



**CE-580**

**COMPUTATIONAL TECHNIQUES  
FOR  
FLUID DYNAMICS**

**HOMEWORK #3**

**Turbulent Flow Between  
Parallel Plates**

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## Calculations

### Assumptions

- Steady flow
- Incompressible flow
- Fully developed turbulent flow
- 1-D uniform flow

### RANS equation in x- direction

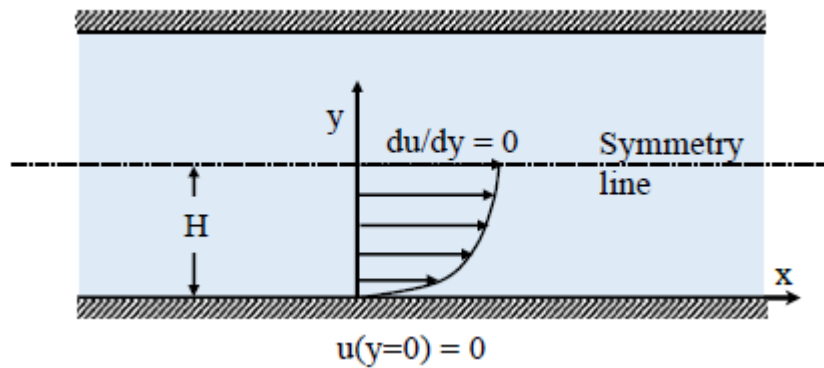


Figure 1: Problem Domain

$$\rho \left( \frac{\partial u}{\partial t} + \frac{\partial u^2}{\partial x} + \frac{\partial \tilde{u}'^2}{\partial x} + \frac{\partial uv}{\partial y} + \frac{\partial \tilde{u}'\tilde{v}'}{\partial y} + \frac{\partial uw}{\partial z} + \frac{\partial \tilde{u}'\tilde{w}'}{\partial z} \right) = -\frac{\partial p}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right)$$

(1)    (2)    (3)    (4)    (5)    (6)    (7)    (8)    (9)    (10)    (11)

Simplify above equation according to our domain and assumptions

- Term (1) drops due to steady flow assumption
- Fully developed so there are no u gradients in x- direction term (2), (3),(9) drops
- Uniform flow, term (4) drops
- 2-D domain term (6),(7),(11) drops

Simplified equation is

$$\rho \frac{\partial \tilde{u}'\tilde{v}'}{\partial y} = -\frac{\partial p}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial y^2} \right)$$

Partial derivatives can be replaced by total derivatives since each dependent variable is dependent one space parameter

$$\rho \frac{d}{dy} \left[ \tilde{u}'\tilde{v}' - \mu \frac{du}{dy} \right] = -\frac{dp}{dx}$$

Using Boussinesq hypothesis for turbulent stress

$$\rho \tilde{u}'\tilde{v}' = \mu_t \frac{du}{dy}$$

Combining the viscous and turbulence resistance terms

$$\frac{d}{dy} \left[ (\mu + \mu_t) \frac{du}{dy} \right] = \frac{dp}{dx}$$

Introducing  $\mu_e$  for effective viscosity and  $C_p$  for pressure gradient

$$\frac{d}{dy} \left[ \mu_e \frac{du}{dy} \right] = C_p$$

Direct integration of above term is not possible due to nonlinearity introduced with viscosity term

### Grid

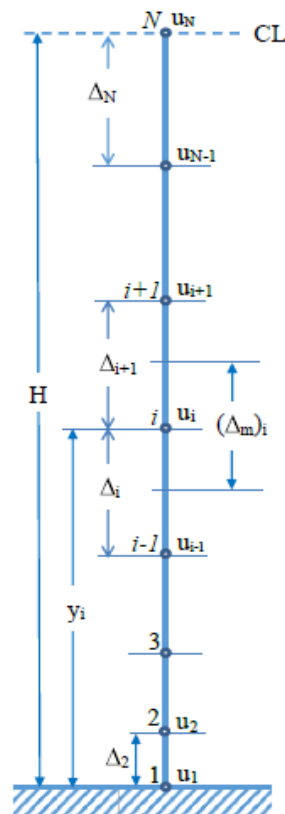


Figure 2: Grid

Grid is generated using constant ratio method

Ratio between any two neighbouring meshes is constant

$$\Delta Y_i = \beta \Delta_{i+1}$$

Distance between the last two mesh point can be calculated as

$$\Delta Y_N = \frac{H}{\sum_{i=0}^{N-2} \beta^i}$$

Mesh distribution can be completed marching down from N to 1

### Discretization

Governing equation can be discretized on domain as

$$\frac{\left(\mu_e \frac{du}{dy}\right)_{i+\frac{1}{2}} - \left(\mu_e \frac{du}{dy}\right)_{i-\frac{1}{2}}}{\frac{\Delta_{i+1} + \Delta_i}{2}} = Cp$$
$$\frac{(\mu_e)_{i+\frac{1}{2}} \frac{u_{i+1} - u_i}{\Delta_{i+1}} - (\mu_e)_{i-\frac{1}{2}} \frac{u_i - u_{i-1}}{\Delta_i}}{\frac{\Delta_{i+1} + \Delta_i}{2}} = Cp$$

**Note** = half term convention is  $i+1/2 \Rightarrow i+1$  ,  $i-1/2 \Rightarrow i$

Collecting the coefficients of the unknown velocity terms

$$\left[ \frac{(\mu_e)_{i+\frac{1}{2}}}{\frac{\Delta_{i+1}(\Delta_{i+1} + \Delta_i)}{2}} \right] u_{i-1} - \left[ \frac{(\mu_e)_{i+\frac{1}{2}}}{\frac{\Delta_{i+1}(\Delta_{i+1} + \Delta_i)}{2}} + \frac{(\mu_e)_{i-\frac{1}{2}}}{\frac{\Delta_i(\Delta_{i+1} + \Delta_i)}{2}} \right] u_i + \left[ \frac{(\mu_e)_{i-\frac{1}{2}}}{\frac{\Delta_i(\Delta_{i+1} + \Delta_i)}{2}} \right] u_{i+1} = Cp$$

We can write above equation such as

$$-A_i u_{i-1} + B_i u_i - C_i u_{i+1} = D_i$$

### Turbulence Model

To obtain turbulent viscosity can be obtained from mixing length theory

$$\mu_t = \rho \nu_t = \rho l_m^2 \frac{du}{dy}$$

Where

$$l_m = H \left[ 0.14 - 0.08 \left( 1 - \frac{y}{H} \right)^2 - 0.06 \left( 1 - \frac{y}{H} \right)^4 \right] f_\mu$$

$$f_\mu = 1 - \exp \left( - \frac{y^+}{A^+} \right) \quad A^+ = 26$$

$$y^+ = \frac{y u_*}{\nu} \quad u_* = \sqrt{\frac{\tau_w}{\rho}}$$

$\tau_w$  can be calculated by integrating the momentum equation

$$\frac{d}{dy} \left[ \mu_e \frac{du}{dy} \right] = Cp, \quad \frac{d\tau}{dy} = Cp, \quad \tau = C_p \times y + C_1$$

Applying the boundary condition at  $y=H, \tau = 0$  ,  $C_1$  calculated as

$$C_1 = -C_p H$$

$$\tau = C_p (y - H)$$

$$\tau = \tau_w = -C_p H$$

### Boundary Conditions

A different approach from previous homework is taken to impose boundary conditions.

Consider below equation

$$-\left[\frac{(\mu_e)_{i-\frac{1}{2}}}{\frac{\Delta_i(\Delta_{i+1} + \Delta_i)}{2}}\right]u_{i-1} + \left[\frac{(\mu_e)_{i+\frac{1}{2}}}{\frac{\Delta_{i+1}(\Delta_{i+1} + \Delta_i)}{2}} + \frac{(\mu_e)_{i-\frac{1}{2}}}{\frac{\Delta_i(\Delta_{i+1} + \Delta_i)}{2}}\right]u_i - \left[\frac{(\mu_e)_{i+\frac{1}{2}}}{\frac{\Delta_{i+1}(\Delta_{i+1} + \Delta_i)}{2}}\right]u_{i+1} = -C_p$$
$$A_i u_{i-1} + B_i u_i + C_i u_{i+1} = D_i$$

At  $i=2$ , impose no slip boundary condition

$$A_2 u_1 + B_2 u_2 + C_2 u_3 = D_2$$

$u_1 = 0$  at the wall, so first term drops

$$B_2 = \frac{(\mu_e)_3}{\frac{\Delta_3(\Delta_3 + \Delta_2)}{2}} + \frac{(\mu_e)_2}{\frac{\Delta_2(\Delta_3 + \Delta_2)}{2}}$$
$$C_2 = \left[\frac{(\mu_e)_3}{\frac{\Delta_2(\Delta_3 + \Delta_2)}{2}}\right]$$
$$D_2 = -C_p$$

At  $i=N-1$ , impose symmetry line, meaning  $u_N = u_{N-1}$ . This makes turbulent velocity zero and cancels the coefficients in B term

$$A_{N-1} u_{N-2} + B_{N-1} u_{N-1} + C_{N-1} u_{N-1} = D_{N-1}$$
$$\mu_e = \mu$$
$$A_{N-1} = -\left[\frac{(\mu_e)_{N-1}}{\frac{\Delta_{N-1}(\Delta_N + \Delta_{N-1})}{2}}\right]$$
$$B_{N-1} - C_{N-1} = B_{N-1} = \left[\frac{(\mu_e)_{N-1}}{\frac{\Delta_{N-1}(\Delta_N + \Delta_{N-1})}{2}}\right]$$
$$D_{N-1} = -C_p$$

When three diagonal system of equations are solved with modified coefficients at  $i=2$  and  $i=N-1$ , boundary conditions will be imposed automatically.

### Error Calculations

Errors are calculated at each iterations as

$$Error = \frac{1}{(N-1)U_N} \sum_{i=2}^N |u_i^n - u_i^{n+1}|$$

### ***Solution Parameters***

$$H = 0.02 \text{ m}, \quad C_p = -100., -1000., -10000. \frac{N}{m^3},$$

$$B = 0.96, 0.94, 0.93 \text{ (for respective pressure gradients)}$$

$$\mu = 0.001 \text{ N} \cdot \frac{s}{m^2}, \quad \rho = 1000 \frac{kg}{m^3}, \quad N = 101$$

### ***Solution algorithm***

1. Initialize solution
  - a. Generate grid
  - b. Calculate laminar boundary layer
2. Start iteration 1 to 100
  - a. Calculate turbulent stresses
    - i. From laminar profile at first iteration
    - ii. From previous solution as the solution advances
  - b. Calculate coefficients of discretized equation with boundary conditions
  - c. Call Thomas algorithm to get a solution from three diagonal system of equations
  - d. Extract the solution from Thomas subroutine
  - e. Calculate error after the first iteration using new and old values of velocity distributions
  - f. Output error at every iteration and final velocity profile
3. Stop iterations
4. Print out  $C_p$  and maximum velocity

## Results and Discussion

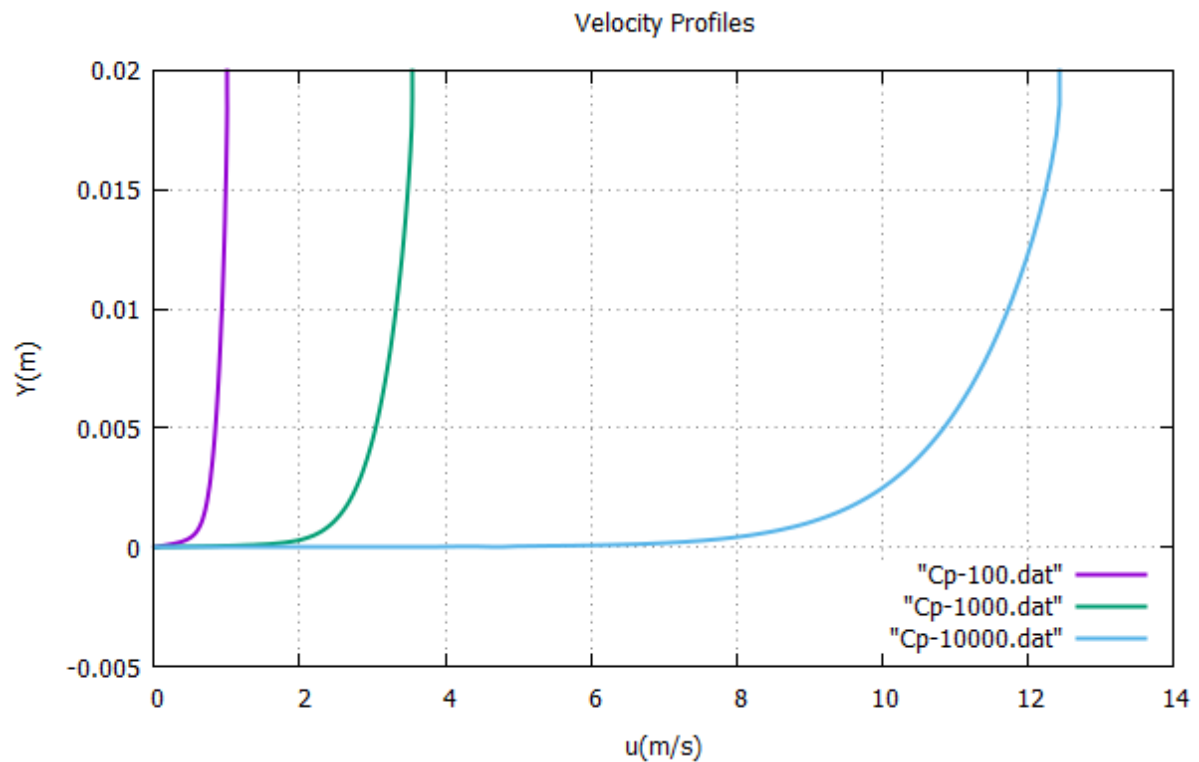


Figure 3: Velocity Profiles

Cp	U_max (m/s) at centerline
-100	1.004
-1000	3.553
-10000	12.453

As seen from figure 3 velocity profiles are very steep and contains very high gradients near the wall zone. The high gradients near wall is an effect of viscous damping function. To catch the high gradients accurately we need very dense mesh close to wall regions. The velocity profile flattens away from the wall and gradients can be calculated with bigger steps. This is the logic behind the constant ratio mesh sizing.

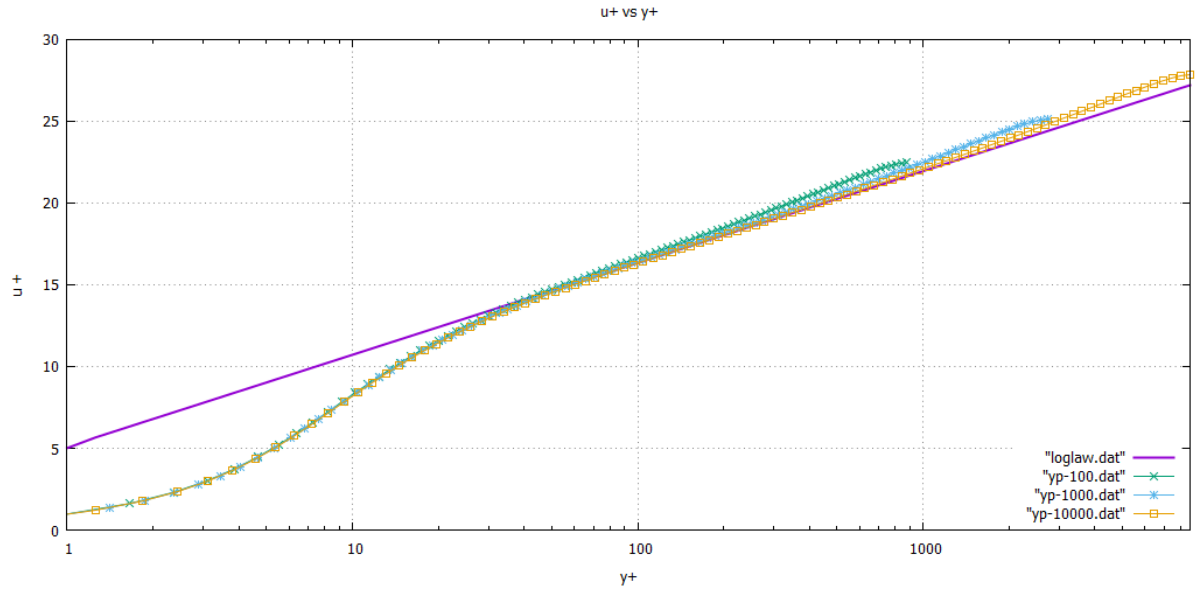


Figure 4:  $u^+$  vs  $y^+$

Figure 4 is another way to view velocity profile using dimensionless quantities. Where

$$u^+ = \frac{u}{u_*} \quad u_* = \text{shear velocity}$$

$$y^+ = \frac{yu_*}{\nu}$$

Log-Law also called Law of the Wall is calculated

$$u^+ = \frac{1}{0.41} \times \ln(y^+) + 5.1$$

The variation of  $u^+$  and  $y^+$  fits very well the log-law starting from  $y^+ \approx 30$ , Meaning above equation can be used to calculate velocities without finite differences up to the so called outer layer. But as  $y^+$  gets smaller there is buffer and viscous sublayer where this correlation does not hold. Thus we need very refined mesh to catch the solution in viscous sublayer. This requires first grid point is at where  $y^+ < 5$ , ideally at  $y^+ = 1$ . In our cases, grid is well refined at the close wall regions and viscous sublayer is well solved



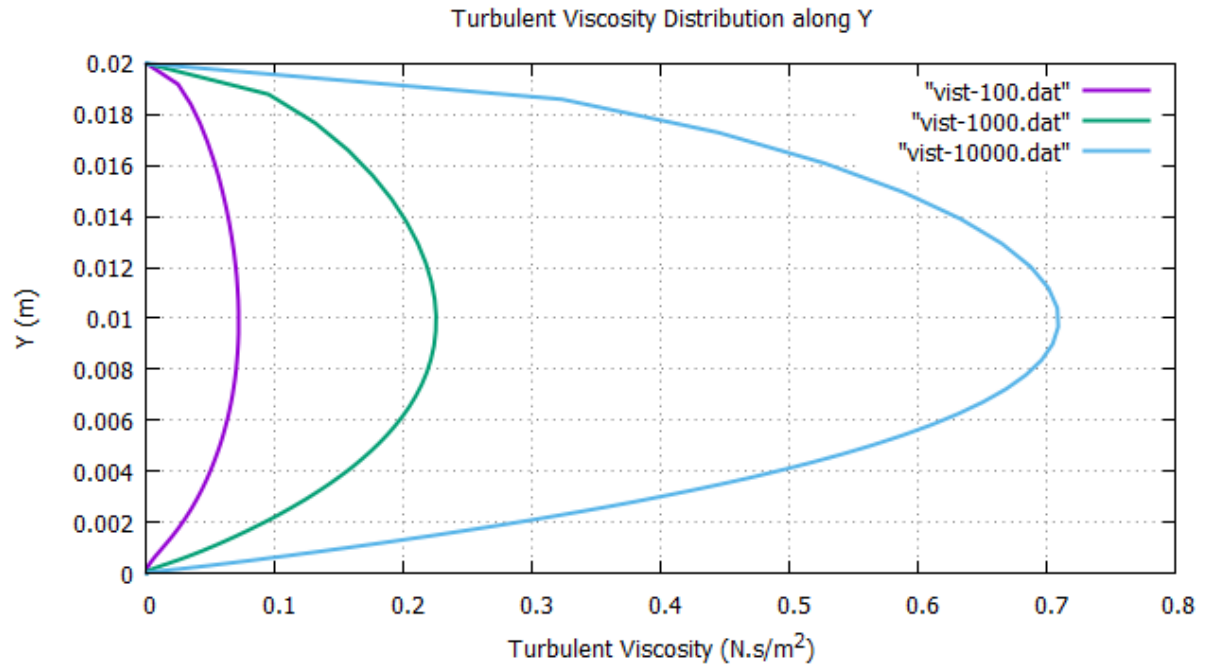


Figure 5: Turbulent Viscosity Distribution with Different Cp

Considering velocity gradients close to the wall zone, turbulent viscosities are expected to be large close to wall. But it is not the case as figure 5 shows. The reason of the small turbulent viscosities is the turbulence model. Viscous damping function and mixing length makes necessary corrections and region close to the wall turbulent stresses are damped. As going away from the wall, velocity gradients become smaller but  $f_\mu$ ,  $l_m$  gets into account so that we get such a turbulent viscosity distribution.

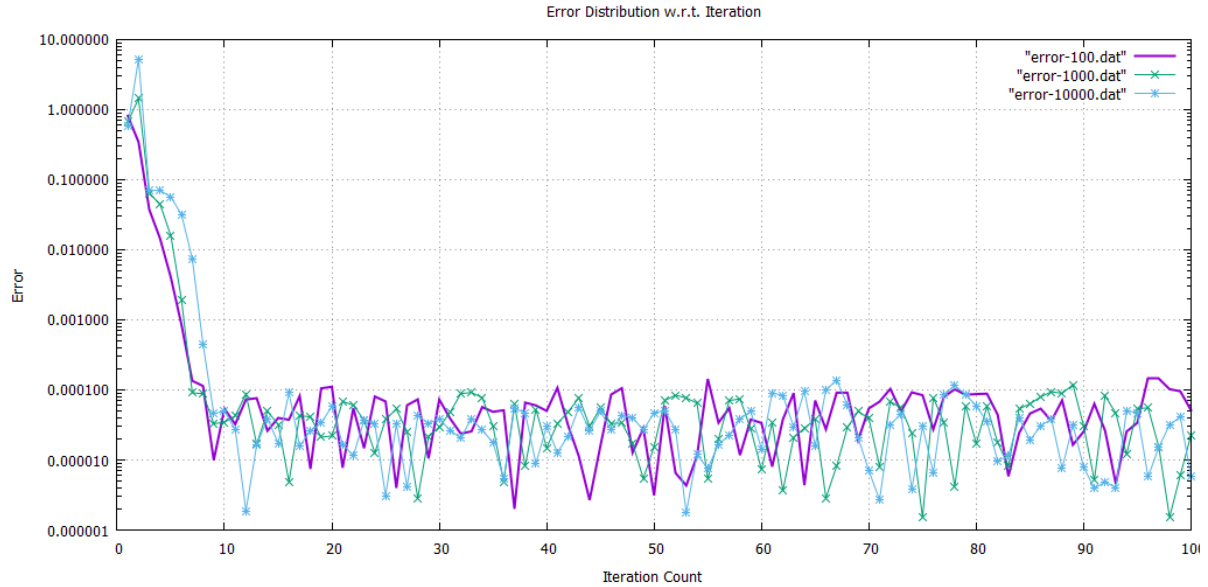


Figure 6: Error Variation Along Iteration Steps

A laminar velocity profile is used for starting point of iterations. But in laminar flow the velocity gradients in y- direction is much greater than a turbulent flow. The result is very high turbulent stresses to start with. High stresses cause almost constant and small velocities and next turbulent stresses will be very low. This will cause oscillations and convergence will be slow. To overcome this problem at first iteration get half of the laminar gradients to start with. And take average of turbulent viscous stresses with the old ones.

Thanks to averaging error magnitude drops to  $1E-6$  levels after 10 iteration as seen in figure 6. As solution marches there is fluctuations in error magnitude caused by round off and truncation errors.

## Fortran Code

```
program TurbulentBetweenParallelPlates
c Taha Yaşar Demir / 1881978
c CE-580 - Homework #3
    parameter (mx=101)
    common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
    common/grid/ Beta, y(mx),yc(mx),dy(mx),N
    common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau
    common/coef/ A(mx), B(mx), C(mx), D(mx)
    common/error/u_old(mx),error(mx)
c  sml = mixinglength , fu = viscous damping function, yp=y+, Ap=A+
c  us = u_star, rnu = kinematic viscosity, tau = wall shear

    open(1,file='error.dat')
    open(2,file='yplus.dat')
    open(3,file='velpr.dat')
    open(4,file='loglaw.dat')
    open(6,file='vist.dat')

    call init
    do l=1,100
        call stress(l) ! Evaluate stresses
        call coefficients ! Get Three diagonal system coefficients
        call THOMAS(2,N-1,A,B,C,D) ! Returns the solution in D vector
        u(1) = 0. ! no-slip Boundary Condition
        u(N) = D(N-1) ! Neumann Boundary at centerline
        do k = 2,N-1 ! extract the solution
            u(k) = D(k)
        enddo
        if (i.gt.1) call error_cal(l) ! After first iteration calculate error
        do k=1,N
            u_old(k) = u(k)
        enddo
        call output(l) ! write solution out in relative files
    enddo

    print*, Cp, u(N) ! get max velocity at centerline

close(1)
close(2)
close(3)
close(4)
close(6)

stop
end
```

Main Program

subroutine init

```
parameter (mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N
common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau
common/coef/ A(mx), B(mx), C(mx), D(mx)
```

```
H = 0.02 ! m
vis = 0.001 ! N.s/m^2
rho = 1000.0 ! kg/m^3
N = mx
print*, "Enter the pressure coef. Cp : "
read(*,*) Cp
if (Cp.eq.-100.) then
    Beta = 0.96
elseif(Cp.eq.-1000.) then
    Beta = 0.94
else
    Beta = 0.93
endif
call makegrid
call analytic

return
end
```

Initialize Solution Parameters and Get Grid and Laminar Solution

subroutine makegrid

```
parameter (mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N
common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau
common/coef/ A(mx), B(mx), C(mx), D(mx)
```

```
real sum
sum = 0.0
do i=0,N-2
    sum = sum+ Beta**i
enddo
dy(N) = H/sum
y(N) = H
```

```
do i=N,2,-1
    y(i-1) = y(i)-dy(i)
    dy(i-1) = Beta*dy(i)
    yc(i) = (y(i)+y(i-1))/2
enddo
```

```
return
end
```

Calculate Grid Spacing and Coordinates

```

subroutine analytic
parameter (mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N

do i=1,N
    ua(i) = -500*y(i)**2 + 200*y(i)
enddo
ua(1) = 0.

return
end

```

Laminar Solution

```

subroutine stress(iter)
parameter (mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N
common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau

Ap = 26
rnu = vis/rho
tau = -Cp*H
us = sqrt(tau/rho)
do k=2,N
    yp(k) = yc(k)*us/rnu
    fu(k) = 1-exp(-yp(k)/Ap)
    sml(k) = H*(0.14-0.08*(1-(yc(k)/H))**2
    &      -0.06*(1-(yc(k)/H))**4)*fu(k)
    if (iter.eq.1) then
        vist(k) = 0.5*(rho*sml(k)**2)*(ua(k)-ua(k-1))/dy(k)
    else
        vist(k) = 0.5*(vist(k)+(rho*sml(k)**2)*(u(k)-u(k-1))/dy(k))
    endif
    vise(k) = (vist(k)+vis)
enddo

return
end

```

Effective Stress Calculation

```

subroutine coefficients
parameter (mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N
common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau
common/coef/ A(mx), B(mx), C(mx), D(mx)
do i=2,N-1
if (i.eq.2) then ! no slip Boundary Conditions
    A(i) = 0
    C(i) = -vise(i+1)/(dy(i+1)*(dy(i+1)+dy(i))/2)
    B(i) = (vise(i)/(dy(i)*(dy(i+1)+dy(i))/2)) - C(i)
    D(i) = -Cp
elseif (i.eq.N-1) then ! Neumann Boundary Condition u_n=u_n-1
    A(i) = -vise(i)/(dy(i)*(dy(i+1)+dy(i))/2)
    C(i) = 0
    B(i) = -A(i)
    D(i) = -Cp
else
    A(i) = -vise(i)/(dy(i)*(dy(i+1)+dy(i))/2)
    C(i) = -vise(i+1)/(dy(i+1)*(dy(i+1)+dy(i))/2)
    B(i) = -A(i) - C(i)
    D(i) = -Cp
endif
enddo
return
end

```

3-Diagonal System coefficients

```

subroutine error_cal(ite)
parameter (mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N
common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau
common/coef/ A(mx), B(mx), C(mx), D(mx)
common/error/u_old(mx),error(mx)

error(ite) = 0.
do i=1,N
error(ite) = error(ite) + (1/((N-1)*u(N)))*abs(u_old(i)-u(i))
enddo

return
end

```

Error Calculation

```

subroutine output(it)
parameter(mx=101)
common/flow/ rho, H, Cp, vis, vist(mx), vise(mx), u(mx), ua(mx)
common/grid/ Beta, y(mx),yc(mx),dy(mx),N
common/turb/ sml(mx), fu(mx), yp(mx), Ap, us, rnu, tau
common/error/u_old(mx),error(mx)
real law(mx),up(mx)

write(1,*) it,error(it)
if(it.eq.1) write(3,*) u(1),y(1),ua(1)
if(it.eq.100) then
    do i=2,N
        law(i) = (1/0.41)*log(yp(i))+5.1
        up(i) = 0.5*(u(i)+u(i-1))/us
        write(2,*) yp(i),up(i)
        write(3,*) u(i),y(i),ua(i)
        write(4,*) yp(i),law(i)
        write(6,*) vist(i),y(i)
    enddo
endif
return
end

```

Write out the Solution

```

subroutine THOMAS(il,iu,aa,bb,cc,ff)
C.....
c Solution of a tridiagonal system of n equations of the form
c  $A(i)*x(i-1) + Bb(i)*x(i) + C(i)*x(i+1) = R(i)$  for i=il,iu
c the solution X(i) is stored in F(i)
c A(il-1) and C(iu+1) are not used
c A,Bb,C,R are arrays to bbe provided bby the user
C.....
parameter (mx=101)
dimension aa(mx),bb(mx),cc(mx),ff(mx),tmp(mx)

tmp(il)=cc(il)/bb(il)
ff(il)=ff(il)/bb(il)
ilp1 = il+1
do i=ilp1,iu
    z=1./(bb(i)-aa(i)*tmp(i-1))
    tmp(i)=cc(i)*Z
    ff(i)=(ff(i)-aa(i)*ff(i-1))*z
enddo
iupil=iu+il
do ii=ilp1,iu
    i=iupil-ii
    ff(i)=ff(i)-tmp(i)*ff(i+1)
enddo
return
end

```

Thomas Algorithm

