

Computational Finance with C++

Finite Difference Methods for Partial Differential Equations

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Finite Difference Methods for Partial Differential Equations

- From the Black and Scholes Equation to the Heat Equation
- Explicit finite difference methods
- Stability of the explicit finite difference scheme
- Implicit Schemes
- Boundary Conditions
- Finite Difference vs Monte Carlo

Model Problem: The BSE PDE

The Black and Scholes Equation:

$$\frac{\partial V}{\partial t} + \frac{\sigma^2}{2} S^2 \frac{\partial^2 V}{\partial S^2} + rS \frac{\partial V}{\partial S} - rV = 0$$

$$V(S, T) = \Phi(S, T)$$

where Φ is the pay-off function e.g. $\Phi(S, T) = \max[S - K, 0]$

Risk Neutral Dynamics:

$$dS(t) = rS(t)dt + \sigma S(t)dW(t) \quad S(0) = s$$

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Risk Neutral Dynamics:

$$dS(t) = rS(t)dt + \sigma S(t)dW(t) \quad S(0) = s$$

A good idea when dealing with PDE is to attempt to simplify it

BSE happens to be related to the Heat Equation with a change of variables

From the BSE PDE to the Heat Equation

The Black and Scholes Equation:
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$$V(S, T) = \Phi(S, T)$$

Following the change of variables (derived on the board), we have:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$
$$u(x, 0) = e^{-\alpha x} \Phi(e^x, T)$$

← The Heat Equation

From the BSE PDE to the Heat Equation

The Black and Scholes Equation:
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← The Heat Equation

Why analyse the heat equation?

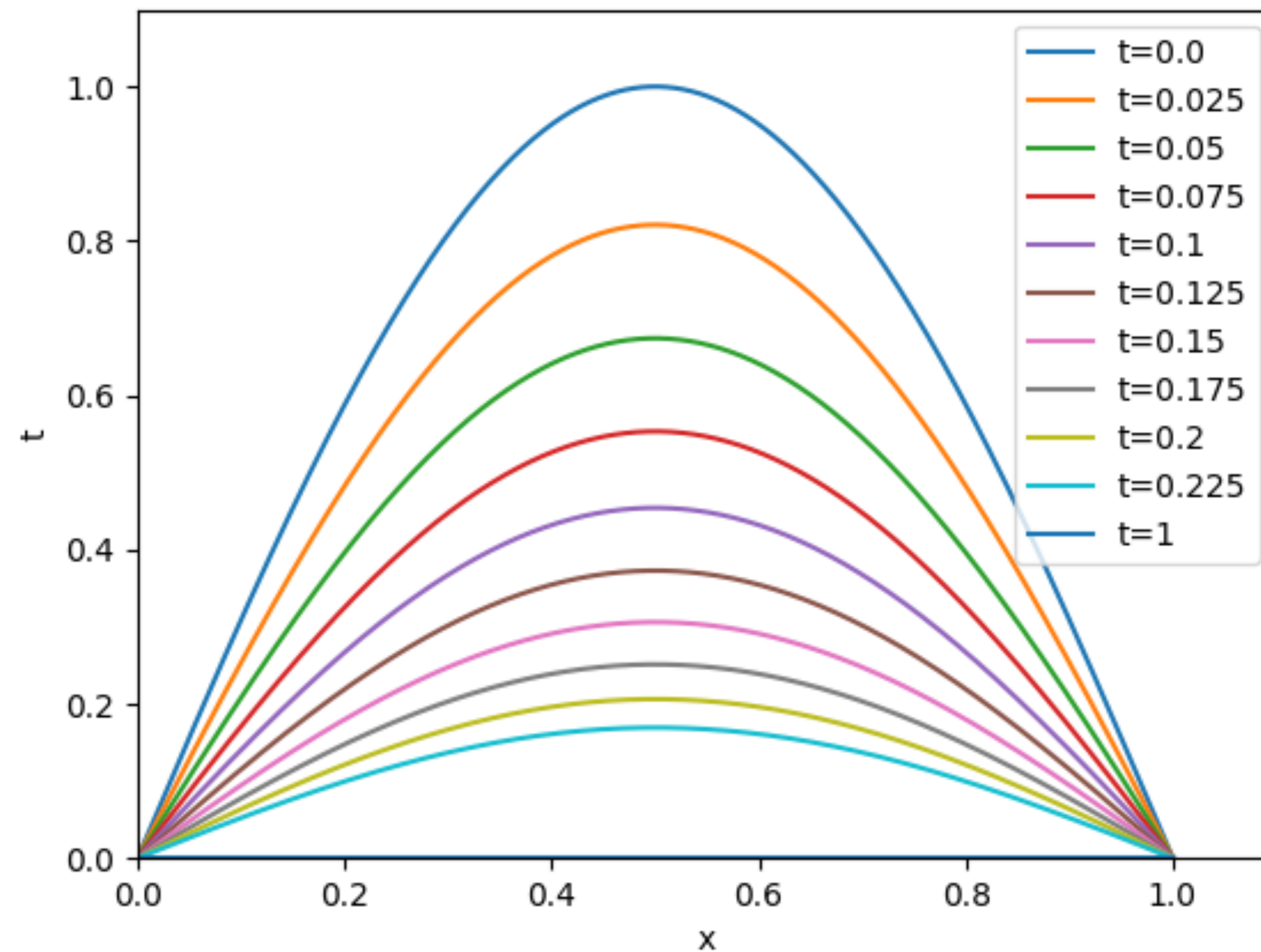
- The finite difference method can (in principle) be applied to the BSE directly
- The **two equations are equivalent**
- But **numerically they are not equivalent**, the heat equation behaves much better!
- **Theoretically it is easier to analyse the heat equation**

Example

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

Boundary conditions: $u(x,0) = \sin(\pi x)$, $u(0,t) = 0$, $u(1,t) = 0$

Exact Solution: $u(x,t) = e^{-\pi^2 t} \sin(\pi x)$



Foundations of Finite Difference Methods

Mean Value Theorem: $f(x + h) = f(x) + f'(x)h + \frac{1}{2}f''(\xi)h^2$, where $\xi \in [x, x + h]$

Rearrange: $f'(x) = \frac{f(x + h) - f(x)}{h} - \frac{1}{2}f''(\xi)h$, where $\xi \in [x, x + h]$

Create equidistant grid: $\dots < x_{i-1} < x_i < \dots$

Constant mesh size: h , use notation $f_i = f(x_i)$, and write

$$f'(x_i) = \frac{f_{i+1} - f_i}{h} + O(h)$$

Basic idea: Use two meshes sizes one for **space** and one for **time**
 Δx Δt

Foundations of Finite Difference Methods

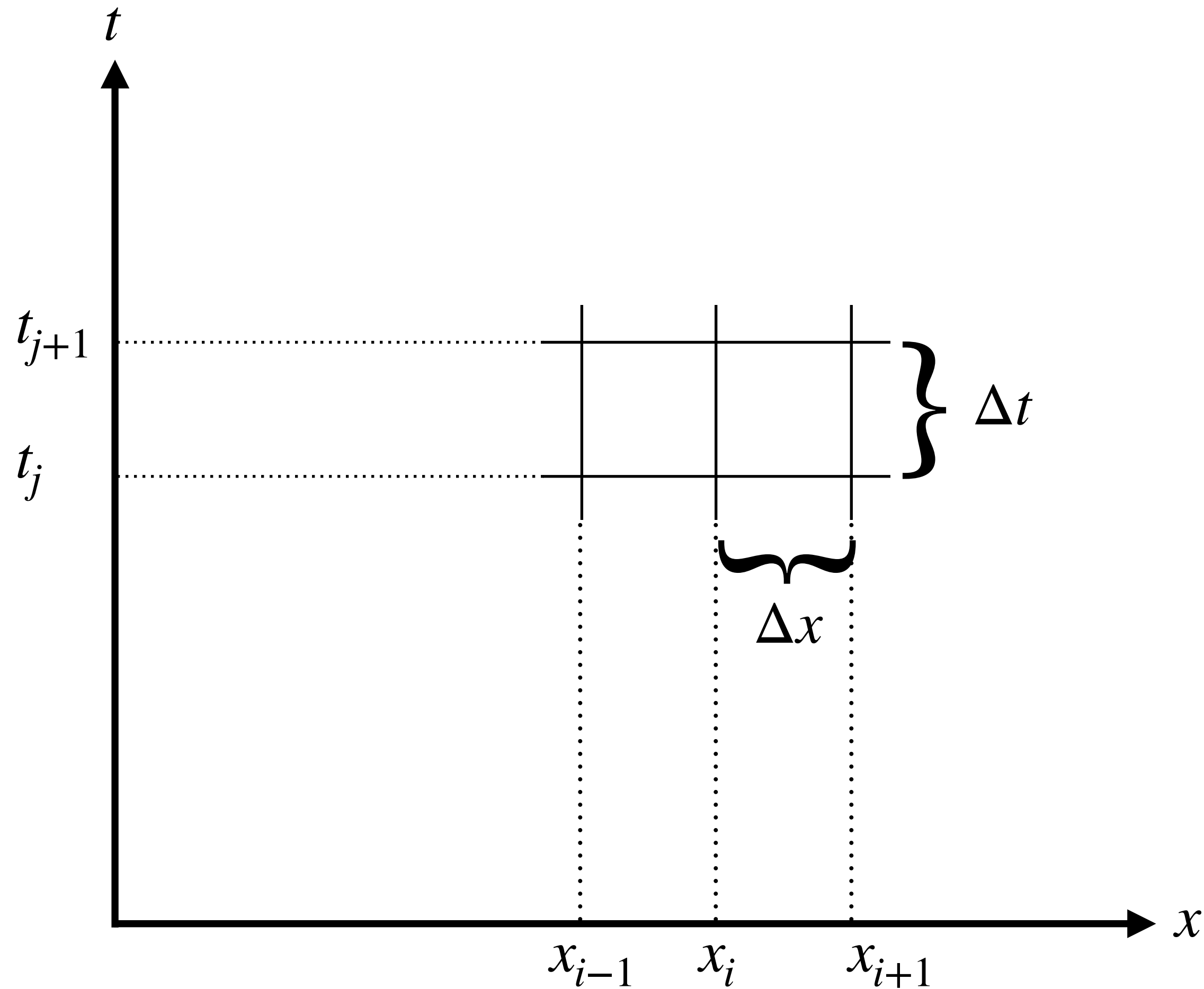
Forward Difference: $f'(x_i) = \frac{f_{i+1} - f_{i-1}}{2h} + O(h)$

Central Difference: $f'(x_i) = \frac{f_{i+1} - f_{i-1}}{2h} + O(h)$

Central Difference:
(second derivative) $f''(x_i) = \frac{f_{i+1} - 2f_i + f_{i-1}}{h^2} + O(h)$

Similar formulas can be used for the time derivative (but lead to very different schemes)

The Finite Difference Grid



Boundary Conditions:

Time Domain: $0 \leq t \leq \frac{\sigma^2 T}{2}$

Space Domain:

Select $0 \leq S_{\min} < S_{\max}$ and set
 $x_{\min} = \ln S_{\min}$, $x_{\max} = \ln S_{\max}$

Explicit Finite Difference Scheme

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad \longrightarrow \quad \frac{\partial u}{\partial t} \approx \frac{u_{i,j+1} - u_{i,j}}{\Delta t} \quad \text{(Forward Difference)}$$

$$\frac{\partial^2 u}{\partial x^2} \approx \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} \quad \text{(Central Difference)}$$

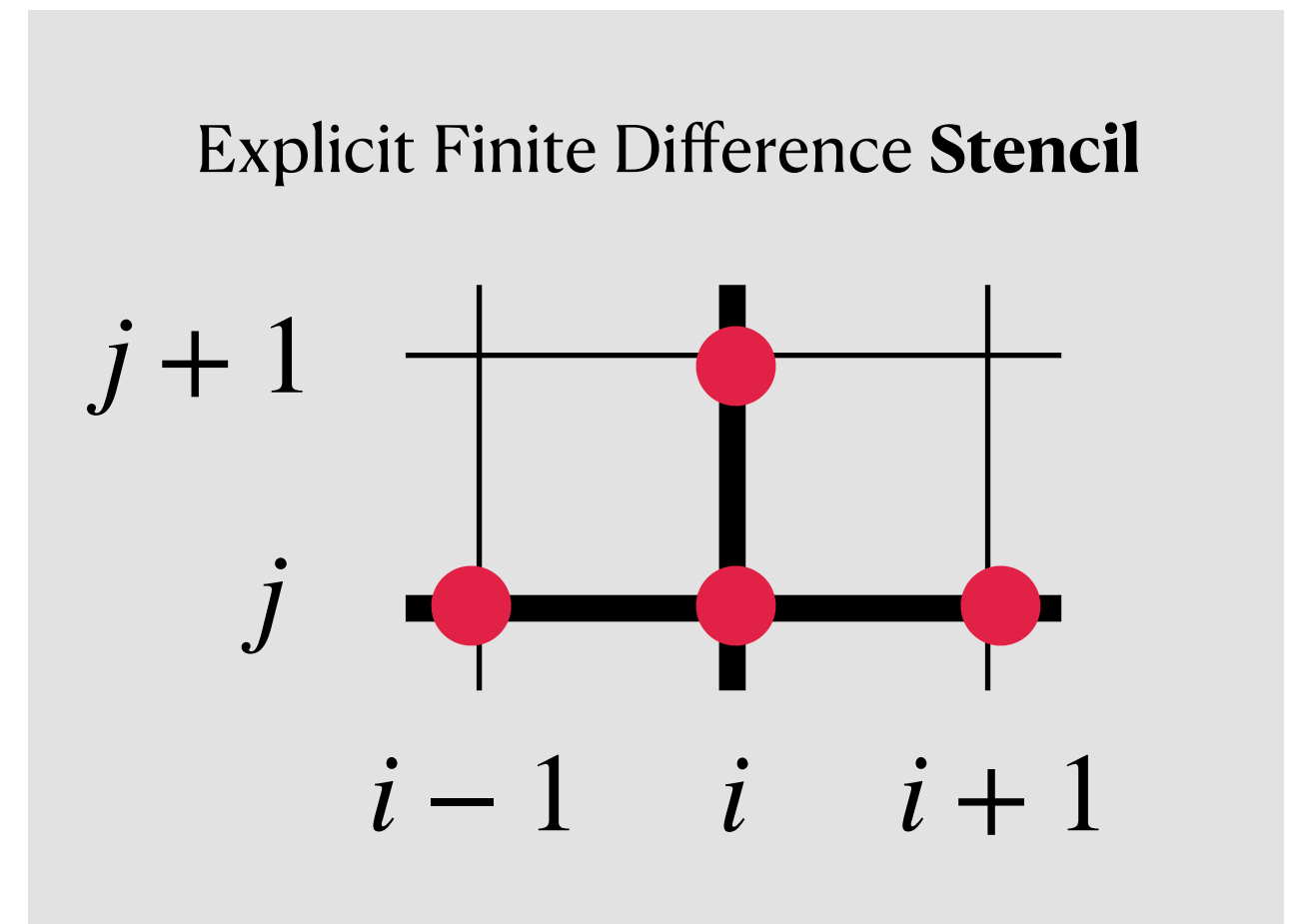


$$\frac{u_{i,j+1} - u_{i,j}}{\Delta t} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2}$$

$$u_{i,j+1} = u_{i,j} + \underbrace{\frac{\Delta t}{\Delta x^2}}_{\lambda} (u_{i+1,j} - 2u_{i,j} + u_{i-1,j})$$



$$u_{i,j+1} = \lambda u_{i-1,j} + (1 - 2\lambda)u_{i,j} + \lambda u_{i+1,j}$$



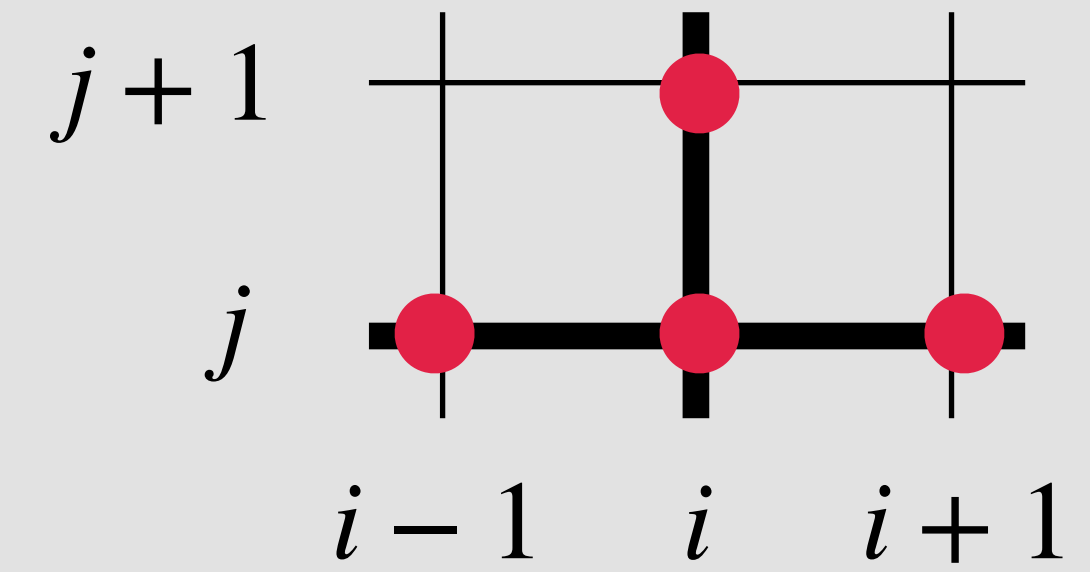
Explicit Finite Difference Scheme

$$u_{i,j+1} = \lambda u_{i-1,j} + (1 - 2\lambda)u_{i,j} + \lambda u_{i+1,j}$$

Vector Form:

$$u^j = \begin{bmatrix} u_{1,j} \\ u_{2,j} \\ \vdots \\ u_{m-1,j} \end{bmatrix}, u^0 \text{ is given, } A = \begin{pmatrix} 1 - 2\lambda & \lambda & 0 & \dots & 0 \\ \lambda & 1 - 2\lambda & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \lambda \\ 0 & \dots & 0 & \lambda & 1 - 2\lambda \end{pmatrix},$$

Explicit Finite Difference **Stencil**



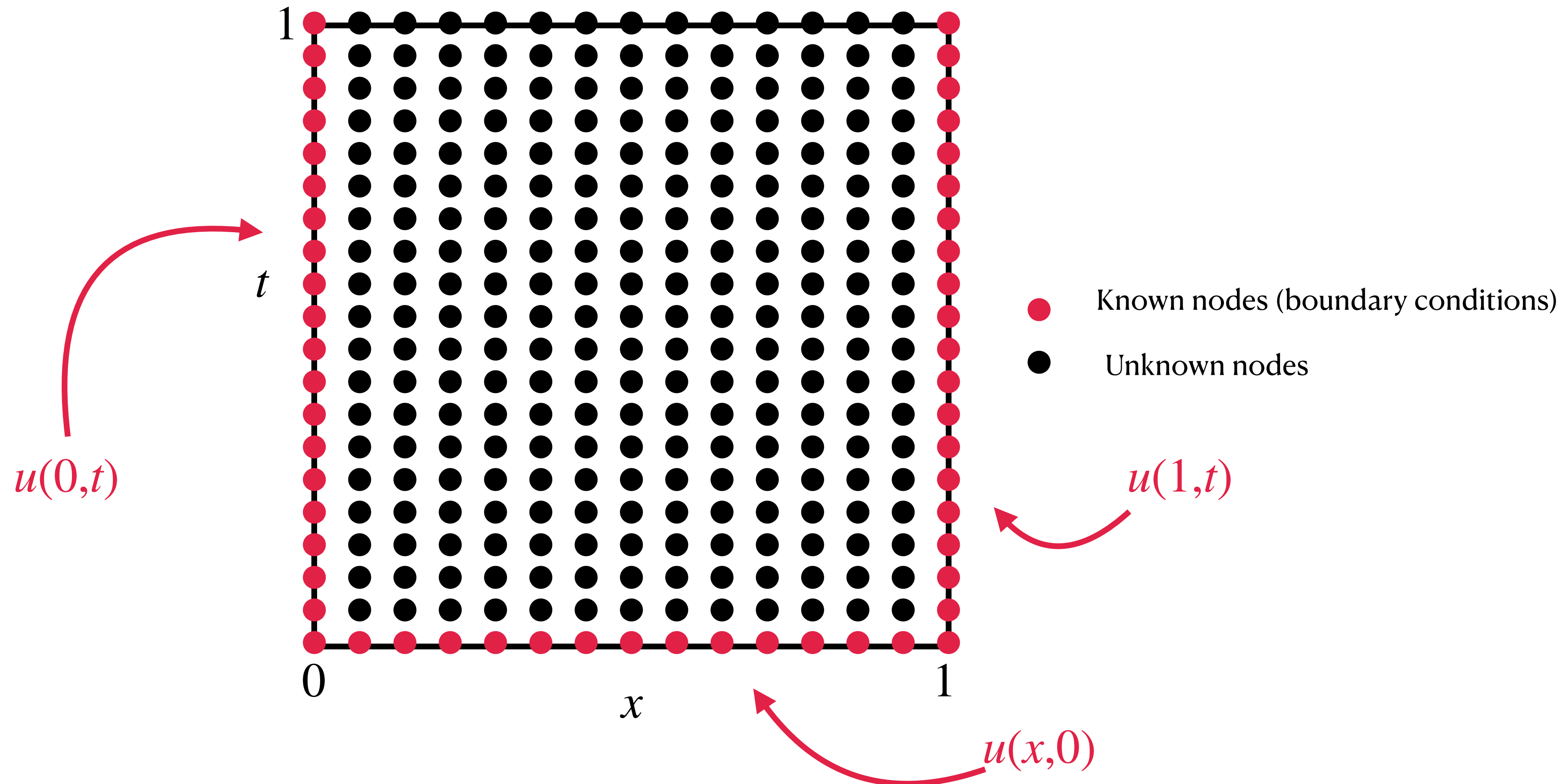
$$u^{j+1} = Au^j$$

Example

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

Boundary conditions: $u(x,0) = \sin(\pi x)$, $u(0,t) = 0$, $u(1,t) = 0$

Exact Solution: $u(x,t) = e^{-\pi^2 t} \sin(\pi x)$

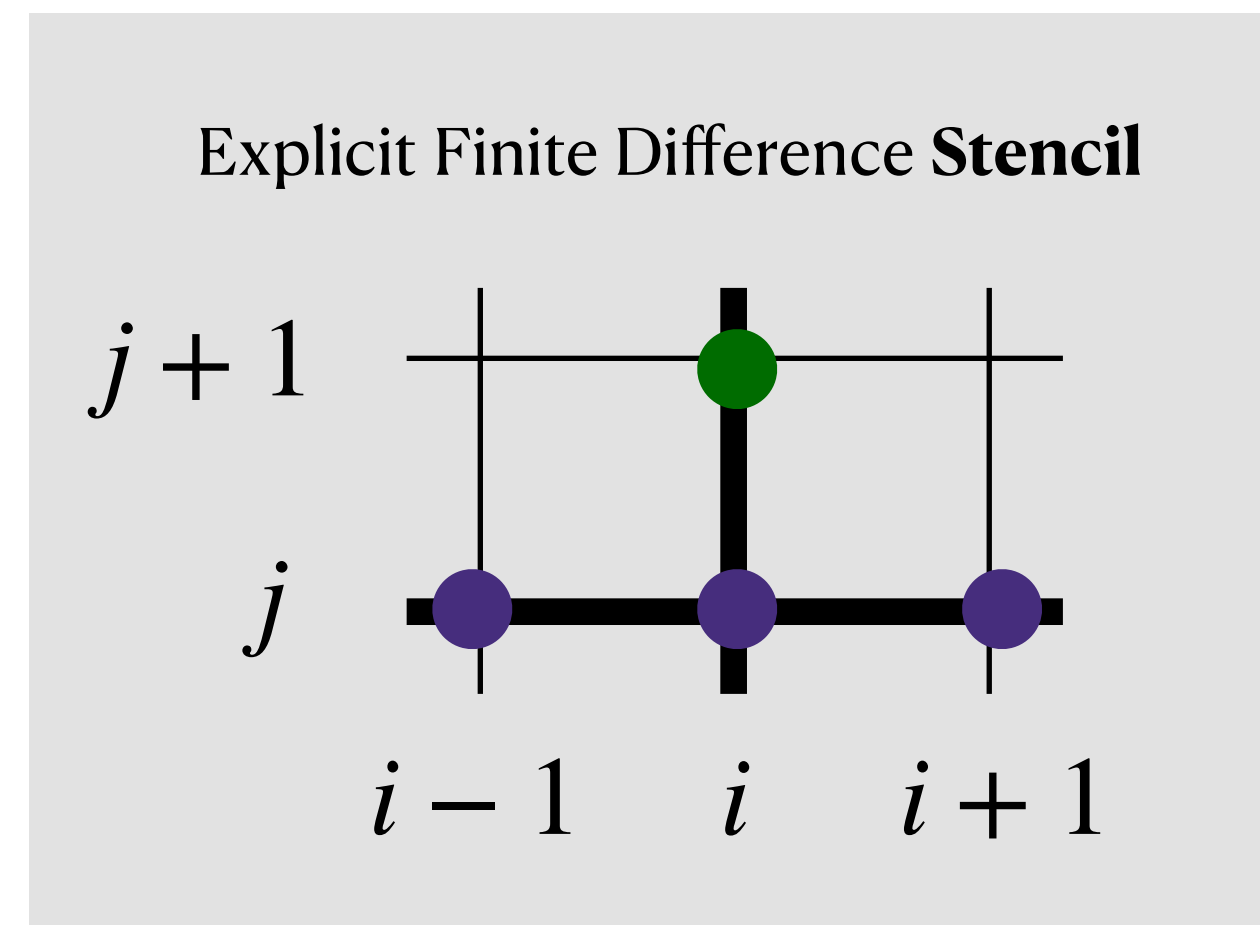
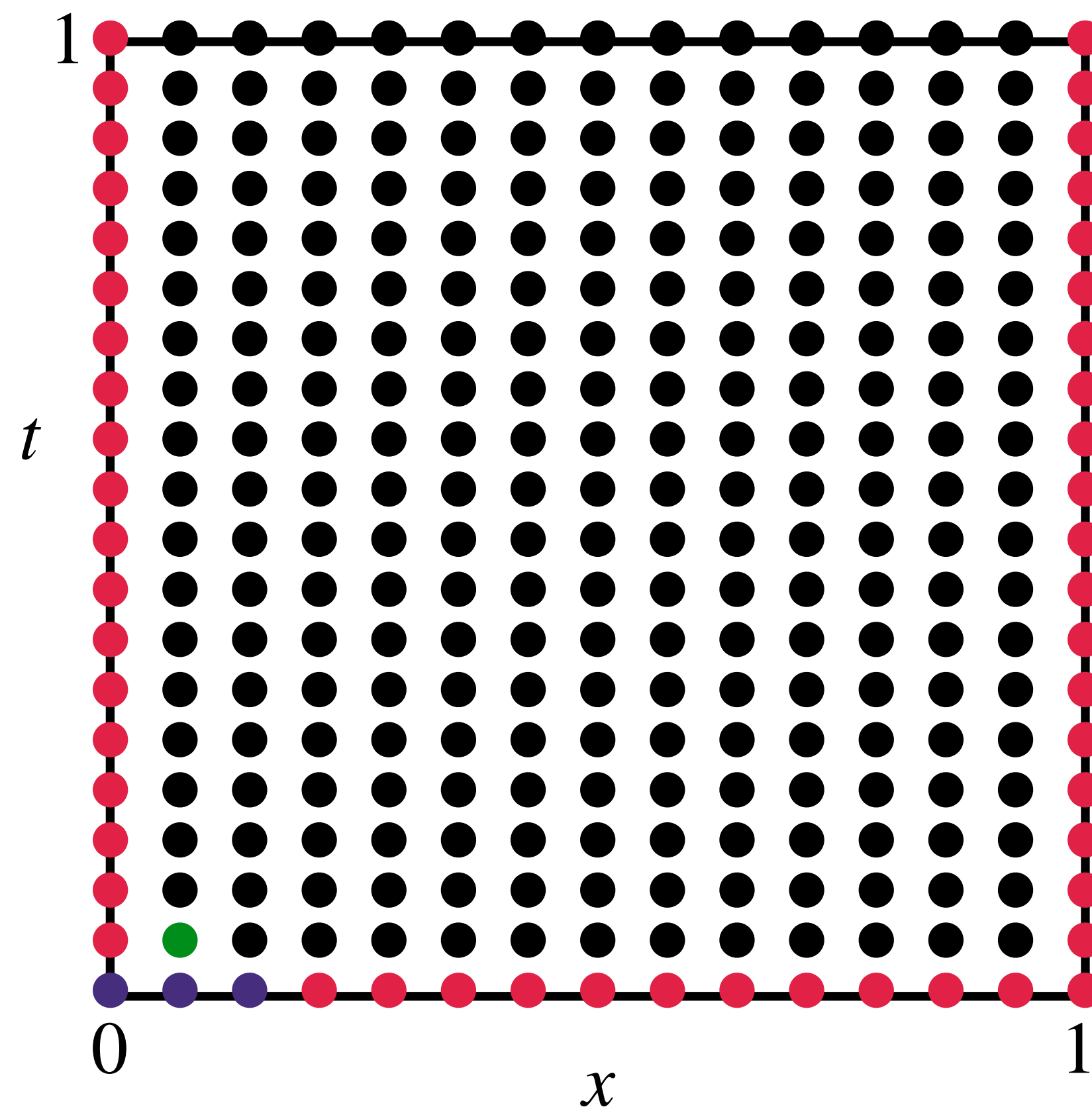


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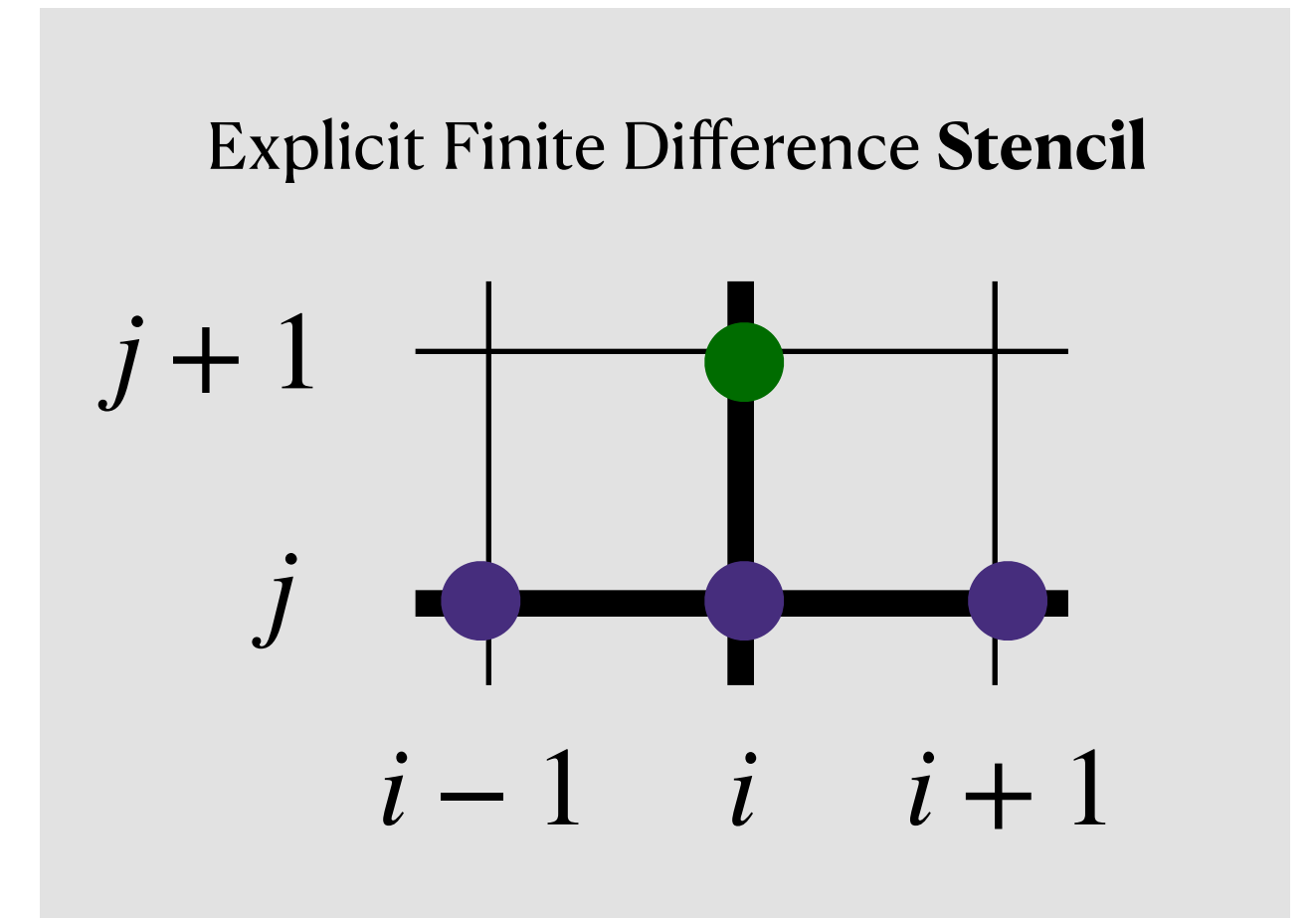
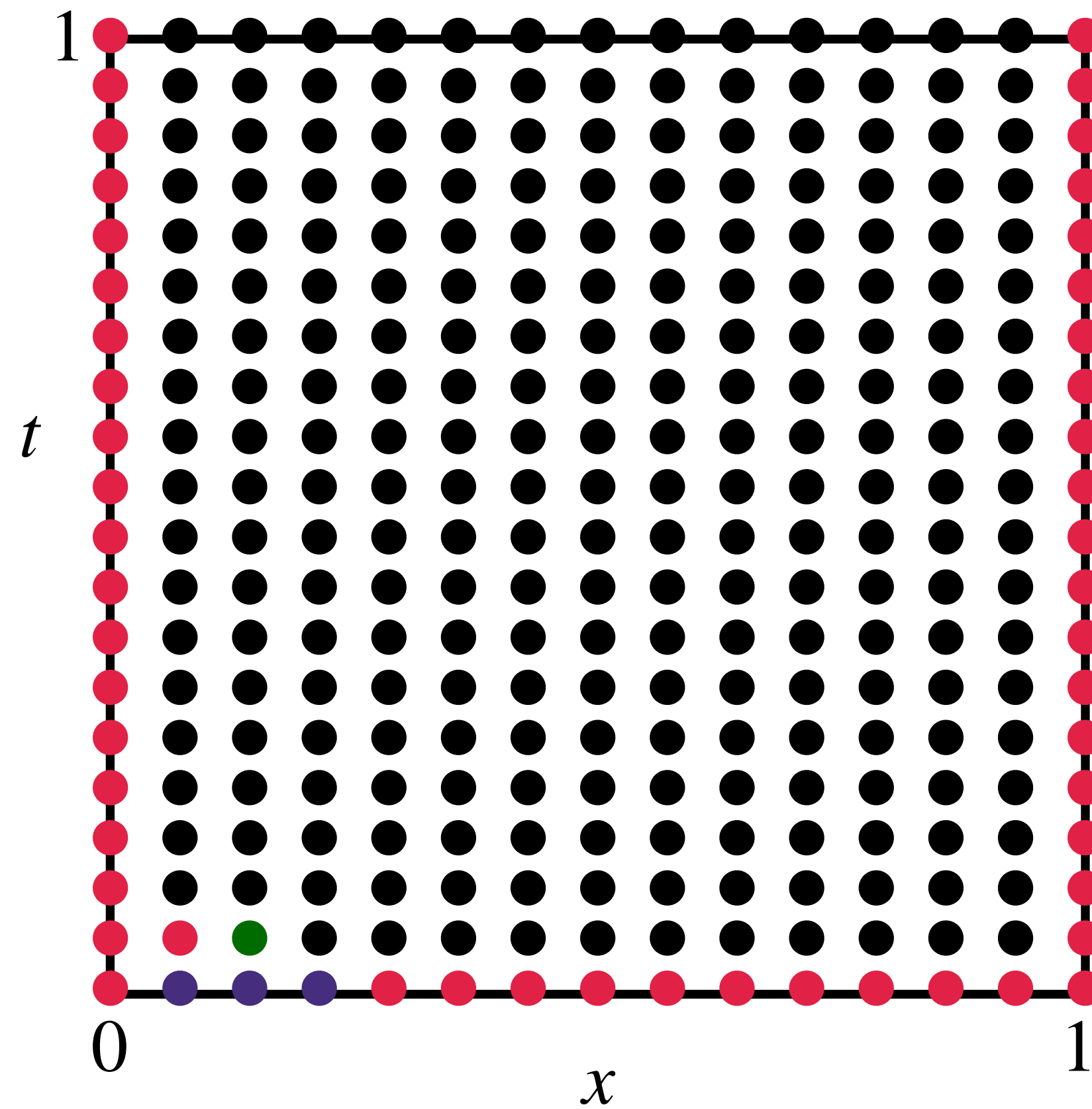
- Known nodes (boundary conditions)
- Unknown nodes
- $u(1,1) = \lambda u_{0,0} + (1 - 2\lambda)u_{1,0} + \lambda u_{2,0}$

Example

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

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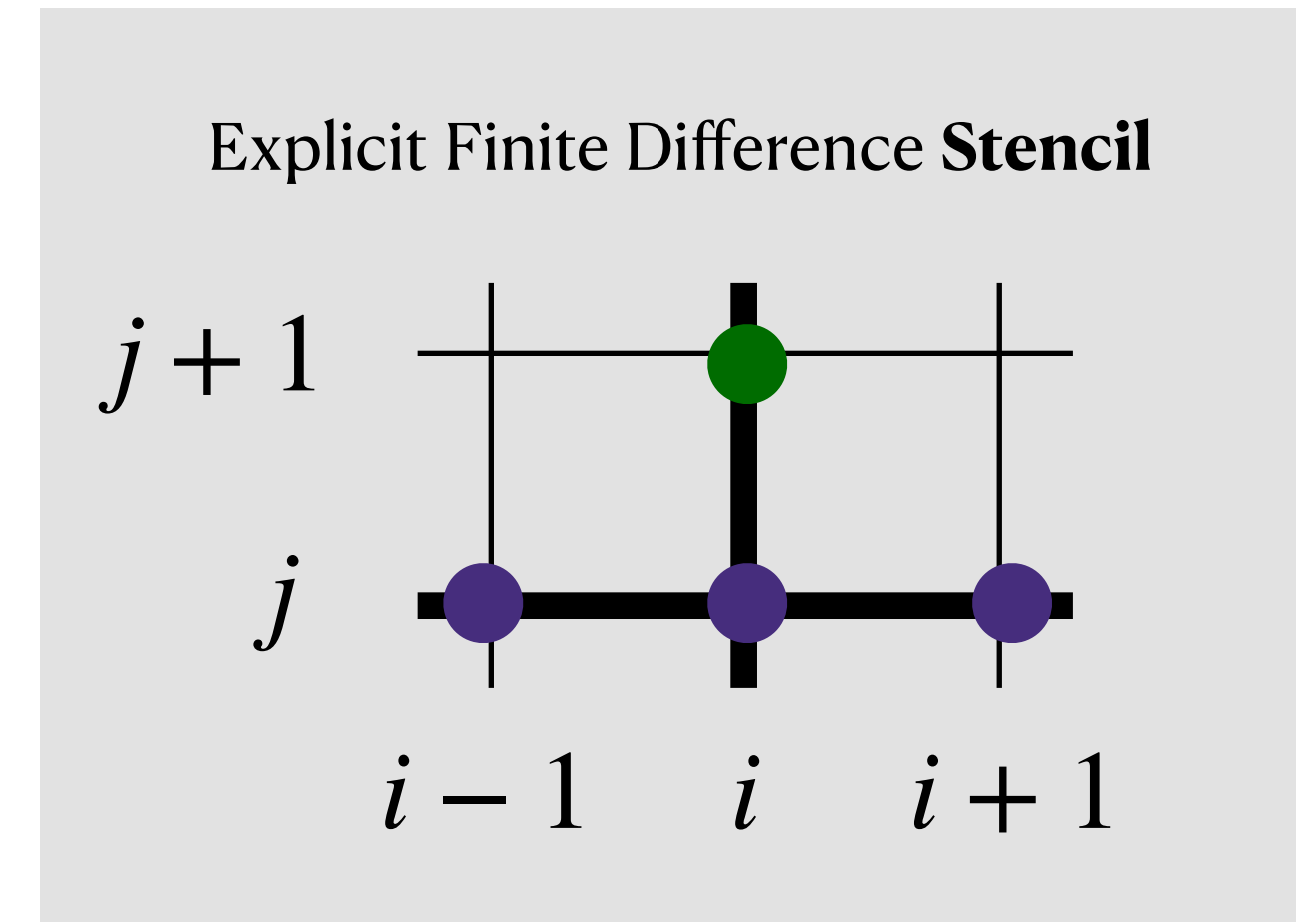
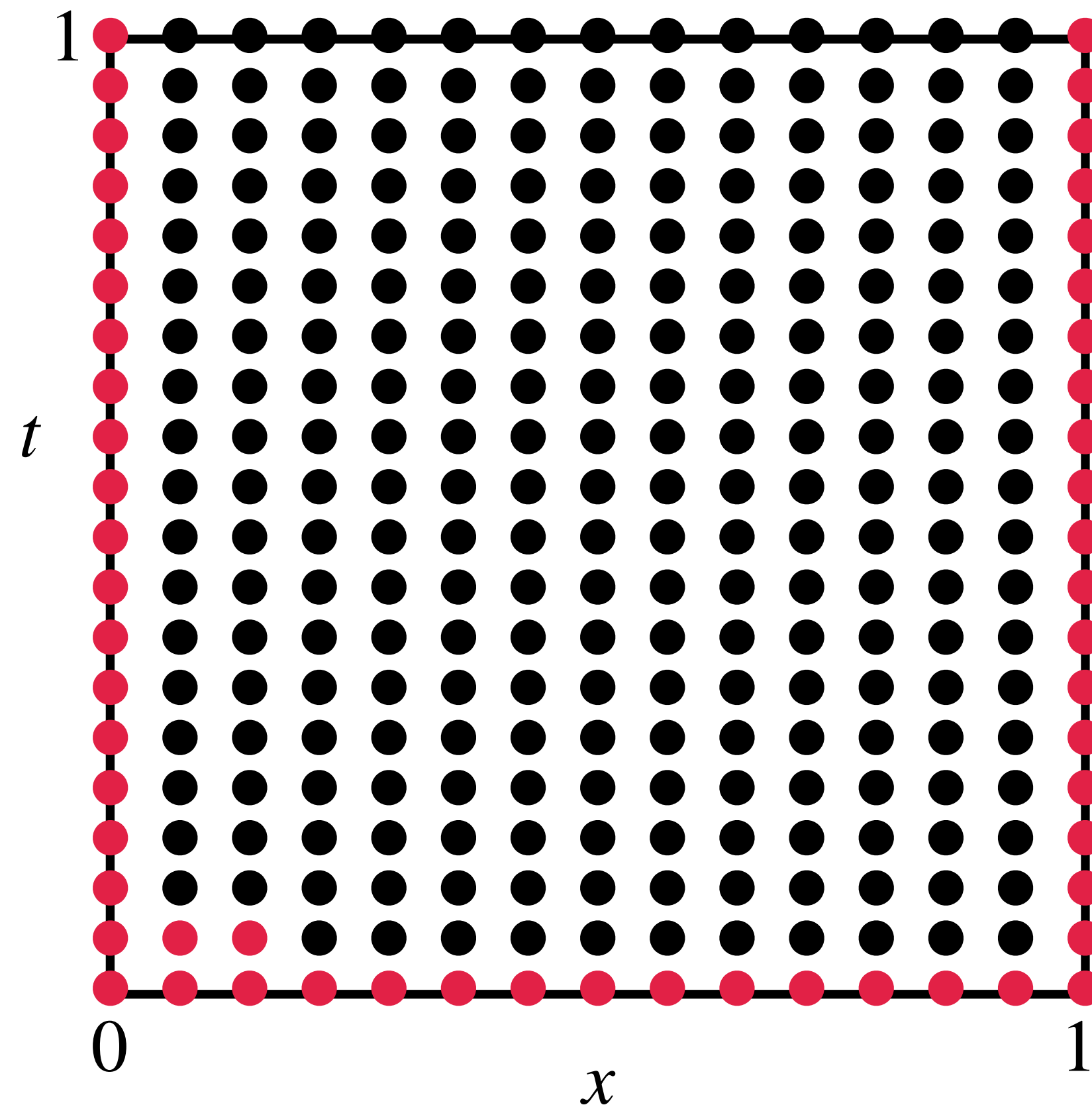
- Known nodes (boundary conditions)
- Unknown nodes
- $u(2,1) = \lambda u_{1,0} + (1 - 2\lambda)u_{2,0} + \lambda u_{3,0}$

Example

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

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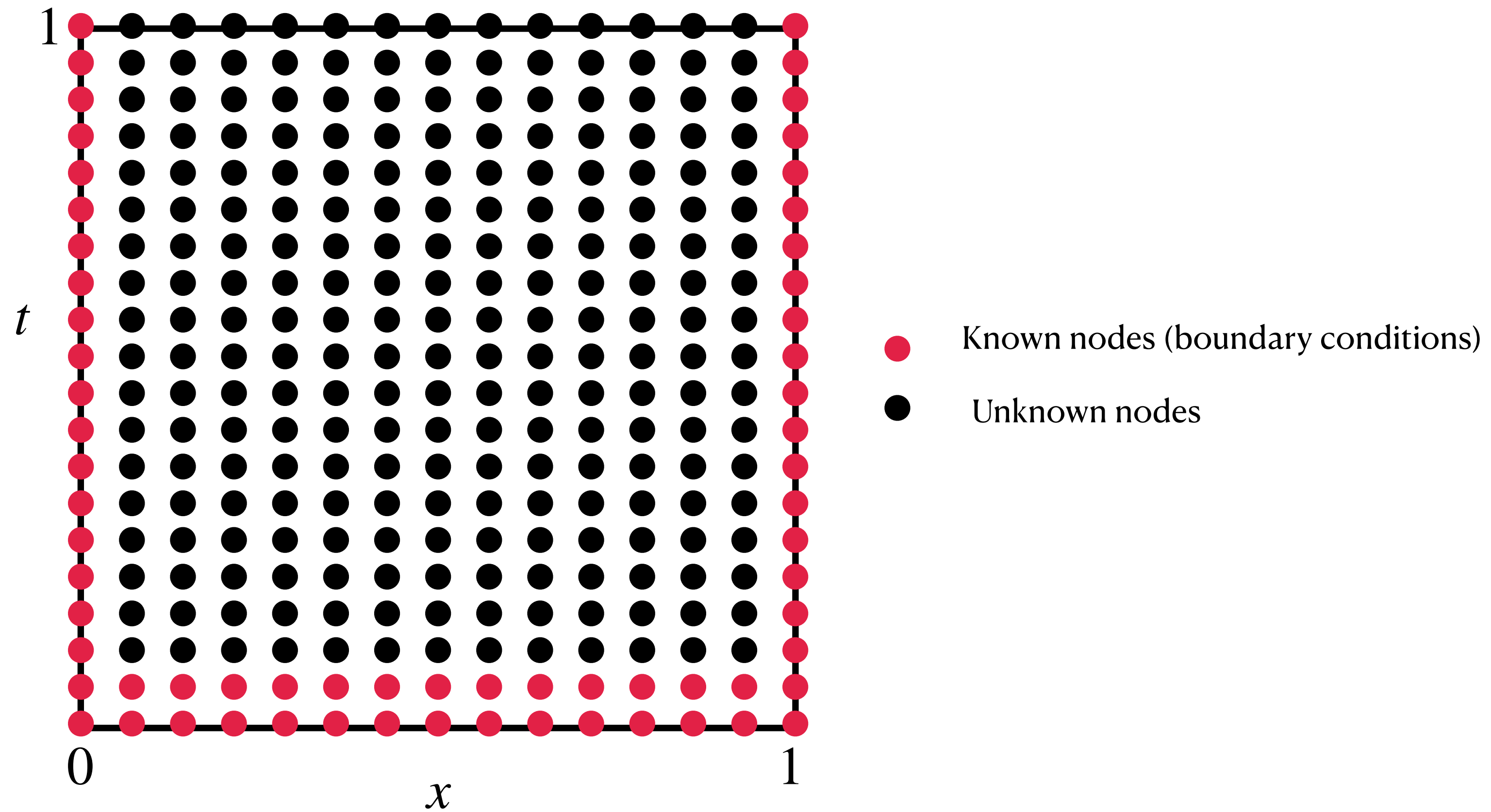
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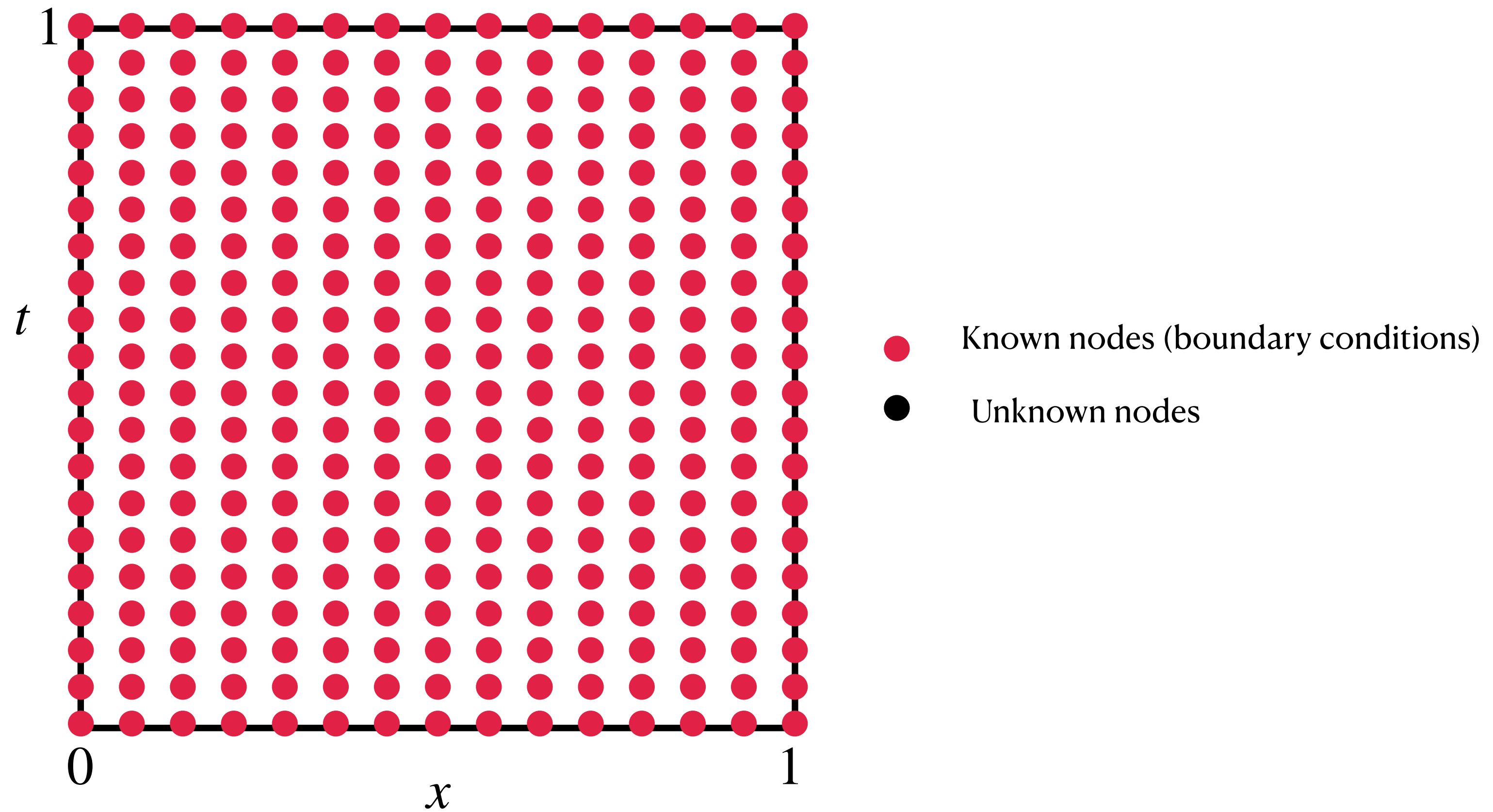


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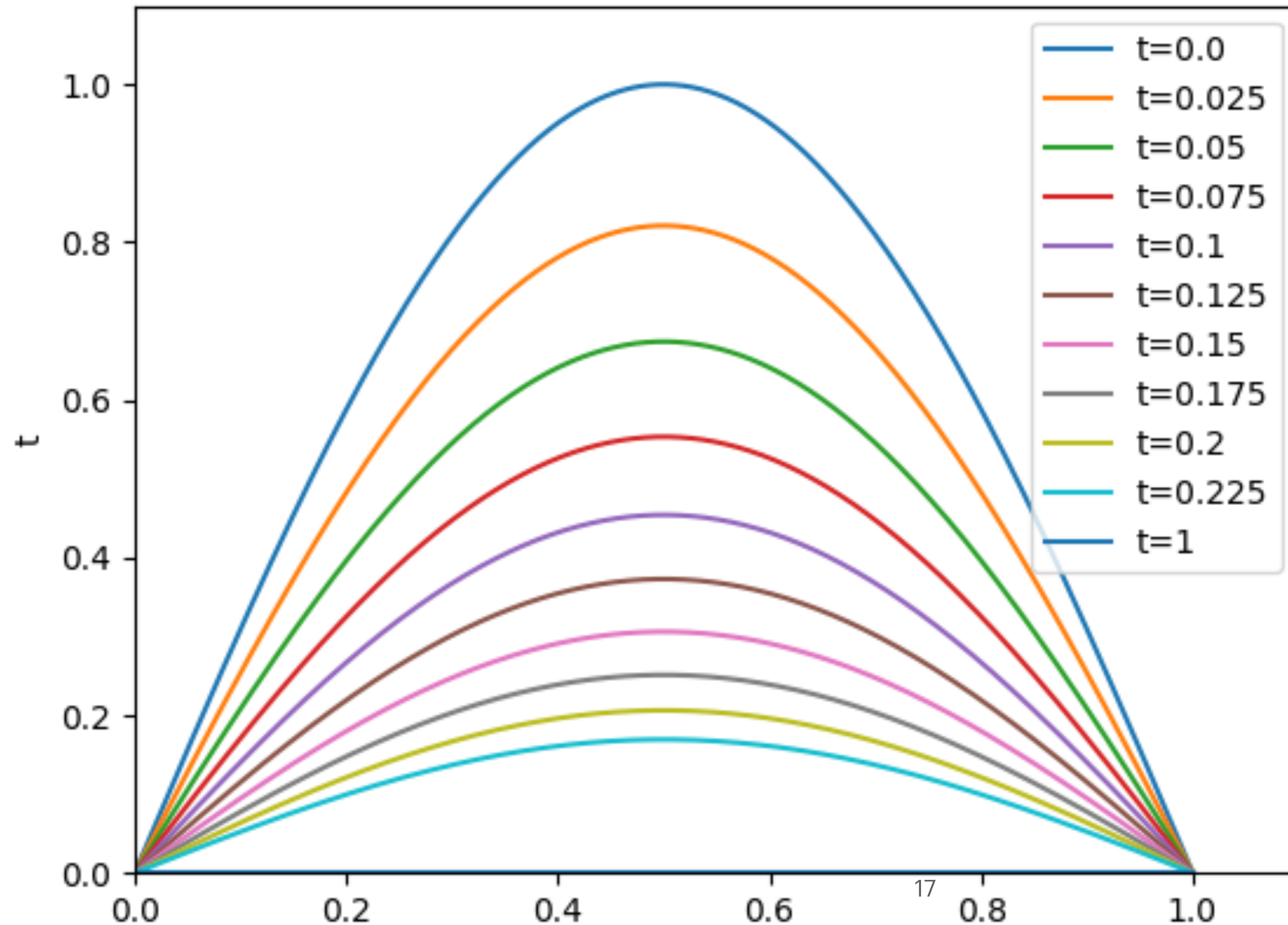
Exact Solution: $u(x,t) = e^{-\pi^2 t} \sin(\pi x)$



$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

Does the method converge?

Boundary conditions: $u(x,0) = \sin(\pi x)$, $u(0,t) = 0$, $u(1,t) = 0$, Solution: $u(x,t) = e^{-\pi^2 t} \sin(\pi x)$



Discrete Scheme:

$$u_{i,j+1} = \lambda u_{i-1,j} + (1 - 2\lambda)u_{i,j} + \lambda u_{i+1,j}$$

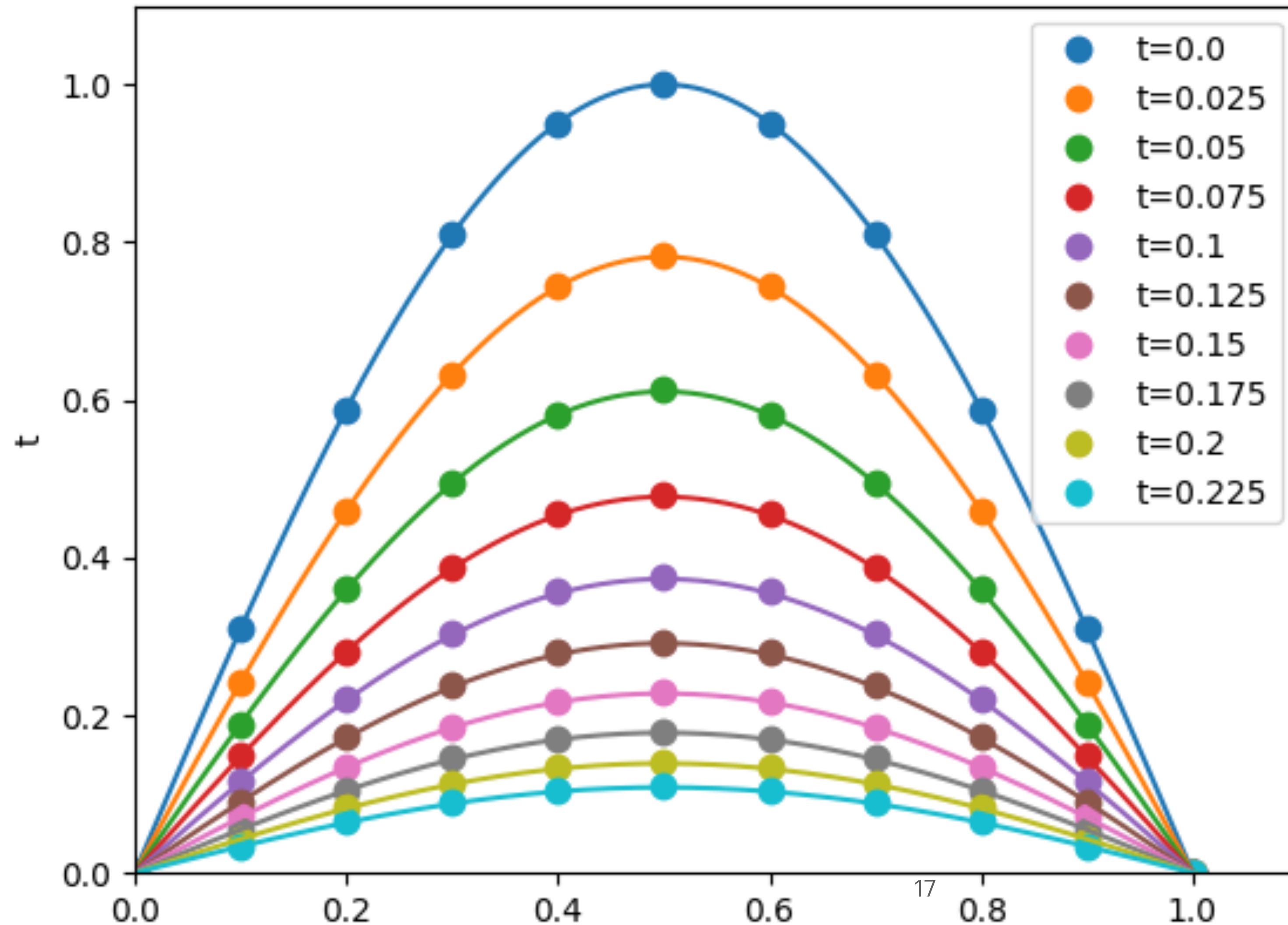
$$\lambda = \frac{\Delta t}{\Delta x^2}$$

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

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$\delta t = 0.0005$, $\delta x = 0.1$, $\lambda = 0.05$



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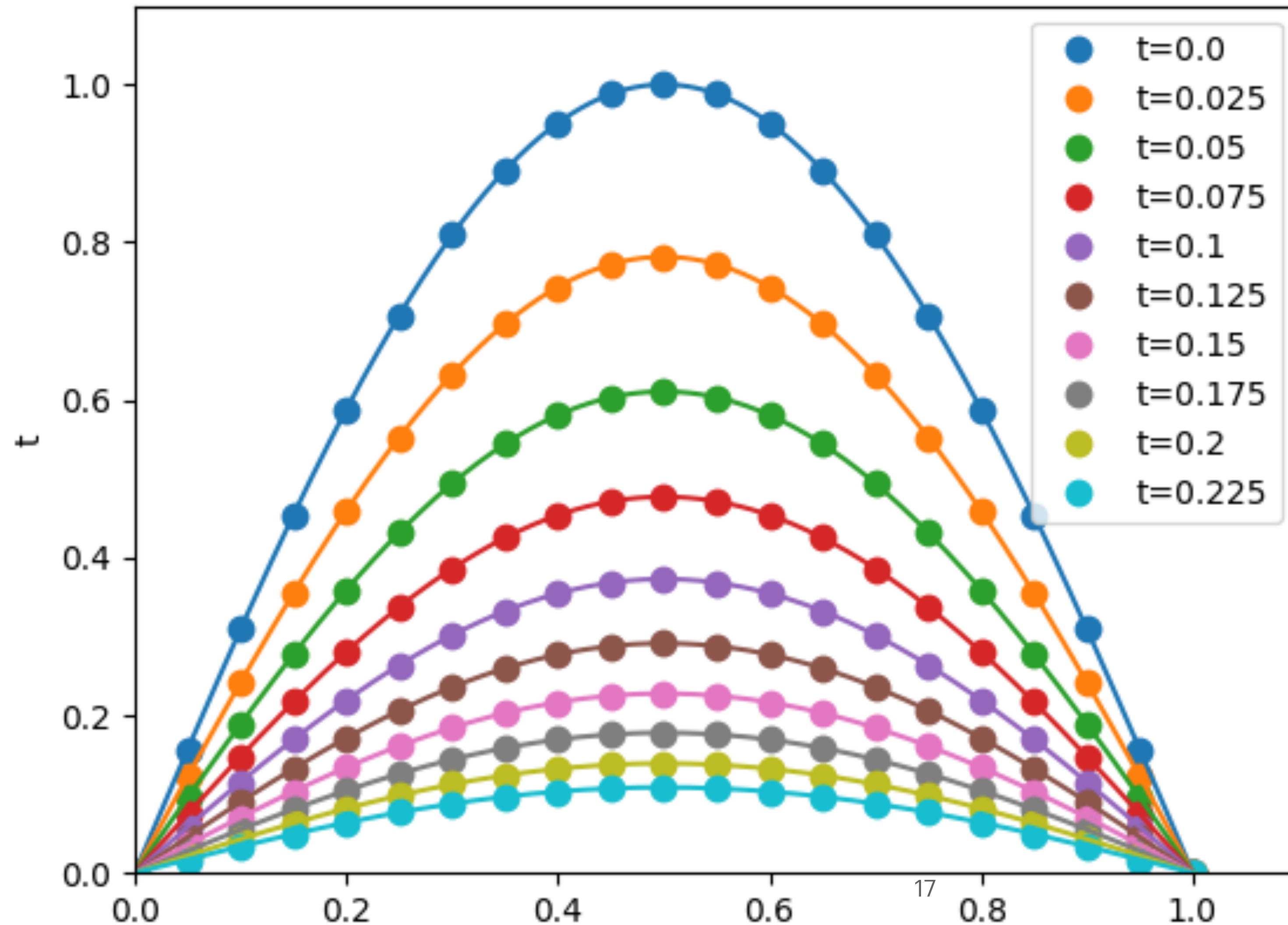
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$\delta t = 0.0005$, $\delta x = 0.05$, $\lambda = 0.2$



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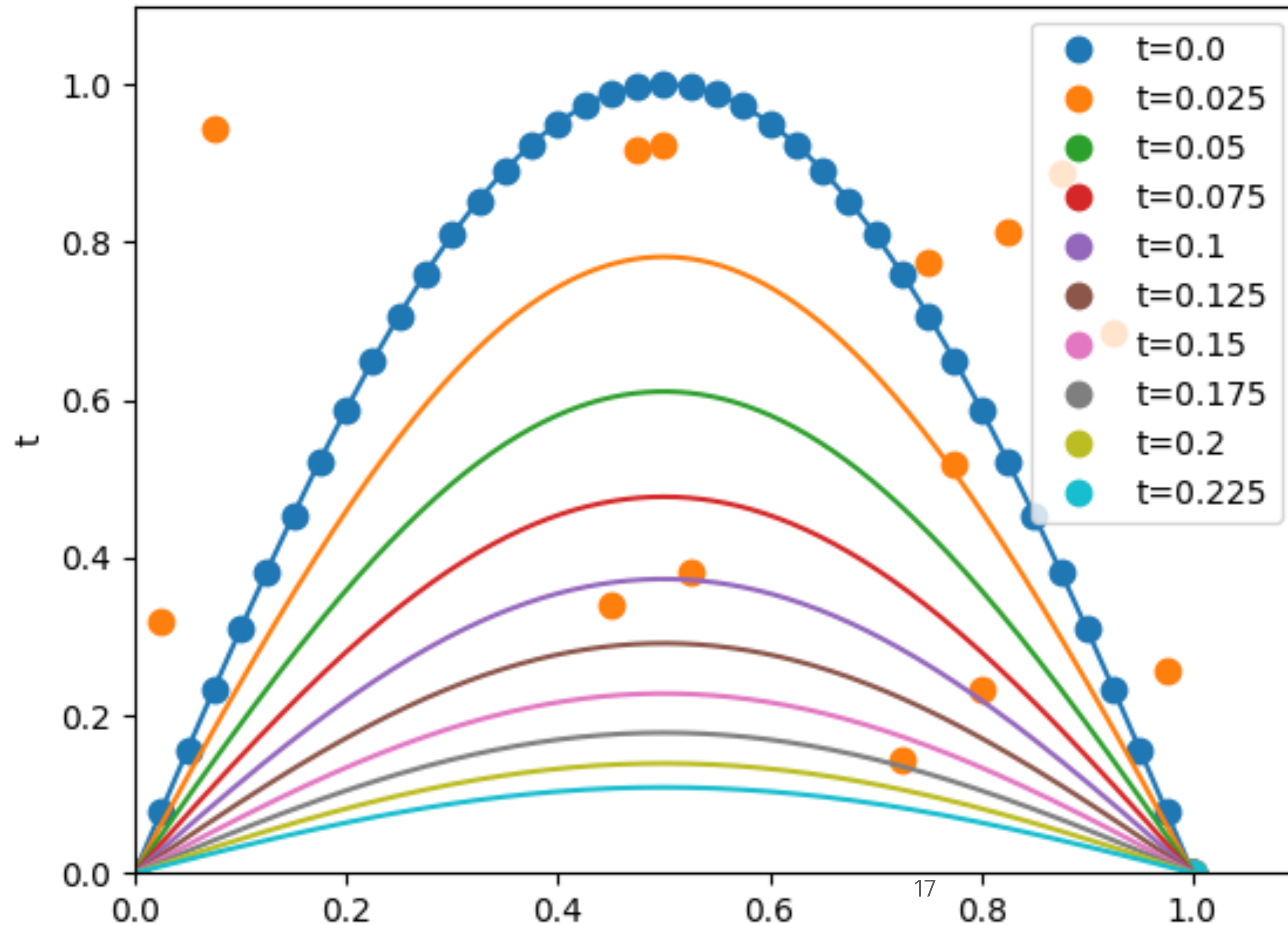
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$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

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$\delta t = 0.0005$, $\delta x = 0.025$, $\lambda = 0.8$



Discrete Scheme:

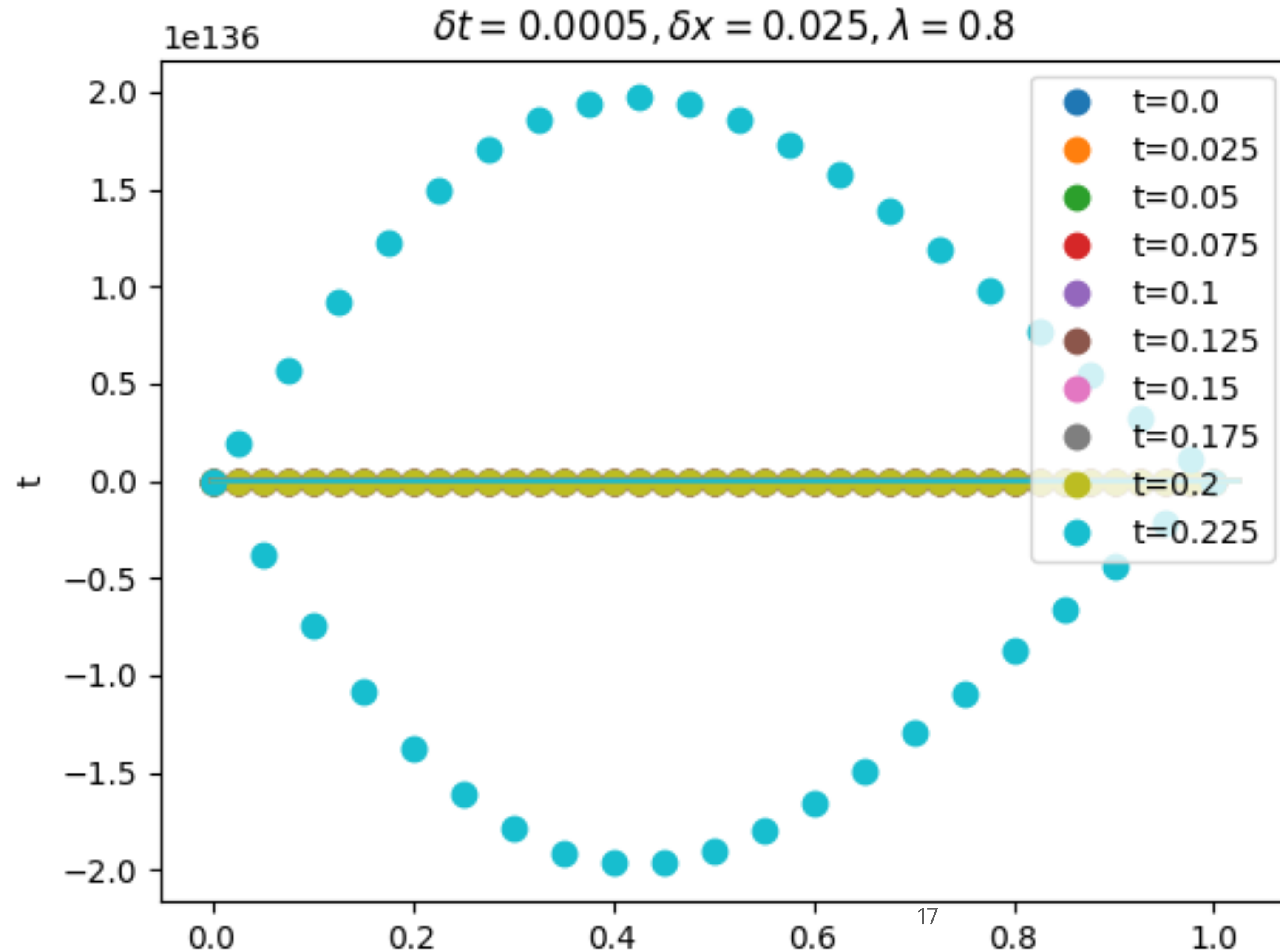
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$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

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What happened?

In numerical analysis of PDEs stability and convergence are equivalent (Lax equivalence theorem)

Stability of PDEs is beyond the scope of this course (ask for references)

A stable method must not increase errors!

Error Propagation in Explicit Finite Difference

A stable method must not increase errors! i.e.

Vector form: $u^{j+1} = Au^j$

$$u^j = \begin{bmatrix} u_{1,j} \\ u_{2,j} \\ \vdots \\ u_{m-1,j} \end{bmatrix}, u_s(j) = \begin{bmatrix} u(\Delta x, j\Delta t) \\ u(2\Delta x, j\Delta t) \\ \vdots \\ u((m-1)\Delta x, j\Delta t) \end{bmatrix}, A = \begin{pmatrix} 1-2\lambda & \lambda & 0 & \dots & 0 \\ \lambda & 1-2\lambda & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \lambda \\ 0 & \dots & 0 & \lambda & 1-2\lambda \end{pmatrix},$$

Where u^j is finite difference approximation and $u_s(j)$ is solution with infinite precision

Error: $e^j = u^j - u(j\Delta t)$ $Ae^j = A(u^j - u(j\Delta t)) = e^{j+1}$ i.e. $e^j = A^j e^0$

A stable method must not increase errors, means:

$$\lim_{j \rightarrow \infty} A^j e^0 = 0 \iff \rho(A) < 1 \iff \lim_{j \rightarrow \infty} A^j[k, l] = 0 \iff \lambda \leq \frac{1}{2}$$

Courant–
Friedrichs–Lewy
condition
(CFL)

Backward Finite Difference Scheme

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \longrightarrow \frac{\partial u}{\partial t} \approx \frac{u_{i,j} - u_{i,j-1}}{\Delta t} \quad \text{(Backward Difference)}$$

$$\frac{\partial^2 u}{\partial x^2} \approx \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} \quad \text{(Central Difference)}$$

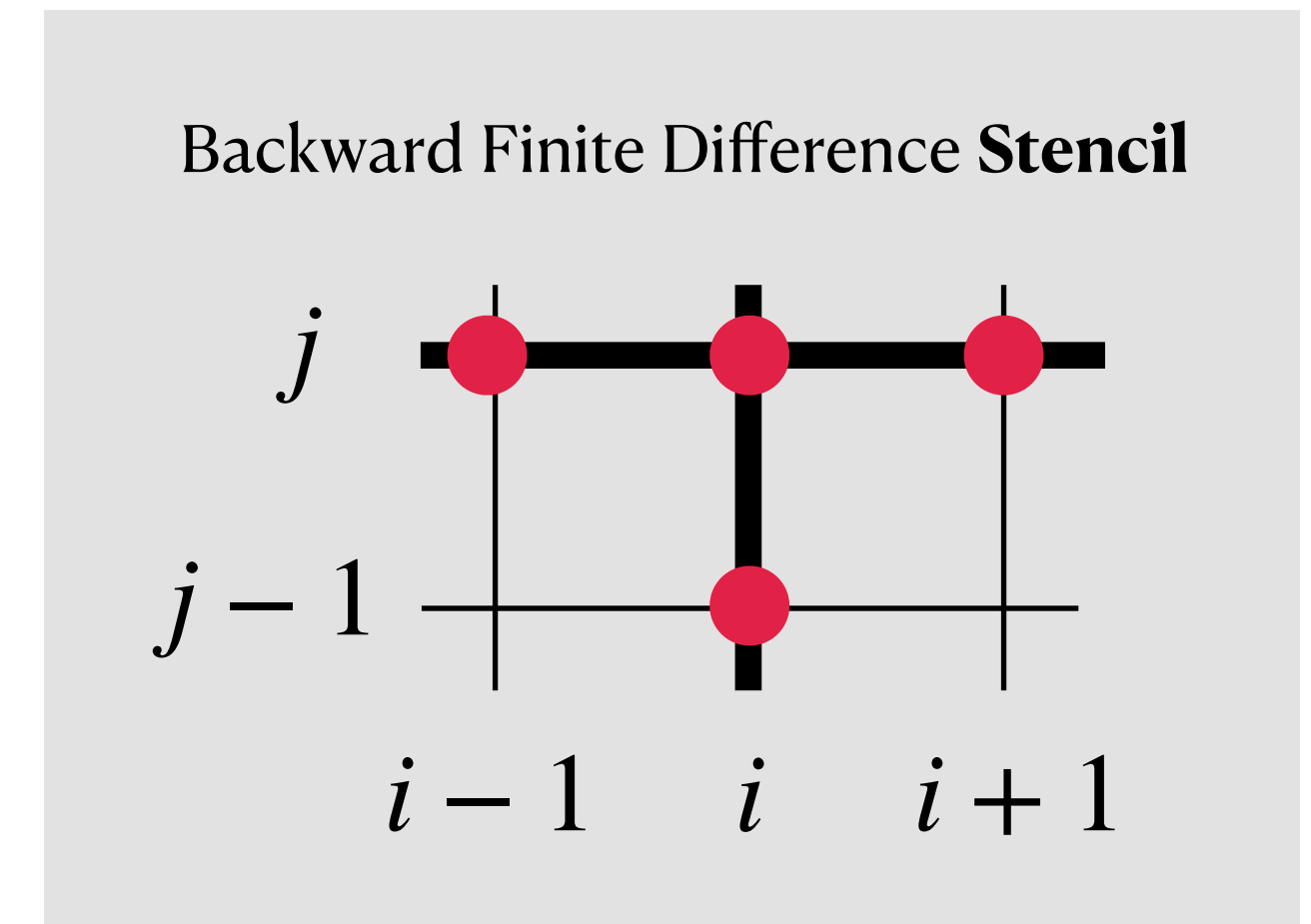


$$\frac{u_{i,j} - u_{i,j-1}}{\Delta t} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2}$$



$$u_{i,j-1} = -\lambda u_{i+1,j} + (1 + 2\lambda)u_{i,j} - \lambda u_{i-1,j}$$

$$\lambda = \frac{\Delta t}{\Delta x^2}$$

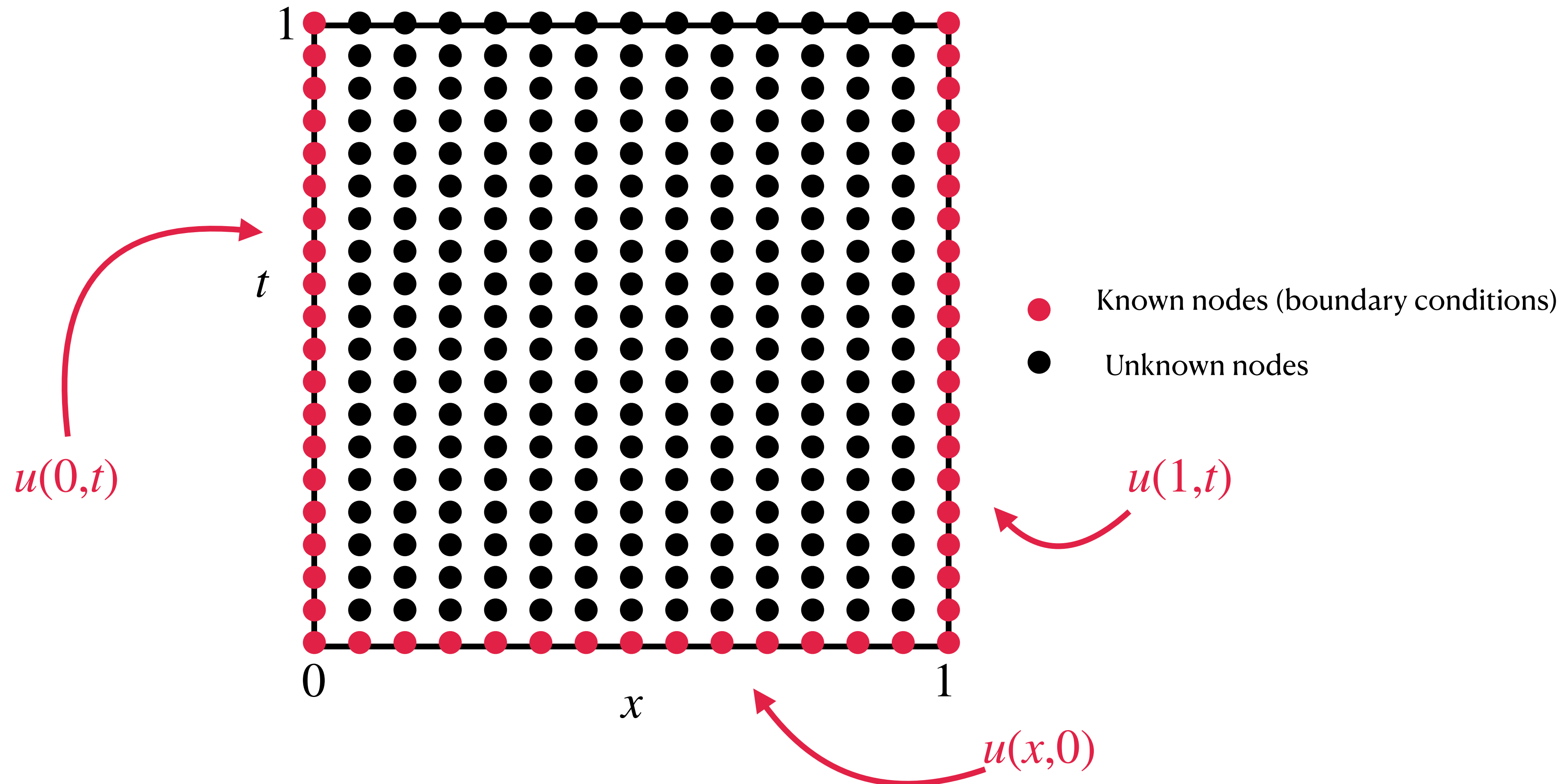


Example

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Exact Solution: $u(x,t) = e^{-\pi^2 t} \sin(\pi x)$

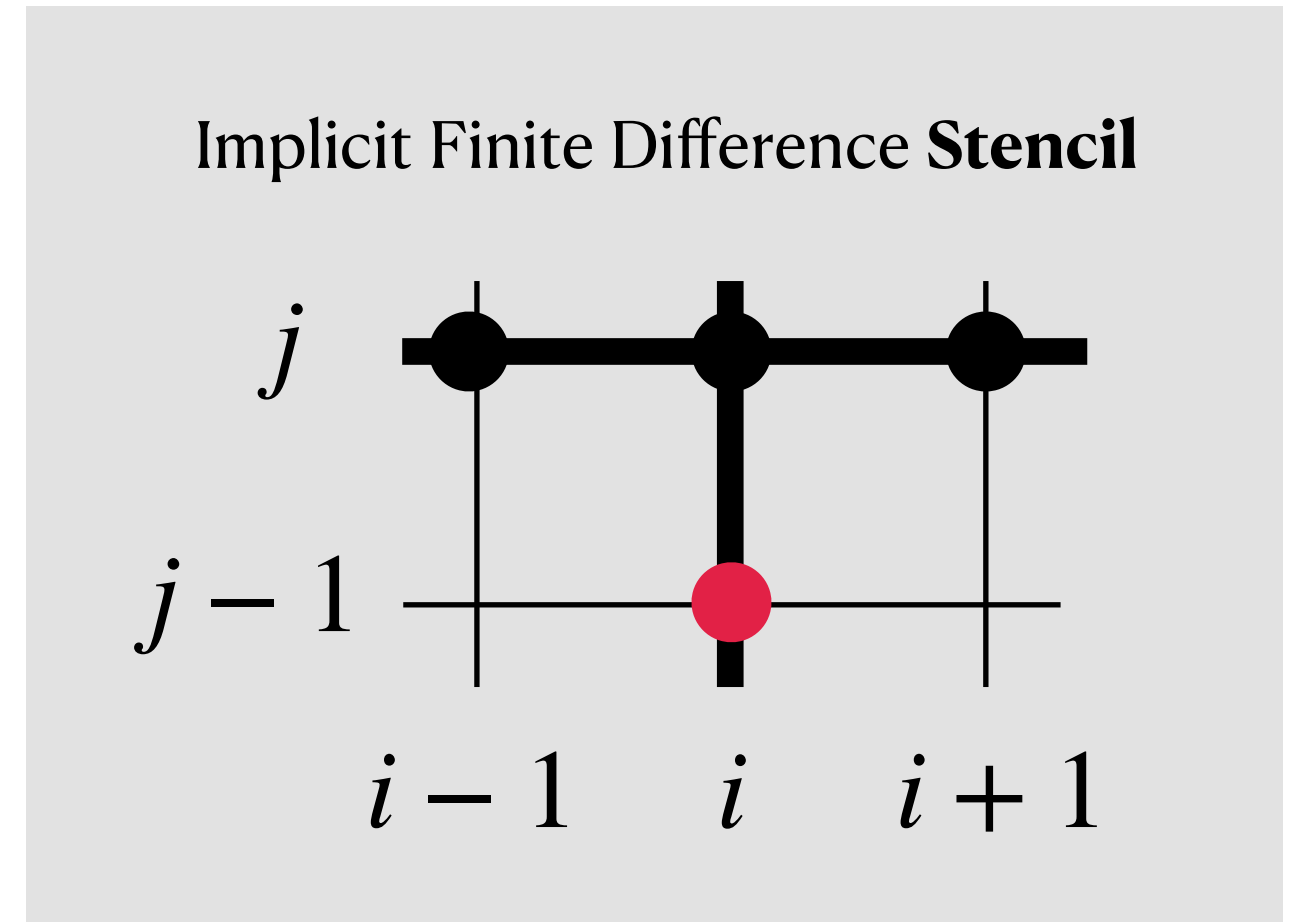
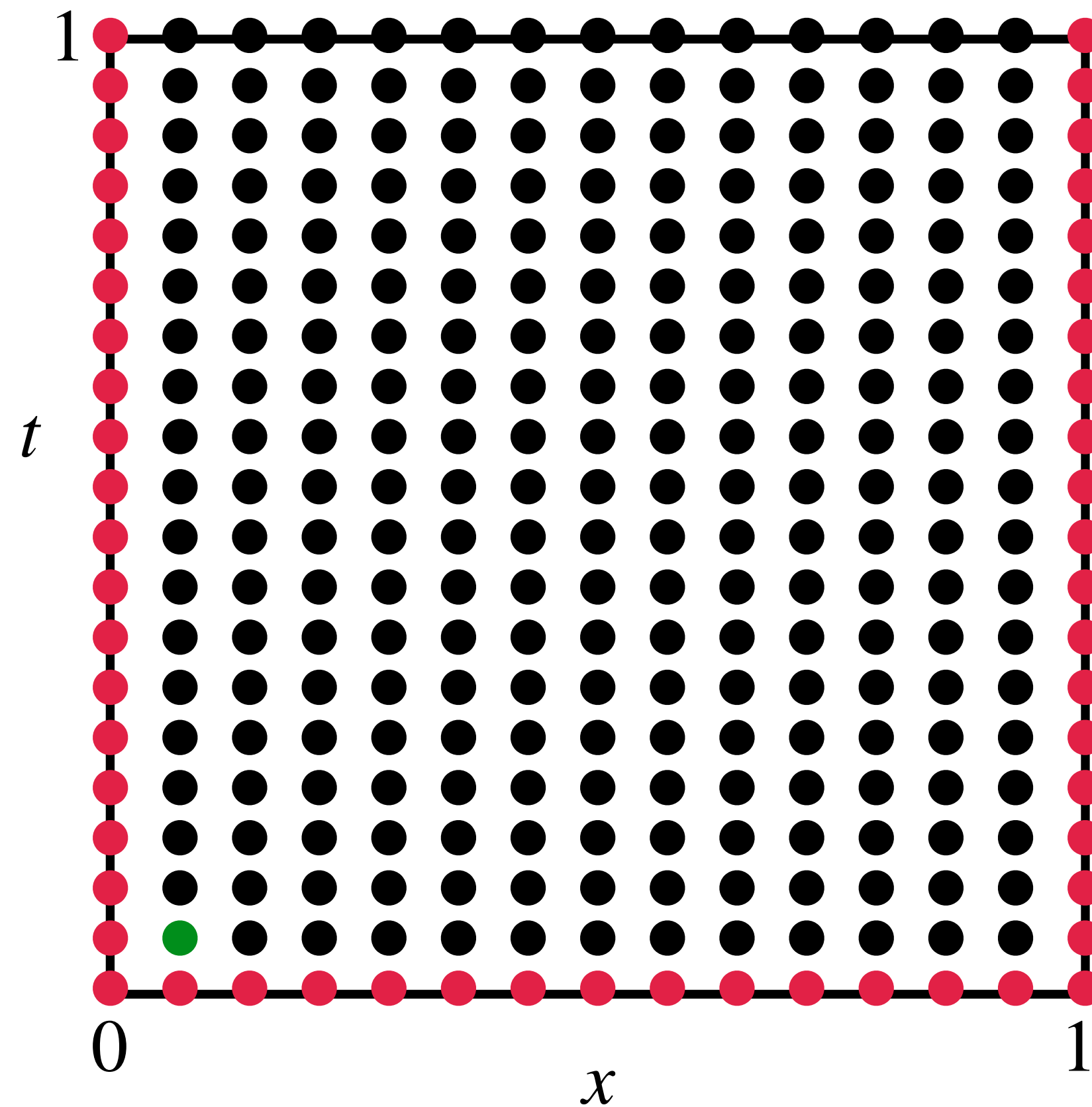


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● Known nodes (boundary conditions)

● Unknown nodes

● $u(1,1) = ??$

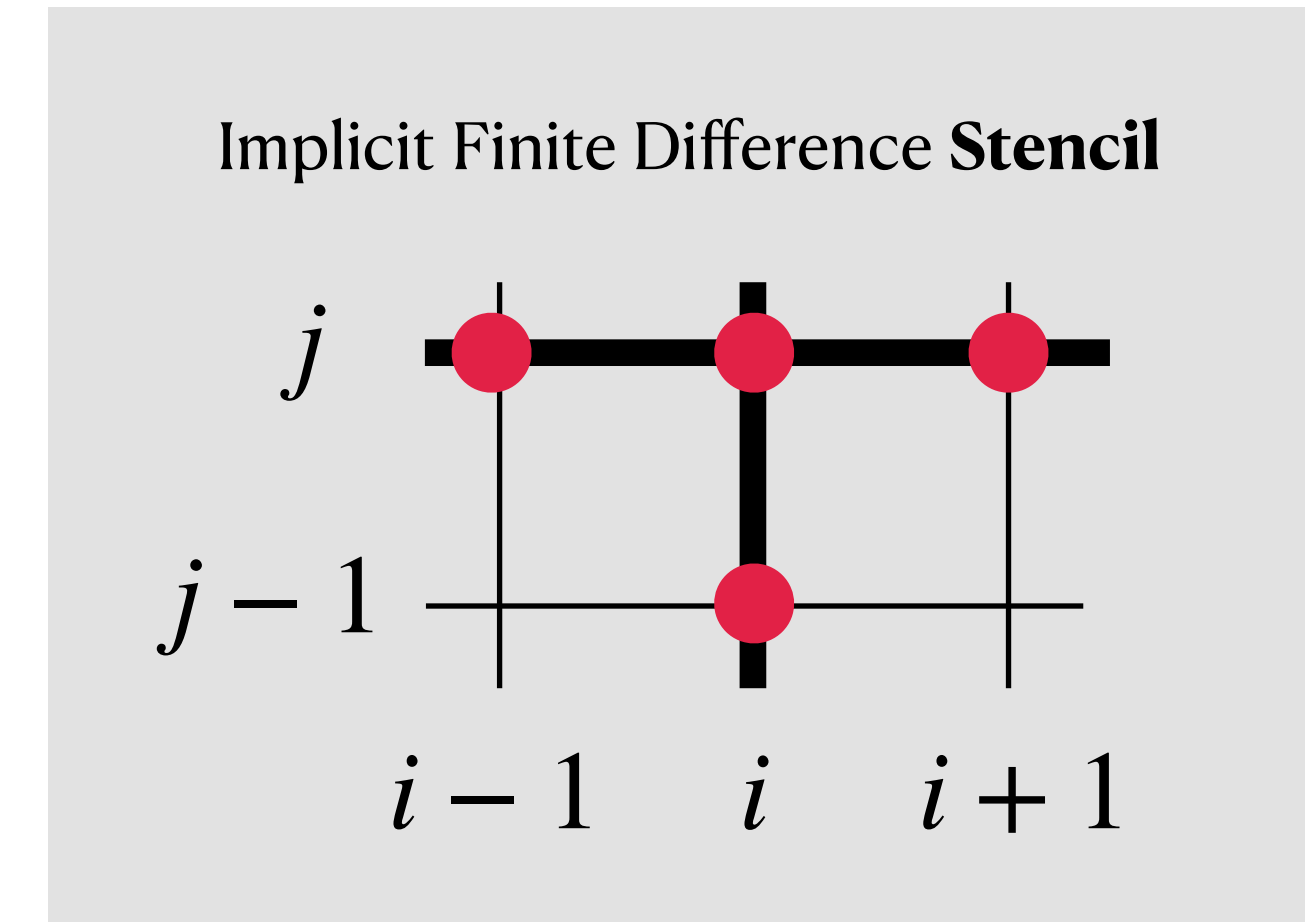
no closed form expression!

Backward finite difference is an example of an **implicit** finite difference scheme

Backward Finite Difference Scheme

$$u_{i,j-1} = -\lambda u_{i+1,j} + (1 + 2\lambda)u_{i,j} - \lambda u_{i-1,j} \quad \lambda = \frac{\Delta t}{\Delta x^2}$$

$$u^j = \begin{bmatrix} u_{1,j} \\ u_{2,j} \\ \vdots \\ u_{m-1,j} \end{bmatrix}, A = \begin{pmatrix} 1 + 2\lambda & -\lambda & 0 & \dots & 0 \\ -\lambda & 1 + 2\lambda & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & -\lambda \\ 0 & \dots & 0 & -\lambda & 1 + 2\lambda \end{pmatrix},$$



Solution must be obtained by solving a (tri-diagonal) linear system of equations:

$$Au^{j+1} = u^j \quad \text{efficient algorithms exist for such sparse matrices}$$

It can be shown that this scheme is unconditionally stable i.e

it converges for all $\Delta x > 0$, $\Delta t > 0$

Some textbooks call this the implicit scheme, but there are many implicit schemes!

Backward Time Cantered Space (BTCS) is the exact name of the method

The Crank-Nicolson Scheme

Combine the implicit and explicit schemes for improved accuracy

$$\frac{u_{i,j+1} - u_{i,j}}{\Delta t} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} \quad \text{Forward } j \text{ (explicit scheme)}$$

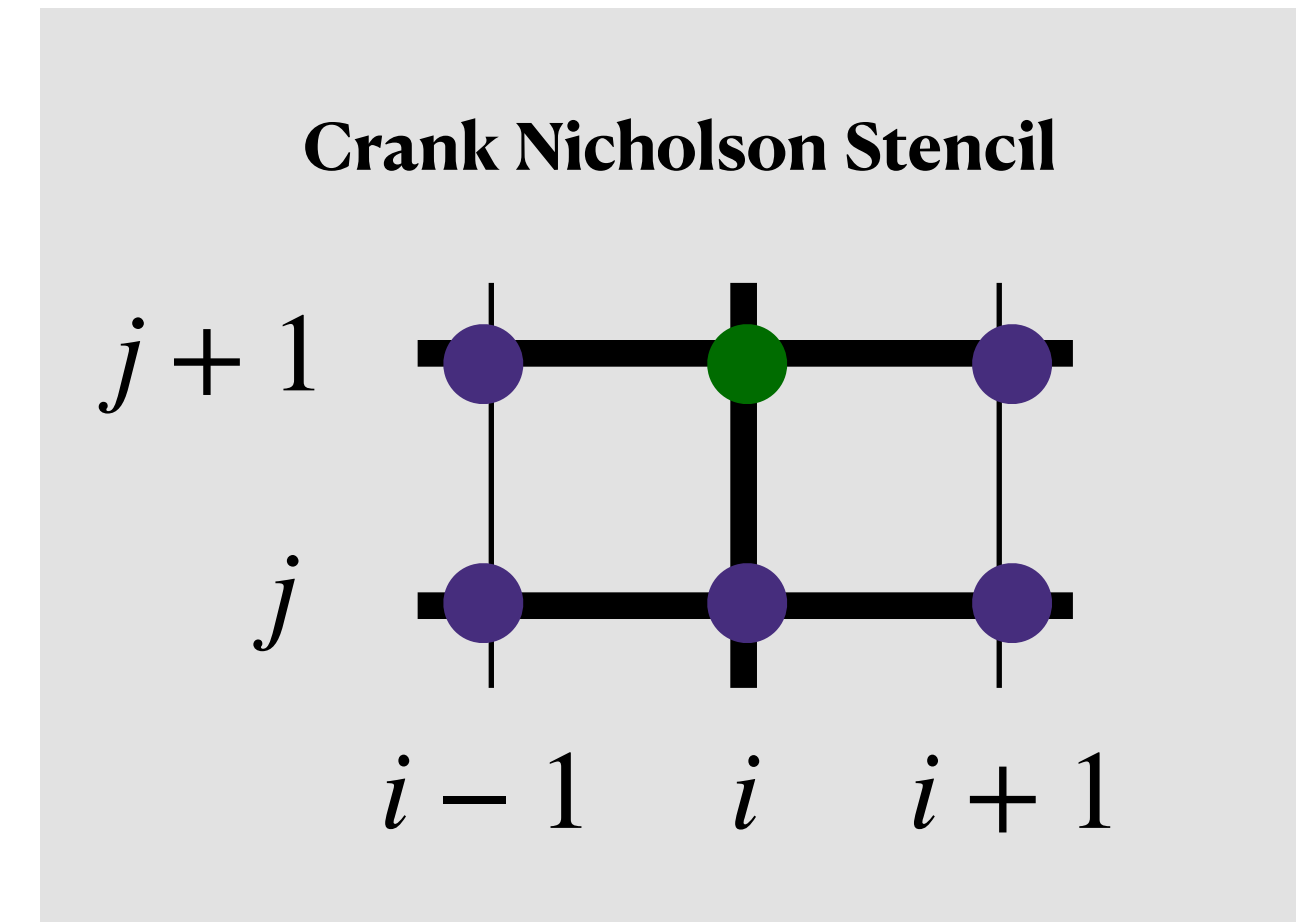
$$\frac{u_{i,j+1} - u_{i,j}}{\Delta t} = \frac{u_{i+1,j+1} - 2u_{i,j+1} + u_{i-1,j+1}}{\Delta x^2} \quad \text{Backward } j+1 \text{ (implicit scheme)}$$

Adding the two:

$$\frac{u_{i,j+1} - u_{i,j}}{\Delta t} = \frac{1}{2\Delta x^2} \left(u_{i+1,j} - 2u_{i,j} + u_{i-1,j} + u_{i+1,j+1} - 2u_{i,j+1} + u_{i-1,j+1} \right)$$

The Crank-Nicolson Scheme

$$\frac{u_{i,j+1} - u_{i,j}}{\Delta t} = \frac{1}{2\Delta x^2} \left(u_{i+1,j} - 2u_{i,j} + u_{i-1,j} + u_{i+1,j+1} - 2u_{i,j+1} + u_{i-1,j+1} \right)$$



- Stability holds for all $\Delta t > 0$
- The convergence order is $O(\Delta t^2) + O(\Delta x^2)$ which is much better than methods studied so far
- Implicit scheme that requires solution of a linear system of equations

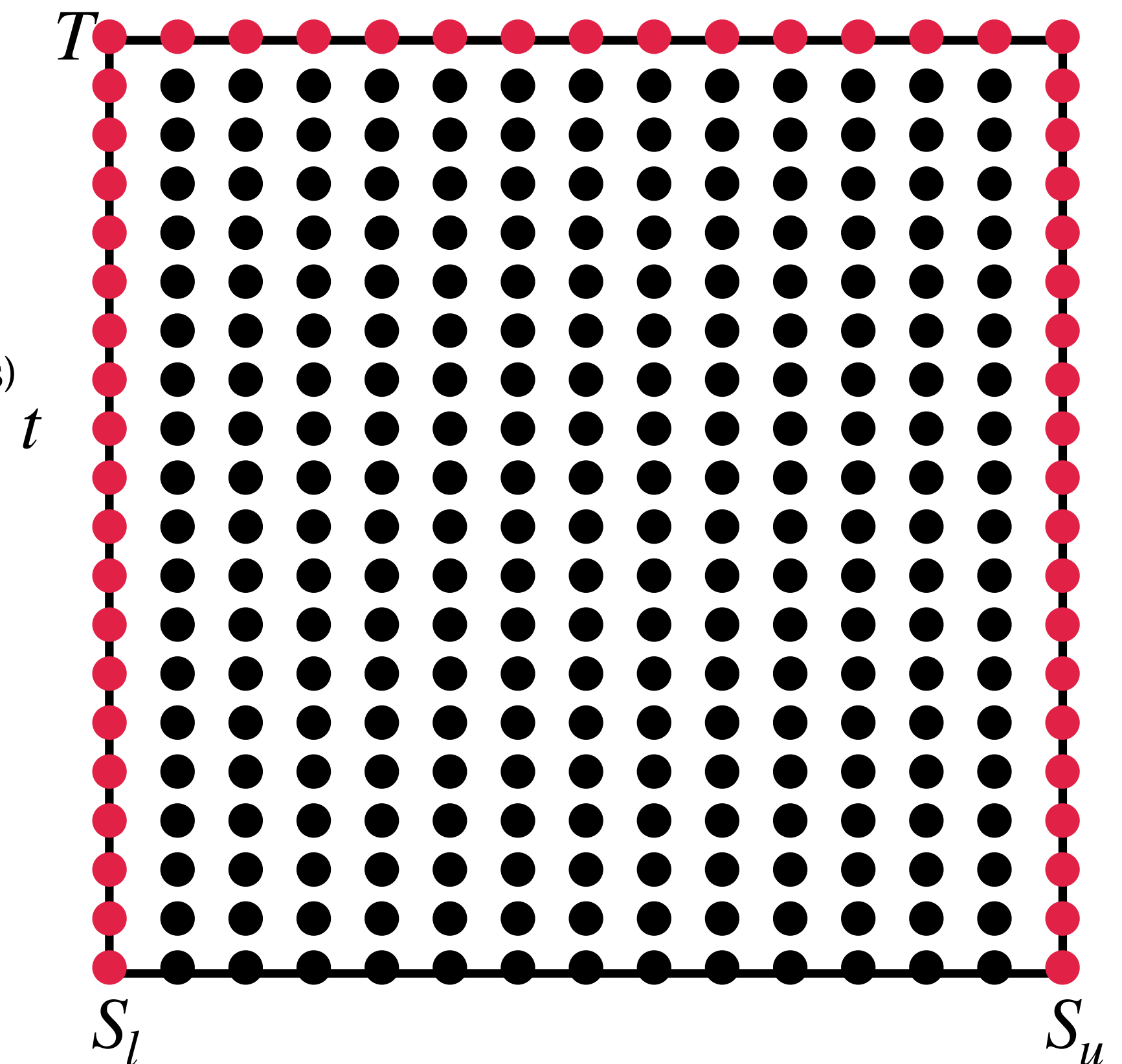
Boundary Conditions

- Finite difference methods need a grid with lower and upper bounds
- Not all boundary conditions will be in contract
- Product knowledge is required here, there are no unique solutions!
- It is easier to apply the boundary conditions on the BSE and then convert it to the heat equation

$$S_l \leq S \leq S_u$$
$$0 \leq t \leq T$$

● Known nodes (boundary conditions)

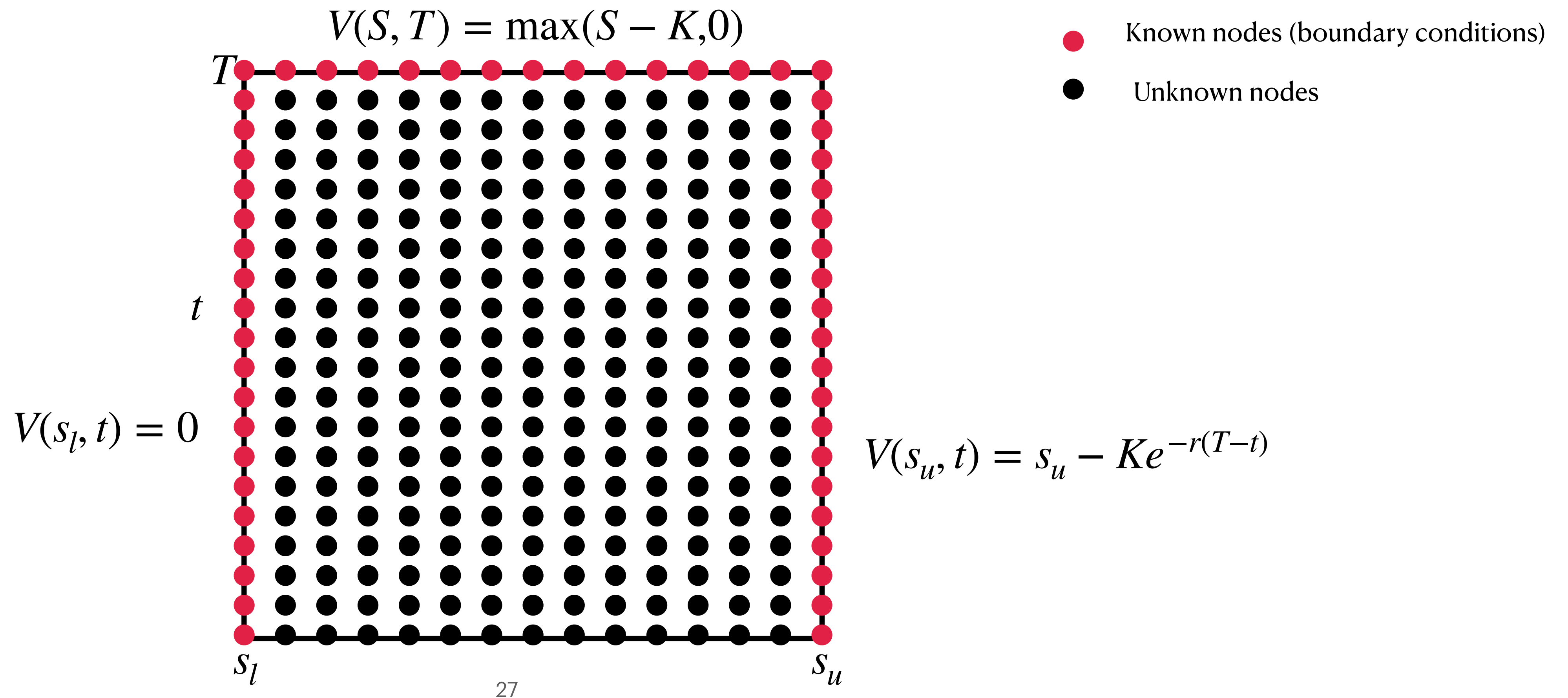
● Unknown nodes



Note: The grid on this slide is for the BSE and is backward in time!

Example: Boundary Conditions Euro Call

- Boundary conditions at expiry are easy: $V(S, T) = \max(S - K, 0)$
- If stock is too low i.e. $S \rightarrow 0$ then option is worthless: $V(s_l, t) \rightarrow 0$
- If stock is too high i.e. $S \rightarrow \infty$ then option will be exercised: $V(s_u, t) = s_u - Ke^{-r(T-t)}$



Finite Difference vs Monte Carlo

- As the dimension increases the FD method suffers from the curse of dimensionality

d dimensions \implies grid of size Δx along each dimension \implies error $\approx (\Delta x)^2$

With a budget of N FD error is $\approx \frac{1}{N^{2d}} = N^{-2d}$

- With a budget of N Monte Carlo error is $\approx 1/\sqrt{N} = N^{-\frac{1}{2}}$ independant of dimension

Summary:^{*}

^{*} This is valid for the methods we encountered in this class, reality can be more complicated but this is a good rule of thumb

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Summary:^{*}

For $d = 1$ FD is better

For $d = 2$ methods are the same

For $d > 2$ MC is better

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