

## Problem Sheet 7

for the tutorial on June 20th, 2025

## Quantum Mechanics II

Summer term 2025

Sheet handed out on June 10th, 2025; to be handed in on June 17th, 2025 until 2 pm

## Exercise 7.1: Laboratory and center-of-mass systems

[2 + 2 + 1 + 4 + 2 P.]

We consider a non-relativistic collision between a projectile particle A of mass  $m_A$  and a target particle of mass  $m_B$  like in the lecture. The laboratory system L is the frame in which the target particle B is at rest before the collision. The center-of-mass system CM is the coordinate system in which the center of mass of the composite system (A+B) is always at rest. In that system the projectile A and target particle B move initially with respect to the center of mass C with equal and opposite momenta,  $\vec{p}_A = -\vec{p}_B = \vec{p}$ , as illustrated in Fig. 1. With respect to the laboratory frame, the center of mass of the two particles moves throughout the collision with a constant velocity  $\vec{v}_c$  along the direction of incidence, with  $\vec{v}_c = \vec{q}_A/(m_A + m_B)$ , where  $\vec{q}_A$  is the momentum of particle A before the collision in the laboratory system.

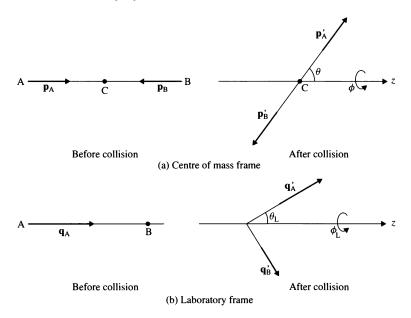


Figure 1: Elastic scattering of a projectile A by target B.

a) Using a Galilean transformation, find the relation between the momenta  $\vec{q}_A$  and  $\vec{p}_A$  before the collision.

- b) After the collision the two particles A and B emerge with equal and opposite momenta  $\vec{p'}_A = -\vec{p'}_B = \vec{p'}$  in the CM frame. In the following we consider an elastic collision, such that the magnitude of the momenta of each particle remain the same p' = p. Relate the components of the momentum of projectile A along the direction of incidence in the two coordinate systems using the two scattering angles  $\theta_L$  and  $\theta$  as illustrated in Fig. 1.
- c) Show that the relation between the two scattering angles is given by

$$\tan \theta_L = \frac{\sin \theta}{\cos \theta + \frac{m_A}{m_B}}.$$
 (1)

- d) Using the equation above, find the relation between the angular differential cross sections  $\frac{d\sigma}{d\Omega}$  in the laboratory and the center-of-mass frames.
- e) Let us now consider a numerical example:

Two beams of protons intersect collinearly. If the kinetic energy of the protons is 5 keV in both beams, calculate:

- i) the magnitude of the relative velocity of a proton in one beam with respect to a proton in the other one,
- ii) the energy in the centre-of-mass system.

## Exercise 7.2: Partial waves and phase shifts

[6 + 8 P.]

We consider in the following the scattering by a central potential V(r) such that the system is completely symmetrical about the direction of incidence, which we choose to be the z-axis. In this case both the wave function  $\psi_k$  and the scattering amplitude f do not depend on the azimuthal angle  $\varphi$ . We then expand them in a series of Legendre polynomials, which form a complete set in the interval  $-1 \le \cos \theta \le 1$ ,

$$\psi_{\mathbf{k}}(r,\theta) = \sum_{l=0}^{\infty} R_l(k,r) P_l(\cos\theta), \qquad (2)$$

$$f(k,\theta) = \sum_{l=0}^{\infty} f_l(k) P_l(\cos \theta).$$
 (3)

Each term in the series is known as a partial wave and is a simultaneous eigenfunction of the operators  $\vec{L}^2$  and  $L_z$  belonging to eigenvalues  $l(l+1)\hbar^2$  and zero, respectively. The radial wave function for the far region where the potential can be neglected is given by a linear combination of Bessel and Neumann functions  $j_l(kr)$  and  $n_l(kr)$ 

$$R_l(k,r) = B_l(k)j_l(kr) + C_l(k)n_l(kr)$$

$$\tag{4}$$

with coefficients B(k) and C(k). Using the asymptotic expressions for the Bessel and Neumann functions given in the lecture, this leads to

$$R_l(k,r) \stackrel{r \to \infty}{\approx} \frac{1}{kr} \left[ B_l(k) \sin\left(kr - \frac{l\pi}{2}\right) - C_l(k) \cos\left(kr - \frac{l\pi}{2}\right) \right].$$
 (5)

a) In the lecture, we have considered a particular choice for the coefficients  $B_l$  and  $C_l$ . Here we consider a more general case. To this end, it is convenient to rewrite the expression above as

$$R_l(k,r) \stackrel{r \to \infty}{\approx} A_l(k) \frac{1}{kr} \sin\left(kr - \frac{l\pi}{2} + \delta_l(k)\right).$$
 (6)

Determine the expressions of the amplitudes  $A_l(k)$  and the phase shifts  $\delta_l(k)$  introduced above as a function of  $B_l$  and  $C_l$ . The phase shifts  $\delta_l(k)$  are real quantities and characterize the strength of the scattering in the lth partial wave by the potential V(r) at the energy  $E = \hbar^2 k^2/(2m)$ .

b) We would like now to relate the phase shifts  $\delta_l(k)$  to the partial wave amplitudes  $f_l(k)$  and to the scattering amplitude  $f(k,\theta)$  in Eq. (3). Use the radial wave function determined above and the general relation between wave function and scattering amplitude

$$\psi_{\mathbf{k}}(r) \stackrel{r \to \infty}{\approx} e^{i\mathbf{k} \cdot \mathbf{r}} + f(k, \theta) \frac{e^{ikr}}{r}$$
 (7)

to determine the partial wave amplitudes  $f_l(k)$ . Write then the expression of the scattering amplitude  $f(k, \theta)$  as a function of the phase shifts.

**Hint:** Use the plane wave expansion

$$e^{i\mathbf{k}\cdot\mathbf{r}} = \sum_{l=0}^{\infty} (2l+1)i^l j_l(kr) P_l(\cos\theta).$$
 (8)

