

### 3.1

解：由于对称性，电场强度 $\vec{E}$ 具有 $r$ 方向分量

由 $\nabla \cdot \vec{E} = \rho / \epsilon_0$ 得

$$E = \rho_0 / 3\epsilon_0 * r + C_1 \quad (0 < r < a)$$

$$C_2 / r^2 \quad (r > a)$$

当 $r=0$ 时， $E=0$

$$\text{当 } r=a \text{ 时， } \rho_0 / 3\epsilon_0 * r + C_1 = C_2 / r^2$$

$$\text{解得 } c_1 = 0, \quad c_2 = \rho_0 a^3 / 3\epsilon_0$$

$$\text{故 } E = \rho_0 r / 3\epsilon_0 \quad (0 < r < a)$$

$$\rho_0 a^3 / 3\epsilon_0 r^2 \quad (r > a)$$

运用积分形式

$$\text{当 } 0 < r < a \text{ 时 } E * 4\pi * r^2 = \frac{4}{3}\pi r^3 \rho_0 / \epsilon_0 \text{ 得到 } E = \rho_0 r / 3\epsilon_0$$

$$\text{当 } r > a \text{ 时 } E * 4\pi * r^2 = \frac{4}{3}\pi a^3 \rho_0 / \epsilon_0 \text{ 得到 } E = \rho_0 a^3 / 3\epsilon_0 r^2$$

验证了结论

### 3.2

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### 3.3

解：1)  $\nabla^2 \varphi = 0 \quad (a_1 < \rho < a_2)$   $\varphi$ 有 $\rho$ 方向上的分量

$$\text{即 } \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho \frac{\partial \varphi}{\partial \rho}) = 0 \text{ 故有 } \varphi = C_1 \ln \rho + C_2 \quad (a_1 < \rho < a_2)$$

$$\text{当 } \rho = a_1 \text{ 时 } \varphi = U_0$$

$$\text{当 } \rho = a_2 \text{ 时 } \varphi = 0$$

$$\text{故 } \varphi = \frac{U_0}{\ln a_2 / a_1} \ln a_2 / \rho \quad (a_1 \leq \rho \leq a_2)$$

$$2) E = -\nabla \varphi = U_0 / \ln \frac{a_2}{a_1} * \frac{1}{\rho}$$

$$\text{当 } \rho = a_1 \text{ 时 } E_{max} = \frac{U_0}{\ln a_2 / a_1} * \frac{1}{a_1}$$

对 $a_1 * \ln a_2 / a_1$ 求导得到  $a_1 = a_2 / e$  时  $E_{max}$  最小为  $e U_0 / a_2$

3.4

解：可认为圆柱形容器为无限长，即不考虑其边缘效应，同样设电容器内外导体单位长度上的电量为为  $+\tau$  和  $-\tau$

$$\text{由 } U_1 = U_2 \Rightarrow r_2 = r_c \left( \frac{r_c}{r_1} \right)^{\varepsilon_{r2}/\varepsilon_{r1}} = 1.966(\text{cm})$$

故：内容介质厚度： $d_1 = r_c - r_1 = 0.5(\text{cm})$

则根据高斯定理，电容器介质中的电位移  $\vec{D} = \frac{\tau}{2\pi r} \vec{r}^0$

$$\text{则内层电介质中电场 } \vec{E}_1 = \frac{1}{\varepsilon_1} \vec{D} = \frac{\tau}{2\pi\varepsilon_0\varepsilon_{r1}r} \vec{r}^0$$

$$\text{外层电介质中电场为 } \vec{E}_2 = \frac{1}{\varepsilon_2} \vec{D} = \frac{\tau}{2\pi\varepsilon_0\varepsilon_{r2}r} \vec{r}^0$$

由于电场强度  $E$  正比于是  $1/r$ ，故在该两层介质分界面处半径为  $r_c$  时有：

$$E_{1\max} = \frac{\tau}{2\pi\varepsilon_0\varepsilon_{r1}r_1}, \quad E_{2\max} = \frac{\tau}{2\pi\varepsilon_0\varepsilon_{r2}r_2}$$

$$U_1 = \int_{r_1}^{r_c} \vec{E}_1 \cdot d\vec{r} = \frac{\tau}{2\pi\varepsilon_0\varepsilon_1} \int_{r_1}^{r_c} \frac{1}{r} dr = \frac{\tau}{2\pi\varepsilon_0\varepsilon_1} \ln \frac{r_c}{r_1}$$

$$U_2 = \int_{r_c}^{r_2} \vec{E}_2 \cdot d\vec{r} = \frac{\tau}{2\pi\varepsilon_0\varepsilon_1} \ln \frac{r_2}{r_c}$$

$$\text{根据 } E_{1\max} = E_{2\max} \Rightarrow r_c = \frac{\varepsilon_{r1}}{\varepsilon_{r2}} r_1 = 1.5(\text{cm})$$

$$\text{由 } U_1 = U_2 \Rightarrow r_2 = r_c \left( \frac{r_c}{r_1} \right)^{\varepsilon_{r2}/\varepsilon_{r1}} = 1.966(\text{cm})$$

故：内容介质厚度： $d_1 = r_c - r_1 = 0.5(\text{cm})$

外层介质厚度： $d_2 = r_2 - r_c = 0.466(\text{cm})$

### 3.5

解: 空气的击穿场强决定了在金属球表面与空气的交界面上的电场的最大值, 另外注意金属导体表面的电场只有法向分量, 金属导体球表面上有自由面电荷 $\delta$ , 其中最大值可由导体与介质界面的衔接条件决定

1) 由于面电荷均匀分布特性, 可求得金属球上最大的电量

$$q_{\max} = \oint_S \sigma_{\max} ds = \sigma_{\max} \cdot 4\pi r^2 = 7.507 \times 10^{-6} (C)$$

2) 只需要计算长圆柱的单位长度电荷总量

$$q_1 = \int_S \sigma_{\max} ds = \sigma_{\max} \cdot 2\pi r \cdot 1 = 2.502 \times 10^{-5} (C)$$

### 3.6

解:

1) (a) 当导体与内套管的电压  $U_1$  称内套管与外套管的电压  $U_2$  相等即:  $U_1 = U_2$

$$\text{此时导体与内套管之间的电场 } E_1 = \frac{U_1}{r \ln \frac{c}{a}}$$

$$\text{于是内套管和外内套管之间的电场 } E_2 = \frac{U_2}{r \ln \frac{b}{c}}$$

于是内外套管上的电荷密度为

$$\sigma_c = +\varepsilon_0 E_1 \Big|_{r=c} = \frac{\varepsilon_0 U_1}{c \ln \frac{c}{a}}$$

$$\sigma_b = +\varepsilon_0 E_2 \Big|_{r=b} = \frac{\varepsilon_0 U_2}{b \ln \frac{b}{c}}$$

由于内外套管上内表面上的总电荷相同, 即

$$2\pi c l \sigma_c = 2\pi b l \sigma_b \Rightarrow L = \frac{c l \sigma_c}{b \sigma_b}$$

$$L=155.13(\text{mm})$$

(b)导体与内套管间的最大场强为导体表面处的场强  $E_{1\max} = \frac{U_1}{a \ln \frac{c}{a}}$

内外套管间的最大场强为导体表面处的场强  $E_{2\max} = \frac{U_2}{b \ln \frac{c}{b}}$

由于  $E_{1\max} = E_{2\max}$

又由于  $2\pi cl\sigma_c = 2\pi bL\sigma_b$  可得

$$\frac{lU_1}{\ln \frac{c}{a}} = \frac{LU_2}{\ln \frac{b}{c}} \Rightarrow \frac{U_1}{U_2} = \frac{L \ln \frac{c}{a}}{l \ln \frac{b}{c}}$$

由上两式可得出  $L = \frac{a}{c}l = 150mm$

$$(2) \quad \begin{aligned} E_{1\max} &= \frac{U_1}{a \ln \frac{c}{a}} = 115.87U_1 \\ E_{2\max} &= \frac{U_2}{b \ln \frac{b}{c}} = 112.04U_2 \end{aligned}$$

当  $U_1 = U_2$  时  $E_{1\max} > E_{2\max}$  因此此时系统中最大场强在导体表面即空气的击穿在此处若保

证气隙不击穿  $E_{1\max} < \text{空气的击穿场强 } E_m$

3.7

解:列写方程组

$$\nabla^2 \varphi_1 = 0 \quad (1)$$

$$\nabla^2 \varphi_2 = 0 \quad (2)$$

$$\varphi_1|_{x=d_1} = \varphi_2|_{x=d_2} \quad (3)$$

$$\varphi_1|_{x=0} = 0 \quad (4)$$

$$\varphi_2|_{x=d_1+d_2} = U_0 \quad (5)$$

$$-\varepsilon_1 \frac{\partial \varphi_1}{\partial n} = -\varepsilon_2 \frac{\partial \varphi_2}{\partial n} \quad (x=d_1) \quad (6)$$

解得  $\varphi_1 = C_1 + C_2x \quad \varphi_2 = C_3 + C_4x$

$$C_1 = 0 \quad C_2 = \frac{\varepsilon_2 U_0}{\varepsilon_2 d_1 + \varepsilon_1 d_2} \quad C_3 = \frac{(\varepsilon_2 - \varepsilon_1) d_1 U_0}{\varepsilon_2 d_1 + \varepsilon_1 d_2} \quad C_4 = \frac{\varepsilon_1 U_0}{\varepsilon_2 d_1 + \varepsilon_1 d_2}$$

介质中的电场

$$E_1 = -\nabla\varphi_1 = -\frac{\partial\varphi_1}{\partial x} = -\frac{\varepsilon_2 U_0}{\varepsilon_2 d_1 + \varepsilon_1 d_2}$$

$$E_1 = -\nabla\varphi_1 = -\frac{\partial\varphi_2}{\partial x} = -\frac{\varepsilon_1 U_0}{\varepsilon_2 d_1 + \varepsilon_1 d_2}$$

3.8

解：用球坐标系写出其拉普拉斯方程及边值问题得

$$\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\varphi}{\partial\theta}\right) = 0 \quad (1)$$

$$\varphi|_{\theta=\theta_0} = U_0 \quad (2)$$

$$\varphi|_{\theta=\frac{\pi}{2}} = 0 \quad (3)$$

$$\text{解得}\varphi = C_1 \ln \tan \frac{\theta}{2} + C_2 \quad C_1 = \frac{U_0}{\ln \tan \frac{\theta_0}{2}} \quad C_2 = 0$$

$$\text{电场}\vec{E} = -\nabla\varphi = -\frac{U_0}{r\sin\theta \ln\left(\tan\frac{\theta_0}{2}\right)} \vec{e}_\theta$$

3.9

解：

1) 嵌入金属壳后，球形电容器的电场有变化，原因如下：

当嵌入金属壳 S 后将球形电容器内电介质分成两部分，S 成为这两部分电介质的分界面。

在该分界面上电位处的相等，即 S 面是等势面。类似地，我们也可把 (a) 图中的介质分成如 (b) 图一样的两部分，显然在 (a) 图中，S 对应的分界面不再是等势面，因此 (a) (b) 两图中的边缘问题的边界条件不相同，其解必不同。

2) 若金属球壳与球形电容器同心，则 S 的嵌入不改变边缘问题的边界条件，因此不改变电场的分布，因为球形电容器中的等势面正好是与球心同心一簇同心球面。

3.10

解：

$$\nabla^2\varphi_1 = 0$$

$$\nabla^2\varphi_2 = 0$$

$$\varphi_1|_{s_{12}} = \varphi_2|_{s_{12}}$$

$$\varepsilon_1 \frac{\partial\varphi_1}{\partial n}|_{s_{12}} = \varepsilon_2 \frac{\partial\varphi_2}{\partial n}|_{s_{12}}$$

$$\varphi_1|_{s_1} = U_1$$

$$\varphi_1|_{s_2} = U_{c2}$$

$$\oiint_{s_2} \left( -\varepsilon_1 \frac{\partial \varphi_1}{\partial n} \right) dS = q_2$$

$$\varphi_1|_{s_3'} = \varphi_2|_{s_3''} = U_{c3}$$

$$\oiint_{s_3'} \left( -\varepsilon_1 \frac{\partial \varphi_1}{\partial n} \right) dS + \oiint_{s_3''} \left( -\varepsilon_2 \frac{\partial \varphi_2}{\partial n} \right) dS = q_3$$

$$\varphi_1|_{s_0'} = \varphi_2|_{s_0''} = 0$$

3.11

解：小带电体可视为一点电荷  $q$ ，它所受静电力，来自导体平板的感应电荷，也就是镜像电荷  $q'$  (平面上方  $h$  处， $q' = -q$ ) 对它的作用力

$$f_e = -\frac{q^2}{4\pi\varepsilon_0(2h)^2}$$

令  $f_e$  与重力  $mg$  大小相等，有

$$\frac{q^2}{4\pi\varepsilon_0(2h)^2} = mg$$

$$\text{即 } q = 4h\sqrt{\pi\varepsilon_0 mg}$$

3.12

解：设  $q_1$  受电场力为  $F_1$ ， $F_1$  是  $q_2$  和分界面上产生感应电荷共同产生的电场对  $q_1$  产生的电场力。同理  $q_2$  受电场力为  $F_2$ ，是  $q_1$  和分界面上产生的感应电荷共同产生的电场对  $q_2$  的电场力。

现对本题情况进行镜像分析可知

$$F_1 = \frac{q_1 \times (q_1' + q_2'')}{4\pi\varepsilon_1 \times (2h)^2} = \frac{q_1}{16\pi\varepsilon_1 h^2} \times \frac{1}{\varepsilon_1 + \varepsilon_2} [(\varepsilon_1 - \varepsilon_2)q_1 + 2\varepsilon_1 q_2]$$

$$F_2 = \frac{q_2 \times (q_2' + q_1'')}{4\pi\varepsilon_2 \times (2h)^2} = \frac{q_2}{16\pi\varepsilon_2 h^2} \times \frac{1}{\varepsilon_1 + \varepsilon_2} [(\varepsilon_2 - \varepsilon_1)q_2 + 2\varepsilon_2 q_1]$$

故两个力大小不相等

3.13

解：

(1)  $Z > 0$  区域的电场由  $q$  和分界面上的极化电荷共同产生

因而  $Z>0$  区域中任意点 P 处的电位

$$\varphi = \frac{q}{4\pi\epsilon_1 r_1} + \frac{q'}{4\pi\epsilon_1 r_1'} = \frac{q}{4\pi\epsilon_0 r_1} + \frac{1}{4\pi\epsilon_0 r_1'} \frac{1-\epsilon_r}{1+\epsilon_r} q$$

$$\text{P 处的电场 } \vec{E}_1 = \frac{q}{4\pi\epsilon_1 r_1^2} \vec{r}_1^0 + \frac{q'}{4\pi\epsilon_1 r_1'^2} \vec{r}_1'^0 = \frac{q}{4\pi\epsilon_0 r_1^2} \vec{r}_1^0 + \frac{1}{4\pi\epsilon_0 r_1'^2} \frac{1-\epsilon_r}{1+\epsilon_r} q \vec{r}_1'^0$$

其中  $\vec{r}_1^0$  和  $\vec{r}_1'^0$  分别为  $q$  和  $q'$  指向 P 点的单位矢量。

(2)  $Z<0$  区域的电场由如图所示的镜像电荷  $q''$  确定

$$q'' = \frac{2\epsilon_2}{\epsilon_1 + \epsilon_2} q = \frac{2\epsilon_r}{1 + \epsilon_r} q$$

$$Z<0 \text{ 区域中任意点 P 处的电位 } \varphi = \frac{q''}{4\pi\epsilon_2 r_2} = \frac{1}{4\pi\epsilon_1 \epsilon_2 r_2} \frac{2\epsilon_r}{1 + \epsilon_r} q$$

$$\text{电场 } \vec{E}_2 = \frac{q''}{4\pi\epsilon_2 r_2^2} \vec{r}_2^0 = \frac{1}{4\pi\epsilon_1 \epsilon_0 r_2^2} \frac{2\epsilon_r}{1 + \epsilon_r} q \vec{r}_2^0$$

(3) 显然如图 Q 点的场强是分界面最大场强

$$\begin{aligned} \vec{E}_Q &= \vec{E}_1|_{Q\text{点}} = \frac{q}{4\pi\epsilon_0 h^2} \vec{r}_1^0 + \frac{1}{4\pi\epsilon_0 h^2} \frac{1-\epsilon_r}{1+\epsilon_r} \vec{r}_1'^0 \\ &= \frac{q}{4\pi\epsilon_0 h^2} \left( 1 - \frac{1-\epsilon_r}{1+\epsilon_r} \right) \vec{r}_1^0 \\ &= \frac{q}{4\pi\epsilon_0 h^2} \frac{2\epsilon_r}{1+\epsilon_r} \vec{r}_1^0 \end{aligned}$$

$\vec{r}_1^0$  的方向是垂直分界面指向介质 2。上面计算的是分界面上方的电场强度下方的电场强度

应当用  $E_2$  来计算

$$\begin{aligned} \vec{E}_2|_{Q\text{点}} &= \frac{q}{4\pi\epsilon_0 h^2} \frac{2\epsilon_r}{1+\epsilon_r} q \vec{r}_2^0 \\ &= \frac{q}{4\pi\epsilon_0 h^2} \times \frac{2}{1+\epsilon_r} \vec{r}_1^0 < \vec{E}_1|_{Q\text{点}} \end{aligned}$$

故最大场强仍为  $E_Q$ 。

(4) 分界面上的极化电荷由  $\sigma_P = P_{1n} - P_{2n}$  来计算不过  $P_{1n} = 0$

$$\sigma_P = -P_{2n} = -(D_{2n} - \varepsilon_0 E_{2n}) = -(\varepsilon_r \varepsilon_0 - \varepsilon_0) E_{2n}$$

由于 Q 点的场强最大, Q 且只有法向分量, 如 (3) 所计算的方法, 可知其大小为

$$\bar{E}_2 \Big|_{Q点} = \frac{q}{4\pi\varepsilon_0 h^2} \frac{2}{1+\varepsilon_r} \text{ 故 } E_{2n} = E_2 \Big|_{Q点}$$

$$\text{故 } \sigma_P = -(\varepsilon_r \varepsilon_0 - \varepsilon_0) E_{2n} = -(\varepsilon_r \varepsilon_0 - \varepsilon_0) \frac{q}{4\pi\varepsilon_0 h^2} \times \frac{2}{1+\varepsilon_r} = \frac{2q}{4\pi h^2} \times \frac{1-\varepsilon_r}{1+\varepsilon_r}$$

3.14

解: 采用镜像法, 在距离球心  $b = \frac{R^2}{a}$  处放置镜像电荷  $q' = -\frac{R}{a}q$ , 在球心放置镜像电荷  $q''$ ,

$$q' + q'' = Q$$

球外电场由  $q, q', q''$  所建立

因此点电荷  $q$  所受电场力为

$$f = \frac{q}{4\pi\varepsilon_0} \left[ \frac{dQ+Rq}{d^3} - \frac{Rdq}{(d^2-R^2)^2} \right]$$

3.15

解: 对于球壳内的电场由  $q$  和  $q'$  共同产生。球壳外层将均匀分布感应电荷, 其电荷总量为  $q$ ,

这部分的感应电荷对壳内不产生吸引力。球壳内表面的电量为  $-q$  外表面为  $+q$  (静电平衡,

导体内电场为 0) 用  $q'$  代替  $-q$

$$\text{其中} \quad \begin{cases} q = -\frac{R}{D} q' \\ d = \frac{R^2}{D} \end{cases} \Rightarrow \begin{cases} q' = -\frac{R}{d} q \\ D = \frac{R^2}{d} \end{cases}$$

$$\text{因此: A 点的电场 } \bar{E}_A = \frac{q}{4\pi\varepsilon_0 (R-d)^2} (-\bar{n}) + \frac{q'}{4\pi\varepsilon_0 (D-R)^2} (+\bar{n})$$

注意: A、B 两点  $\bar{n}$  的方向均指向球心.

则 A 点感应电荷密度

$$\sigma_A = \bar{D}_A \bullet \bar{n} = \varepsilon_0 \bar{E}_A \bar{n} = -\varepsilon_0 \left[ \frac{q}{4\pi\varepsilon_0 (R-d)^2} - \frac{q'}{4\pi\varepsilon_0 (D-R)^2} \right] (C/m^2)$$

$$\text{B 点电场 } \bar{E}_B = \frac{q}{4\pi\varepsilon_0 (R+d)^2} (-\bar{n}) + \frac{q'}{4\pi\varepsilon_0 (D+R)^2} (-\bar{n})$$



$$\sigma_B = \vec{D}_B \cdot \vec{n} = -\epsilon_0 \vec{E}_B \cdot \vec{n} = -\epsilon_0 \left[ \frac{q}{4\pi\epsilon_0(R+d)^2} + \frac{q'}{4\pi\epsilon_0(D+R)^2} \right] (C/m^2)$$

(2) 点 q 受到的电场力是  $q'$  所施加的

$$\text{故 } f = \frac{q'q}{4\pi\epsilon_0(D-d)^2} = \frac{-\frac{R}{d}q^2}{4\pi\epsilon_0(\frac{R^2}{d} - d)^2} = -\frac{dRq^2}{4\pi\epsilon_0(R^2 - d^2)^2} \text{ 其中负表示吸引力}$$

以上作答与导体球壳接地无关，与球壳带电无关。有关 A、B 两点电荷密度的正负号可以根据两种方法进行判断，第一种方法就是解题中应用矢量运算的方法，第二种方法可以根据物理规律直接进行判断：即同种电荷相互排斥，异种电荷相互吸引。

3.16

解：用电轴法计算其电轴位置

$$b^2 = h^2 - a^2 \Rightarrow b = \sqrt{h^2 - a^2} = 8cm$$

(1) 设电轴上的电荷密度分别为  $-\tau$  和  $+\tau$  则空间中任意点 P 处的电位

$$\varphi = \frac{\tau}{2\pi\epsilon_0} \ln \frac{r_2}{r_1} \text{ 式中 } r_1, r_2 \text{ 分别是 P 点到电轴 } +\tau \text{ 和 } -\tau \text{ 的距离，其电位参考点选在 } y \text{ 轴上}$$

于是 A 点的电位

$$\varphi_A = \frac{\tau}{2\pi\epsilon_0} \ln \frac{b+(h-a)}{b-(h-a)} = \frac{\tau}{2\pi\epsilon_0} \ln 3$$

B 点电位

$$\varphi_B = \frac{\tau}{2\pi\epsilon_0} \ln \frac{b-(h-a)}{b+(h-a)} = -\frac{\tau}{2\pi\epsilon_0} \ln 3$$

由题意知导体间电压为 1000V 即  $U_{AB} = \varphi_A - \varphi_B = 1000V$

$$\text{于是 } \frac{\tau}{2\pi\epsilon_0} \ln 3 - (-\frac{\tau}{2\pi\epsilon_0} \ln 3) = 1000V$$

$$\therefore \tau = \frac{1000}{2 \ln 3} \times 2\pi\epsilon_0 \text{ 于是空间中任意点 P 处的电位:}$$

$$\varphi = \frac{1000}{2 \ln 3} \times 2\pi\epsilon_0 \ln \frac{r_2}{r_1} = 455.12 \ln \frac{r_2}{r_1}$$

$$\text{其中 } r_2 = \sqrt{(b+x)^2 + y^2}, r_1 = \sqrt{(b-x)^2 + y^2}$$

$$\varphi = 228 \ln \frac{(0.08+x)^2 + y^2}{(0.08-x)^2 + y^2} (\text{V})$$

电场强度 (P 点电场强度)

$$\vec{E} = \frac{\tau}{2\pi\epsilon_0 r_1} \vec{r}_1^0 + \frac{-\tau}{2\pi\epsilon_0 r_2} \vec{r}_2^0 = \frac{\tau}{2\pi\epsilon_0 r_1} \vec{r}_1^0 - \frac{\tau}{2\pi\epsilon_0 r_2} \vec{r}_2^0 = 455.12 \left( \frac{\vec{r}_1^0}{r_1} - \frac{\vec{r}_2^0}{r_2} \right)$$

其中  $\vec{r}_1^0$  和  $\vec{r}_2^0$  分别是  $+\tau$  电轴和  $-\tau$  电轴横向 P 点的单位矢量

(2) 显然对于带正电的导线上面电荷密度最大值和最小值分别出现在 A 点和 C 点  
对于 A 点

$$\sigma_{\max} = \sigma_A = D_A = \epsilon_0 E_A = \epsilon_0 \times 455.12 \left( \frac{1}{b-(h-a)} + \frac{1}{b+(h-a)} \right) = 1.34 \times 10^{-7} (\text{C/m}^2)$$

对于 C 点

$$\sigma_{\min} = \sigma_C = D_C = \epsilon_0 E_C = \epsilon_0 \times 455.12 \left( \frac{1}{a+(h-b)} - \frac{1}{a+(b+h)} \right) = 1.678 \times 10^{-8} (\text{C/m}^2)$$

应为0.336乘十的负七次方

(3) 导体单位长度受到的电场力可用电轴法计算该力可认为是其中一根轴对另一根电轴的

$$\text{单位长度电荷产生的作用力设电轴 } +\tau \text{ 产生的电场为 } E \text{ 则 } E = \frac{\tau}{2\pi\epsilon_0 r}$$

则在该电场中的另一根电轴单位长度电荷  $\tau$  所受到的作用力

$$f = qE = \tau E = \frac{\tau^2}{2\pi\epsilon_0 \times 2b} = 7.19 \times 10^{-5} \text{ N}$$

3.17

解：用镜像法计算镜像线电荷有两个分别为  $\tau'$  和  $\tau''$  且

$$\begin{cases} +\tau = -\frac{a}{D}\tau' \\ -\tau = -\frac{a}{D}\tau'' \\ d = \frac{a^2}{D} \end{cases} \Rightarrow \begin{cases} \tau' = -\frac{D}{a}\tau, \tau'' = \frac{D}{a}\tau \\ D = \frac{a^2}{d} \end{cases}$$

(1) 圆管内电场任意 P 的电场等于 4 根线电荷电场的矢量和

$$\begin{aligned} \vec{E} &= \vec{E}_1 + \vec{E}_2 + \vec{E}_3 + \vec{E}_4 = \frac{\tau''}{2\pi\epsilon_0 r_1} \vec{r}_1^0 + \frac{\tau}{2\pi\epsilon_0 r_2} \vec{r}_2^0 + \frac{\tau}{2\pi\epsilon_0 r_3} \vec{r}_3^0 + \frac{\tau'}{2\pi\epsilon_0 r_4} \vec{r}_4^0 \\ &= \frac{\tau}{2\pi\epsilon_0 r_2} \left[ \frac{D}{r_1} \vec{r}_1^0 - \frac{1}{r_2} \vec{r}_2^0 + \frac{1}{r_3} \vec{r}_3^0 - \frac{1}{r_4} \vec{r}_4^0 \right] \end{aligned}$$

(2) 两根线电荷之间作用力为 0, 即  $\tau''$ ,  $-\tau$ ,  $\tau'$  共同产生的电场在  $+\tau$  处场强为 0。设  $+\tau$  处

为 A 点则  $\vec{E}_A = 0$

$$\text{而 } \vec{E}_A = (\vec{E}_1 + \vec{E}_2 + \vec{E}_4) \Big|_A = \frac{\tau''}{2\pi\epsilon_0 r_1} \vec{r}_1^0 + \frac{-\tau}{2\pi\epsilon_0 r_2} \vec{r}_2^0 + \frac{\tau'}{2\pi\epsilon_0 r_4} \vec{r}_4^0$$

$$\text{其中 } r_1 = D + d = \frac{a^2}{d} + d, r_2 = 2d, r_4 = D - d = \frac{a^2}{d} - d$$

$$\vec{E}_A = \frac{\tau''}{2\pi\epsilon_0 r_1} \vec{r}_1^0 - \frac{\tau}{2\pi\epsilon_0 r_2} \vec{r}_2^0 - \frac{\tau'}{2\pi\epsilon_0 r_4} \vec{r}_4^0$$

$$E_A = \frac{\tau}{2\pi\epsilon_0 \left( \frac{a^2}{d} + d \right)} - \frac{\tau}{2\pi\epsilon_0 2d} + \frac{\tau'}{2\pi\epsilon_0 r_4} = 0$$

将  $D = \frac{a^2}{d}$  代入可求得 d 然后两线荷间距离为 2d

3.18

解:

$$a_1=1\text{cm}, \quad a_2=2\text{cm}, \quad d=0.5\text{cm}$$

$$\text{用电轴法计算有 } \begin{cases} a_1^2 = (h_1 + b)(h_1 - b) \\ a_2^2 = (h_2 + b)(h_2 - b) \\ d = h_2 - h_1 \end{cases}$$

$$\text{得 } \begin{cases} h_1 = 2.75\text{cm} \\ h_2 = 3.25\text{cm} \\ b = 2.56\text{cm} \end{cases}$$

设电缆两导体间的电压为 U 即  $U = \varphi_A - \varphi_B$  以 y 轴为参考点则任意点 P 的电位

$$\varphi = \frac{\tau}{2\pi\epsilon} \ln \frac{r_1}{r_2} \text{ 其中 } r_2, r_1 \text{ 分别为 } -\tau \text{ 和 } +\tau \text{ 到 P 点的距离}$$

$$\text{因此对于 A 点 } \varphi_A = \frac{\tau}{2\pi\epsilon} \ln \frac{b + (h_1 - a_1)}{b - (h_1 - a_1)}$$

$$B \text{ 点 } \varphi_B = \frac{\tau}{2\pi\epsilon} \ln \frac{b+(h_2-a_2)}{b-(h_2-a_2)}$$

于是可以得出

$$\frac{\tau}{2\pi\epsilon} = \frac{U}{\ln\left(\frac{b+(h_1-a_1)}{b-(h_1-a_1)} \times \frac{b-(h_2-a_2)}{b+(h_2-a_2)}\right)}$$

$$= 1.65u$$

电缆中最大场强处是 A 点，当 A 点场强达到击穿场强时，电缆电压不能连续升高。

$$\begin{aligned} \vec{E}_A &= \frac{\tau}{2\pi\epsilon r_1} \vec{r}_1^0 + \frac{-\tau}{2\pi\epsilon r_2} \vec{r}_2^0 \\ &= \left( \frac{\tau}{2\pi\epsilon r_1} + \frac{\tau}{2\pi\epsilon r_2} \right) \vec{r}_1^0 \\ &= \frac{\tau}{2\pi\epsilon} \left( \frac{1}{r_1} + \frac{1}{r_2} \right) \vec{r}_1^0 \end{aligned}$$

$$\text{其中 } r_1 = b - (h_1 - a_1), r_2 = b + (h_2 - a_2)$$

$$\text{由题意: } E_A = 50 \times 10^5 \text{ 另外前面求出 } \frac{\tau}{2\pi\epsilon} = 1.65u \text{ 所以可解出 } u = 20.24\text{kv}$$

3.19

(a) 可以

(b) 由  $\alpha = 120^\circ$  不满足  $\alpha = \frac{\pi}{n}$ ,  $n$  为正整数的条件故 (b) 图电场不能用镜像法求出

(c) 可以其镜像法如右图其中  $q' = -\frac{R}{d}q$  位置  $b = \frac{R^2}{d}$

(d) 不可以，由于  $q$  在平面以下，其对应于平面的镜像应当在平面以上，如右图，仍然在求解区域中这与镜像电荷不能处于求解区域的要求不符

(d) 可以

对于  $\epsilon_1$  介质的  $1/4$  空间而言

对于  $\epsilon_2$  介质的  $1/4$  空间而言

3.20

解：设内外导体沿轴线方向单位长度所带电量分别为  $+\tau$  和  $-\tau$ 。由高斯定理，电缆中电介质的电位移  $\vec{D} = \frac{\tau}{2\pi r} \vec{r}^0$

则  $\xi_1$  介质中

$$\vec{E}_1 = \frac{\vec{D}}{\xi_1} = \frac{\tau}{2\pi\xi_1 r} \vec{r}^0 (R_1 < r < R_2)$$

$\xi_2$  介质中

$$\vec{E}_2 = \frac{\vec{D}}{\xi_2} = \frac{\tau}{2\pi\xi_2 r} \vec{r}^0 (R_2 < r < R_3)$$

于是内外导体间的电压

$$U = \int_{R_1}^{R_2} \vec{E}_1 d\vec{r} + \int_{R_2}^{R_3} \vec{E}_2 d\vec{r} = \int_{R_1}^{R_2} E_1 dr + \int_{R_2}^{R_3} E_2 dr = \frac{\tau}{2\pi\xi_1} \ln \frac{R_2}{R_1} + \frac{\tau}{2\pi\xi_2} \ln \frac{R_3}{R_2}$$

$$\text{于是电容 } C = \frac{q}{U} = \frac{\tau}{U} = \frac{\tau}{\frac{\tau}{2\pi\xi_1} \ln \frac{R_2}{R_1} + \frac{\tau}{2\pi\xi_2} \ln \frac{R_3}{R_2}} = \frac{2\pi\xi_1\xi_2}{\xi_2 \ln \frac{R_2}{R_1} + \xi_1 \ln \frac{R_3}{R_2}}$$

若电缆中介质常数为  $\xi_1$ ，则可得：

$$U' = \int_{R_1}^{R_3} E_1 dr = \frac{\tau}{2\rho\xi_1} \ln \frac{R_3}{R_1}$$

$$\therefore C' = \frac{\tau}{U'} = \frac{2\pi\xi_1}{\ln \frac{R_3}{R_1}} = \frac{2\pi\xi_1\xi_2}{\xi_2 \ln \frac{R_3}{R_1}} = \frac{2\pi\xi_1\xi_2}{\xi_2 \ln \frac{R_2}{R_1} + \xi_2 \ln \frac{R_3}{R_2}}$$

$C'$ 与 $C$ 进行对比， $C' > C$ 。

故可以判断：当  $R_3$  趋向无穷大时  $C$  也趋向无穷大。

3.21

解：  $R=1.5 \times 10^{-3} \text{m}$ ,  $h=14 \text{m}$

$$\text{电轴位置 } b^2 = h^2 - R^2 \Rightarrow b = \sqrt{h^2 - R^2} = h$$

以大地为参考点则

$$\varphi_A = \frac{\tau}{2\pi\xi_0} \ln \frac{b + (h - R)}{b - (h - R)} = \frac{\tau}{2\pi\xi_0} \ln \frac{2h}{R}$$

$$\text{即天线与大地之间的电压 } U = \varphi_A \frac{\tau}{2\pi\xi_0} \ln \frac{2h}{R}$$

则单位长度的天线电容

$$C_0 = \frac{\tau}{U} = \frac{2\pi\epsilon_0}{\ln \frac{2h}{R}} = \frac{2 \times 3.14 \times 8.85 \times 10^{-12}}{\ln(\frac{30}{1.5 \times 10^{-3}})} = 5.6 \times 10^{-12} F/m$$

则整个天线的总电容

$$C = C_0 L = 5.6 \times 10^{-12} \times 10 = 56 pF$$

3.22

解:

$$R_1 = 1 \times 10^{-2} m, \quad R_2 = 0.5 \times 10^{-2} m, \quad d = 2 \times 10^{-2} m$$

$$\begin{cases} b^2 = h_1^2 - R_1^2 \\ b^2 = h_2^2 - R_2^2 \\ d = h_1 + h_2 \end{cases} \Rightarrow \begin{cases} h_1 = 1.1875 cm \\ h_2 = 0.8125 cm \\ b = 0.64 cm \end{cases}$$

$$\varphi_A = \frac{\tau}{2\pi\epsilon_0} \ln \frac{b + (h_1 - R_1)}{b - (h_1 - R_1)} = \frac{\tau}{2\pi\epsilon_0} \ln \frac{0.8275}{0.4225}$$

$$\varphi_B = \frac{\tau}{2\pi\epsilon_0} \ln \frac{b - (h_2 - R_2)}{b + (h_2 - R_2)} = \frac{\tau}{2\pi\epsilon_0} \ln \frac{0.3275}{0.9525}$$

两轴间的电压

$$U = \varphi_A - \varphi_B = \frac{\tau}{2\pi\epsilon_0} \ln \frac{0.8275 \times 0.9525}{0.4225 \times 0.3275}$$

$$\text{则单位长度的电容 } C_0 = \frac{\tau}{U} = \frac{2\pi\epsilon_0}{\ln \frac{0.8275 \times 0.9525}{0.4225 \times 0.3275}} = 31.94 pF/m$$

3.23

解:

$$a_1 = 4 \times 10^{-2} m, \quad a_2 = 12 \times 10^{-2} m, \quad d = 5 \times 10^{-2} m$$

$$\begin{cases} b^2 = h_1^2 - R_1^2 \\ b^2 = h_2^2 - R_2^2 \\ d = h_1 + h_2 \end{cases} \Rightarrow \begin{cases} h_1 = 19.5 cm \\ h_2 = 10.3 cm \\ b = 15.3 cm \end{cases}$$

$$\varphi_A = \frac{\tau}{2\pi\epsilon_0} \ln \frac{b + (h_1 - a_1)}{b - (h_1 - a_1)}$$

$$\varphi_B = \frac{\tau}{2\pi\epsilon} \ln \frac{b + (h_2 - a_2)}{b - (h_2 - a_2)}$$

两导体间电压

$$U = \varphi_A - \varphi_B = \frac{\tau}{2\pi\epsilon} \ln \frac{\frac{b + (h_1 - a_1)}{b - (h_1 - a_1)}}{\frac{b + (h_2 - a_2)}{b - (h_2 - a_2)}}$$

$$\text{单位长度的电容 } C_0 = \frac{\tau}{U} = 318.84 \text{ pF/m}$$

若同轴电缆

$$\vec{E} = \frac{\tau}{2\pi\epsilon r} \vec{r}^0 \quad U' = \int_{a_1}^{a_2} \vec{E} d\vec{r} = \int_{a_1}^{a_2} \frac{\tau}{2\pi\epsilon r} dr = \frac{\tau}{2\pi\epsilon} \ln \frac{a_2}{a_1} = \frac{\tau}{2\pi\epsilon} \ln \frac{12}{4}$$

$$\text{于是 } C'_0 = \frac{\tau}{U'} = \frac{2\pi\epsilon}{\ln 3} = 253 \text{ pF/m}$$

显然  $C_0 > C'_0$

3.24

解:

1. 镜像电荷如图所示

则球  $q_1$  上的电位

$$\begin{cases} \varphi_1 = \frac{q_1}{4\pi\epsilon_0 r_1} + \frac{q_2}{4\pi\epsilon_0 (h_1 - h_2 - r_1)} + \frac{-q_2}{4\pi\epsilon_0 (h_1 + h_2 - r_1)} + \frac{-q_1}{4\pi\epsilon_0 (2h_1 - r_1)} \\ \varphi_2 = \frac{q_2}{4\pi\epsilon_0 r_2} + \frac{q_1}{4\pi\epsilon_0 (h_1 - h_2 - r_2)} + \frac{-q_1}{4\pi\epsilon_0 (h_1 + h_2 - r_2)} + \frac{-q_2}{4\pi\epsilon_0 (2h_1 - r_2)} \end{cases}$$

考虑到,  $h_1, h_2 \gg r_1, r_2$ , 故:

$$\begin{cases} \varphi_1 = \frac{q_1}{4\pi\epsilon_0 r_1} + \frac{q_2}{4\pi\epsilon_0 (h_1 - h_2)} + \frac{-q_2}{4\pi\epsilon_0 (h_1 + h_2)} + \frac{-q_1}{4\pi\epsilon_0 (2h_1)} \dots\dots\dots(1) \end{cases}$$

$$\begin{cases} \varphi_2 = \frac{q_2}{4\pi\epsilon_0 r_2} + \frac{q_1}{4\pi\epsilon_0 (h_1 - h_2)} + \frac{-q_1}{4\pi\epsilon_0 (h_1 + h_2)} + \frac{-q_2}{4\pi\epsilon_0 (2h_1)} \dots\dots\dots(2) \end{cases}$$

2. 由带电导体球 1 和 2 以及大地 0 组成的静电系统, 由

$$\begin{cases} \varphi_1 = \alpha_{11}q_1 + \alpha_{12}q_2 \dots\dots\dots(3) \end{cases}$$

$$\begin{cases} \varphi_2 = \alpha_{21}q_1 + \alpha_{22}q_2 \dots\dots\dots(4) \end{cases}$$

将 (1) (2) 与 (3) (4) 对照可得电位系数

$$\begin{cases} \alpha_{11} = \frac{1}{4\pi\epsilon_0 r_1} - \frac{1}{4\pi\epsilon_0 (2h_1)} \\ \alpha_{12} = \frac{1}{4\pi\epsilon_0 (h_1 - h_2)} - \frac{1}{4\pi\epsilon_0 (h_1 + h_2)} \\ \alpha_{21} = \frac{1}{4\pi\epsilon_0 (h_1 - h_2)} - \frac{1}{4\pi\epsilon_0 (h_1 + h_2)} \\ \alpha_{22} = \frac{1}{4\pi\epsilon_0 r_2} - \frac{1}{4\pi\epsilon_0 (2h_1)} \end{cases}$$

(3) (4) 解得:

$$\begin{cases} q_1 = \frac{\alpha_{22}}{\alpha_{11}\alpha_{22} - \alpha_{12}\alpha_{21}}\varphi_1 - \frac{\alpha_{12}}{\alpha_{11}\alpha_{22} - \alpha_{12}\alpha_{21}}\varphi_2 \\ q_2 = \frac{\alpha_{21}}{\alpha_{11}\alpha_{21} - \alpha_{11}\alpha_{22}}\varphi_1 - \frac{\alpha_{11}}{\alpha_{12}\alpha_{21} - \alpha_{11}\alpha_{22}}\varphi_2 \end{cases}$$

于是该静电系统的感应系数

$$\beta_{11} = \frac{\alpha_{22}}{\alpha_{11}\alpha_{22} - \alpha_{12}\alpha_{21}}, \quad \beta_{12} = -\frac{\alpha_{12}}{\alpha_{11}\alpha_{22} - \alpha_{12}\alpha_{21}}$$

$$\beta_{21} = \frac{\alpha_{21}}{\alpha_{11}\alpha_{21} - \alpha_{11}\alpha_{22}}, \quad \beta_{22} = -\frac{\alpha_{11}}{\alpha_{12}\alpha_{21} - \alpha_{11}\alpha_{22}}$$

$$\text{于是由} \begin{cases} C_{i0} = \beta_{i1} + \beta_{i2} + \dots + \beta_{im} \\ C_{ij} = -\beta_{ij} \end{cases}$$

得到系统各个分布电容:

$$\begin{cases} C_{10} = \beta_{11} + \beta_{12} = \frac{\alpha_{22} - \alpha_{12}}{\alpha_{11}\alpha_{22} - \alpha_{12}\alpha_{21}} \\ C_{20} = \beta_{21} + \beta_{22} = \frac{\alpha_{21} - \alpha_{11}}{\alpha_{12}\alpha_{21} - \alpha_{11}\alpha_{22}} \\ C_{12} = -\beta_{12} = \frac{\alpha_{12}}{\alpha_{11}\alpha_{22} - \alpha_{12}\alpha_{21}} \end{cases}$$

系统工作电容:

$$C_p = C_{12} + \frac{C_{10}C_{20}}{C_{10} + C_{20}} = \frac{4\pi\epsilon_0}{\frac{1}{r_1} + \frac{1}{r_2} - \frac{1}{2h_1} - \frac{1}{2h_2} - \frac{2}{h_1 - h_2} + \frac{2}{h_1 + h_2}}$$

3.25

解:

$$r_0 = 0.005m, \quad h_1 = 10m, \quad h_2 = 8m, \quad d = 2m$$



设导线 1, 2 单位长度电荷量为  $\tau_1$  和  $\tau_2$ , 根据镜像法, 考虑到地面影响及  $r_0 \ll d, h_1, h_2$

可近似认为电轴与导线的几何轴线重合, 则导线 1, 2 电势分别为

$$\varphi_1 = \frac{\tau_1}{2\pi\epsilon_0} \ln \frac{2h_1}{r_0} + \frac{\tau_2}{2\pi\epsilon_0} \ln \frac{\sqrt{d^2 + (h_1 + h_2)^2}}{\sqrt{d^2 + (h_1 - h_2)^2}}$$

$$\varphi_2 = \frac{\tau_2}{2\pi\epsilon_0} \ln \frac{2h_2}{r_0} + \frac{\tau_1}{2\pi\epsilon_0} \ln \frac{\sqrt{d^2 + (h_1 + h_2)^2}}{\sqrt{d^2 + (h_1 - h_2)^2}}$$

$$\text{故由 (*)} \begin{cases} \varphi_1 = \alpha_{11}\tau_1 + \alpha_{12}\tau_2 \\ \varphi_2 = \alpha_{21}\tau_1 + \alpha_{22}\tau_2 \end{cases}$$

得到导线系统电势系数

$$\alpha_{11} = \frac{1}{2\pi\epsilon_0} \ln \frac{2h_1}{r_0}$$

$$\alpha_{12} = \alpha_{21} = \frac{1}{2\pi\epsilon_0} \ln \frac{\sqrt{d^2 + (h_1 + h_2)^2}}{\sqrt{d^2 + (h_1 - h_2)^2}}$$

$$\alpha_{22} = \frac{1}{2\pi\epsilon_0} \ln \frac{2h_2}{r_0}$$

由  $[\varphi] = [\alpha][\tau] \Rightarrow [\tau] = [\alpha]^{-1}[\varphi] = [\beta][\varphi]$  可得出感应系数  $[\beta]$

同样解方程组 (\*) 可得:

$$\begin{cases} \tau_1 = \frac{1}{\alpha_{11}\alpha_{12} - \alpha_{12}\alpha_{21}} (\alpha_{22}\varphi_1 - \alpha_{12}\varphi_2) \\ \tau_2 = \frac{1}{\alpha_{12}\alpha_{21} - \alpha_{11}\alpha_{22}} (\alpha_{21}\varphi_1 - \alpha_{11}\varphi_2) \end{cases}$$

于是由

$$\begin{cases} \tau_1 = \beta_{11}\varphi_1 + \beta_{12}\varphi_2 \\ \tau_2 = \beta_{21}\varphi_1 + \beta_{22}\varphi_2 \end{cases}$$

$$\text{得 } \beta_{11} = \frac{\alpha_{22}}{\alpha_{11}\alpha_{12} - \alpha_{12}\alpha_{21}}, \quad \beta_{12} = \frac{-\alpha_{12}}{\alpha_{11}\alpha_{12} - \alpha_{12}\alpha_{21}}$$

$$\beta_{21} = \frac{\alpha_{21}}{\alpha_{12}\alpha_{22} - \alpha_{11}\alpha_{22}}, \quad \beta_{22} = \frac{-\alpha_{11}}{\alpha_{12}\alpha_{22} - \alpha_{11}\alpha_{22}}$$

$$\begin{cases} \tau_1 = (\beta_{11} + \beta_{12})\varphi_1 - \beta_{12}(\varphi_1 - \varphi_2) = C_{10}U_{10} + C_{12}U_{12} \\ \tau_2 = -\beta_{21}(\varphi_2 - \varphi_1) + (\beta_{21} + \beta_{22})\varphi_2 = C_{21}U_{21} + C_{20}U_{20} \end{cases}$$


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$$\text{于是} \begin{cases} C_{10} = \beta_{11} + \beta_{12} \\ C_{12} = -\beta_{12} \\ C_{21} = -\beta_{21} \\ C_{20} = \beta_{21} + \beta_{22} \end{cases}$$

下面其工作电容:

$$C_p = C_{12} + \frac{C_{10}C_{20}}{C_{10} + C_{20}}$$

得  $C_{10} = 5.64 \text{ pF/m}$ ,  $C_{12} = C_{21} = 1.63 \text{ pF/m}$ ,  $C_{10} = 5.44 \text{ pF/m}$ ,  $C_p = 4.4 \text{ pF/m}$

3.26

解:

$$\begin{cases} w_e = \frac{1}{2} \sum \varphi_k q_k \\ w_e = \frac{1}{2} \int_{v'} \varphi \rho dv + \frac{1}{2} \int_{s'} \varphi \sigma ds \\ w_e = \frac{1}{2} \int_v \vec{E} \cdot \vec{D} dv \end{cases}$$

介电常数为  $\epsilon_0$  的均匀带电球,  $\rho = \frac{q}{\frac{4}{3}\pi a^3}$ . 球内有电场, 电场仍然球对称应用高

$$\text{斯定理可得电场} \begin{cases} \vec{E} = \frac{r\rho}{3\epsilon_0} \vec{r}^0 \dots\dots\dots (r < a) \\ \vec{E} = \frac{a^3 q}{3\epsilon_0 r^2} \vec{r}^0 \dots\dots\dots (r > a) \end{cases}$$

$$\text{电位} \begin{cases} \varphi = \frac{a^3 \rho}{3\epsilon_0 r} \dots\dots\dots (r > a) \\ \varphi = \frac{\rho}{2\epsilon_0} (a^2 - \frac{r^2}{3}) \dots\dots\dots (r < a) \end{cases}$$

方法 1:

$$\begin{aligned} w_e &= \frac{1}{2} \int_{v'} \varphi \rho dv = \frac{\rho}{2\epsilon_0} (a^2 - \frac{r^2}{3}) = \frac{\rho^2}{4\epsilon_0} \iiint_V (a^2 - \frac{r^2}{3}) r^2 \sin \theta d\theta d\varphi dr \\ &= \frac{\rho^2}{4\epsilon_0} \int_0^\pi \sin \theta d\theta \int_{-\pi}^\pi d\varphi \int_0^a (a^2 - \frac{r^2}{3}) r^2 dr = \frac{\rho^2}{4\epsilon_0} \cdot 4\pi \int_0^a (a^2 - \frac{r^2}{3}) r^2 dr \\ &= \frac{\pi \rho^2}{\epsilon_0} [\int_0^a a^2 r^2 dr - \int_0^a \frac{r^4}{3} dr] = \frac{\pi \rho^2}{\epsilon_0} [a^2 \times \frac{1}{3} r^3 \Big|_0^a - \frac{1}{15} r^5 \Big|_0^a] \end{aligned}$$

$$= \frac{4\pi\rho^2 a^5}{15\epsilon_0} = \frac{4\pi(\frac{3q}{4\pi a^3})^2 a^5}{15\epsilon_0} = \frac{3q^2}{20\pi\epsilon_0 a}$$

方法 2:

$$\begin{aligned} w_e &= \frac{1}{2} \int_v \vec{E} \cdot \vec{D} dv \\ &= \frac{1}{2} \int_{\text{球内}} \frac{r\rho}{3} \cdot \frac{r\rho}{3\epsilon_0} dr + \frac{1}{2} \int_{\text{球外}} \frac{a^3\rho}{3r^2} \cdot \frac{a^3\rho}{3\epsilon_0 r^2} dr \\ &= \frac{1}{2} \iiint_{\text{球内}} \frac{r^2\rho^2}{9\epsilon_0} r^2 \sin\theta d\theta d\varphi dr + \frac{1}{2} \iiint_{\text{球外}} \frac{a^6\rho^2}{9\epsilon_0 r^4} r^2 \sin\theta d\theta d\varphi dr \\ &= \frac{1}{2} \int_0^\pi \sin\theta d\theta \int_{-\pi}^\pi d\varphi \int_0^a \frac{r^2\rho^2}{9\epsilon_0} r^2 dr + \frac{1}{2} \int_0^\pi \sin\theta d\theta \int_{-\pi}^\pi d\varphi \int_a^\infty \frac{a^6\rho^2}{9\epsilon_0 r^4} r^2 dr \\ &= \frac{4a^5\rho^2}{15\epsilon_0} = \frac{3q^2}{20\pi\epsilon_0 a} \end{aligned}$$

3.27

解: 如图所示, 由导体 1, 2 和大地构成的静电系统中的储能壳视为右图中三个电容中的储能之和。显然对于  $C_{10}$  其电压为  $U_{10} = \varphi_1$ , 对于  $C_{20}$  其电压为  $U_{20} = \varphi_2$ , 而  $C_{12}$  上的电压为  $U_{12} = \varphi_1 - \varphi_2$ 。

因此, 整个系统中静电能

$$w_e = \frac{1}{2} C_{10} U_{10}^2 + \frac{1}{2} C_{20} U_{20}^2 + \frac{1}{2} C_{12} U_{12}^2 = \frac{1}{2} C_{10} \varphi_1^2 + \frac{1}{2} C_{20} \varphi_2^2 + \frac{1}{2} C_{12} (\varphi_1 - \varphi_2)^2$$

即证。

3.28

解: 设同轴电缆内外两导体上每单位长度的电荷为  $\pm\tau$ , 则电介质中的电场强度

$$\vec{E} = \frac{\tau}{2\pi\epsilon_0 r} \vec{r}^0, \quad \text{电位差 } \vec{D} = \epsilon_0 \vec{E} = \frac{\tau}{2\pi} \vec{r}^0$$

假定内外导体半径分别为  $a, b$  则  $a=5\text{mm}, b=10\text{mm}$

则电介质中场强最大在内导体表面  $E_{\max} = \frac{\tau}{2\pi\epsilon_0 a}$ , 令  $E_{\max} = E_{\text{击穿}}$ , 则  $\tau = 2\pi\epsilon_0 a E_{\text{击穿}}$

将上式带入  $\vec{E}$  和  $\vec{D}$  的表达式中, 并计算 1 公里长电缆静电能量

$$w_e = \frac{1}{2} \int_v \vec{E} \cdot \vec{D} dv = \frac{1}{2} \int_a^b (\vec{E} \cdot \vec{D}) \cdot 2\pi r L dr = \frac{1}{2} \int_a^b \left( \frac{\tau}{2\pi\epsilon_0 r} \cdot \frac{\tau}{2\pi} \right) \cdot 2\pi r L dr = \frac{\tau^2 L}{4\pi\epsilon_0} \ln \frac{b}{a}$$

将  $a = 5 \times 10^{-3} \text{ m}$  ,  $b = 10 \times 10^{-3} \text{ m}$  ,  $E_{\text{击穿}} = 200 \times 10^3 \times 10^3 \text{ V/m}$

$$\xi = 5\xi_0 = 5 \times 8.85 \times 10^{-12} \text{ F/m}, \quad L = 1 \text{ km} = 10^3 \text{ m}$$

代入的:  $w_e = 965 \text{ J}$

3.29

解: 在介质电轴出之前, 电容器中的静电能为

$$w_e = \frac{1}{2} CU^2 = \frac{1}{2} \times 1 \times 10^{-12} \times (10^3)^2 = 0.5 \times 10^{-6} \text{ (J)}$$

电容器的极板上的电荷电量

$$Q = CU = 1 \times 10^{-12} \times 10^3 = 10^{-9} \text{ (C)}$$

(1) 在电源断开的情况下, 抽出电介质, 极板上的电量不变, 但电压会变化。

抽出电介质电后, 电容器的电容为

$$C' = \frac{\xi_0 S}{d} = \frac{1}{5} \frac{\xi S}{d} = \frac{1}{5} C \quad (\text{其中 } c \text{ 为原来的电容})$$

$$\therefore C' = 0.2 \times 10^{-12} \text{ F}$$

因此, 在电源断开情况下, 抽出电介质电, 电容器中的静电变为

$$w'_e = \frac{1}{2} \frac{Q^2}{C'} = \frac{1}{2} \times \frac{(10^{-9})^2}{0.2 \times 10^{-12}} = 2.5 \times 10^{-6} \text{ J}$$

在此过程中电容器的静电能增加了

$$\Delta w'_e = w'_e - w_e = 2 \times 10^{-6} \text{ J}$$

由于外电源不提供能量, 因此电容器中静电能增加说明了外力作正功, 而电场力作负功。所得外力做功

$$w_{\text{外}} = \Delta w'_e = 2 \times 10^{-6} \text{ J}$$

(2) 由于电源一直连在电容器上, 因而电容器的电压不变, 当电介质电抽出之后, 电容器中静电能

$$w''_e = \frac{1}{2} C' U^2 = \frac{1}{2} \times 0.2 \times 10^{-12} \times (10^3)^2 = 0.1 \times 10^{-6} \text{ (J)}$$

电容器静电能增加了

$$\Delta w''_e = w''_e - w_e = -0.4 \times 10^{-6} \text{ (J)}$$

其增加量为负。

我们知道在电源一直连在电容器上时, 电源提供的能量一半用来作电场储能的增量, 另一半用于电场力对外做功。因此可以推断, 本小题中电源从电容器中吸收了  $2 \times 0.4 \times 10^6 \text{ J}$  的能量, 其中一半是电容器提供的, 另一半是外力做功提供的。即  $w'_{\text{外}} = 0.4 \times 10^6 \text{ J}$ , 电场力

作正功，外力作负功。

3.30

解：设单位长度的内外导体圆柱上的电量为 $\pm\tau$ ，则两导体间的电压

$$U = \int_{R_1}^{R_2} \vec{E} \cdot d\vec{r} = \int_{R_1}^{R_2} \frac{\tau}{2\pi\epsilon_0 r} dr = \frac{\tau}{2\pi\epsilon_0} \ln \frac{R_2}{R_1}$$

则两导体间单位长度上的电容

$$C_0 = \frac{\tau}{U} = \frac{2\pi\epsilon_0}{\ln \frac{R_2}{R_1}}$$

则当电容器总长为 $L$ 时，其总电容为 $C = C_0 L = \frac{2\pi L \epsilon_0}{\ln \frac{R_2}{R_1}}$

$$\text{电容器中的静电能 } w_e = \frac{1}{2} C U^2 = \frac{\pi L \epsilon_0}{\ln \frac{R_2}{R_1}} U^2$$

小圆柱所受到的轴向吸力

$$\begin{aligned} f &= + \frac{\partial w_e}{\partial L} \Big|_{\varphi = \cos t} = + \frac{\partial}{\partial L} \left( \frac{\pi L \epsilon_0}{\ln \frac{R_2}{R_1}} U^2 \right) = + \frac{\pi \epsilon_0}{\ln \frac{R_2}{R_1}} U^2 \\ &= \frac{\pi \times 8.85 \times 10^{-12}}{\ln \left( \frac{6 \times 10^{-2}}{5 \times 10^{-2}} \right)} \times (10^3)^2 = 152.49 \times 10^{-6} N = 1.5249 \times 10^{-4} N \end{aligned}$$

3.31

解：选 $\theta$ 角为广义坐标，则其对应的广义力为轴动力矩， $n=5$ ，说明总电容为10个小电容并联而成

$$C = 10 \times C_1 = 10 \times \frac{\epsilon_0 S}{d} = 10 \times \frac{\epsilon_0 \times \frac{1}{2} \theta R^2}{\delta} = 5 \epsilon_0 \theta R^2 / \delta$$

其总静电储能

$$w_e = \frac{1}{2} C U^2 = \frac{1}{2} \times 5 \times \frac{\epsilon_0 \theta R^2}{\delta} U^2$$

$$\text{轴动力矩: } M = f = + \frac{\partial w_e}{\partial \theta} \Big|_{\varphi = \cos t} = 2.5 \times \frac{\epsilon_0 R^2}{\delta} U^2 = 1.593 \times 10^{-4} (N \cdot m) \quad M \text{ 为正, 有}$$

使 $\theta$ 增大的标量

3.32

解：液体电介质上升是由于电场力的作用，并且该电场力与液面上的电介质液体的重力和平衡，求电场力用坐标平移法。

$$f = \frac{\partial w_e}{\partial q} \Big|_{q_k - c_k} = - \frac{\partial w_e}{\partial q} \Big|_{\varphi_k - c_k}$$

先求  $w_e$

$$w_e = \frac{1}{2} C U_0^2 \Rightarrow C = ?$$

电容器的电容：

$$C = \frac{\xi_0 S_1}{d} + \frac{\xi S_2}{d} = \frac{\xi_0 a(l-x)}{d} + \frac{\xi ax}{d}$$

$$w_e = \frac{1}{2} C U_0^2 = \frac{1}{2} \times \left[ \frac{\xi_0 a(l-x)}{d} + \frac{\xi ax}{d} \right] \times U_0^2$$

因此：

$$= \frac{1}{2} \frac{a U_0^2}{d} [\xi_0(l-x) + \xi x]$$

选广义坐标为  $x$ ，则液体受到沿  $x$  方向的电场力为：

$$f_x = \frac{\partial w_e}{\partial x} = \frac{\partial}{\partial x} \left[ \frac{1}{2} \frac{a U_0^2}{d} (\xi_0(l-x) + \xi x) \right] = \frac{U_0^2}{2} \cdot \frac{a}{d} (\xi - \xi_0)$$

此力与水平面以上液体重力和平衡，即：

$$f_x = \frac{a U_0^2}{2d} (\xi - \xi_0) = \rho_m g a d h \Rightarrow \frac{U_0^2}{2d} (\xi - \xi_0) = \rho_m g d h$$

$$\text{于是 } \xi = \xi_0 + \frac{2g\rho_m h d^2}{U_0^2}$$