

Quantum field theory in the solid state, Exercise sheet 2

Corrections: Monday 12th of May

The periodic Anderson and Kondo lattice models: symmetries and strong coupling limit.

The periodic Anderson model (PAM) describes the physics of a wide band that hybridizes with a narrow band of nearly localized electrons. To define the model we will consider a hyper-cubic lattice of linear length L in d dimensions with 2 orbitals per unit cell, one extended, the c -orbital, and one localized, the f -orbital. Since the f -orbital is localized the Coulomb repulsion will have to be taken into account. The periodic $SU(N)$ generalization of Anderson model describes this situation and reads:

$$\hat{H}_{\text{PAM}} = -t \sum_{\langle i,j \rangle} \sum_{\sigma=1}^N \left(\hat{c}_{i,\sigma}^\dagger \hat{c}_{j,\sigma} + \text{H.c.} \right) + V \sum_i \sum_{\sigma=1}^N \left(\hat{c}_{i,\sigma}^\dagger \hat{f}_{i,\sigma} + \text{H.c.} \right) + \frac{U_f}{N} \sum_i \left(\hat{n}_i^f - \frac{N}{2} \right)^2. \quad (1)$$

Here $\hat{n}_i^f = \sum_{\sigma=1}^N \hat{f}_{i,\sigma}^\dagger \hat{f}_{i,\sigma}$.

1. Symmetries

(a) Show that the PAM enjoys the same global symmetries, particle number, total spin, total momentum, as the Hubbard model that we discussed in class.

(b) Let $\hat{c}_{\mathbf{k},\sigma}^\dagger = \frac{1}{\sqrt{L^d}} \sum_{\mathbf{i}} e^{i\mathbf{k} \cdot \mathbf{i}} \hat{c}_{\mathbf{i},\sigma}^\dagger$ and an equivalent form holds for the f -operator. For the $SU(2)$ case, under which conditions does

$$\langle \hat{c}_{\mathbf{k}_1,\sigma_1}^\dagger \hat{c}_{\mathbf{k}_2,\sigma_2} \hat{c}_{\mathbf{k}_3,\sigma_3}^\dagger \hat{f}_{\mathbf{k}_4,\sigma_4} \rangle \quad (2)$$

not vanish?, Here, for an operator \hat{O} , $\langle \hat{O} \rangle = \frac{\text{Tr}(e^{-\beta \hat{H}_{\text{PAM}}} \hat{O})}{\text{Tr}(e^{-\beta \hat{H}_{\text{PAM}}})}$ and β is the inverse temperature.

2. Non-interacting limit

Consider the specific case of $U_f = 0$, and $d=1$. Can you diagonalize the Hamiltonian and plot the band structure.

3. Strong coupling limit

(a) Consider a single site periodic Anderson model in the Hilbert with a total of two particles for the $SU(2)$ case. Use symmetries to diagonalize the Hamiltonian.

(b) Analyze the strong coupling limit, $U_f \gg t$ and show that the low energy spectrum is

the same as that of the single site Kondo lattice model:

$$\hat{H}_{KLM} = -t \sum_{\langle i,j \rangle} \left(\hat{\mathbf{c}}_i^\dagger \hat{\mathbf{c}}_j + \text{H.c.} \right) + \frac{2J_K}{N} \sum_i \mathbf{c}_i^\dagger \frac{\boldsymbol{\sigma}}{2} \mathbf{c}_i \cdot \hat{\mathbf{S}}_i \quad (3)$$

with $\hat{\mathbf{S}}_i = \frac{1}{2} \hat{\mathbf{f}}_i^\dagger \boldsymbol{\sigma} \hat{\mathbf{f}}_i$ and

$$\boldsymbol{\sigma} = \left(\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \right). \quad (4)$$

Here we have used a spinor notation $\hat{\mathbf{c}}_i^\dagger = (\hat{c}_{i,\uparrow}^\dagger, \hat{c}_{i,\downarrow}^\dagger)$ and $\hat{\mathbf{f}}_i^\dagger = (\hat{f}_{i,\uparrow}^\dagger, \hat{f}_{i,\downarrow}^\dagger)$. The Kondo lattice model is defined in the Hilbert space where $\hat{\mathbf{f}}_i^\dagger \hat{\mathbf{f}}_i = 1$.