



TP PROJET  
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## Meta - lenses

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## 1 Introduction

A meta-surface is composed of nanostructures smaller than the wavelength of light, enabling the creation of ultra-thin functional devices. Research using electron beam lithography has shown that a wide range of optical devices, including lenses, spectral filters, polarization band-pass filters, and beam shapes, can be effectively replaced by 'flat' meta-surface optics. These meta-surfaces have reached a level of advancement suitable for large-scale production. Semiconductor companies are gearing up for a significant increase in the meta-surface technology market in the coming years. To support this transition, accurate simulations of optical meta-surfaces and innovative design approaches are essential.

Traditional simulation techniques, like Finite-Difference Time-Domain (FDTD), can be challenging to implement due to the relatively large dimensions (mm) of these optical systems and the sub-wavelength diffraction patterns involved. Therefore, alternative approaches based on simplified models have proven effective, facilitating the simulation of large meta-surfaces.

The primary objective of these projects is to explore the potential of these new meta-surfaces and gradually develop the expertise to design meta-lenses.

Figure 1 depicts a simplified diagram of the diffraction structure. Incident light propagates from the left side and passes through the layers of the meta-surface. After reaching the near-field (NF) plane, the wave continues to propagate through free space towards the far-field (FF).

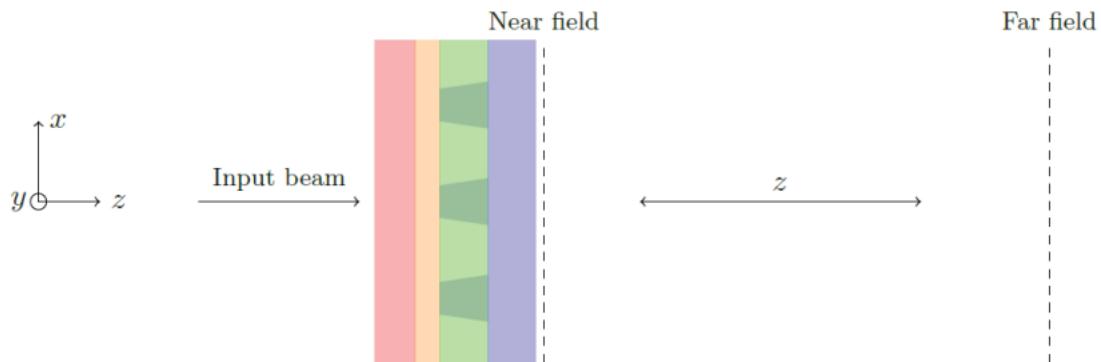


FIGURE 1 – When a laser illuminates a meta surface composed of pillars, it generates a distinct near-field phase pattern. After propagating a distance  $z$ , we measure the electromagnetic field's amplitude in the so-called far field.

Our initial focus will be on understanding the phase shift of the wave as it passes through a specific lens and its connection to the optical characteristics of the propagating waves. The phase of the signal contains significant information about the resulting image once it has propagated a distance  $z$ . We aim to acquire the knowledge necessary to project the Supoptique logo in the far field.<sup>1</sup>

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1. Denis Rideau's document served as the model for this introduction. For the source, see [1].

## 2 Step 1 : Copying a moth's eye

In this part, we will draw inspiration from a moth's eye to craft our first meta-surface, which aims to enhance the transparency of an interface (to allow light to pass through).

### 2.1 The relationship between the dielectric constant and wavelength

As an initial step in optical modeling, engineers need to obtain the required data tables for all materials used in the simulation. To do that, we use the data from the website [refractive index](#) and extract the silicon SiO<sub>2</sub> material that we would like to compare with.

We follow the following steps :

1. Shelf = "MAIN"
2. Book = "Si" (in Oxygen and oxides) or "SiO2"
3. Page = "Popova"
4. Go in the section called "Optical constants of SiO<sub>2</sub> (Silicon dioxide, Silica, Quartz)" and download the DATA in "TXT - Tab separated".
5. Copied and baste the data in a Excel with two columns : one for the wavelength (wl) and one for the index (n)

Now that we have the data, we would like to generate figures showing n and k according to the wavelength in two different units, eV and  $\mu\text{m}$  (see Figures 3 and 4). To do this and for practical modeling purposes, it's essential to understand the relationship between the pair (n, k) and the material's dielectric constants ( $\epsilon_r, \epsilon_i$ ). Similarly, it's crucial to correlate the beam energy E with its wavelength,  $\lambda$  (let us recall that  $E(\text{eV}) = \frac{hc}{\lambda e}$  with  $e = 1,602.10^{-19}\text{J}$ ). We use the code that's presented below to make four figures.

$$E = \frac{n h c N}{\lambda} \quad \text{with } N \text{ the number of photon} \quad (1)$$

$$[n'_r(w) + ik(w)]^2 = \epsilon_r + i\epsilon_i \quad (2)$$

```

close all
% Constantes
h = 6.63e-34;
c = 3.00e8;
e = 1.6e-19;
%chargement des données
load('Si_Palik.txt','ascii');%changer Si_Palik avec SiO2 quand
% nécessaire
lambda=Si_Palik(:,1);%nm
E=h.*c./e./lambda;%eV
n=Si_Palik(:,2);
k=Si_Palik(:,3);
figure;
plot(lambda/1000, Si_Palik(:,2))%micromètre
hold on
plot(lambda/1000, Si_Palik(:,3))%micromètre
xlabel('Lambda (um)'), ylabel('n, k[1]');
legend('n', 'k');
figure;
plot(E, Si_Palik(:,2))%eV
hold on
plot(E, Si_Palik(:,3))%eV
xlabel('Photon Energy (eV)'), ylabel('n, k[1]');
legend('n', 'k');

```

FIGURE 2 – Code to generate Figures 3 and 4

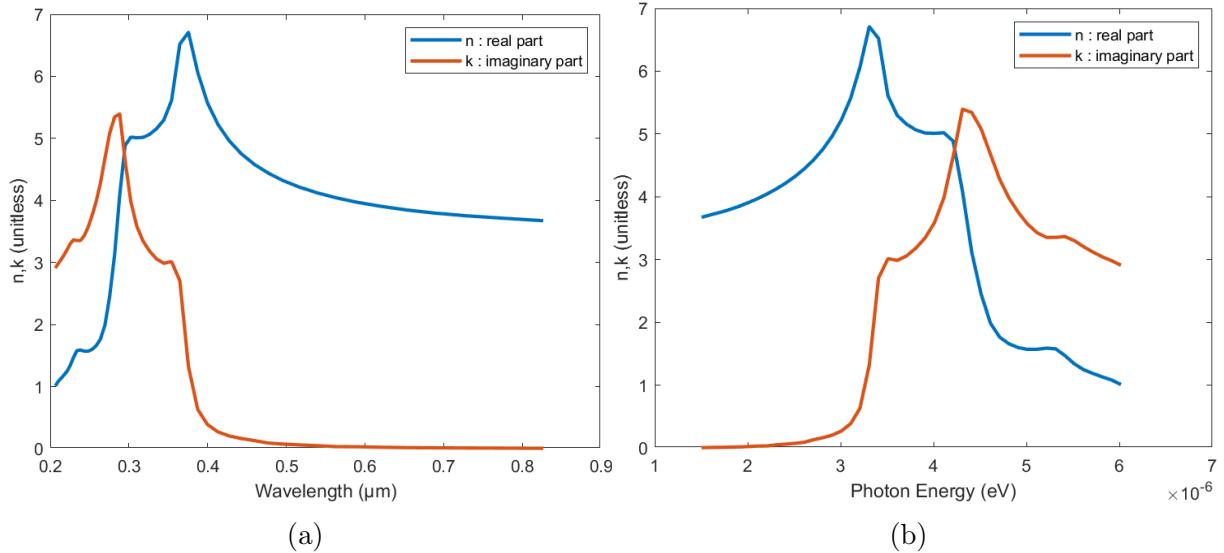


FIGURE 3 – Real and imaginary parts of Si index function of wavelength (a) and photon energy (b)

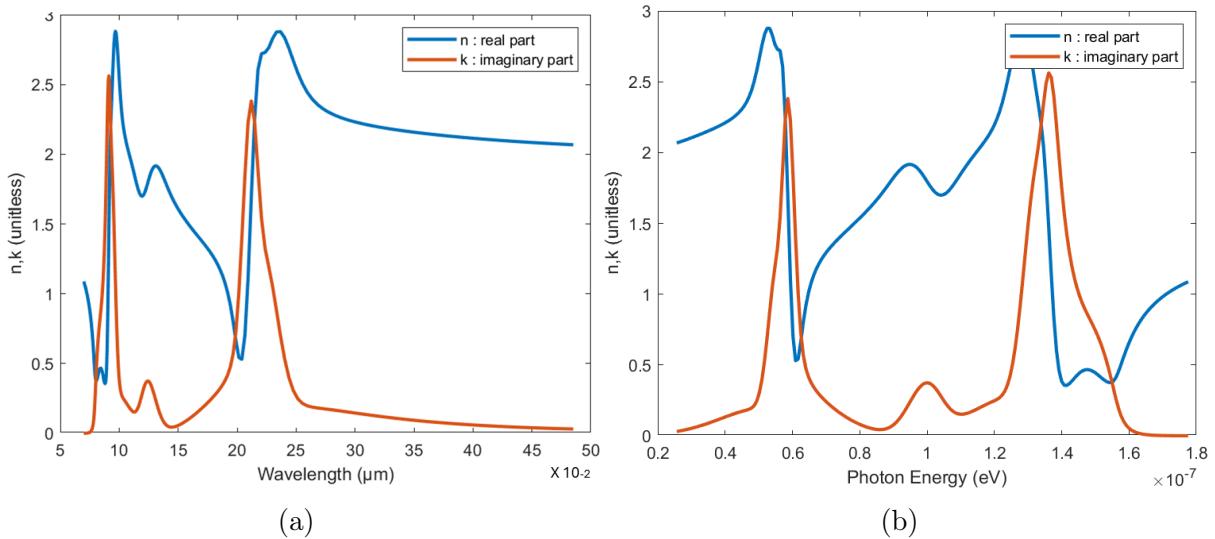


FIGURE 4 – Real and imaginary parts of SiO<sub>2</sub> index function of wavelength (a) and photon energy (b)

Another important quantity related to the imaginary part of  $\epsilon$  is the optical absorption coefficient  $\alpha$ . This value is a measure of how strongly the material absorbs light at a particular wavelength.

$$\alpha = \frac{4\pi k}{\lambda} \quad (3)$$

The Beer-Lambert absorption law is :

$$I = I_0 e^{-\alpha d} \quad (4)$$

With  $I_0$  the initial intensity of the beam before entering the material,  $I$  is the transmitted intensity after passing through the material,  $\alpha$  is the absorption coefficient, and  $d$  is the thickness of the material through which the beam is passing.

So, the Beer-Lambert Law states that the intensity  $I$  of light passing through an absorbing medium decreases exponentially with the thickness of the medium. This relationship helps in determining properties such as the absorption coefficient, which can provide insights into the characteristics of the material, such as its composition and structure. It is also useful in various applications, including the analysis of chemical concentrations, understanding of light-matter interactions, and designing devices like optical filters and sensors.

Since  $k$  represents the imaginary portion, Figures 3 and 4 enable us to understand that the material can be transparent or absorbent at a given wavelength (visible at 0.4 to 0.8  $\mu\text{m}$  and infrared at 1  $\mu\text{m}$ ). By examining the curves, we can observe that Si is somewhat transparent in the visible spectrum and rather absorbant in the UV. In the  $\text{SiO}_2$ ,  $k$  is nearly zero, making it absorbent between 250 and 500 nm.

## 2.2 Transmission, Absorption and Reflection coefficients

When a light beam enters a material, part of it is reflected back at the interface while the rest is transmitted (or refracted) into the material. Augustin-Jean Fresnel described this phenomenon, explaining that the amount of light reflected is quantified by the reflectance,  $R$ , and the amount transmitted is quantified by the transmittance,  $T$ . This occurs because light waves experience reflection at points where there is a difference in optical permittivity (also known as impedance) between adjacent materials. This difference is often referred to as an impedance mismatch. The degree of reflection of the light's energy at the boundary between two materials is directly proportional to the severity of the impedance mismatch ; thus, a greater mismatch leads to a larger portion of the light's energy being reflected.

The community focused on optical devices has long been engaged in researching Anti-reflective coatings (ARCs) to reduce effective reflectance and glare. A well-known approach involves depositing multiple layers on the surface of the material to create a gradual adjustment in impedance.

The expression of the (amplitude) reflection  $r$  and transmission  $t$  coefficient at normal incidence are :

$$r = \frac{n_1 - n_2}{n_1 + n_2} \quad t = \frac{2n_1}{n_1 + n_2} \quad (5)$$

With  $n_1$  is the refractive index of the initial medium (from which the light is coming) and  $n_2$  is the refractive index of the second medium (into which the light is entering).  $r$  measures the fraction of the incident wave's amplitude that is reflected back into the first medium, and  $t$  measures the fraction of the incident wave's amplitude that is transmitted into the second medium.

We implement these equations for Air/Si and Air/ML/Si systems and obtain the Figures 5 with  $T = |t|^2$ ,  $R = |r|^2$  and  $A + R + T = 1$  with the code from Annexe A. These figures are implemented with a Rigorous Coupled-wave Analysis (RCWA). This code solves the Maxwell's equations, decomposing the incident light into planar waves of different directions, and taking into account both the amplitude and phase of the light wave. With that code one propagates the incident wave into a small portion of the coating (a square of 2.5  $\mu\text{m}$ ) and can calculate numerically, for each wavelength, what portion of the incident light is transmitted, reflected or absorbed.

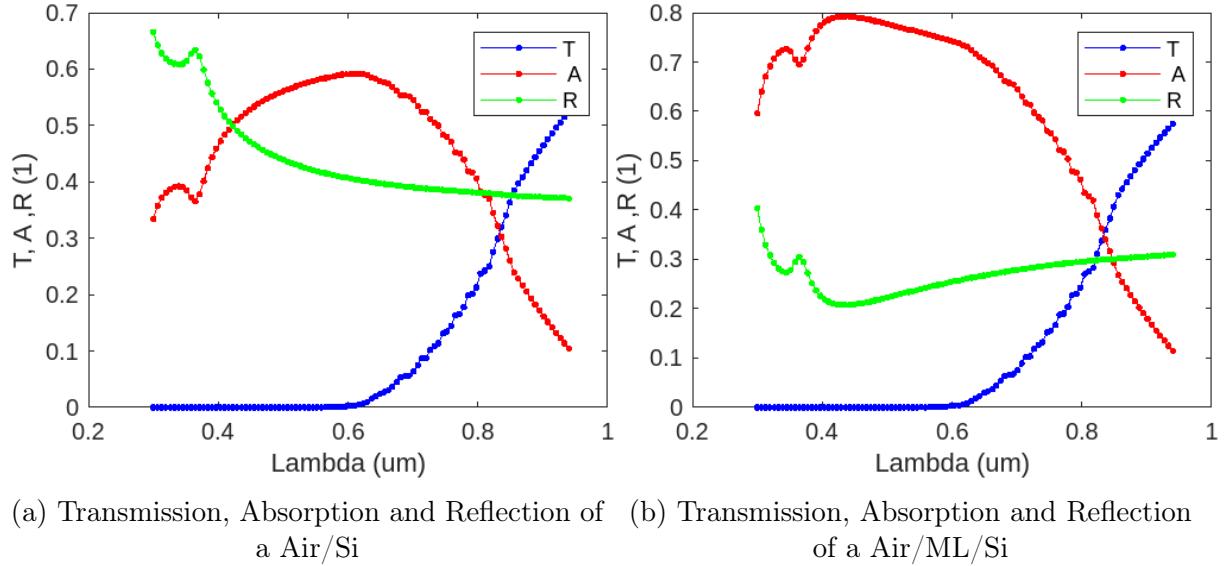


FIGURE 5 – T, A and R coefficients for different layers at normal incidence, calculated with RCWA

As we can see, the absorbent coefficient increases with the number of layers. The reflexion and transmission coefficients are weakened as a result. However, this data was at a typical incidence.

When dealing with light that is incident at an angle other than normal (perpendicular) to the boundary between two media, the expressions for the reflection coefficient (R) and the transmission coefficient (T) become dependent on the angle of incidence as well as the indices of refraction of the two media. These equations differ for two polarizations : perpendicular (s-polarization) and parallel (p-polarization) to the plane of incidence. So we write the equations for at an interface between two media with indices of refraction  $n_1$  and  $n_2$ , and with an angle of incidence  $\theta^2$ .

$$R_s(\theta) = \frac{[\sqrt{a+z} - \sqrt{2}\cos(\theta)]^2 + a - z}{[\sqrt{a+z} + \sqrt{2}\cos(\theta)]^2 + a - z} \quad \text{and} \quad T_s(\theta) = 1 - R_s(\theta) \quad (6)$$

$$R_p(\theta) = \frac{[\sqrt{a+z} - \sqrt{2}\sin(\theta)\tan(\theta)]^2 + a - z}{[\sqrt{a+z} + \sqrt{2}\sin(\theta)\tan(\theta)]^2 + a - z} \quad \text{and} \quad T_p(\theta) = 1 - R_p(\theta) \quad (7)$$

2. For the source, see [7].

With,  $z = n^2 - K^2 - \sin(\theta_1)^2$  and  $a = \sqrt{z^2 + 4n^2K^2}$ .

For unpolarized light, the average reflection and transmission coefficients can be approximated as :

$$R = \frac{R_s + R_p}{2}, \quad T = \frac{T_s + T_p}{2} \quad (8)$$

We obtain two figures with two different angles for the same system, Air/Si (see Figure 6). It is seen that the blue of Rs lies beneath the red of Rp in the figure (a).

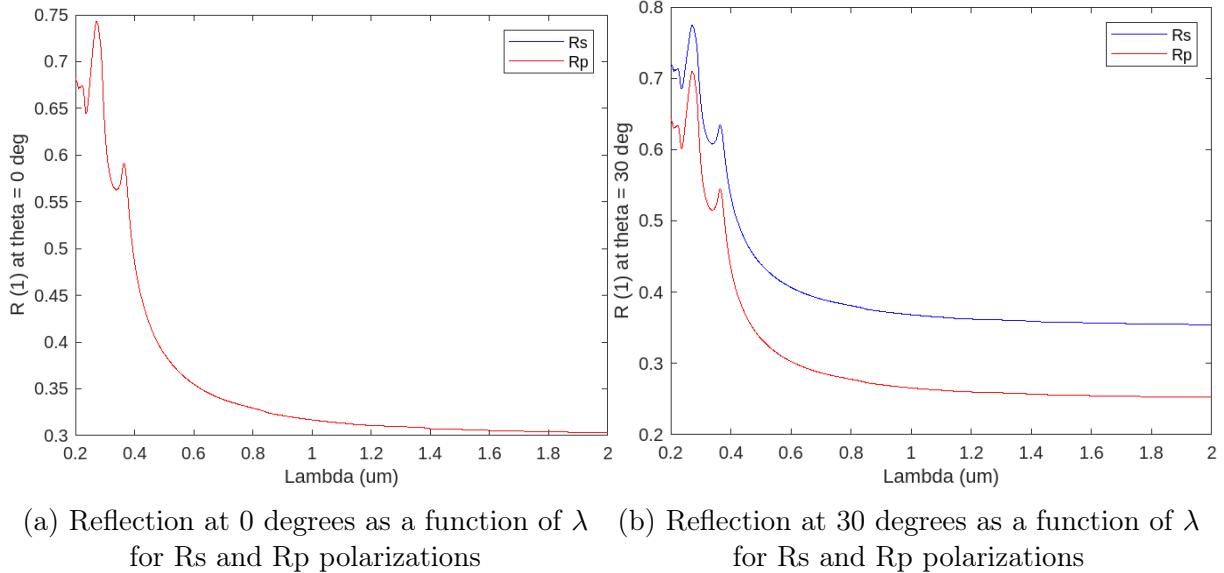


FIGURE 6 – Varying degrees of reflection

We observe that Rs and Rp are not reflected in the same way when the angle is larger. Figure 7 illustrates the evolution.

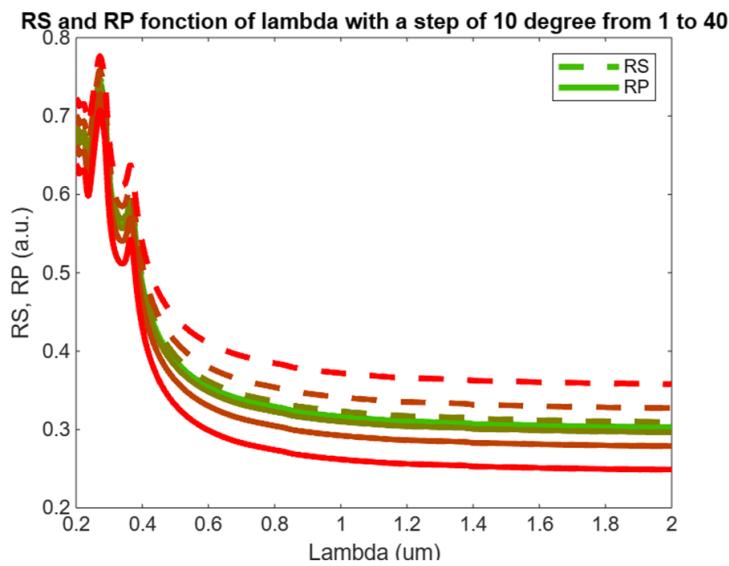


FIGURE 7 – Rs and Rp as a function of  $\lambda$  with a step of 10 degrees

To better describe, we can plot the R and T coefficients for angles between 0 and 90° and for TE and TM polarization. This is what is displayed in Figure 8. We can see the reflections are equal at normal incidence, and then become different with one always increasing and the other one slightly decreasing and reaching 0 before increasing.

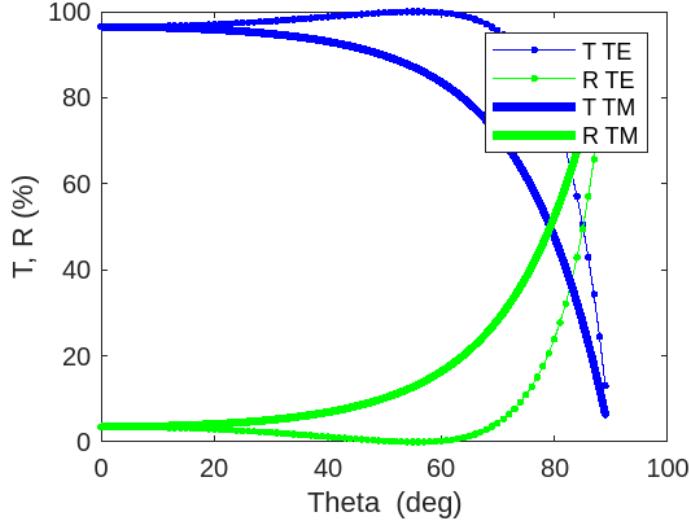


FIGURE 8 – Transmission and reflection of a Air/SiO<sub>2</sub> system as a function of incident angle for two polarizations

We see that a given angle there is the **Critical Angle**. This angle ( $\theta_c$ ) is the smallest angle of incidence at which total internal reflection occurs. For angles of incidence greater than the critical angle, all the incident light is reflected, and none is transmitted through the boundary. The critical angle can be calculated using Snell's law :

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$

where :

- $n_1$  is the index of refraction of the denser medium (higher index of refraction)
- $n_2$  is the index of refraction of the less dense medium (lower index of refraction)
- $\theta_1$  is the angle of incidence in the denser medium
- $\theta_2$  is the angle of refraction in the less dense medium

For total internal reflection to occur, the angle of refraction ( $\theta_2$ ) would reach 90 degrees, since beyond this angle, the refracted ray does not enter the second medium but skims along the interface. Setting  $\theta_2 = 90$  in Snell's law gives :

$$\theta_c = \arcsin\left(\frac{n_2}{n_1}\right) \quad (9)$$

For the critical angle to exist :

- Light must be traveling from a medium of higher refractive index to one of lower refractive index
- $\theta_c$  is only defined for  $n_1 > n_2$  otherwise, the arcsinus function does not yield a real number since the sine of an angle cannot be greater than one.

### 2.3 ARC with a moth's eye

The objective is to develop an Anti-Reflective Coating (ARC) by structuring the silicon layer to resemble the pattern of a moth's eye, as depicted in Figure 9 and discussed in reference [3].

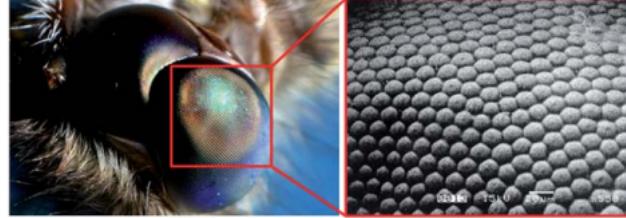


FIGURE 9 – The image on the left shows the eye of the moth, while the SEM image on the right displays the nano bumps or protrusions on the outer surface of the corneal lens.

Using RCWA, we model the absorption and reflection as a function of wavelength and simulate the effect of a surface structuration on light reflection on the Air/Si system. The surface structuration is represented by Figure 10, and the output of the RCWA code is shown in Figure 11. This time, exercise caution as the coefficients' units are percentages.

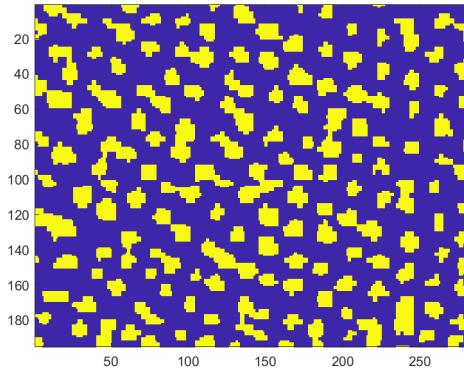


FIGURE 10 – Structure used for the anti-reflective metasurface simulation

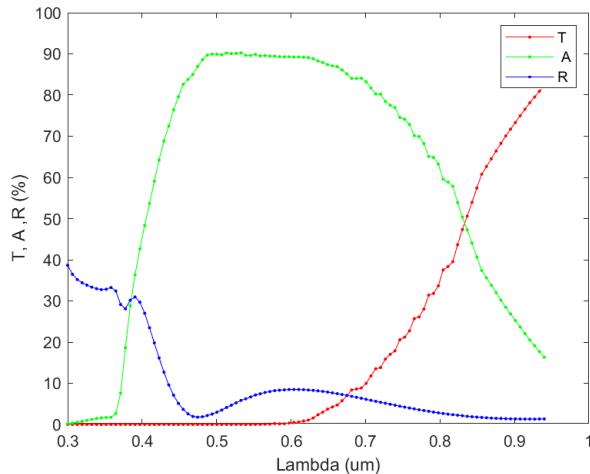


FIGURE 11 – T, R and A of the Air/Si system the surface of which is structured with a random pattern mimicking the moth's eye

Results : when compared to all other examples previously analyzed, we can observe a relatively rapid decline in the R coefficient. Therefore, we may infer that, in comparison to the single and multilayer coatings previously researched, the utilization of a nano-sized metastructure is an excellent solution for our anti-reflective coating.

## 2.4 Wave front phase shift

We characterize the maximum phase shift as the discrepancy in phase between a wave traversing through a reference material of uniform composition, having a thickness of  $h_0$  and a refractive index of  $n_1$ , and a wave traversing through a non-uniform layer also with a thickness of  $h_0$ , consisting of two materials with refractive indices  $n_1$  and  $n_2$  respectively.

The source [1], page 5, 6 and 7, helps us to find the thickness of this non-homogeneous layer to varies the phase shift up to to 2 in the second material with respect to the reference one. We can write that :

$$h_0 = \frac{\lambda_0}{n_1 - n_2} \quad (10)$$

## 3 Step 2 : Propagation of light

We shall discuss the close field and the far field in this section. We also describe the Canonical transformation, which has a proportionality to the angle  $\alpha_z$  and a variable change.

### 3.1 Definition of the Fourier transform

The Fourier transform and the inverse Fourier transform are defined as follows :

$$F[H(x, y)](k_x, k_y) = \iint_{-\infty}^{\infty} H(x', y') e^{-i(k_x x' + k_y y')} dx' dy' \quad (11)$$

$$F^*[F(k_x, k_y)](x, y) = \iint_{-\infty}^{\infty} F(k_x, k_y) e^{i(k_x x + k_y y)} dk_x dk_y \quad (12)$$

In practical applications, the real space is finite, with dimensions  $L_x$  and  $L_y$ , and discretized into  $N_x$  and  $N_y$  points. It's important to note that the FFT yields a set of Fourier coefficients  $F_{nx,ny}$ , with  $nx$  ranging from  $-\frac{N_x}{2}$  to  $\frac{N_x-1}{2}$  and  $ny$  ranging from  $-\frac{N_y}{2}$  to  $\frac{N_y-1}{2}$ , implicitly defining the Fourier space as :

$$\left\{ \frac{2\pi nx}{L_x} \mid nx \in \left[ -\frac{N_x}{2}, \frac{N_x-1}{2} \right] \right\} \quad (13)$$

$$\left\{ \frac{2\pi ny}{L_y} \mid ny \in \left[ -\frac{N_y}{2}, \frac{N_y-1}{2} \right] \right\} \quad (14)$$

It is possible to implement a "brute" force FT. This is a good exercise to understand the basics of Fourier analysis. We can start with a 1D model and then implement a 2D model. Here are some explanation about this two models :

- In a 1D model, you have a signal  $f(x)$  where  $x$  is the independent variable. The Fourier Transform of  $f(x)$ , denoted as  $F(k)$ , is given by :

$$F(k) = \int_{-\infty}^{\infty} f(x) \cdot e^{-2\pi i k x} dx \quad (15)$$

For a discrete implementation, we approximate this integral by a sum :

$$F(k) \approx \sum_{n=0}^{N-1} f(n\Delta x) \cdot e^{-2\pi i k n \Delta x} \quad (16)$$

Where  $\Delta x$  is the spacing between the sample points and  $N$  is the total number of sample points.

- For the 2D case, where you have a function  $f(x,y)$ , the Fourier Transform  $F(k_x, k_y)$  is given by :

$$F(k_x, k_y) = \iint_{-\infty}^{\infty} f(x, y) \cdot e^{-2\pi i (k_x x + k_y y)} dx dy \quad (17)$$

For a discrete implementation, similar to the 1D case, you'll have :

$$F(k_x, k_y) \approx \sum_{m=0}^{M-1} \sum_{n=0}^{N-1} f(m\Delta x, n\Delta y) \cdot e^{-2\pi i (k_x m \Delta x + k_y n \Delta y)} \quad (18)$$

One can make a comparison between an FFT-based solution and a basic 2D example. Regretfully, our algorithms did not function as intended ; however, the FFT-based solution outperforms the brute-force method in terms of speed, particularly for bigger datasets.

## 3.2 Near Field and Far Field

A scalar approximation of the vectorial electromagnetic wave can be derived using the Fresnel transform. Assuming uniform wave propagation (with a wavelength of  $\lambda = \sqrt{\epsilon_r} \lambda_0$ ), the fields at any given position can be inferred from the field at a location situated at  $z=0$ .

$$E(x, y, z) = \frac{1}{i\lambda} \iint_{-\infty}^{\infty} E(x', y', z=0) \frac{e^{ikr} z}{r^2} dx' dy' \quad (19)$$

Where  $E(x', y', z=0)$  represents the near field, and  $r = \sqrt{(x-x')^2 + (y-y')^2 + z^2}$ . Part 3.6 illustrates how this equation is put into practice.

## 3.3 Fresnel integrals as a convolution

Given that the Near Field region, where the electric field is non-zero, has dimensions considerably smaller than  $z$ , the Fresnel approximation  $r \approx z + \frac{(x-x')^2 + (y-y')^2}{2z}$  can be employed for the exponential term (and  $r \approx z$  in the  $\frac{1}{r}$  denominators) :

$$E(x, y, z) = e^{ikz} \frac{i}{\lambda z} \iint_{-\infty}^{\infty} E(x', y', z=0) e^{ik \frac{(x-x')^2 + (y-y')^2}{2z}} dx' dy' \quad (20)$$

This equation represents a convolution that can be effectively computed in reciprocal space by multiplying the Fourier transforms of the two terms and then performing an inverse Fourier transform. For efficiency purposes, one can utilize the Fourier transform of the impulse response of the free space propagation kernel.

$$H_{Fresnel}(x, y) = e^{i \frac{\pi}{\lambda z} (x^2 + y^2)} \quad (21)$$

$$\mathcal{F}\{e^{i \frac{\pi}{\lambda z} (x^2 + y^2)}\}(k_x, k_y) = i \lambda z e^{-i \frac{\lambda z}{4\pi} (k_x^2 + k_y^2)} \quad (22)$$

When considering the reciprocal space defined with the  $(2\pi)$  factors, by redefining the pre-exponential factor in the kernel, one can express the Fourier transform of Equation 20 as :

$$F[E(x, y, z)](k_x, k_y) = e^{iz \left[ \frac{2\pi}{\lambda} + \frac{\lambda(k_x^2 + k_y^2)}{4\pi} \right]} \times F[E(x, y, z=0)](k_x, k_y) \quad (23)$$

### 3.4 Helmholtz as a convolution

It's worth noting that this kernel closely resembles another more precise kernel, with which the convolution satisfies the Helmholtz equation. We can write the Equation 23 as :

$$F[E(x, y, z)](k_x, k_y) = e^{i \frac{z}{2\pi} \sqrt{(\frac{2\pi}{\lambda})^2 - k_x^2 - k_y^2}} \times F[E(x, y, z=0)](k_x, k_y) \quad (24)$$

### 3.5 Fresnel integrals with a canonical transformation

An alternative way of evaluating the Fresnel integral can be performed using a Fourier transform and a canonical transformation. Using :

$$(x - x')^2 = x^2 + x'^2 - 2xx' \quad (25)$$

$$(y - y')^2 = y^2 + y'^2 - 2yy' \quad (26)$$

and  $k = \frac{2\pi}{\lambda}$ , the expression becomes :

$$E(x, y, z) = \frac{e^{ikz}}{i\lambda z} e^{i \frac{\pi}{\lambda z} (x^2 + y^2)} \iint_{-\infty}^{\infty} E(x', y', z=0) e^{i \frac{\pi}{\lambda z} (x'^2 + y'^2)} e^{-i 2\pi \left( \frac{x}{\lambda z} x' + \frac{y}{\lambda z} y' \right)} dx' dy' \quad (27)$$

We recognize a Fourier transform with wave numbers  $k_x = \frac{2\pi x}{\lambda z}$  and  $k_y = \frac{2\pi y}{\lambda z}$ . It ends with :

$$E(x, y, z) = \frac{e^{i \frac{2\pi}{\lambda} z}}{i\lambda z} H(x, y) F [E(x', y', z=0) H(x', y')] (k_x, k_y) \quad (28)$$

One notices the implicit dependence of  $x$  and  $y$  in  $k_x$  and  $k_y$ .

### 3.6 Pinhole and Far Field

To further proceed, let's define two planes, namely the Near Field plane at  $z = 0$  and the Far Field plane at a given  $z$ . We have observed that the far field plane and the reciprocal Fourier space are linked by a canonical transformation. Recalling the definition of the reciprocal wave vectors, one can therefore define the set of points in the Far field space  $(x^*, y^*) := (\lambda z k_{x,y})$ . Considering the Fourier space definition, the Far Field space  $(x_n^*) := \frac{n\lambda z}{L_x}$ , where  $n \in [-\frac{N_x}{2}, \frac{N_x-1}{2}]$ , differs from the Near field space  $(x_n) := \frac{nL_x}{N_x}$ , where  $n \in [-\frac{N_x}{2}, \frac{N_x-1}{2}]$ . One also notes that the size along the X direction of the far field ( $x^*$ ) is smaller than the near field one  $L_x$  when  $z < \frac{N_x^2}{L_x \lambda}$ . Three Figures are shown in Figure 12, which we obtained by using the code from section B.

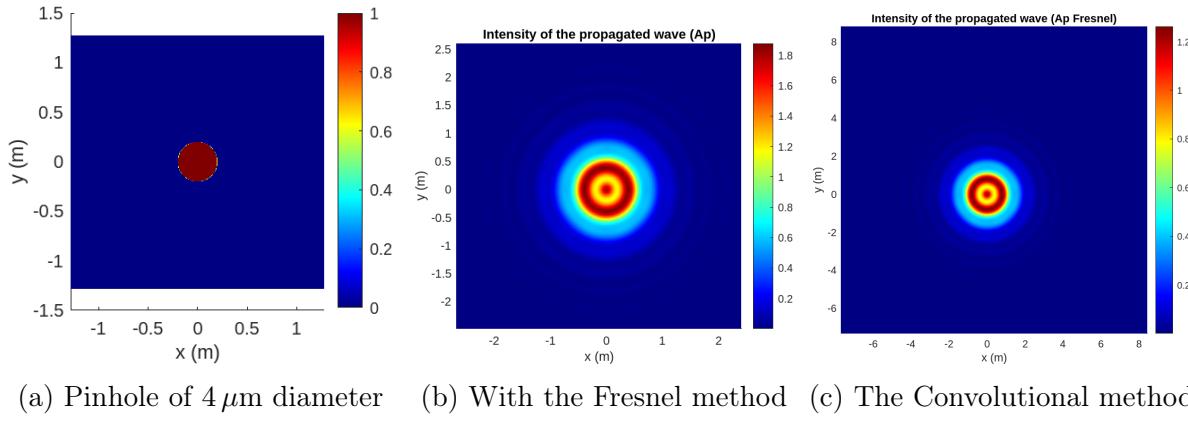


FIGURE 12 – Propagation of a pinhole at  $z = 10 \mu\text{m}$

The boundaries for the Fourier transform integration are theoretically infinite. However, in practice, padding with zeros outside the near field plane can be employed to effectively increase its size. Figures 13 provide a comparison between convolution and canonical transform, highlighting the impact of padding.

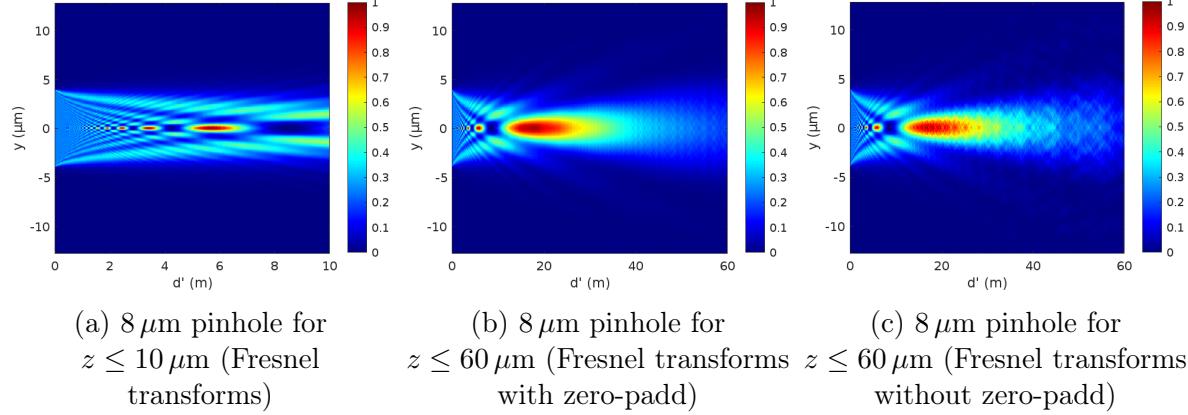


FIGURE 13 – Near field domain with a pinhole, vertical cut using convolution with FFT

The Figure [14] is quite different because we still had to implement the interpolation for the variable change but it still seems correct.

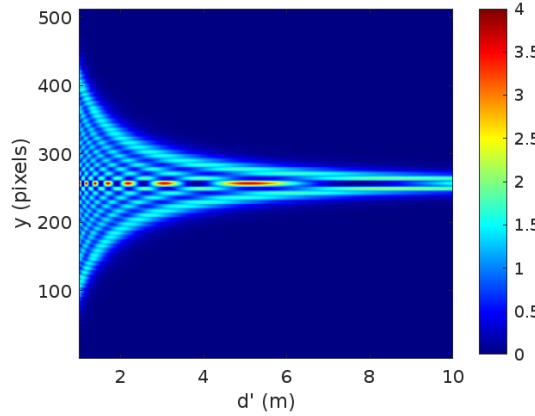


FIGURE 14 – Near field domain with a pinhole, vertical cut using canonical transform with FFT

## 4 Step 3 : Design of a lens

We begin by revisiting the theoretical lens equation to solidify our understanding of how lenses function to manipulate light paths. This formula is crucial for any practical application involving lens design :

$$\frac{1}{f'} = \frac{1}{x'} + \frac{1}{x} \quad (29)$$

Following this, we will simulate the in-plane phase shift that occurs just after light passes through the lens. This involves plotting how the phase of the light changes across the plane of the lens, providing a visual representation of the optical effect induced by the lens :

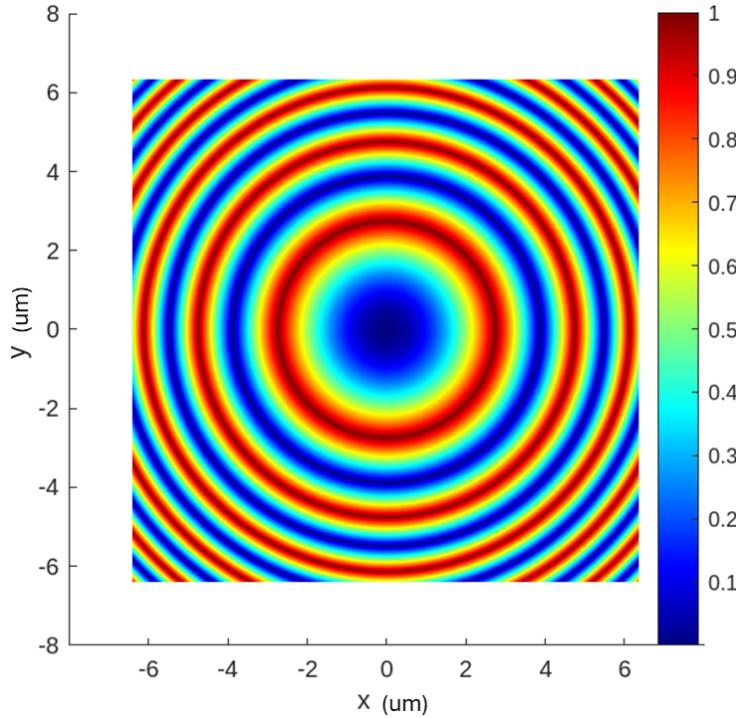


FIGURE 15 – In plane phase shift after the lens

Then we examine the vertical central-cut of the propagating fields. This involves plotting both the modulus (or magnitude) and phase of the light field as it moves through space. This plot will help us identify the focal point of the lens, which is where the light rays converge, creating a sharp focus. An example of these plots can be found in Figure [16].

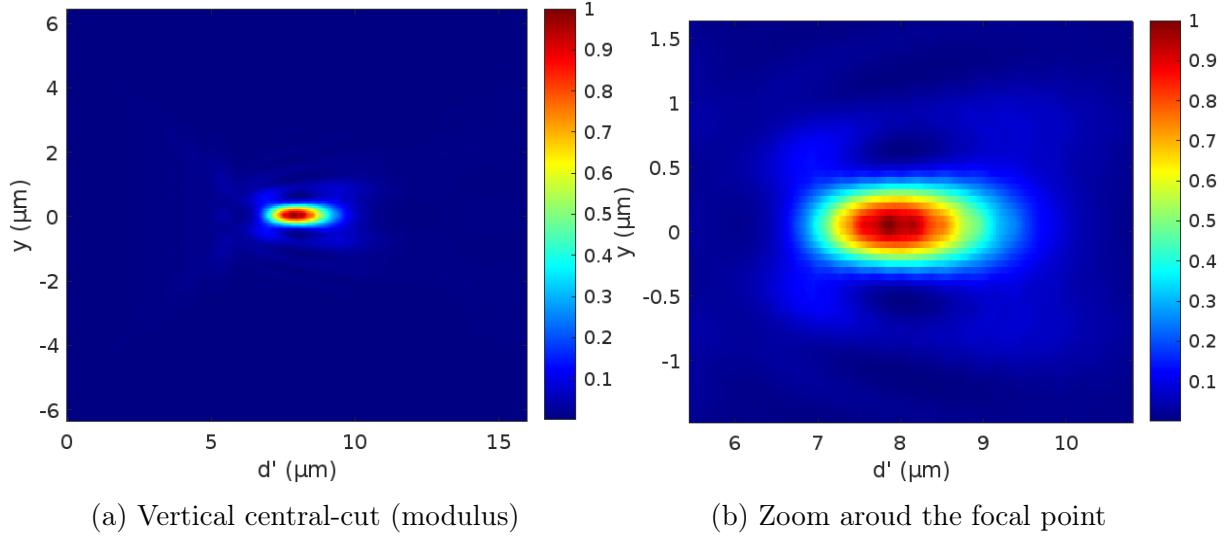


FIGURE 16 – Vertical central-cut of the propagating fields (modulus), focal length of  $f = 8 \mu\text{m}$  at 940 nm

However, we can notice that the focal point is not well resolved : we can measure  $2 \mu\text{m}$  on the Figure[16b]. We can resolve this by increasing the number of pixels on the image. Here is the result of it after going from 256 to 1000 pixels on each side :

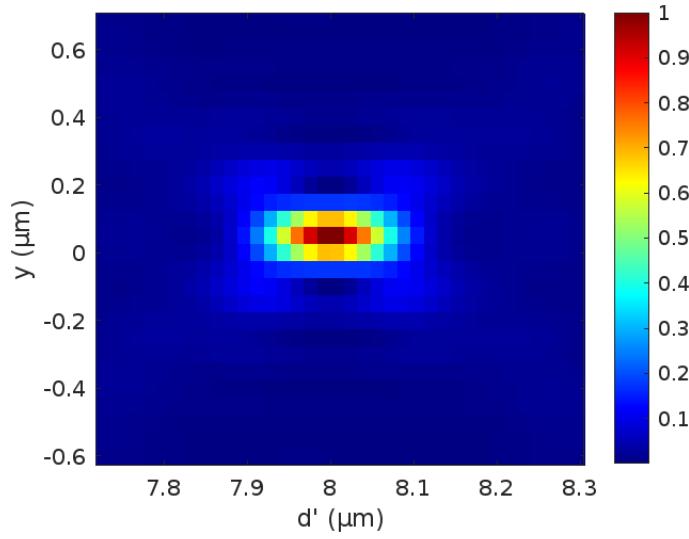


FIGURE 17 – Vertical central-cut of the propagating fields (modulus) with a better resolution

Here is what we obtain for the phase of the propagating fields : we can also notice that that at  $z = 8 \mu\text{m}$ , we have a symmetry for the converging incident wave.

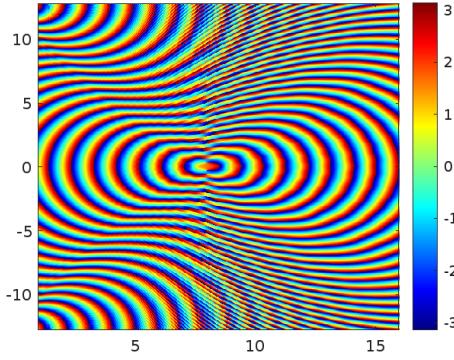


FIGURE 18 – Vertical central-cut of the propagating fields (phase)

Once we are comfortable with these concepts and the basic lens equation, we proceed to study the propagation of light after passing through such a lens using the Helmholtz equation. This step involves a deeper dive into wave optics, allowing us to predict and visualize how light behaves in more complex scenarios. We will conduct simulations to propagate light at various positions beyond the lens and display 2D-cut planes at positions that are deemed relevant based on the optical setup and the specific experiment objectives. Additionally, we plot a vertical cut perpendicular to the lens to examine the cross-sectional distribution of the light field :

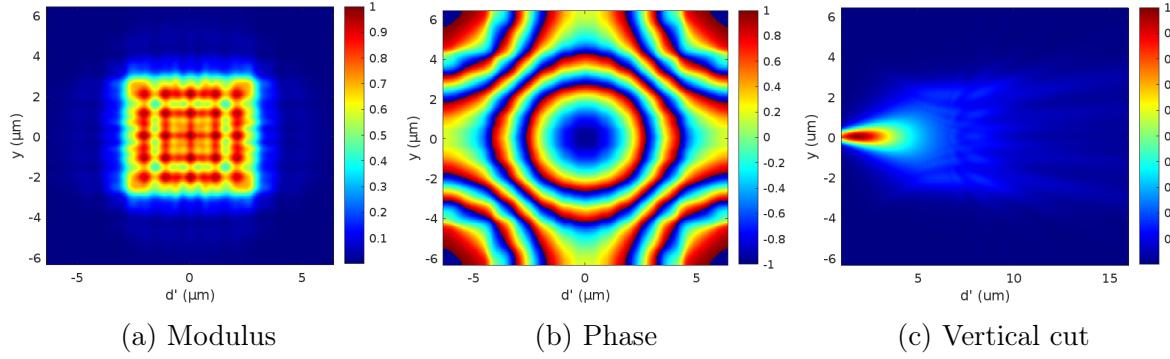


FIGURE 19 – Helmholtz method : Vertical central-cut of the propagating fields at the focal point, focal length of  $f = 8 \mu\text{m}$

Finally, we will vary the focal length and the lateral size of the lens to explore how these changes affect the focusing properties and the overall light distribution :

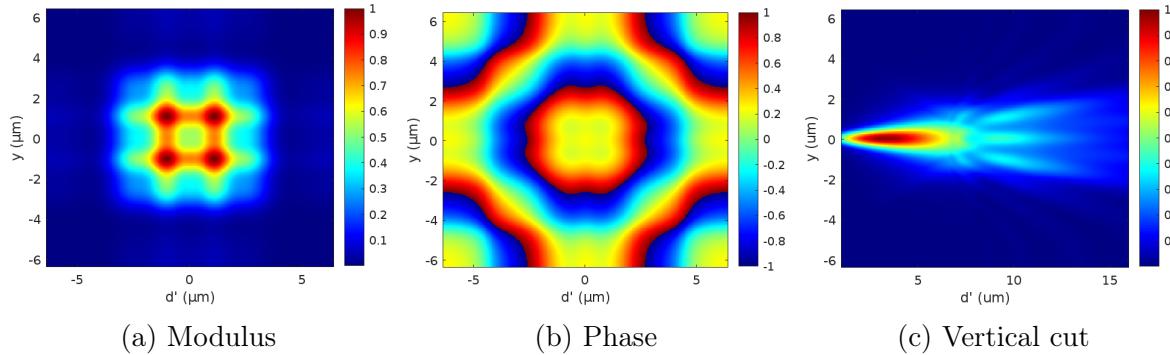


FIGURE 20 – Helmholtz method : Vertical central-cut of the propagating fields at the focal point, focal length of  $f = 20 \mu\text{m}$

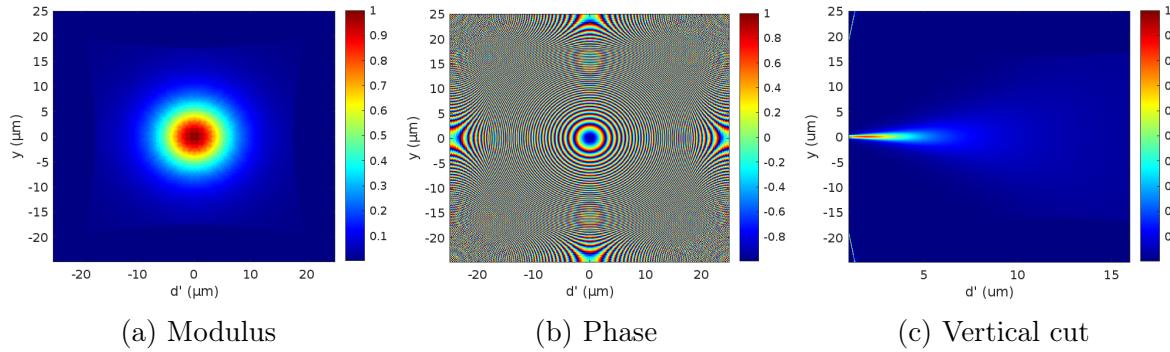


FIGURE 21 – Helmholtz method : Vertical central-cut of the propagating fields at the focal point, focal length of  $f = 8 \mu\text{m}$  and lateral size of  $50 \mu\text{m}$

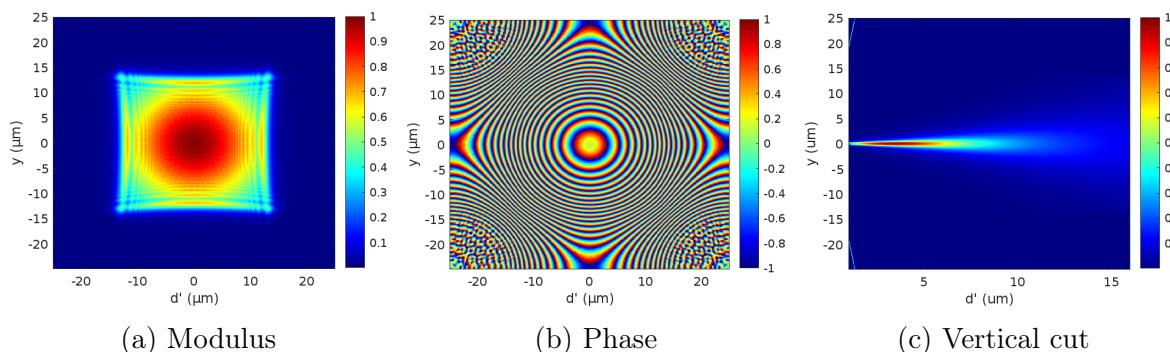


FIGURE 22 – Helmholtz method : Vertical central-cut of the propagating fields at the focal point, focal length of  $f = 20 \mu\text{m}$  and lateral size of  $50 \mu\text{m}$

We can easily notice here that the focal length  $f$  and grid size  $L$  significantly impact the results. The focal length  $f$  determines the curvature of the phase profile in the complex transmittance function, affecting how sharply the wavefront focuses; a shorter  $f$  enhances focusing power, leading to a more pronounced convergence, while a longer  $f$  produces a flatter wavefront that converges over greater distances. Meanwhile, the grid size  $L$  influences both the resolution and the field of view of the simulation, with larger values providing better spatial sampling that can capture more detailed wavefront behaviors and reduce aliasing in Fourier transforms. However, increased  $L$  also raises computational demands, impacting processing time for example. These parameters should be carefully selected based on the specific requirements of the simulation to balance accuracy and computational efficiency.

However, we could still improve these results by adjusting the zero-padding.

Finally, we can combine these lens configurations with a pinhole aperture to study the effects of aperture size on light diffraction and focus quality (see Figure 23).

Smaller apertures seem to have a higher tendency to show pronounced diffraction effects due to their closer approach to the wavelength of light used. The diffracted waves interfere more significantly, potentially creating complex patterns such as a central minimum or "hole" in the center of the main diffraction spot.

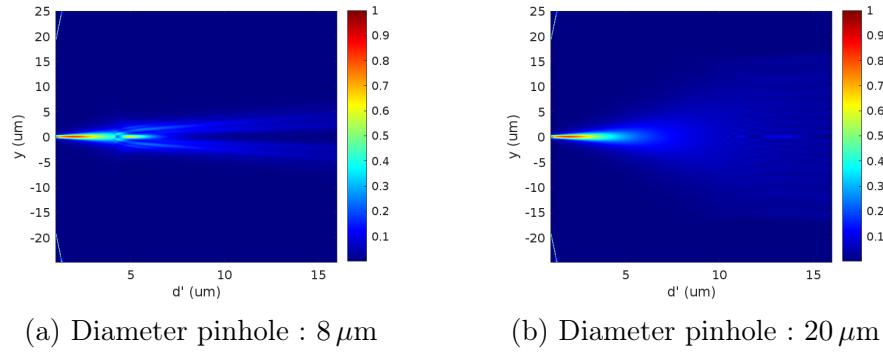


FIGURE 23 – Helmholtz method : Vertical central-cut of the propagating fields at the focal point, focal length of  $f = 8 \mu\text{m}$  for different apertures

## 5 Step 4 : Toward a metasurface

There are three sections here. The first describes how to depict a three-dimensional meta surface; the second describes the characteristics of the phase and transmission as a function of a pillar width; and the third and final section examines Supoptique's logo.

## 5.1 Three-dimensional meta surface

Using a metasurface to implement the previously designed optical function is the last step in this course. As seen in Figure 24, where the silicon tube—which is incapable of floating in the air—is represented in black and the SiO<sub>2</sub> is depicted in blue—we usually start by generating a 3D image. This figure represents a gaussian light-illuminated pillar. The metasurface, which turned into a mirror, is this pillar. When the metasurface is fully lit, resonance occurs.

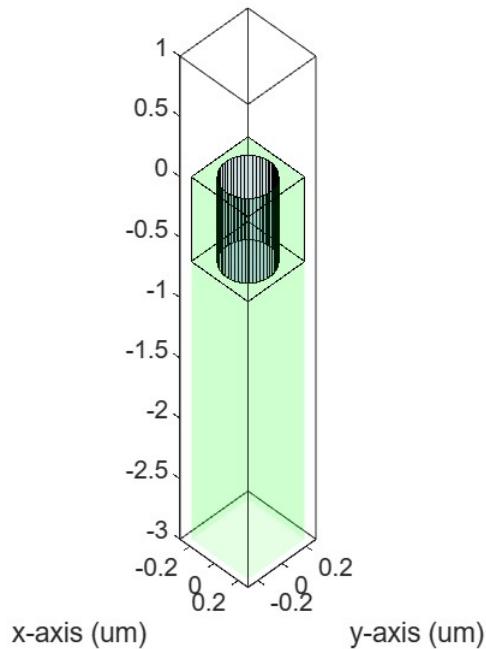


FIGURE 24 – Eight-mode 3D depiction of silicon in SiO<sub>2</sub>

## 5.2 Phase and transmission as a function of a pillar width

As seen in Figures 25 and 26, we can have a look-up table, also called a library, that shows the phase shift versus the width of the diffracting element (in this case, a pillar). We apply the code from Annexe C to obtain these Figures. The codes employ functions to alter the pillar's diameter dimension and to depict the transmission and phase as a function of diameter. The interval between each pillar can also be altered ; in Figure 25, it is  $0,47 \mu\text{m}$ , while in Figure 26, it is  $0,6 \mu\text{m}$ .

Alternatively, the mode can be altered by using the "OPTION.halfnpw = " 8 or 0 in the code included in the C Annexe. Two curves are depicted in Figure 25 : one with mode 0, which is the Fundamental of FFT and corresponds to a basic average of the structure permittivity, and the other with eight modes.

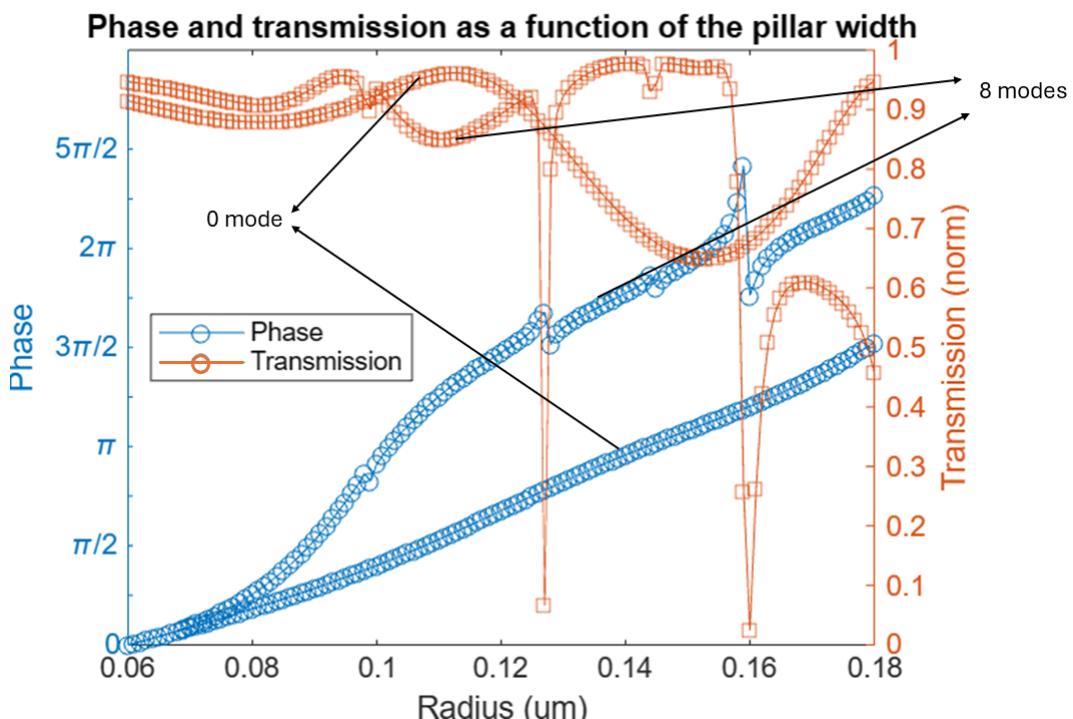


FIGURE 25 – Phase and Transmission : 0 and 8 modes

Explanation : both modalities have the same commencement of the phase and the transmission. After that, we can observe that the phase doubles with the addition of modes. That can be explained by the fact that the optical path grows (doubles in our case) as the wave bounds on the surface. This explains why the phase curve for the eight modes increases quickly.

We may observe that there are no points in the transmission where the fundamental is 0. This is a good news because it is challenging if we need to add modes if the wave failed to pass for the fundamental mode. The value of the transmission for eight modes decreases to zero for certain levels. It indicates that the surface is reflective. Finding the ideal radius for a suitable transmission is therefore feasible.

We are now altering the interval between the pillars. We entered  $0.6 \mu\text{m}$  and put 8 modes. Figure 26 illustrates that a large number of points in the transmission are null. It is evident that there is no transmission between  $\pi$  and  $2\pi$ . This is the plage where the pillar is shown (see Figure 24), which makes it problematic. Thus, the separation between each pillar is crucial.

These numbers are displayed following 215 seconds. Therefore, it will be ideal to optimize the code in order to get the results faster.

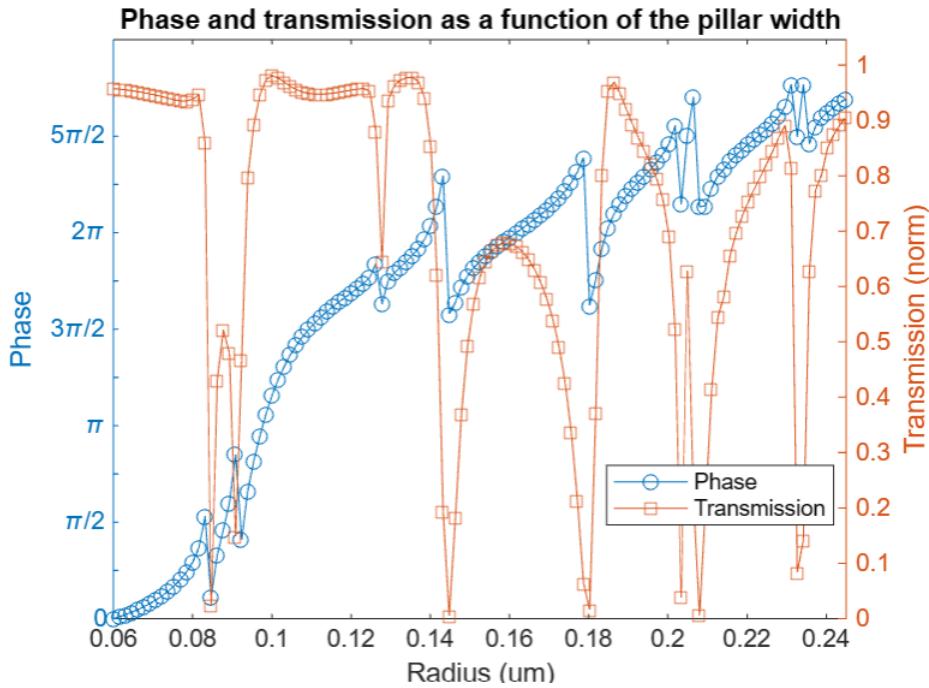


FIGURE 26 – Period of 0.6, instead of 0.47, between the pillar

### 5.3 Examination of Supoptique Logo

Projecting the Supoptique logo in the far field is the ultimate objective (see the logo in Figure 27).

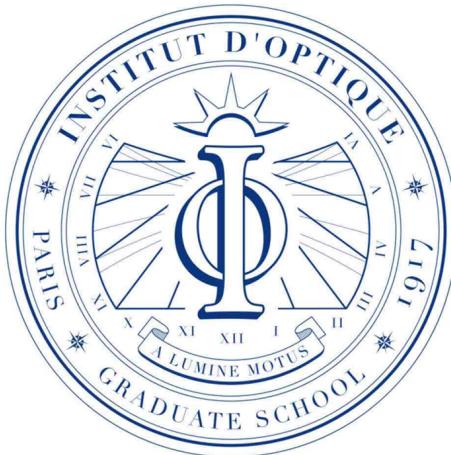


FIGURE 27 – Supoptique Logo

Here is the images we obtain by using the [C] code :

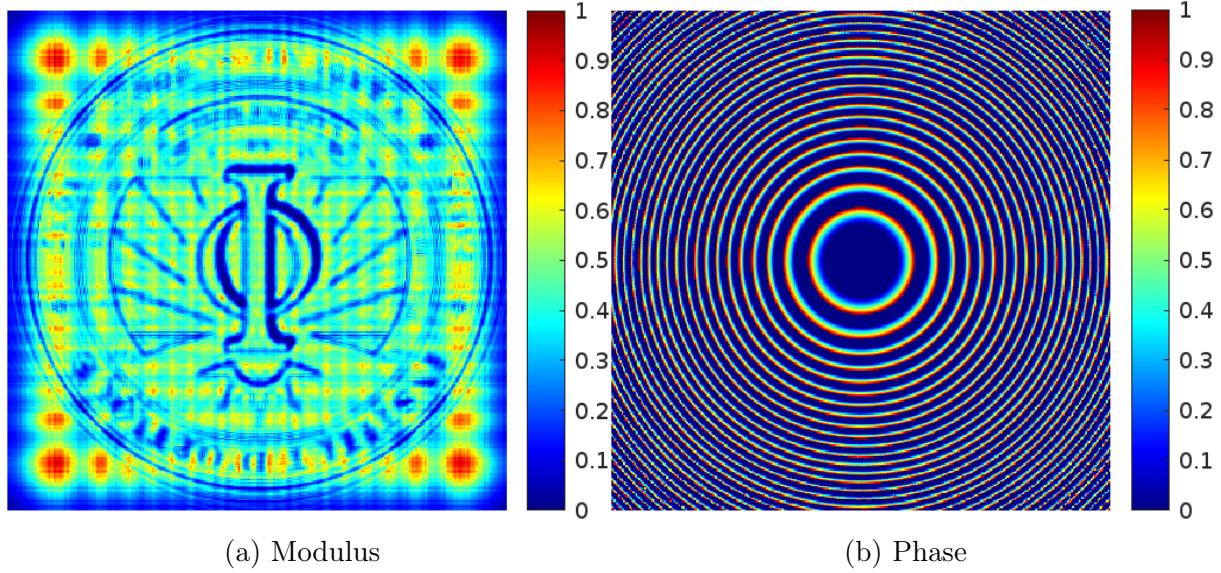


FIGURE 28 – Fresnel transforms : far field propagating fields, lens with a focal length of  $f = 8 \mu\text{m}$  at  $16 \mu\text{m}$  from both the object and the image

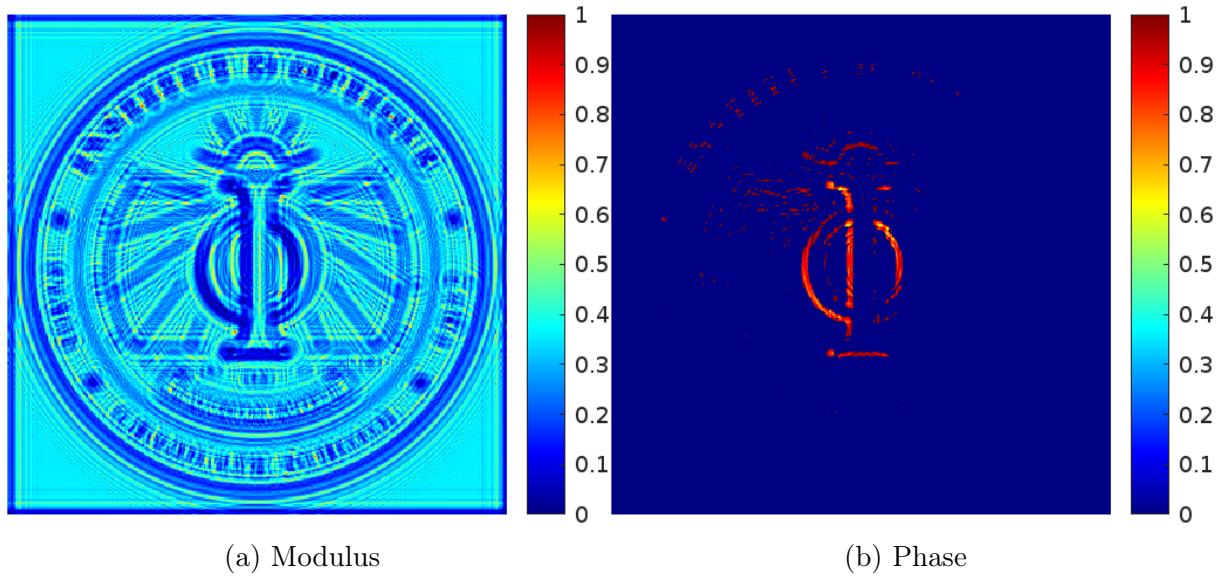


FIGURE 29 – Canonical transforms : far field propagating fields, lens with a focal length of  $f = 8 \mu\text{m}$  at  $16 \mu\text{m}$  from both the object and the image

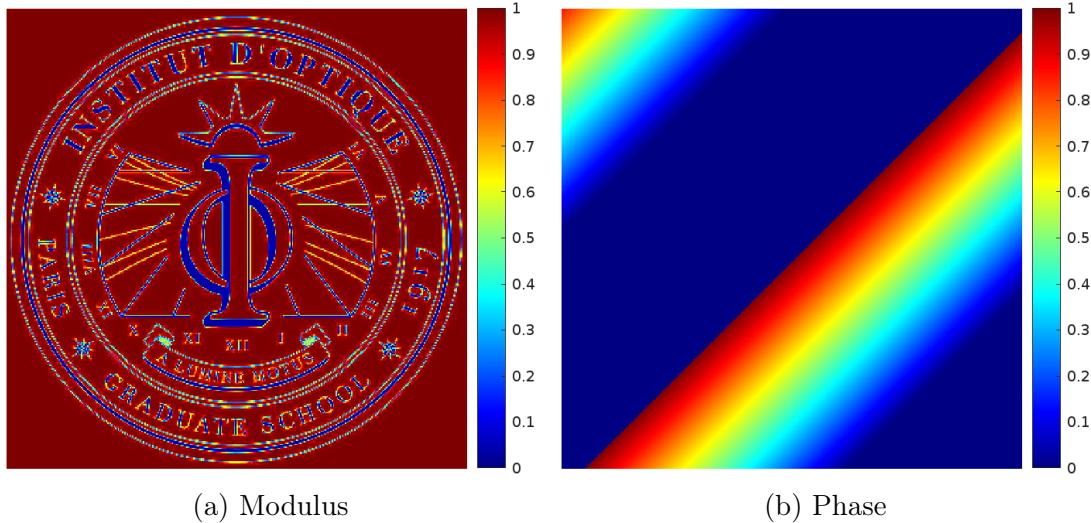


FIGURE 30 – Fresnel propagation using Helmholtz method : far field propagating fields, lens with a focal length of  $f = 8 \mu\text{m}$  at  $16 \mu\text{m}$  from both the object and the image

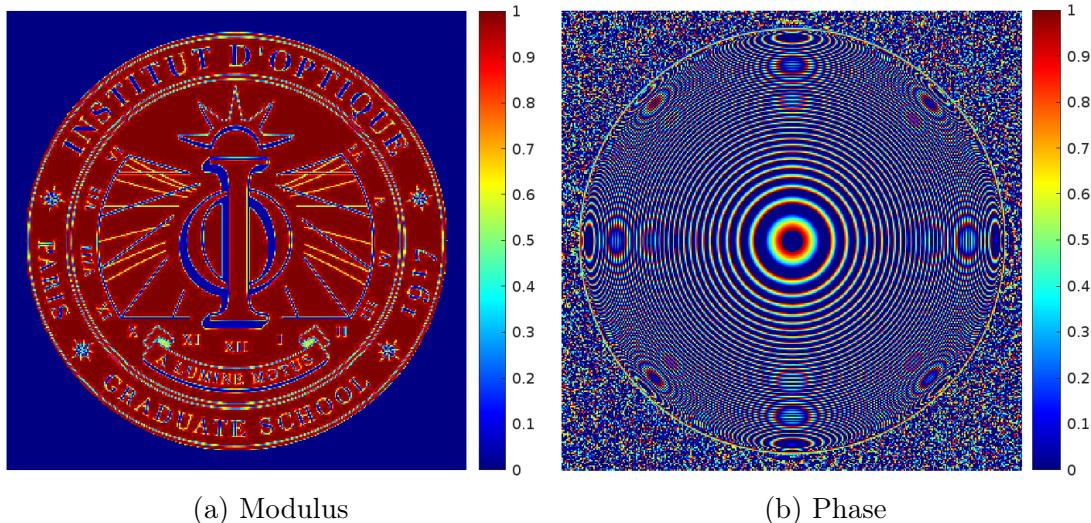


FIGURE 31 – Fraunhofer propagation using Helmholtz method : far field propagating fields, lens with a focal length of  $f = 8 \mu\text{m}$  at  $16 \mu\text{m}$  from both the object and the image

In summary, the choice of propagation method and the precise positioning of the lens relative to the object and image planes critically influence the clarity, resolution, and fidelity of the projected images. The Helmholtz method, particularly with optimal lens placement, shows superior performance in maintaining the integrity and sharpness of the Supoptique logo in the far field projection.

## 6 Conclusion

In summary, this project allowed us to discover meta surfaces and further specialize on meta lenses. Every step was required to comprehend the work we did, and the end product truly satisfied us. It was beneficial to work with Bruno and Benjamin at the beginning to avoid starting off alone. Making a lot of Matlab code at last was a great task.

## A Matlab codes Step 1

```

1 % Création de figures reliant les index n et k à la longueur d'onde ou
2 % à l'énergie pour des angles d'incidences pouvant changer
3 close all
4 % Constantes
5 h = 6.63e-34;
6 c = 3.00e8;
7 e = 1.6e-19;
8 %chargement des données
9 load('Si_Palik.txt','ascii');%changer Si_Palik avec Si02 quand
    nécessaire
10 lambda=Si_Palik(:,1);%nm
11 E=h.*c./lambda;%eV
12 n=Si_Palik(:,2);
13 k=Si_Palik(:,3);
14 figure;
15 plot(lambda/1000,Si_Palik(:,2))%micromètre
16 hold on
17 plot(lambda/1000,Si_Palik(:,3))%micromètre
18 xlabel('Lambda (um)'),ylabel('n, k[1]');
19 legend('n','k');
20 figure;
21 plot(E,Si_Palik(:,2))%eV
22 hold on
23 plot(E,Si_Palik(:,3))%eV
24 xlabel('Photon Energy (eV)'),ylabel('n, k[1]');
25 legend('n','k');
26 theta_=deg2rad(1:10:40)
27 icol=0;for theta=theta_;icol=icol+1 %degré
28 z=n.^2-k.^2-sin(theta).^2;
29 a=sqrt(z.^2+4*n.^2.*k.^2);
30
31 rs=((sqrt(a+z)-sqrt(2)*cos(theta)).^2+a-z)./((sqrt(a+z)+sqrt(2)*cos(
    theta)).^2+a-z);
32 rp=rs.*((sqrt(a+z)-sqrt(2)*sin(theta)*tan(theta)).^2+a-z)./((sqrt(a+z)+
    sqrt(2)*sin(theta)*tan(theta)).^2+a-z);
33
34 figure(3300)
35 h=plot(lambda,rs,'--');set(h,'linewidth',3)
36 set(h,'color',[icol/max(size(theta_)) 1-icol/max(size(theta_)) 0])
37 hold on
38 h=plot(lambda,RP,'-');set(h,'linewidth',3)
39 set(h,'color',[icol/max(size(theta_)) 1-icol/max(size(theta_)) 0])
40
41 xlabel('Lambda (um)'),ylabel('RS, RP (a.u.)');
42 legend('RS','RP');
43 title('RS and RP fonction of lambda with a step of 10 degree from 1 to
    40')
44
45 end
46
47 %%%%%%%%%%%%%%
48
49 % This test aims to launch RCWA 3D solver on simple layered system
50 format longE;
51 clear;

```

```

52 addpath(genpath('/MATLAB Drive/pcode'));
53 % Number of mode in fourier space (MANDATORY)
54 % Recommended : 18
55 OPTION.halfnpw = 18;
56 % TM or TE mode
57 % TM=0 -> TE mode.
58 % TM=1 -> TM mode.
59 OPTION.TM = 1;
60 % Transverse resolution
61 OPTION.FIELDnx = 100; % Number of dots in x direction, to reconstruct E
   and H vector fields
62 OPTION.FIELDny = 100; % Number of dots in y direction, to reconstruct E
   and H vector fields
63 OPTION.FIELD = 0; % Plotting fields
64
65 % Number of frequency (you can define om as a frequency range)
66
67 numom = 1;
68 om = 1e3*1.24/940; % photon energy (hbar omega) in eV
69 numom = 100;
70 if OPTION.halfnpw>0;numom = 10;end;
71 om = linspace(1e3*1.24/1000,1e3*1.24/250 ,numom); % photon energy (hbar
   omega) in eV
72 lambdaall = linspace(.3,.94 ,numom); % photon energy (hbar omega) in eV
73
74 % monitor: layer index to measure the phase (with 0 being the lower one)
75
76 monitor = 0;
77
78 % More parameters
79 % Incident light angle (NEVER SET ZERO)
80
81 theta = 0; % Angle with z-axis
82 phi = 0; % Angle with x-axis
83
84 % Visualization options
85 OPTION.SlideType = [ 'x' , 'z' ]; % Type of slices.
86 OPTION.SlideValue = [0 , -0.05]; % value of slices.
87 OPTION.display_colormap = 1; % 1 -> Display colormap on plots of fields
88
89 %OPTION.BEAM.TYPE='PW';
90 %OPTION.BEAM.TYPE='GA';OPTION.OPTION.BEAM.PARAM=[.05 0 0];
91 disp(" ");
92 disp("STARTING calculation");
93 disp(" ");
94 warning off
95 Nb_sim=1
96 % Cycle over angles
97 for j = 1:Nb_sim
98   disp(' ')
99   disp(['Step: ', num2str(j), '/', num2str(Nb_sim)])
100  %parfor stuff
101  Lambda_all =zeros(1,numom);RR = zeros(1,numom);TT = zeros(1,numom);
102  AA = zeros(1,numom);
103  % RCWA Solving : cycle over frequency
104  %%%
105  for i = 1:size(lambdaall, 2)
106    ARG= struct;

```

```

104 % Wavelength
105 ARG.lambda =lambdaall(i);% 1.24/om(i); % (um)
106 ARG.t = j;
107 % Size of the simulated region
108 ARG.ax = 1.7;
109 % ARG.ax = 2.36;
110 ARG.ax = 2.5;
111 % ARG.ax = 20;
112 ARG.Material_Sub='silicon';
113 % ARG.Material_Sub='SiO2';
114 % Multi coating
115 %ARG.d_ML1=0.06; ARG.Mat_ML1='SiO2_Palik';
116 %ARG.d_ML1=0.08; ARG.Mat_ML1='NewTiO2';
117 %ARG.d_ML2=0.03; ARG.Mat_ML2='NewTiO2';
118 %ARG.d_ML3=0.013; ARG.Mat_ML3='Al2O3_refractiveindex';
119 %DSA
120 ARG.DSA='y';ARG.DSA_Thickness=0.07;
121 % DEVICE creation
122 DEVICE = struct_LayerCoating(ARG);
123 % Visualization
124 % OPTION.SlideType = [ 'z']; % Type of slices.
125 % OPTION.SlideValue = [ -0.05]; % value of slices.
126 % visualization_structure3D(DEVICE, OPTION);
127 if 0==1
128     figure
129         OPTION_3D=OPTION;
130         OPTION_3D.SlideType=[ ]; % type of slides.
131         OPTION_3D.SlideValue=[ ]; % value of slides.
132         OPTION_3D.PlotEdge='no';
133         OPTION_3D.PlotEdge='yes'; OPTION_3D.Plotfast='on';
134         OPTION_3D.epsilon_Visible='no';
135         visualization_structure3D_Perpective(DEVICE, OPTION_3D,
136         get(gcf,'number'));
137         set(gca,'ZLim',[ -1 8])
138         %remove air Visualisation3D_Remove(get(gcf,'number'),[
139             1 0 0]);set(gca,'zlim',[ -3 1])
140 end
141 % RCWA3D on frequency om(i)
142 SOL = RCWA_3D_singlecore(DEVICE, ARG.lambda, theta, phi, OPTION);
143 ;
144 if OPTION.FIELD==1
145     OPTIONPF=OPTION; OPTIONPF.Type='n'
146     PlotField3D(SOL, DEVICE, OPTIONPF);
147 end
148 % ModalPhase(j, :) = SOL.ModalPhase;
149 disp([' Step duration: ' num2str(round(SOL.Time0_All/60, 2)) ' min'])
150 % Retrieving result
151 Lambda_all(i) =ARG.lambda;
152 RR(i) = SOL.RR;
153 TT(i) = SOL.TT;
154 AA_= SOL.AA;
155 AA(i) = AA_(end);
156 end
157 end
158 figure(2)
159 hold on
160 h=plot(Lambda_all,100*TT,'.-b')

```

```

157 h=plot(Lambda_all,100*AA,'.-r')
158 %set(h,'DisplayName',[': ' num2str(OPTION.halfnpw)]);
159 hold on
160 h=plot(Lambda_all,100*RR,'.-g')
161 xlabel('Lambda (um)')
162 ylabel('T, A ,R (%)')
163 legend('T','A','R')
164 box on
165 set(gcf,'color',[1 1 1])
166 figure(20)
167 hold on
168 h=plot(Lambda_all,TT,'.-b')
169 h=plot(Lambda_all,AA,'.-r')
170 %set(h,'DisplayName',[': ' num2str(OPTION.halfnpw)]);
171 hold on
172 h=plot(Lambda_all,RR,'.-g')
173 xlabel('Lambda (um)')
174 ylabel('T, A ,R (1)')
175 legend('T','A','R')
176 box on
177 set(gcf,'color',[1 1 1])
178 %%
179 % Air to Glass vs angle
180 clear Theta_all TT AA RR ARG
181 for TM = 0:1;
182     OPTION.TM=TM;
183 i=0;for theta=0:90;i=i+1;
184     ARG.lambda = 0.5; % (um)
185     ARG.t = j;
186     % Size of the simulated region
187     ARG.ax = 10;%ARG.lambda * 1 / 2; % (um)
188     ARG.Material_Sub='silicon';
189 % ARG.Material_Sub='siO2';
190 % DEVICE creation
191 DEVICE = struct_LayerCoating(ARG);
192 % Visualization
193 % visualization_structure3D(DEVICE, OPTION);
194 % RCWA3D on frequency om(i)
195 %OPTION.BEAM.TYPE='GA'
196 OPTION.BEAM.TYPE='PW';
197 OPTION.OPTION.BEAM.PARAM=[.1 0 0];
198 SOL = RCWA_3D_singlecore(DEVICE, ARG.lambda, theta, phi, OPTION);
199 % Retrieving result
200 Theta_all(i) =theta;
201 RR(i) = SOL.RR; TT(i) = SOL.TT; AA_ = SOL.AA;
202 % AAall(i) = sum(AA_(2:end));
203 % AA(i) = AA_(L);
204 end
205 figure(3)
206 hold on
207 h=plot(Theta_all,100*TT,'.-b')
208 if OPTION.TM==1;set(h,'LineWidth',3);end
209 h=plot(Theta_all,100*RR,'.-g')
210 if OPTION.TM==1;set(h,'linewidth',3);end
211 end
212 xlabel('Theta (deg)')
213 ylabel('T, R (%)')
214 legend('T TE','R TE','T TM','R TM')

```

```
215 box on
216 set(gcf, 'color', [1 1 1])
```

## B Matlab codes Step 2 and 3

```

1 %% Conception of the pinhole
2 diameter_pinhole = 4; % um
3 lambda = 0.940e-6/sqrt(12.95); % m,
4 pp = 0.05*1e-6; % m
5 L = 512; C = 512;
6 Dp = diameter_pinhole * 1e-6 / pp;
7 U0 = zeros(L, C);
8 [x, y] = meshgrid((-C/2):(C/2-1), (-L/2):(L/2-1));
9 rho = sqrt(x.^2 + y.^2);
10 U0(rho < (Dp / 2)) = 1;
11 x_m = x * pp;
12 y_m = y * pp;
13 figure
14 hsurfaces = surface(x_m, y_m, U0);
15 set(hsurfaces, 'FaceColor', 'interp', 'EdgeColor', 'none');
16 set(gcf, 'color', [1 1 1]);
17 daspect([1 1 1]);
18 colorbar;
19 colormap(jet);
20 xlabel('x (m)');
21 ylabel('y (m)');
22 title('Image of a pinhole');
23
24 %% Propagation at z=10 um
25 d_prime = 10e-6;
26 % Using the Fresnel method :
27 Ap_fresnel = Propagation_zero_padd(U0, lambda, d_prime, pp, 1, 0);
28 imshow(abs(Ap_fresnel).^2);
29 % Using the convolutional method :
30 Ap_fr_conv = Propagation_zero_padd_conv(U0, lambda, d_prime, pp);
31 imshow(abs(Ap_fr_conv).^2);
32
33 %% Near field domain
34 % First, we redo the pinhole with the right diameter:
35 diameter_pinhole = 8; % um
36
37 lambda = 0.940e-6; % m,
38 pp = 0.05*1e-6; % m
39 w = 0;
40 L = 512/2; C = L;
41 Dp = diameter_pinhole * 1e-6 / pp;
42 % pinhole
43 U0 = zeros(L, C);
44 [x, y] = meshgrid((-C/2):(C/2-1), (-L/2):(L/2-1));
45 rho = sqrt(x.^2 + y.^2);
46 U0(rho < (Dp / 2)) = 1;
47 t_a_prime = U0;
48 % For z <= 10 m
49 d_prime_values = linspace(1e-6, 10e-6, 1000);
50 % For z <= 60 m
51 %d_prime_values = linspace(1e-6, 60e-6, 1000);
```

```

52 phase_circles = ones(L, length(d_prime_values));
53 for i = 1:length(d_prime_values)
54     d_prime = d_prime_values(i);
55         % With zero-padd
56     Ap = Propagation_zero_padd(t_a_prime, lambda, d_prime, pp, 1, 0);
57         % Without zero-padd
58     %Ap = Propagation_zero_padd(t_a_prime, lambda, d_prime, pp, 0, 0);
59     phase_circles(:, i) = abs(Ap(L/2, :)).^2;
60 end
61 phase_circles = phase_circles/max(phase_circles,[],"all");
62 y_pixels = ((1:L) - L/2) * pp * 1e6;
63 figure;
64 imagesc(d_prime_values * 1e6, y_pixels, phase_circles);
65 colormap(jet);
66 colorbar;
67 axis xy;
68 xlabel('d' '(um)');
69 ylabel('y (um)');
70
71 %% In plane phase shift after the lens
72 [x, y] = meshgrid((-C/2):(C/2-1), (-L/2):(L/2-1));
73 x_m = x * pp;
74 y_m = y * pp;
75 f = 8e-6;
76 t_a_prime=complex_transmittance(lambda,L, C, pp,f);
77 d_prime = 0;
78 Ap = Propagation_zero_padd(t_a_prime, lambda, d_prime, pp, 1, 0);
79 phase_circles = angle(Ap);
80 phase_circles = phase_circles/max(phase_circles,[],"all");
81 % imshow(phase_circles);
82 % colormap(jet);
83 figure
84 hsurfaces = surface(x_m, y_m, phase_circles);
85 set(hsurfaces, 'FaceColor', 'interp', 'EdgeColor', 'none');
86 set(gcf, 'color', [1 1 1]);
87 daspect([1 1 1]);
88 colorbar;
89 colormap(jet);
90 xlabel('x (m)');
91 ylabel('y (m)');
92 zlabel('Intensity');
93
94 %% Vertical central-cut of the propagating fields, focal length of f=8um
95 at 940nm
95 L = 512; C = L;
96 U0 = ones(L, C);
97 [x, y] = meshgrid((-C/2):(C/2-1), (-L/2):(L/2-1));
98 x_m = x * pp;
99 y_m = y * pp;
100 t_a_prime = complex_transmittance(lambda, L, C, pp, f);
101 d_prime_values = linspace(1e-6, 16e-6, 1000); % Range of d_prime values
102 phase_circles = ones(L, length(d_prime_values));
103 for i = 1:length(d_prime_values)
104     d_prime = d_prime_values(i);
105     Ap = Propagation_zero_padd(t_a_prime, lambda, d_prime, pp, 1, 0);
106         % Modulus :
107     phase_circles(:, i) = abs(Ap(L/2, :)).^2;
108         % Phase :

```

```

109     % phase_circles(:, i) = angle(Ap(L/2, :));
110 end
111 phase_circles = phase_circles/max(phase_circles,[],"all");
112 y_pixels = ((1:L) - L/2) * pp * 1e6;
113 figure;
114 imagesc(d_prime_values * 1e6, y_pixels, phase_circles);
115 colormap(jet);
116 colorbar;
117 axis xy;
118 xlabel('d', '(um)');
119 ylabel('y (um)');
120
121 %% Helmholtz method: vertical central-cut at the focal point
122 clear all
123 lambda = 0.940e-6; %m
124 C = 256;
125 L = C;
126 k = 2*pi/lambda;
127 pp = 0.05e-6; %m
128 [x, y] = meshgrid((-C/2):(C/2-1), (-L/2):(L/2-1));
129 x_m = x * pp;
130 y_m = y * pp;
131 f = 8e-6;
132 t_a_prime=complex_transmittance(lambda,L, C, pp,f);
133 d_prime = f;
134 [UH_RS,UH_F,UH_Fraun] = propagation_H(t_a_prime,lambda,d_prime,pp);
135 phase_circles = angle(UH_F);
136 phase_circles = phase_circles/max(phase_circles,[],"all");
137 figure
138 y_pixels = ((1:L) - L/2) * pp * 1e6; % Conversion de l'index des pixels
    en micromètres
139 figure;
140 imagesc(y_pixels, y_pixels, phase_circles); % Conversion de m en m
    pour l'affichage
141 colormap(jet);
142 colorbar;
143 axis xy;
144 xlabel('d', '( m )');
145 ylabel('y ( m )');

146
147
148 %%
149 %%
150 %%
151 %%
152 %% Functions
153     % Fresnel propagation
154 function Ap=Propagation_zero_padd(tim,lambda,z,pp,zero_padd,w)
155 % Usage :Ap=PropagationZP(tim,lambda,z,pp,zero_padd,w);
156 % Goal : Simulation of free propagation of coherent light from the plane
    1 of transmittance tim to a plane 2 distant of
157 % z. Depending on the z value the program uses impulse response hz or
    transfer function Hz
158 % Note : the uniform phase term exp(-i*2*pi*z:lambda) can be considered
    to be equal to 1 without changing the results in some cases.
159 % input parameters :
160 % tim : transmittance of the plane
161 % lambda : wavelength

```

```

162 % z : propagation distance along the optical axis
163 % pp : pixel pitch
164 % zero_padd
165 % if zero_pad=0 no zero padding
166 % if zero_pad=1 zero padding and output image of size equal to the input
167 % image
168 % if zero_pad=2 zero padding and output image of size equal to two times
169 % the
170 % input image size
171 % w=0, if we do not want warnings to be written in the command window
172 % w=1, warnings are written in the command window
173
174 [L,C]=size(tim);
175 if L~=C
176     if w==1,fprintf('columns number is not equal to rows number \n
177     Warning : to select the method "impulse response" or "transfer
178     function" the number of rows is used\n'),end
179 end
180 zNS=L*(pp^2)/lambda;% z0 is computed without zero padding because in
181 % that case what interests us is the in field of view reconstruction
182
183 if zero_padd==0
184     Hz=FresnelFunctionFT(L,C,z,lambda,pp);% Matlab coordinate origin
185     if w>=1,fprintf('use of the transfer function for the
186     propagation computation\n'), end
187     fftt=fft2(tim);% % Matlab coordinate origin
188     Ap=ifft2(Hz.*fftt);
189 else % (zero_padd>=1)
190
191     tim2=ones(2*L,2*C)*mean(tim(:));tim2(L-L/2:L+L/2-1,C-C/2:C+C/2-1)=
192     tim;
193     if w==2,figure;imshow(tim2,[],'XData',[1:size(tim2,2)]*1000*pp,
194     'YData',[1:size(tim2,1)]*1000*pp);colorbar;title(' 0 padded image');
195     axis on;xlabel('x(mm)');ylabel('y(mm)'), end
196     Hz=FresnelFunctionFT(2*L,2*C,z,lambda,pp);
197     if w>=1,fprintf('use of the transfer function for the
198     propagation computation\n'), end
199 end
200 %%%%%%%%
201 function hz=FresnelFunction(L,C,z,lambda,pp);
202 %% Fresnel function
203 % input parameters :
204 % L,C : number of lines and columns of the Fresnel function (
205 % output parameter)
206 % z : distance of propagation in meter
207 % lambda : wavelength in meter
208 % pp : pixel pitch in meter
209 % output parameter :
210 % hz : matrix of L lines and C columns of complexe values of the

```

```

    Fresnel function

210
211     %% Usage : hz=FresnelFunction(L,C,z,lambda,pp);
212     hz=zeros(L,C);
213     [Ki,Et]=meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1)); % units in pix-1
214     Ki=Ki*pp;% units in m
215     Et=Et*pp;% units in m
216     c1=1/lambda/z/1i;
217     c2=1i*pi/lambda/z;
218     hz=ifftshift(exp(1i*2*pi*z/lambda)*c1*exp(c2*( Ki.^2+Et.^2) ) );
219     % the phase term exp(i*2*pi*z/lambda ) can be omitted in some cases
220     % the phase c2*( Ki.^2+Et.^2) is positive , it is the sign
221     convention considered in the lectures
222
223 end
224 %%%%%%
225 function Hz=FresnelFunctionFT(L,C,z,lambda,pp);
226     %% Fresnel function Fourier transform
227     % input parameters :
228     % L,C : number of lines and columns of the Fresnel function (
229     output parameter)
230     % z : distance of propagation in meter
231     % lambda : wavelength in meter
232     % pp : pixel pitch in meter
233     % output parameter :
234     % Hz : matrix of L lines and C columns of complexe values of the
235     Fresnel
236     % function Fourier Transform
237
238     %% Usage : Hz=FonctionDeFresnelFT(L,C,z,lambda,pp);

239 Hz=zeros(L,C);
240 [u,v]=meshgrid((-floor(C/2):ceil(C/2)-1)/C,(-floor(L/2):ceil(L/2)-1)/L);
241     % for images with an even number of columns and rows, the
242     geometrical center is C/2+1,L/2+1
243     % for images with an odd number of columns and rows, the
244     geometrical center
245     % is ceil (C/2+1),ceil(L/2+1)
246     % The frequencies varies from -0.5 to 0.5-frequency sampling
247 step, There are LxC
248     % samples. The unit is pix-1
249 u=u/pp; % unit : m-1
250 v=v/pp; % unit : m-1

251 ipilambdaz=-i*pi*lambda*z;
252 Hz=ifftshift(exp(1i*2*pi*z/lambda)*exp(ipilambdaz*( u.^2+v.^2) ))
253 );
254 end
255 end
256
257 %% Fresnel propagation by using the convolutional method
258 function Ap_Fc = Propagation_zero_padd_conv(tim,lambda,z,pp);
259 % Propagation selon la méthode de Fresnel sous sa forme canonique
260
261 [L,C] = size(tim);
262 Ap_Fc = zeros(L,C);

```

```

257 [X,Y] = meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1));
258 Xp = X/(C);
259 Yp = Y/(C);
260 X = X.*pp;
261 Y = Y.*pp;
262 Xp = Xp.*pp.*lambda.*z;
263 Yp = Yp.*pp.*lambda.*z;
264 c1 = (1/lambda/z/1i).*exp(1i*2*pi*z/lambda);
265 c2 = 1i.*pi/lambda/z;
266 %Ap_Fc = c1.*exp(c2.*(Xp.^2 + Yp.^2)).*fftshift(fft2(tim.*exp(c2.*(X.^2
267 + Y.^2)))); 
267 Ap_Fc = ifftshift(c1 .* exp(c2 .* (X.^2 + Y.^2))) .* fft2(tim .* exp(c2
268 .* (X.^2 + Y.^2)));
268 Ap_Fc = ifft2(Ap_Fc);
269 end
270
271 % Lens
272 function tL = complex_transmittance(lambda, L, C, pp, f)
273 X = -C/2:C/2-1;
274 Y = -L/2:L/2-1;
275 [x, y] = meshgrid(X, Y);
276 Delta = (x.^2 + y.^2) .* (pp^2 / (2*f));
277 Phi = (2*pi/lambda) .* Delta;
278 tL = exp(-1i .* Phi);
279 end
280
281 %% Propagation using Helmholtz method :
282 function [ M ] = MaskDisk( L,C,Dp,pp );
283 % Usage :[ M ] = MaskDisk( L,C,Dp,pp );
284 % Goal : Create a binary mask of size LxC with a disk of diameter D.
285 % Its center is located in (L/2+1,C/2+1) if L,C are even.
286 % Input Parameters :
287 % - Dp : Disk diameter (type : double, scalar, unit=meter)
288 % - L,C : size of the mask (type : double, scalar, unit=pixel)
289 % Output Parameters :
290 % - M : Mask (type : double, size : LxC)
291 % pp : pixel pitch
292
293 x=-C/2:C/2-1;
294 y=-L/2:L/2-1;
295 x=x.*pp;y=y.*pp;
296 [X,Y]=meshgrid(x,y);
297 rho=sqrt(X.^2+Y.^2);
298 M=double(rho<(Dp/2));
299 end
300
301 function [Ap_Fc]=propagation_Fc(tim,lambda,z,pp);
302 % Propagation selon la mÃ thode de Fresnel sous sa forme canonique
303
304 [L,C] = size(tim);
305 Ap_Fc = zeros(L,C);
306 [X,Y] = meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1));
307 Xp = X/(C);
308 Yp = Y/(C);
309 X = X.*pp;
310 Y = Y.*pp;
311 Xp = Xp.*pp.*lambda.*z;
312 Yp = Yp.*pp.*lambda.*z;

```

```

313 c1 = (1/lambda/z/1i).*exp(1i*2*pi*z/lambda);
314 c2 = 1i.*pi/lambda/z;
315 Ap_Fc = c1.*exp(c2.*(Xp.^2 + Yp.^2)).*fftshift(fft2(tim.*exp(c2.*(X.^2 +
316 Y.^2)))); 
316 end
317 function [Ap_H_RS,Ap_H_F,Ap_H_Fraun]=propagation_H(tim,lambda,z,pp)
318 % Propagation selon la méthode de Helmholtz
319
320 [L,C] = size(tim);
321 Ap_H_RS = zeros(L,C);
322 Ap_H_F = zeros(L,C);
323 Ap_H_Fraun = zeros(L,C);
324 [X,Y] = meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1));
325 X = X.*pp;
326 Y = Y.*pp;
327 Xp = X/(C);
328 Yp = Y/(C);
329 Xp = Xp.*pp.*lambda.*z;
330 Yp = Yp.*pp.*lambda.*z;
331
332 % Définition de la solution des équations d'Helmholtz
333
334 k0 = 2*pi/lambda;
335 kx = (2*pi/lambda).*(X./z);
336 ky = (2*pi/lambda).*(Y./z);
337 EH = tim.*exp(1i.*z.*sqrt(k0.^2 - kx.^2 - ky.^2));
338
339 % propagation suivant les différentes méthodes :
340
341 % Rayleigh Sommerfeld
342
343 R = sqrt(X.^2+Y.^2+z.^2);
344 c1 = z./(lambda*1i*(R.^2));
345 c2 = 1i*2*pi/lambda;
346 hzRS = ifftshift(c1.*exp(c2.*R));
347 TF_RS = fft2(hzRS);
348 Ap_H_RS = ifft2(fft2(EH).*TF_RS);
349
350 % Fresnel
351
352 % c3 = (1 / (lambda * z * 1i)) .* exp(1i * 2 * pi * z / lambda);
353 % c4 = 1i * pi / (lambda * z);
354 % pre_transform = tim .* exp(c4 .* (X.^2 + Y.^2));
355 % hzF = fftshift(fft2(pre_transform));
356 % Ap_H_F = c3 .* ifft2(hzF);
357 c3 = (1 / (lambda * z * 1i)) .* exp(1i * 2 * pi * z / lambda); %
Common phase factor
358 c4 = 1i * pi / (lambda * z); % Scaling for the quadratic phase factor
359 pre_transform = EH .* exp(c4 .* (X.^2 + Y.^2)); % Modifying the input
field with quadratic phase
360 hzF = fftshift(fft2(pre_transform)); % Applying fft to the phase-
modified field
361 % Ap_H_F = c3 .* ifft2(hzF); % Applying inverse fft and scaling with c3
362 Ap_H_F = ifft2(hzF); % Applying inverse fft and scaling with c3
363
364 % Fraunhofer
365
366 c5 = (1/lambda/z/1i).*exp(1i*2*pi*z/lambda);

```

```

367 c6 = 1i.*pi/lambda/z;
368 Ap_H_Fraun = ifftshift(c5.*exp(c6.*Xp.^2 + Yp.^2)).*fftshift(fft2(EH));
369 ;
370 Ap_H_Fraun = ifft2(Ap_H_Fraun);
371 end
372
372 function [Ap_F]=propagation_F(tim,lambda,z,pp);
373 % Propagation selon la mÃ thode de Fresnel (approximation de la mÃ
373 thode de Rayleigh Sommerfeld)
374
375 [L,C] = size(tim);
376 hzF = zeros(L,C);
377 [X,Y] = meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1));
378 X = X.*pp;
379 Y = Y.*pp;
380 c1 = (1/lambda/z/1i).*exp(1i*2*pi*z/lambda);
381 c2 = 1i.*pi/lambda/z;
382 hzF = ifftshift(c1.*exp(c2.*(X.^2 + Y.^2)));
383 TF_F = fft2(hzF);
384 Ap_F = ifft2(fft2(tim).*TF_F);
385 end
386 function [Ap_Fraun]=propagation_Fraun(tim,lambda,z,pp)
387
388 [L,C] = size(tim);
389
390 Ap_Fraun = zeros(L,C);
391 [X,Y] = meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1));
392 Xp = X/(C);
393 Yp = Y/(C);
394 X = X.*pp;
395 Y = Y.*pp;
396 Xp = Xp.*pp.*lambda.*z;
397 Yp = Yp.*pp.*lambda.*z;
398 c1 = (1/lambda/z/1i).*exp(1i*2*pi*z/lambda);
399 c2 = 1i.*pi/lambda/z;
400 Ap_Fraun = c1.*exp(c2.*(Xp.^2 + Yp.^2)).*fftshift(fft2(tim));
401 end
402
403 function [Ap_RS]=propagation_RS(tim,lambda,z,pp);
404 % Propagation selon la mÃ thode de Rayleigh Sommerfeld
405
406 [L,C] = size(tim);
407 hzRS = zeros(L,C);
408 [X,Y] = meshgrid((-floor(C/2):ceil(C/2)-1),(-floor(L/2):ceil(L/2)-1));
409 X = X.*pp;
410 Y = Y.*pp;
411 R = sqrt(X.^2+Y.^2+z.^2);
412 c1 = z./(lambda*1i*(R.^2));
413 c2 = 1i*2*pi/lambda;
414 hzRS = ifftshift(c1.*exp(c2.*R));
415 TF_RS = fft2(hzRS);
416 Ap_RS = ifft2(fft2(tim).*TF_RS);
417 end

```

## C Matlab codes Step 4

```

1 %% Propa logo supop
2
3 clear all
4 close all
5
6 lambda = 0.940e-6; %m
7 k = 2*pi/lambda;
8 pp = 0.05e-6; %m
9
10
11 tim = imread('Logo_Supop.png');
12 tim = rgb2gray(tim);
13 % imagesc(tim)
14
15 tim = imresize(tim, [699 699]);
16 imshow(tim);
17
18 [L,C] = size(tim);
19 [x, y] = meshgrid((-C/2):(C/2-1), (-L/2):(L/2-1));
20 x_m = x * pp;
21 y_m = y * pp;
22 f = 8e-6;
23 d_prime = 2*f;
24 Dp = 600;
25
26
27 tim_propa=Propagation_zero_padd(tim, lambda, d_prime, pp, 1, 0);
28 % tim_propa = Propagation_zero_padd_conv(tim, lambda, d_prime, pp);
29 tim = double(tim);
30
31 tL=complex_transmittance(lambda,L, C, pp,f);
32 t_lens_transmited=tim_propa.*tL;
33 t_lens=t_lens_transmited;
34
35 Ap_propa=Propagation_zero_padd(t_lens, lambda, d_prime, pp, 1, 0);
36 % Ap_propa = Propagation_zero_padd_conv(t_lens, lambda, d_prime, pp);
37
38 phase_circles = abs(Ap_propa).^2;
39 % phase_circles = angle(Ap_propa);
40 phase_circles = phase_circles/max(phase_circles,[],"all");
41
42 imshow(phase_circles);
43 colormap(jet); colorbar;
44
45
46%%%%%%%%%%%%%%%
47
48 function [Layers, current_z, ind_current_layer,nz]=ADD_SStructured_Layer2
    (Layers, current_z, ind_current_layer,nz,epsSi,epsSi02,ALLcenter,
     ALLcentery,ARG)
49 % A layer of structuration
50 % Parameters:
51 % into ARG
52 % Mandatory PARAMETERS:
53 % % definition of the centers of all pillars as rectanglura grid:
54 %ALLcenter=[-1 0 1] ;
55 %ALLcenter=[ 0 ] ;
56 %ALLcentery= ALLcenter ;

```

```

57 %OR like a list !
58
59 Nduplicate=(size(ALLcenter,1));
60
61 if ~isempty(ALLcentery)
62
63     Nduplicate=max(size(ALLcenter));
64     Nduplicatey=max(size(ALLcentery)); %ALLcentery is a 1 vector. We are
65     dealing with rectangular pillar grids
66     for icenterx= 1:Nduplicate
67         for icentery= 1:Nduplicatey
68             %      centre cercle=[x_min + (x_max - x_min)/2, y_min + (y_max -
69             % y_min)/2 ]+[ALLcenter(icenterx), ALLcenter(icentery)];
70             centre cercle(icentery,icenterx,:)=[ALLcenter(icenterx),
71             ALLcenter(icentery)];
72         end
73     end
74 else
75     Nduplicatey=1; % provide pillar position and other parameters as a
76     list
77     for icenterx= 1:Nduplicate
78         %      centre cercle=[x_min + (x_max - x_min)/2, y_min + (y_max -
79         % y_min)/2 ]+[ALLcenter(icenterx), ALLcenter(icentery)];
80         centre cercle(1,icenterx,:)=[ALLcenter(icenterx,1), ALLcenter(
81             icenterx,2)];
82     end
83 end
84
85 if isfield(ARG,'Pillar_Rotation'); Pillar_Rotation=ARG.Pillar_Rotation;
86 else; Pillar_Rotation=zeros(Nduplicatey,Nduplicate);
87 end %rotate the pillars (matrix of rotation)
88
89 if isfield(ARG,'FINERGRID'); FINERGRID=ARG.FINERGRID;
90 else; FINERGRID=1;
91 end %not used but do not remove
92
93 %other parameters (default if not specified)
94 % Pillar width of the bottom circle of the pillar (the largest). The
95 % light is comming from the top
96 % if Pillar width2 not specified it is equal to width
97
98 if isfield(ARG,'PillarWidth'); PillarWidth=ARG.PillarWidth; else;
99     PillarWidth=zeros(Nduplicatey,Nduplicate)+.2; end
100 if isfield(ARG,'PillarWidth2'); PillarWidth2=ARG.PillarWidth2; else;
101     PillarWidth2=zeros(Nduplicatey,Nduplicate)+.2; end
102
103 % SPECIFIC Pillars
104
105 if isfield(ARG,'Height_pillar'); Height_pillar=ARG.Height_pillar; else;
106     Height_pillar=0.5;      end
107 if isfield(ARG,'angle_pillar'); angle_pillar=ARG.angle_pillar; else;
108     angle_pillar=0+zeros(Nduplicate,Nduplicate);      end
109
110 %deg / SLicing vertical des pilliers (number of cuts)
111
112 if isfield(ARG,'Npillar');
113     Npillar=ARG.Npillar;
114 else;

```



```

156     OUTPUT=-1;
157     return;
158 end
159
160 if max(size(PillarWidth2,1)) ~=Nduplicatey
161     disp('size(PillarWidth2) does not match with ALLcenter
definition')
162     OUTPUT=-1;
163     return;
164 end
165 end
166
167 %create some variable for the parfor
168
169 current_z_(1)=current_z;ind_current_layer_0=ind_current_layer;for
iNpillar=1:(1+Npillar)
170
171 if iNpillar<=N_overetch/2
    d_(iNpillar) = Top_Pillar_Overetch;
173 elseif iNpillar>(1+Npillar-N_overetch/2)
    d_(iNpillar) = Bot_Pillar_Overetch;
175 else
    d_(iNpillar) = (Height_pillar-Top_Pillar_Overetch-
Bot_Pillar_Overetch)/(1+Npillar-N_overetch); %hauteur des pilliers
end
177 ind_current_layer = ind_current_layer + 1; nz=[nz round(100/(1+
Npillar)) ];
179 current_z_(iNpillar+1) = current_z_(iNpillar) -d_(iNpillar);
180 end
181
182 %for rand Pillars
183
184 if ISRAND==1 %reported here for parfor reason
    Widthextracted=[];Widthextracted_bottom = [];Nshape=0; ALLf=[];
    %if strcmp(ARG.Pillar_type, "circleRand")
    N=(1+Npillar);rL=1;% Warning corelation lenth are in rL units
    for icenterx= 1:Nduplicate
        for icentery= 1:Nduplicate
            % R0=PillarWidth(icenterx,icentery)/2; %/2 because it is a
            radius
            %[f,x,y] =rsgene2D_Loumi(N,rL,RMS,Lradial*R0*2*pi,Lvertical*
            Height_pillar,Seeds_maps(icenterx,icentery)); %rescaled correlation
            length
            [f,x,y] =rsgene2D(N,rL,RMS,Lradial,Lvertical*Height_pillar);
            ALLf(icentery,icenterx,:,:)=f;
        end
    end
196
197 Nshape=N; % DO NOT CHANGE Nshape=N !
198 %buid a circle as polynoms (radius is unity)
199 P=[];
200 for thetha_grid=1:Nshape
    P=[P; [ cos(2*pi*thetha_grid/Nshape) sin(2*pi*thetha_grid/Nshape
)]];
201 end
202
203 if PLOTALL==1;figure(66);hold off;end
%Widthextracted=[];Widthextracted_bottom = [];

```

```

206 %end
207 end
208
209 %end rand Pillars
210 Layers_(1)=Layers(1);
211 for iNpillar=1:(1+Npillar)
212
213 %eps_outpillar = epsSi02;%outer material _pillier;
214 if iNpillar<=N_overetch/2
215     eps_outpillar = Top_Pillar_Overetch_eps; %overetch
216 elseif iNpillar>(1+Npillar-N_overetch/2)
217     eps_outpillar = Bot_Pillar_Overetch_eps;%overetch
218 else
219     eps_outpillar = epsSi02;%outer material _pillier;
220 end
221 eps_square = eps_outpillar; %outer square
222 Array_of_circle=[];Array_of_Polygon=[];
223 %Pillar with side angle
224 Array_of_circle=[];    Array_of_Polygon=[];
225 if strcmp(ARG.Pillar_type, "perfectcircle")
226     for icenterx= 1:Nduplicate
227         for icentery= 1:Nduplicatey
228             %centrecercle=[x_min + (x_max - x_min)/2, y_min + (y_max
229 - y_min)/2 ]+[ALLcenter(icenterx), ALLcenter(icentery)];
230             %centrecercle=[ALLcenter(icenterx), ALLcenter(icentery)
231 ];
232             centrecercle_=[centrecercle(icentery,icenterx,1),
233 centrecercle(icentery,icenterx,2)];
234             radiusofcircle=(PillarWidth(icentery,icenterx)-2*
235 Height_pillar*tan(2*pi/360*angle_pillar(icentery,icenterx))*((Npillar
236 -iNpillar+1)/(Npillar+1))/2; %bug
237             Array_of_circle = [Array_of_circle Disk_2D(
238 radiusofcircle, centrecercle_,epsSi)];
239             end
240         end
241     elseif strcmp(ARG.Pillar_type, "none")
242         %do noting
243     else
244         for icenterx= 1:Nduplicate
245             for icentery= 1:Nduplicatey
246                 if strcmp(ARG.Pillar_type, "circle")
247                     Theta=0:(2*pi)/51:2*pi;multfact= 1;
248                 elseif strcmp(ARG.Pillar_type, "losange")
249                     Theta=0:(pi/2):2*pi;multfact= 1;
250                 elseif strcmp(ARG.Pillar_type, "hexagon")
251                     Theta=0:(pi/3):2*pi;multfact= 1;
252                 elseif strcmp(ARG.Pillar_type, "octogon")
253                     Theta=0:(pi/4):2*pi; multfact= 1;
254                 elseif strcmp(ARG.Pillar_type, "circSquare")
255                     Pillar_circSquareangle= ARG.Pillar_circSquareangle(
icentery,icenterx);
256                     Theta_0=(pi/4):(pi/2):(2*pi);
257                     cutangle=-pi/Pillar_circSquareangle:(pi/
258 Pillar_circSquareangle)/11:pi/Pillar_circSquareangle;
259                     Theta=[];for icutangle=1:max(size(Theta_0))
260                         Theta=[Theta Theta_0(icutangle)+cutangle];
261                     end
262                     multfact= 1/cos(pi/4-pi/Pillar_circSquareangle); %

```

```

increase the radius to get the square part at the good dimention
givent by pillarwidth
256    end
257    % centre cercle=[x_min + (x_max - x_min)/2, y_min +
258    % centre cercle=[ALLcenter(icenterx), ALLcenter(icentery)];
259    % centre cercle=[ALLcenter(icenterx), ALLcenter(
260    icentery)];
261    centre cercle_=[centre cercle(icentery,icenterx,1),
262    centre cercle(icentery,icenterx,2)];
263    radiusofcircle=(PillarWidth(icentery,icenterx)-2*
264 Height_pillar*tan(2*pi/360*angle_pillar(icentery,icenterx))*((Npillar-
265 -iNpillar+1)/(Npillar+1))/2; %bug
266    % Building the radius:
267    RRR = [];
268    tmpR = 0;
269    Nb_Period=ARG.Pillar_NbPeriod(icentery,icenterx);
270    if Nb_Period>0
271        PeriodSi = (ARG.PillarInnerWidth(icentery,icenterx)
272 +2*Height_pillar*tan(2*pi/360*angle_pillar(icentery,icenterx))*((Npillar-
273 -iNpillar+1)/(Npillar+1))/2;%bug
274        PeriodSi02 = PillarWidth(icentery,icenterx)/2-
275 PeriodSi-Height_pillar*tan(2*pi/360*angle_pillar(icentery,icenterx))
276 *((Npillar-iNpillar+1)/(Npillar+1));%bug
277        for ir = 1:Nb_Period
278            if ir == 1
279                RRR = [RRR (tmpR +PeriodSi)];
280                tmpR = tmpR + PeriodSi;
281                RRR = [RRR (tmpR + PeriodSi02)];
282                tmpR = tmpR + PeriodSi02;
283            else
284                RRR = [RRR (tmpR +PeriodSi)];
285                tmpR = tmpR + PeriodSi;
286                RRR = [RRR (tmpR + PeriodSi02)];
287                tmpR = tmpR + PeriodSi02;
288            end
289        end
290    else % juste one regular pillar
291        RRR= radiusofcircle;
292    end
293    RRR = multfact*RRR;
294    RRR2 = PillarWidth2(icentery,icenterx)/PillarWidth(
icentery,icenterx)*RRR; %this is a rectangular shape
295    for ir = 1:(max(size(RRR)))
        tmp = centre cercle_+[ RRR(ir)*cos(Theta)', RRR2(ir)*
sin(Theta)'];
        if ISRAND==1
            Pm=zeros(Nshape,2);
            %linear interpolation of the shapes (prior to
distortion)
            [pt,dudt,foft] = interparc(0.1,RRR(ir)*cos(Theta),
RRR2(ir)*sin(Theta),'linear')
            P=foft(linspace(0,1,Nshape));
            %distortion fields
            Rescale_distortion_to_Perimeter=perimeter(polyshape(P));
            if iNpillar ==1 [distortion_field(
icentery,icenterx,:,:)] =rsgene2D(N,rL,RMS,Lradial/
Rescale_distortion_to_Perimeter,Lvertical/Height_pillar);
            end

```

```

296         ishape=1;
297         VP=[P(ishape,:)-P(end,:)];VP=[VP(2) -VP(1)];VP=
298 VP/norm(VP);% normal direction
299         Pm(ishape,:)= P(ishape,:)+ALLf(icentery,icenterx
300 ,iNpillar,ishape)*VP;
301         for ishape=2:Nshape
302             VP=[P(ishape,:)-P(ishape-1,:)];VP=[VP(2) -VP
303 (1)];VP=VP/norm(VP);% normal direction
304             Pm(ishape,:)= P(ishape,:)+distortion_field(
305 icentery,icenterx,iNpillar,ishape)*VP;
306             end
307             Pm(1,:)= 0.5*(Pm(2,:)+Pm(end,:));
308             %plot(P(:,1),P(:,2),'.');hold on
309             %plot(Pm(:,1),Pm(:,2),'.')
310             tmp = centre cercle+Pm;
311             end
312 %rotate pillars:
313 %polygoneC_coordinate = Rot2D(polygoneC_coordinate, pi/2, [0, 0]);
314             if (Pillar_Rotation(icentery,icenterx))>0;
315                 tmp = Rot2D(tmp, Pillar_Rotation(icentery,
316 icenterx), centre cercle_);
317             end
318             if Nb_Period>0
319                 if mod(ir,2)==1 % Starting with eps_outpillar
320                     material = eps_outpillar;
321                 else
322                     material = epsSi ;
323                 end
324             else % juste one regular pillar
325                 material = epsSi ;
326             end
327             Array_of_Polygon=[Array_of_Polygon Polygon2D(tmp,
328 material)];
329             end
330             if Nb_Period>0
331 %OUTPUT.Area_pillar= -area(polyshape(Array_of_Polygon(1).P))+area(
332 polyshape(Array_of_Polygon(2).P));% area of the area of the pillars
333             else
334 % OUTPUT.Area_pillar= area(polyshape(Array_of_Polygon(1).P));% area of
335             the area of the pillars
336             end
337             end
338             % d = .94/(sqrt(real(epsSi))-sqrt(real(eps_square))); %hauteur des
339             pilliers theorique
340             %Bug was d a,d not d_(iNpillar)
341             Layers_(iNpillar) = [layer_2D(d_(iNpillar), eps_square,
342             Array_of_Polygon, Array_of_circle, current_z_(iNpillar)-d_(iNpillar),
343             current_z_(iNpillar))];
344             %ind_current_layer = ind_current_layer + 1; nz=[nz round(100/Npillar
345 ) ]; current_z = current_z - d;
346         end
347
348 Layers(ind_current_layer_0+[0:Npillar]) =Layers_(1+[0:Npillar]);
349 current_z= current_z_(end); %ici peut etre un bug
350
351 %%%%%%%%%%%%%%%%

```

```

342
343 function [ALLcenter, ALLcentery, ax, ay]=Create_Grid_Structuration(ax,ay,
344 ARG);
345 % create a grid for structuration with or without periodic square
346 % lattice
347 if isfield(ARG,'RectangularGrid'); RectangularGrid=ARG.RectangularGrid;
348 else; RectangularGrid=1; end
349 if RectangularGrid==1
350 % STRUCTURATION:
351 % definition of the centers of all pillars
352
353 if isfield(ARG,'ALLcenter'); ALLcenter=ARG.ALLcenter;
354 else
355 %ALLcenter=[-1 0 1] ;
356 ALLcenter=[ 0 ] ;
357 end
358 if isfield(ARG,'ALLcentery'); ALLcentery=ARG.ALLcentery;
359 else
360 ALLcentery= ALLcenter ;
361 end
362 ALLcenter=ALLcenter*ax;
363 ALLcentery=ALLcentery*ay;
364 Nduplicate=max(size(ALLcenter,2));
365 Nduplicatey=max(size(ALLcentery)); %ALLcentery is a 1 vector
366 % the total size of the device
367 ax=Nduplicate*ax;ay=Nduplicatey*ay; %
368 else % list of absolute values for pillar positions
369 if isfield(ARG,'ALLcenter'); ALLcenter=ARG.ALLcenter;
370 else
371 ALLcenter=[ 0 ; 0] ;
372 end
373 ALLcentery=[]; % not used ==> a list of absolute pillar positions
374 end
375 %%%%%%%%%%%%%%
376
377
378 function OUTPUT = struct_LayerCoating(ARG)
379 if isfield(ARG,'ax')
380 ax=ARG.ax;
381 else
382 ax=1;
383 end
384 if isfield(ARG,'ay')
385 ay=ARG.ay;
386 else
387 ay=ax;
388 end
389 if isfield(ARG,'PML')
390 PML=ARG.PML;
391 else
392 PML = 0;
393 end
394 if isfield(ARG,'N_PML')
395 N_PML=ARG.N_PML;
396 else

```

```

397     N_PML = 0;
398
399 if isfield(ARG, 'thickness_PML')
400     thickness_PML = ARG.thickness_PML;
401 else
402     thickness_PML = 1;
403 end
404 if isfield(ARG, 'max_eps_PML')
405     max_eps_PML = ARG.max_eps_PML;
406 else
407     max_eps_PML = 20;
408 end
409 if isfield(ARG, 'DSA')
410     DSA = ARG.DSA;
411 else
412     DSA = 'no';
413 end
414 %DSA Thickness ;
415 if isfield(ARG, 'DSA_Thickness')
416     DSA_Thickness = ARG.DSA_Thickness;
417 else
418     DSA_Thickness = 0.1;
419 end
420 % multilayer
421 NBML=0;for iML=1:10
422 if isfield(ARG, [ 'd_ML' num2str(iML)])
423     NBML=NBML+1;
424     eval(['d_ML' num2str(iML) '=' ARG.d_ML' num2str(iML) ';' ]);
425     eval(['Mat_ML' num2str(iML) '=' ARG.Mat_ML' num2str(iML) ';' ]);
426 else
427     eval(['d_ML' num2str(iML) '= 0;' ]);
428 end
429 end
430 OUTPUT.NBML=NBML;
431 if isfield(ARG, 'Material_Sub')
432     Material_Sub = ARG.Material_Sub;
433 else
434     Material_Sub = 'Silicon';
435 end
436 % Define the structure
437 % Period in x in micrometer
438 x_min = -ax / 2;
439 x_max = x_min + ax;
440 %ay = 0.5; % Periode in y in micrometer
441 y_min = -ax / 2;
442 y_max = y_min + ay;
443 MAXCOORD.xmin = x_min;
444 MAXCOORD xmax = x_max;
445 MAXCOORD.ymin = y_min;
446 MAXCOORD ymax = y_max;
447 % PERMITTIVITY
448 lambda = ARG.lambda;
449 % Palik data were used in fdtd and dgtd simulations.
450 % load('Ta205.txt','ascii') % wl epsilon
451 % Eps_Ta205 = Ta205(:,2)+1j*Ta205(:,3); % Eps_R Eps_I
452 % wl_Ta205 = Ta205(:,1);
453 %load('C:\Users\lecal\OneDrive\Documents\SUPOP\2AS\TP Projet\SiO2_Palik
454 % ','ascii') % wl n k

```

```

454 load('Si_Palik.txt','ascii') % wl n k
455 load('Si02_Palik.txt','ascii') % wl n k
456 wl_Si_Palik = Si_Palik(:,1) ; wl_Si02_Palik = Si02_Palik(:,1);
457 Eps_Si_Palik = (Si_Palik(:,2).^2 - Si_Palik(:,3).^2) + 1i*(2*Si_Palik
    (:,2).*Si_Palik(:,3)); % Eps_R Eps_I
458 Eps_Si02_Palik = (Si02_Palik(:,2).^2 - Si02_Palik(:,3).^2) + 1i*(2*
    Si02_Palik(:,2).*Si02_Palik(:,3)); % Eps_R Eps_I
459 %Eps_Ta205 = (Ta205_refractiveindex(:,2).^2 - Ta205_refractiveindex(:,3)
    .^2) + 1i*(2*Ta205_refractiveindex(:,2).*Ta205_refractiveindex(:,3));
    % Eps_R Eps_I
460 %wl_Ta205 = 1e3*Ta205_refractiveindex(:,1);
461 % interp at wavelength
462 epsSi02 = interp1(wl_Si02_Palik, Eps_Si02_Palik, lambda*1e3);
463 epsSi = interp1(wl_Si_Palik, Eps_Si_Palik, lambda*1e3);
464 %epsTa205 = interp1(wl_Ta205, Eps_Ta205, lambda*1e3);
465 epsAl203=3.02295;
466     epsAir = 1;
467 if strmatch(Material_Sub, 'silicon', 'exact') ;
468     epsMap = [ epsSi, epsAir];
469 else
470     epsMap = [ epsSi02, epsAir];
471 end
472 %multilayer
473 for iML=1:NBML
    eval(['tmpMat=' 'Mat_ML' num2str(iML) ';' ])
474 tmpMattxt=[tmpMat '.txt'];
475 eval(['load(''' tmpMattxt ''', '-ascii');']);
476 eval(['tmpMat =' tmpMattxt(1:end-4) ';' ])
477 tmpepsilon=nk2epsilon(tmpMat(:,2),tmpMat(:,3));
478 supermicronlambda=lambda*(1e3*isempty(findstr('Palik',tmpMattxt))
+1*isempty(findstr('Palik',tmpMattxt)));
479 if supermicronlambda<tmpMat(1,1);supermicronlambda=tmpMat(1,1);end
480 if supermicronlambda>tmpMat(end,1);supermicronlambda=tmpMat(end,1);
481 end
482 tmpepsilon=interp1(tmpMat(:,1), tmpepsilon, supermicronlambda); %
normlisation of the W axis
483 if isnan(tmpepsilon);tmpepsilon=1;end
484 eval(['epsMat_ML' num2str(iML) '=tmpepsilon;'])
485 epsMap = [epsMap tmpepsilon];
486 end
487 epsMap=unique(epsMap);
488 %Attribute the top and bottom material (must be real)
489     epssup = epsAir;
490     epssub = real(epsSi); % super- and sub-strate permeabilities
491     if strmatch(Material_Sub, 'silicon', 'exact') ; epssub = real(
        epsSi); else epssub = real(epsSi02);end
492     MAX_Z = 10; % Max z of your structure
493 %SUP-Strate
494     nz=1;
495     d_sup = [0.01];
496     current_z = MAX_Z;
497     ind_current_layer = 1;
498 %% 4 um Air
499     Array_of_Polygon = [];
500     eps_square = epsAir;
501     Array_of_circle = [];
502     d = 10;
503     Layers(ind_current_layer) = [layer_2D(d, eps_square,

```

```

      Array_of_Polygon, Array_of_circle, current_z-d, current_z)];
504    ind_current_layer = 1 + ind_current_layer;
505    nz = [nz 60];
506    current_z = current_z - d;
507 %ML
508 for iML=1:NBML
509     Array_of_Polygon = [];
510     eval(['eps_square =epsMat_ML' num2str(iML) ';' ])
511     Array_of_circle = [];
512     eval(['d = d_ML' num2str(iML) ';' ])
513     Layers(ind_current_layer) = [layer_2D(d, eps_square,
514     Array_of_Polygon, Array_of_circle, current_z-d, current_z)];
514     ind_current_layer = 1 + ind_current_layer; nz=[nz 10 ]; current_z =
515     current_z -d;
516 end
517 %DSA
518 if DSA(1)=='y'
519 ARG_DSA.PLOT=0; ARG_DSA.DSA_Thickness=DSA_Thickness;
520 %realstructure=imread('selfassembly.PNG');
521 % One level
522 %DSA_Shape=realstructure(:,:,1)>60;
523 %realstructure=imread('Mouche_zoom.PNG');
524 realstructure=imread('Mouche.png');
525 %realstructure=imread('Mouche_Nice.PNG');
526 % One level
527 DSA_Shape=realstructure(:,:,1)>100;
528 %DSA_Shape=DSA_S
529 ape(20:end-20,20:end-20);
530 % figure;imagesc(DSA_Shape)
531 ARG_DSA.DSA_Min_Nb_Edges=5; %nim size of poly
532 ARG_DSA.DSA_Min_Around_Space=0.95; % border guard
533 Layers_DSA=ADD_DSA_Layer(current_z,x_max,x_min,y_max,y_min,epsSi,epsAir,
534 DSA_Shape,ARG_DSA);
535 Layers(ind_current_layer) =Layers_DSA;
536 ind_current_layer = 1 + ind_current_layer; nz=[nz 10 ]; current_z =
537 current_z -ARG_DSA.DSA_Thickness;
538 end
539 %% 10um of Si
540 Array_of_Polygon=[];
541 if strmatch(Material_Sub, 'silicon', 'exact') ; eps_square = epsSi;
542 else eps_square = epsSiO2;end%outer square
543 Array_of_circle = [];
544 d = 10;
545 Layers(ind_current_layer) = [layer_2D(d, eps_square,
546 Array_of_Polygon, Array_of_circle, current_z-d, current_z)];
547 ind_current_layer = 1 + ind_current_layer;
548 nz=[nz 60];
549 current_z = current_z -d;
550 %% SUB-strate
551 d_sub = 0.01;
552 % d_sub = 10;
553 nz = [nz 1];
554 % Constructing usual variable : d and L
555 L = max(size(Layers));
556 d=0; tmp_d = [];
557 for i = 1:max(size(Layers))
558     tmp_d = [tmp_d , Layers(i).d];
559 end

```

```

555 d = [d_sup, tmp_d, d_sub];
556 BOT_Z = MAX_Z - sum(d(2:(end-1)));
557 % disp(Layers)
558 st = dbstack;
559 OUTPUT.name = st.name;
560 OUTPUT.ax = ax;
561 OUTPUT.x_min = x_min;
562 OUTPUT.x_max = x_max;
563 OUTPUT.ay = ay;
564 OUTPUT.y_min = y_min;
565 OUTPUT.y_max = y_max;
566 OUTPUT.epssup = epssup;
567 OUTPUT.epssub = epssub;
568 OUTPUT.MAX_Z = MAX_Z;
569 OUTPUT.BOT_Z = BOT_Z;
570 OUTPUT.epsMap = epsMap;
571 OUTPUT.nz = nz;
572 OUTPUT.d = d;
573 OUTPUT.L = L;
574 OUTPUT.Layers = Layers;
575 end
576
577 %%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
578 clear all
579
580 format longE;
581
582 addpath(genpath('/MATLAB Drive/pcode'));
583 %RCWA Library generation
584 %Generate a phase/radius library using RCWA solver for a single pillar
585 %This is a demo script, where you just have to set up few parameters.
586 %It uses another script to generate the device structure called:
587 %struct_Pilier_ODIF_PML_robin.m
588 %The structure is as follow (in nanometer):
589
590 clear;
591 % TM or TE mode
592 %TM=0 ==> H is along x [0 Ey Ez(theta)] [Hx 0 0]
593 %TM=1 ==> E is along x [Ex 0 0] [0 Hy Hz(theta)]
594
595 OPTION.TM = 0; % sweep in Lambda
596 numom = 1; %nb of points to sweep
597 % if OPTION.halfnpw>0;numom = 16;else;numom =100;end;
598 % om = linspace(1e3*1.24/940,1e3*1.24/300 ,numom); % photon energy (hbar
      omega) in eV
599 % lambdaall = 1.24./om; %micrometer
600 %lambdaall = linspace(.3,.94 ,numom); %micrometer
601 lambdaall = linspace(.94,.94 ,numom); %micrometer
602 % Incident light angle (NEVER SET ZERO)
603 theta = 0; % Angle with z-axis
604 phi = 0; % Angle with x-axis
605 % Number of mode in fourier space (RCWA MANDATORY)
606 % Recommended 18 in case of non uniform layers: 18
607 OPTION.halfnpw = 8; % On change le nombre de mode ici
608 % Visualization options
609 OPTION.SlideType = ['x', 'z']; % Type of slices.
610 OPTION.SlideValue = [0, -0.05]; % value of slices.
611 OPTION.display_colormap = 1; % 1 -> Display colormap on plots of fields

```

```

612 disp(" "); disp("STARTING calculation"); disp(" ");
613
614 warning off
615 % Device
616
617 period_of_Pillar_Grid=0.47 ; % um
618 Nb_sim=121; %mettre bcp de point pour ne pas manquer la résonance
619 Wmin=0.12; Wmax=period_of_Pillar_Grid-0.11;%
620 Width_ = linspace(Wmin, Wmax, Nb_sim);
621 [PX PY]=meshgrid([-1:1]); %in um directly
622 [PX PY]=meshgrid([0]); %in um directly
623 N_grid=size(PX,1);
624     % RCWA Solving : cycle over frequency
625     for i = 1:size(lambdaall, 2)
626         %parfor stuff
627         Lambda_all = zeros(1,numom); RR = zeros(1,numom); TT = zeros(1,numom);
628         AA = zeros(1,numom);
629 tic
630     % Cycle over devices
631         for j = 1:Nb_sim
632             % Create the Device to be simulated
633             ARG= struct;
634             % Wavelength
635             ARG.lambda =lambdaall(i);% 1.24/om(i); % (um)
636             % ARG.t = j;
637             % Size of the simulated region
638             ARG.ax = period_of_Pillar_Grid*N_grid;
639             ARG.Material_Sub='silicon';
640             ARG.Material_Sub='SiO2';
641             %pillar grid
642             ARG.STRUCT='y';
643             ARG.ALLcenter=period_of_Pillar_Grid*[PX(:) PY(:)];
644             ARG.ALLcenterY=[];
645             ARG.RectangularGrid=0;
646             ARG.PillarWidth=Width_(j)+zeros(1,(size(ARG.ALLcenter,1)));
647             ARG.PillarWidth2=Width_(j)+zeros(1,(size(ARG.ALLcenter,1)));
648             ARG.PillarInnerWidth=ARG.PillarWidth-.12;
649             ARG.Pillar_type="circSquare" ;
650             ARG.Pillar_type="circle" ; ARG.Pillar_NbPeriod
651             =0+0*ARG.PillarWidth;ARG.Pillar_circSquareangle=10+0*ARG.PillarWidth;
652             %from circle to sqr
653             ARG.Npillar=0;
654             ARG.angle_pillar=0+0*ARG.PillarWidth;
655             % Multi coating
656             %ARG.d_ML1=0.06;
657             ARG.Mat_ML1='SiO2_Palik';
658             %ARG.d_ML2=0.047;ARG.Mat_ML2='Ta205_refractiveindex';
659             %ARG.d_ML3=0.013;ARG.Mat_ML3='Al203_refractiveindex';
660             % DEVICE creation
661             DEVICE = struct_LayerCoating2(ARG);
662             % Visualization at the last frequency (do not work with a parfor)
663             % visualization_structure3D(DEVICE, OPTION);
664             if (j==Nb_sim)
665                 figure(4)
666             % Possibilité de faire run section ici

```

```

667     OPTION_3D.PlotEdge='no';
668     OPTION_3D.PlotEdge='yes';
669     OPTION_3D.Plotfast='on';
670     OPTION_3D.epsilon_Visible='no';
671     visualization_structure3D_Perspective(DEVICE,
672     OPTION_3D,get(gcf,'number'));
673 % Pour visualiser La figure 2 faire run section ici %remove air
674 % Visualisation3D_Remove(get(gcf,'number'),[ 1 0
675 0]);set(gca,'zlim',[ -3 1])
676 drawnow
677 end
678 % RCWA3D Simulation
679 SOL = RCWA_3D_singlecore(DEVICE, ARG.lambda, theta, phi, OPTION);
680 ;
681 ModalPhase(j, :) = SOL.ModalPhase;
682 % Retrieving result
683 RR(j) = SOL.RR;
684 TT(j) = SOL.TT;
685 % AA_ = SOL.AA;
686 % L = DEVICE.L;
687 % AA(j) = AA_(L);
688 end
689 toc
690 ALLPhase( :, :) = ModalPhase;
691 ALLTrans( :) = TT;
692 % Figure of the phase/transmission versus pillar width
693 monitor=0;
694 figure(40)
695 Phh = ModalPhase(:, end-monitor) - ModalPhase(1, end-monitor);
696 %Phh(find(Phh<0)) = Phh(find(Phh<0))+2*pi;
697 Phh=unwrap(Phh);
698 %Phh = abs(Phh-2*pi);
699 yyaxis left
700 title('Phase and transmission as a function of the pillar width')
701 plot(Width_(1:size(Phh,1))/2, Phh, '-o')
702 xlabel('Radius (um)')
703 yticks([0 pi/4 pi/2 3*pi/4 pi 5*pi/4 3*pi/2 7*pi/4 2*pi 9*pi/4 10*pi/4
704 ])
705 yticklabels({'0', ' ', ' ', '\pi/2', ' ', ' ', '\pi', ' ', ' ', '3\pi/2', ' ', ' ', '2\pi',
706 ' ', ' ', '5\pi/2'})
707 axis([Wmin/2 Wmax/2 0 3*pi])
708 ylabel('Phase')
709 hold on
710 yyaxis right
711 plot(Width_(1:size(Phh,1))/2, TT, '-s')
712 legend({'Phase','Transmission'},'Location','west')
713 ylabel('Transmission (norm)')
714 ylim([0 1])
715 set(gcf,'color',[1 1 1])
716 % eval([' savefig( '' NAME_SIMULATION '_halfnpw-' num2str(OPTION.
717 halfnpw) ''')'])
718 end
719 return
720 load /prj/ai4odiff/users/rd10/MPW6_AI_BENCH_2023/
721 DS_Odiff2_etchV4_1_insensitive_220122.mat
722 PHASE=(unwrap(ds.phiT - ds.phiT(1)));
723 PillarW=2*ds.scatterer.scatterer_radius_top_x;

```

```

718 yyaxis left
719 plot(0.5*PillarW,PHASE,'--')
720 legend({'Phase','Transmission','Phase Ref'},'Location','west')
721 NAME_SIMULATION='ODIFF_LIB'
722 eval([' savefig( ''', NAME_SIMULATION '_halfnpw-' num2str(OPTION.halfnpw
    ) '''')'])
723 %% Plot fields after library calculation
724 % Index of the radius to plot
725 idx = Nb_sim / 2;
726 r = radius(idx);
727 SOL = ALL_SOL(idx);
728 if OPTION.FIELD == 1
729     OPTION.PLOTPHASE = 0;
730     PlotField3D(SOL, DEVICE, OPTION);
731     OPTION.PLOTPHASE = 1;
732     PlotField3D(SOL, DEVICE, OPTION);
733 end

```

## D Sources

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