  $b = m/2kT$ . Integrating by parts, we have

$$\int_{v_{\text{esc}}}^{\infty} v^3 e^{-bv^2} = -\frac{v^2}{2b} e^{-bv^2} \Big|_{v_{\text{esc}}}^{\infty} + \frac{1}{b} \int_{v_{\text{esc}}}^{\infty} v e^{-bv^2} = \left( -\frac{v^2}{2b} + \frac{1}{2b^2} \right) e^{-bv^2} \Big|_{v_{\text{esc}}}^{\infty}$$

Evaluating between the two limits and multiplying the  $A/4$  factor back onto the result, we have

$$\phi = \frac{A}{4} \left( \frac{v_{\text{esc}}^2}{2b} + \frac{1}{2b^2} \right) e^{-bv_{\text{esc}}^2}$$

Substituting in the constants for  $A$  and  $b$ , we have

$$\phi = \frac{\pi m^{3/2} n}{(2\pi kT)^{3/2}} \left( \frac{kT v_{\text{esc}}^2}{m} + \frac{2(kT)^2}{m^2} \right) e^{-mv_{\text{esc}}^2/2kT}$$

Finally, in terms of  $v_s$  and  $\lambda_{\text{esc}}$ , we have

$$\phi = \frac{nv_s}{2\sqrt{\pi}} (\lambda_{\text{esc}} + 1) e^{-\lambda_{\text{esc}}}$$

### **Part D**

Over 1 Gyr the Earth would lose  $4.2 \times 10^{42}$  molecular hydrogen molecules, nearly half of its current reservoir of hydrogen.

### **Part E**

Over this same period of time the it is unlikely that even a single oxygen molecule will escape the atmosphere through this mechanism. I calculate that about  $3 \times 10^{-54}$  molecules escape over 1 Gyr. Thus, over time, the ratio of oxygen to hydrogen molecules should increase as hydrogen escapes at oxygen stays put.

### **Part E**

For deuterium-hydrogen molecules, I calculate that  $4 \times 10^{39}$  molecules escape over 1 Gyr, compared to a thousand times more molecules for hydrogen alone. As a result, the D/H ratio should increase over time.

My calculations for parts D-E are on my github: