Instabilities of scalar fields around oscillating stars

Taishi Ikeda w/ Vitor Cardoso, Miguel Zilhao

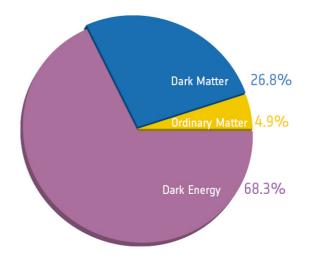






Scalar field beyond GR and SM

- Mystery of our Universe
 - Dark matter ?? Dark energy ??
 - Quantum theory of gravity ??



Scalar-tensor theory is natural extension of general relativity.

$$\mathscr{L} = \frac{R}{16\pi} - \frac{1}{2} g^{\mu\nu} \nabla_{\mu} \Phi \nabla_{\nu} \Phi - V(\Phi) + \cdots$$

The scalar field may have nonminimal coupling to matter.

$$\mathcal{L} \supset \mathcal{L}_{\text{matter}}(\Psi_{\text{m}}, A(\Phi)^2 g_{\mu\nu})$$

- Chameleon screening mechanism
- Spontaneous scalarization in neutron star et al....

Scalar field beyond GR and SM

$$\Box \Phi = V_{\rm eff}'(\Phi,\rho) \qquad V_{\rm eff}(\Phi,\rho) = \frac{1}{2} \left(\mu_0^2 - |\beta| \rho \right) \Phi^2 \qquad \left(A(\Phi) = e^{\frac{\beta}{2} \Phi^2} \right)$$
 threshold for scalarizarion :
$$|\beta_c^*| \sim \left(1 + \left(\mu_0 R_* \right)^2 \right) C_*^{-1} \qquad \beta < 0$$

- Spontaneous scalarization in neutron star (Damour et al (1993))
 - around Sun or WD : $\Phi = 0$

$$|\beta_c^{\text{NS}}| < |\beta| < |\beta_c^{\text{WD}}|$$

• around NS : $\Phi \neq 0$ • scalar hair

$$C_{\odot} \qquad C_{\text{WD}} \qquad C_{\text{NS}}$$

$$10^{-6} \qquad 10^{-3} \qquad 0.3$$

$$C = \frac{2M}{R}$$

$$\Phi = 0 \qquad \Phi \neq 0$$

Scalar field beyond GR and SM

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 threshold for scalarizarion :
$$|\beta_c^*| \sim \left(1 + \left(\mu_0 R_* \right)^2 \right) C_*^{-1} \qquad \beta < 0$$

- No Spontaneous scalarization in all normal stars
 - around all normal stars : $\Phi = 0$ $|\beta| < |\beta_c^{NS}|$
 - How to constraint the parameter region ??

$$C_{\odot} \qquad C_{\text{WD}} \qquad C_{\text{NS}}$$

$$10^{-6} \qquad 10^{-3} \qquad 0.3 \qquad C = \frac{2M}{R}$$

$$\Phi = 0 \qquad \Phi \neq 0$$

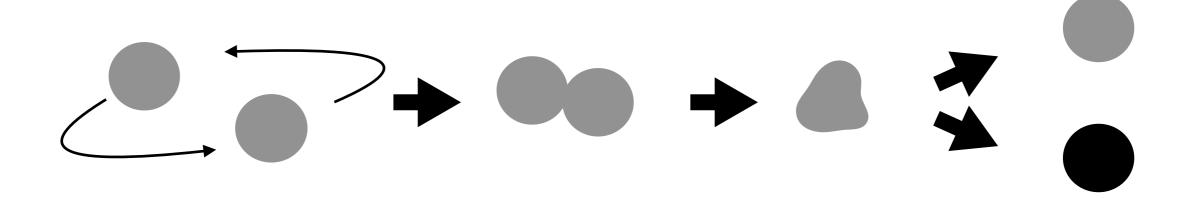
When does the field become relevant??

Our model and parameter region

$$|\beta| < |\beta_c^{NS}|$$

$$\mathcal{L} = \frac{R}{16\pi} - \frac{1}{2}g^{\mu\nu}\nabla_{\mu}\Phi\nabla_{\nu}\Phi - V(\Phi) + \mathcal{L}(A(\Phi)^{2}g, \Psi_{m}) \qquad A(\Phi) = e^{\frac{\beta}{2}\Phi^{2}}$$

- The stationary profile around all normal stars is same as GR. $\Phi = 0$
- The dynamical situations ??
 - eg) Neutron star merger remnants et al

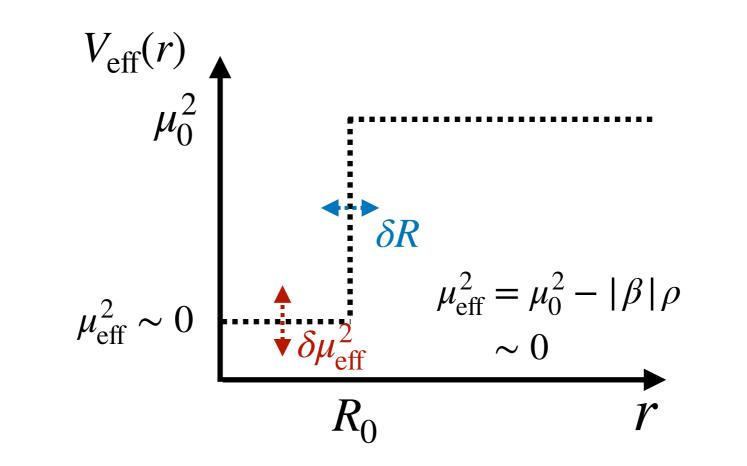


Effective potential

$$\Phi(t, r, \theta, \phi) = \sum_{l,m} \frac{\psi_{lm}(t, r)}{r} Y_{lm}(\theta, \phi)$$

Effective potential of scalar field around relativistic star.

$$-\partial_t^2 \psi - \frac{\alpha^2}{a^2} \partial_r \left(\log \frac{a}{\alpha} \right) \partial_r \psi + \frac{\alpha^2}{a^2} \partial_r^2 \psi - V_{\text{eff}}(r) \psi = 0$$
$$V_{\text{eff}}(r) = \dots + \alpha^2 \left\{ \mu_0^2 - |\beta| \left(\tilde{\rho} - 3\tilde{p} \right) \right\}$$



Effective potential

$$\Phi(t, r, \theta, \phi) = \sum_{l,m} \frac{\psi_{lm}(t, r)}{r} Y_{lm}(\theta, \phi)$$

• Effective potential of scalar field around relativistic star.

$$-\partial_{t}^{2}\psi - \frac{\alpha^{2}}{a^{2}}\partial_{r}\left(\log\frac{a}{\alpha}\right)\partial_{r}\psi + \frac{\alpha^{2}}{a^{2}}\partial_{r}^{2}\psi - V_{\mathrm{eff}}(r)\psi = 0$$

$$V_{\mathrm{eff}}(r) = \cdots + \alpha^{2}\left\{\mu_{0}^{2} - |\beta|\left(\tilde{\rho} - 3\tilde{p}\right)\right\}$$

$$V_{\mathrm{eff}}(r)$$

$$\mu_{0}^{2}$$

$$V_{\mathrm{eff}}(r)$$

$$\mu_{0}^{2}$$

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$$\mu_{0}^{2}$$

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$$V_{\mathrm{eff}}(r)$$

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$$V_{\mathrm{eff}}(r)$$

$$V_{\mathrm{eff}$$

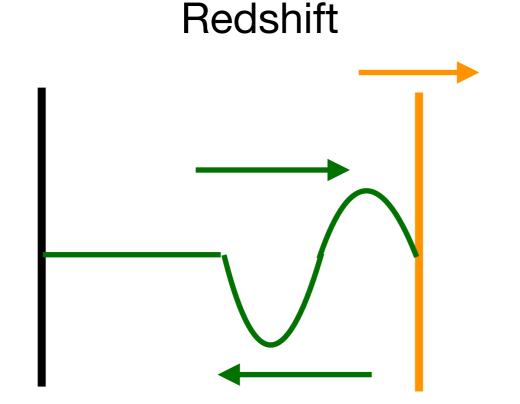
Effect of δR

Blueshift instability (1+1 dim)

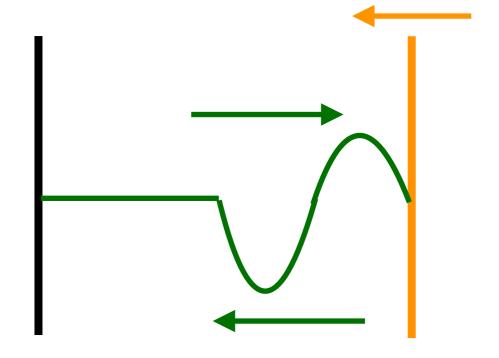
• 1+1 dim. wave eq. with oscillating boundary (Dittrich et al 1994)

$$-\partial_t^2 \Phi + \partial_x^2 \Phi = 0 \quad \text{with} \quad \begin{cases} \Phi(t,0) = 0 \\ \Phi(t,L(t)) = 0 \end{cases}$$

$$L(t) = L_0 + \delta L \sin(\omega t)$$



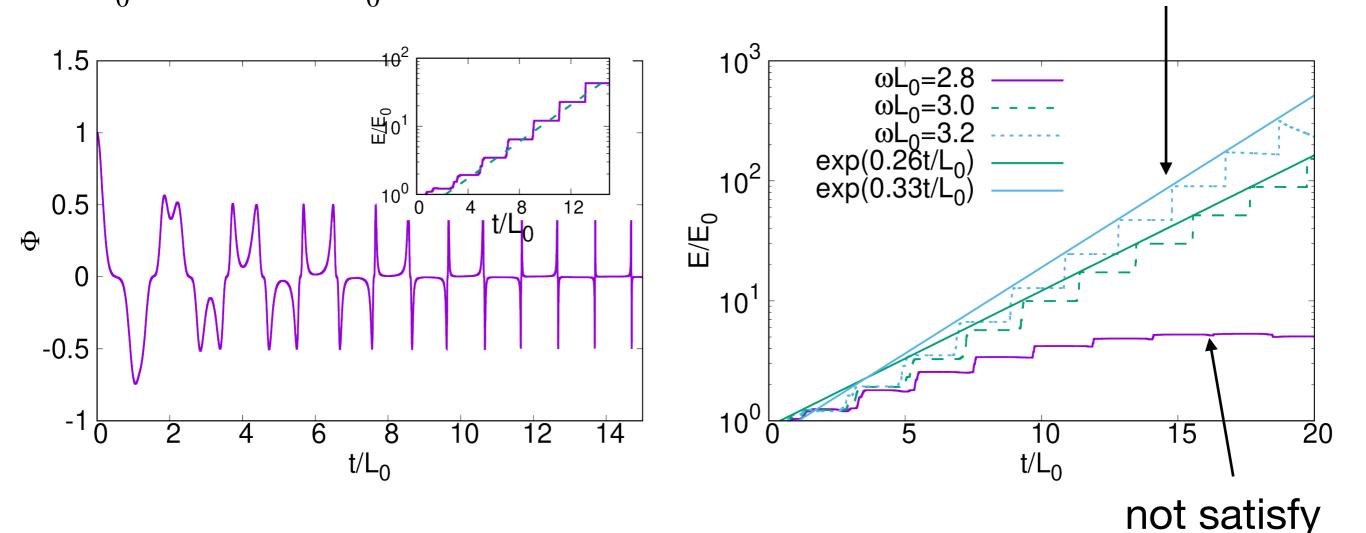
Blueshift



Blueshift instability (1+1 dim)

Condition for the blueshift instability

$$\frac{\pi}{L_0 + \delta L} < \omega < \frac{\pi}{L_0 - \delta L}$$
 The blueshift accumulates. $E \propto e^{\lambda_B(\omega, \delta L)t}$



Blueshift instability (3+1 dim)

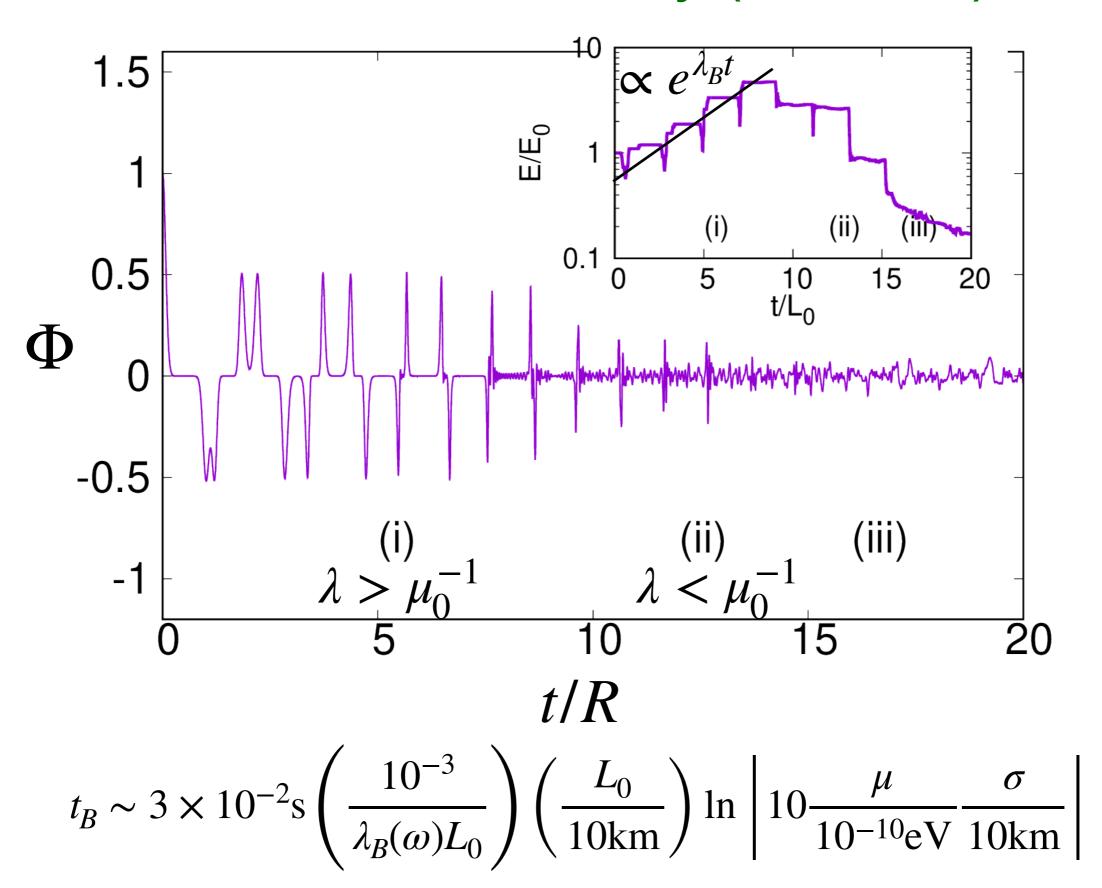
• 3+1 evolution in flat spacetime

$$\Box_{\text{flat}} \Phi = V_{\text{eff}}(r) \Phi \qquad V_{\text{eff}}(r) = \begin{cases} \mu_0^2 & (R + \delta R(t) < r) \\ \mu_{\text{eff}}^2 & (r < R + \delta R(t)) \end{cases}$$

$$\delta R(t) \sim \sin(\omega t)$$

$$\lambda V_{\text{eff}}(r)$$

Blueshift instability (3+1 dim)



Effect of $\delta\mu^2$

Parametric instability (1+1 dim)

1+1 dim. wave eq. with oscillating effective mass

$$-\partial_t^2 \Phi + \partial_x^2 \Phi - \mu^2(t) \Phi = 0$$
$$\mu^2(t) = \bar{\mu}^2 + \delta \mu^2 \sin(\omega t)$$

This equation can be reduced to Mathieu equation.

$$\partial_{\tau}^{2} \phi_{k} + \left(\delta + 2 \epsilon \sin(\tau)\right) \phi_{k} = 0$$

$$\delta = \frac{\kappa^{2} + \bar{\mu}^{2}}{\omega^{2}}$$

Instability condition

$$\delta \sim \frac{n^2}{4}$$

Parametric instability (3+1 dim)

• 3+1 evolution in flat spacetime

$$\square_{\text{flat}} \Phi = V_{\text{eff}}(r) \Phi$$

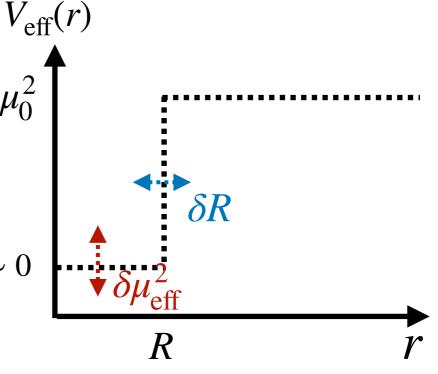
$$V_{\rm eff}(r) = \begin{cases} \mu_0^2 & (R + \delta R < r) \\ \mu_{\rm eff}^2 + \delta \mu_{\rm eff}^2 \sin(\omega t) & (r < R + \delta R) & \mu_{\rm eff}^2 \sim 0 \end{cases}$$

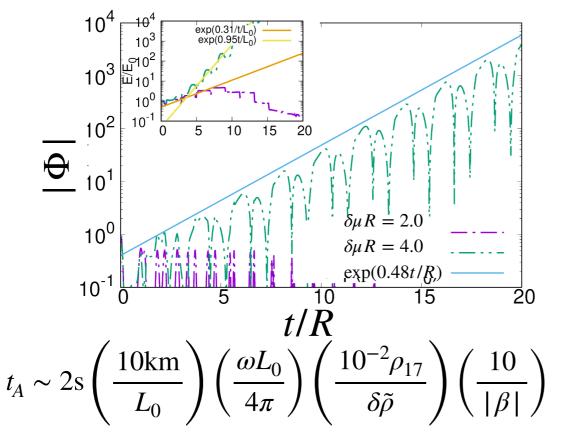
$$\delta R(t) \sim \sin(\omega t)$$

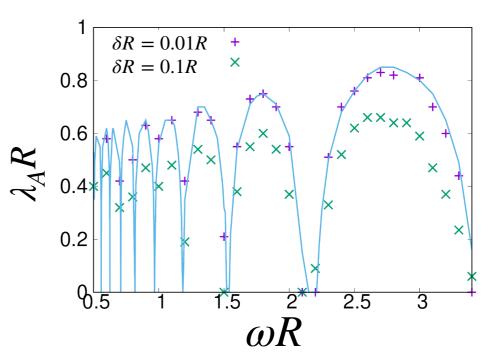
(Wang et al (2011)

$$(R + \delta R < r)$$

$$(r < R + \delta R)$$

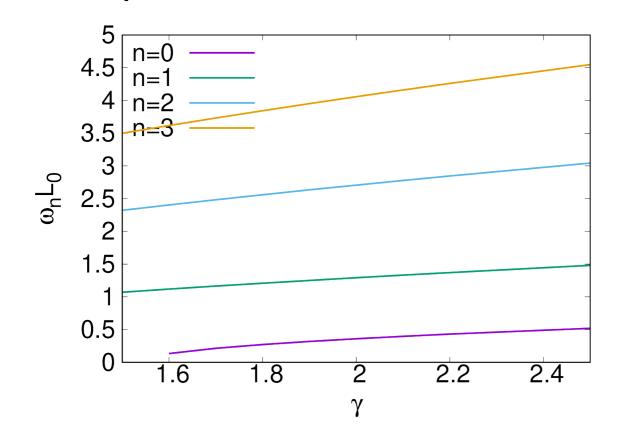


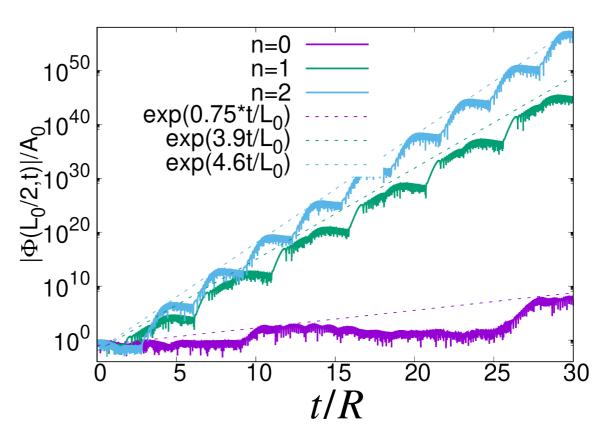




Parametric instability in relativistic stars

- We checked the parametric instability in relativistic constant density star.
 - solve the normal mode of the star
 - solve the scalar field around star with normal mode perturbations.





Summary

$$\mathcal{L} = \frac{R}{16\pi} - \frac{1}{2} g^{\mu\nu} \nabla_{\mu} \Phi \nabla_{\nu} \Phi - V(\Phi) + \mathcal{L}(A(\Phi)^{2}g, \Psi_{m}) \qquad A(\Phi) = e^{\frac{\beta}{2}\Phi^{2}}$$
$$|\beta| < |\beta_{c}^{\text{NS}}|$$

- Two instabilities of the scalar field around oscillating star.
 - blueshift instability

$$t_B \sim 3 \times 10^{-2} \text{s} \left(\frac{10^{-3}}{\lambda_B(\omega) L_0} \right) \left(\frac{L_0}{10 \text{km}} \right) \ln \left| 10 \frac{\mu}{10^{-10} \text{eV}} \frac{\sigma}{10 \text{km}} \right|$$

parametric instability

$$t_A \sim 2s \left(\frac{10 \text{km}}{L_0}\right) \left(\frac{\omega L_0}{4\pi}\right) \left(\frac{10^{-2} \rho_{17}}{\delta \tilde{\rho}}\right) \left(\frac{10}{|\beta|}\right)$$

We can expect the similar effect in other models.

$$\mathcal{L} \supset f(\Phi)\mathcal{G}_{GB}$$

