

## Quantum field theory in the solid state, Exercise sheet 7

Corrections: Monday June 23<sup>rd</sup>

### Berry phase

The Berry phase is a very general concept present in the classical and quantum worlds. It refers to a phase that is acquired when a system is transported adiabatically around a closed path. For example the Aharonov-Bohm effect is a Berry phase. A Dirac cone is characterized by a Berry phase. Here we will discuss two further Berry phases for boson and spin-1/2 systems. For your convenience I have uploaded some notes for general definition of the Berry phase, in the realm of quantum mechanics.

#### 1. Bosons

(a) In Exercise sheet 5, we considered the real time coherent state path integral for the harmonic oscillator. Carry out the Wick rotation to show that the partition reads:

$$Z = \int D\{\alpha\} e^{-\int_0^\beta d\tau \left[ \left( \bar{\alpha} \frac{\partial}{\partial \tau} \alpha - \alpha \frac{\partial}{\partial \tau} \bar{\alpha} \right) + \omega \bar{\alpha} \alpha \right]} \quad (1)$$

with  $\alpha(\beta) = \alpha(0)$ .

(b) The first term is the Berry phase. Show that it can be written as

$$\int_0^\beta d\tau \left[ \bar{\alpha} \frac{\partial}{\partial \tau} \alpha - \alpha \frac{\partial}{\partial \tau} \bar{\alpha} \right] \equiv \int_0^\beta d\tau \left( \langle \alpha(\tau) | \frac{d}{d\tau} | \alpha(\tau) \rangle - \text{H.c.} \right) = i \oint_\gamma \mathbf{A} \cdot d\mathbf{l}. \quad (2)$$

Compute the corresponding *magnetic* field  $\mathbf{B} = \nabla \times \mathbf{A}$ . Hint: express the complex variable  $\alpha$  in terms of a vector in  $\mathbb{R}^2$ .

#### 2. Coherent states for a spin 1/2.

Consider a spin 1/2 degree of freedom where the Hilbert space is spanned by the states  $\{|\uparrow\rangle, |\downarrow\rangle\}$ . The spin operator satisfies the algebra  $[\hat{S}^\alpha, \hat{S}^\beta] = i \sum_{\gamma=1}^3 \epsilon^{\alpha\beta\gamma} \hat{S}^\gamma$  with  $\hat{\mathbf{S}}^2 = s(s+1)$  with  $s = 1/2$ .

(a) Find normalized states,  $|\mathbf{n}\rangle$ , such that:

$$\langle \mathbf{n} | \hat{\mathbf{S}} | \mathbf{n} \rangle = s\mathbf{n}. \quad (3)$$

These states correspond to spin coherent states, and  $\mathbf{n}$  is a vector on the unit sphere. Hint: Consider  $|\mathbf{n}\rangle = z_1 |\uparrow\rangle + z_2 |\downarrow\rangle$  with  $z \in \mathbb{C}$ . With this Ansatz  $\mathbf{n} = \mathbf{z}^\dagger \boldsymbol{\sigma} \mathbf{z}$  with  $\mathbf{z}^\dagger = (\bar{z}_1, \bar{z}_2)$ .

(b) Show that a resolution of unity reads:

$$\int_{S^2} \frac{d^2 \mathbf{n}}{2\pi} |\mathbf{n}\rangle \langle \mathbf{n}| = \hat{1}. \quad (4)$$

Here the integration runs over the unit sphere  $S^2$ .

(c) Consider the Hamiltonian

$$\hat{H} = \mathbf{B}(t) \cdot \hat{S} \quad (5)$$

describing the spin degree of freedom in a magnetic field. Show that the real time evolution can be expressed as:

$$\langle \mathbf{n}_b | \hat{U}(t, 0) | \mathbf{n}_a \rangle = \int \prod_{n=1}^{N-1} \frac{d^2 \mathbf{n}_n}{2\pi} e^{iS(\{\mathbf{n}_n\})} \quad (6)$$

with

$$S = \int_0^t dt' \left[ i \langle \mathbf{n}(t') | \frac{d}{dt'} | \mathbf{n}(t') \rangle - s \mathbf{B}(t') \cdot \mathbf{n}(t') \right] \quad (7)$$

The first term corresponds to the Berry phase:

$$S_B = \int_0^t dt' i \langle \mathbf{n}(t') | \frac{d}{dt'} | \mathbf{n}(t') \rangle. \quad (8)$$

(d) Show that the Berry phase is a gauge dependent object for open paths, but has a well defined meaning for closed ones.

(e) Consider a closed path. Show that

$$S_B = \int_{\Omega} d\Omega \cdot \mathbf{B} \quad \text{with} \quad \mathbf{B}(\mathbf{x}) = \frac{s}{r^2} \hat{n} \quad (9)$$

In the above  $\mathbf{x} = r\mathbf{n}$  with  $\mathbf{n}$  a unit vector, and  $\Omega$  is the domain on the unit sphere enclosed by the path. The above is valid for a general value of  $s$ . For things to be consistent one will require that  $4\pi s = 1$  thus leading to the known quantization of  $s$ .

# 1. Bosons

(a) In Exercise sheet 5, we considered the real time coherent state path integral for the harmonic oscillator. Carry out the Wick rotation to show that the partition reads:

$$Z = \int D\{\alpha\} e^{-\int_0^\beta d\tau \left[ \left( \bar{\alpha} \frac{\partial}{\partial \tau} \alpha - \alpha \frac{\partial}{\partial \tau} \bar{\alpha} \right) + \omega \bar{\alpha} \alpha \right]} \quad (1)$$

with  $\alpha(\beta) = \alpha(0)$ .

(b) The first term is the Berry phase. Show that it can be written as

$$\int_0^\beta d\tau \left[ \bar{\alpha} \frac{\partial}{\partial \tau} \alpha - \alpha \frac{\partial}{\partial \tau} \bar{\alpha} \right] \equiv \int_0^\beta d\tau \left( \langle \alpha(\tau) | \frac{d}{d\tau} | \alpha(\tau) \rangle - \text{H.c.} \right) = i \oint_\gamma \mathbf{A} \cdot d\mathbf{l}. \quad (2)$$

Compute the corresponding magnetic field  $\mathbf{B} = \nabla \times \mathbf{A}$ . Hint: express the complex variable  $\alpha$  in terms of a vector in  $\mathbb{R}^2$ .

$$\langle \alpha_b | e^{-it\hat{H}} | \alpha_a \rangle \propto \int D\{\alpha(t)\} e^{i \int_{t_a}^{t_b} dt L(\alpha, \dot{\alpha}, t)}$$

$$L(\alpha, \dot{\alpha}, t) = \frac{i}{2} (\bar{\alpha} \dot{\alpha} - \alpha \dot{\bar{\alpha}}) - \omega \bar{\alpha} \alpha$$

$$t = -i\tau$$

$$\frac{d}{dt} = i \frac{d}{d\tau}$$

$$i \int_{t_a}^{t_b} \frac{i}{2} (\bar{\alpha} \dot{\alpha} - \alpha \dot{\bar{\alpha}}) - \omega \bar{\alpha} \alpha dt$$

$$= i \int_0^\beta -\frac{1}{2} (\bar{\alpha} \partial_\tau \alpha - \alpha \partial_\tau \bar{\alpha}) - \omega \bar{\alpha} \alpha (-i d\tau)$$

$$= - \int_0^\beta \frac{1}{2} (\bar{\alpha} \partial_\tau \alpha - \alpha \partial_\tau \bar{\alpha}) + \omega \bar{\alpha} \alpha d\tau$$

$$\int_0^\beta (\bar{\alpha} \partial_\tau \alpha - \alpha \partial_\tau \bar{\alpha}) d\tau =$$

$$\langle \alpha | \frac{d}{d\tau} | \alpha \rangle =$$

2. Coherent states for a spin 1/2.

Consider a spin 1/2 degree of freedom where the Hilbert space is spanned by the states  $\{|\uparrow\rangle, |\downarrow\rangle\}$ . The spin operator satisfies the algebra  $[\hat{S}^\alpha, \hat{S}^\beta] = i \sum_{\gamma=1}^3 \epsilon^{\alpha\beta\gamma} \hat{S}^\gamma$  with  $\hat{\mathbf{S}}^2 = s(s+1)$  with  $s = 1/2$ .

(a) Find normalized states,  $|\mathbf{n}\rangle$ , such that:

$$\langle \mathbf{n} | \hat{\mathbf{S}} | \mathbf{n} \rangle = s \mathbf{n}. \tag{3}$$

These states correspond to spin coherent states, and  $\mathbf{n}$  is a vector on the unit sphere. Hint: Consider  $|\mathbf{n}\rangle = z_1 |\uparrow\rangle + z_2 |\downarrow\rangle$  with  $z \in \mathbb{C}$ . With this Ansatz  $\mathbf{n} = \mathbf{z}^\dagger \boldsymbol{\sigma} \mathbf{z}$  with  $\mathbf{z}^\dagger = (\bar{z}_1, \bar{z}_2)$ .

$$|\vec{n}\rangle = z_1 |\uparrow\rangle + z_2 |\downarrow\rangle$$