

Introduction of three-dimensionally expanded wave function

In the beginning, I propose to expand the wave function into three-dimensional. Due to this operation, we can express the width of a photon's orbital in it. The following is the method.

Firstly, we think how to rotate the wave function φ degrees (Fig. 2a). If we can find the value of point ψ' , we can expand the wave function into three-dimensional since the point ψ is the wave function.

$$\psi = R(\cos\theta + i\sin\theta) (\theta = \omega t - k_x x - k_y y - k_z z) \quad (1)$$

$$\left[\begin{array}{l} R: \text{the existence probability of a photon,} \\ \omega: \text{the number of vibrations, } k_x, k_y, k_z: \text{the wave number,} \\ t: \text{the measurable time variable, } x, y, z: \text{the measurable space variables} \end{array} \right]$$

Well, let us draw a perpendicular from ψ' to $O\psi$ (origin at the point O). The length is the \tilde{z} -coordinate of ψ' and the coordinates of the intersection ψ'' are the \tilde{x} - and \tilde{y} -coordinate of ψ' (Fig. 2b, Fig. 2c). Please note, however, that \tilde{x} , \tilde{y} and \tilde{z} are unmeasurable space variables in the spherical coordinate system.

$$\psi' = (\tilde{x}, \tilde{y}, \tilde{z}) = (R\cos\varphi\cos\theta, R\cos\varphi\sin\theta, R\sin\varphi) \quad (2)$$

Secondly, we should find those numbers which take the place of an imaginary number i . In the polar form of a complex number ψ , the imaginary number i means that the existence probability R rotates through θ degrees about the origin. Based on hints from this fact, I propose to use Pauli matrices.

$$\psi' = R\cos\varphi\cos\theta\sigma_{\tilde{x}} + R\cos\varphi\sin\theta\sigma_{\tilde{y}} + R\sin\varphi\sigma_{\tilde{z}} (\varphi = \omega t - k_x x - k_y y - k_z z) \quad (3)$$

$$\left[\theta: \text{the width of a photon's orbital, } \sigma_{\tilde{x}}, \sigma_{\tilde{y}}, \sigma_{\tilde{z}}: \text{the Pauli matrix} \right]$$

For the sake of simplicity, let us suppose that each Pauli matrix is handled as a number and they belong to non-Abelian group. This is hereinafter referred to as the Pauli coordinate system.

$$\begin{aligned} (\sigma_{\tilde{x}})^2 &= 1, \\ (\sigma_{\tilde{y}})^2 &= 1, \\ (\sigma_{\tilde{z}})^2 &= 1 \end{aligned}$$

$$\begin{aligned} \sigma_{\tilde{x}}\sigma_{\tilde{y}} + \sigma_{\tilde{y}}\sigma_{\tilde{x}} &= 0, \\ \sigma_{\tilde{y}}\sigma_{\tilde{z}} + \sigma_{\tilde{z}}\sigma_{\tilde{y}} &= 0, \\ \sigma_{\tilde{x}}\sigma_{\tilde{z}} + \sigma_{\tilde{z}}\sigma_{\tilde{x}} &= 0 \end{aligned}$$

If we calculate with attention to the direction, the square of ψ' named “dice function

$\Xi(R, \varphi, \theta)$ is its own absolute value.

$$\begin{aligned}
\Xi^2(R, \varphi, \theta) &= \left(R \cos \varphi \cos \theta \sigma_{\tilde{x}} + R \cos \varphi \sin \theta \sigma_{\tilde{y}} + R \sin \varphi \sigma_{\tilde{z}} \right)^2 \\
&= R^2 \cos^2 \varphi \cos^2 \theta (\sigma_{\tilde{x}})^2 + R^2 \cos^2 \varphi \sin^2 \theta (\sigma_{\tilde{y}})^2 + R^2 \sin^2 \varphi (\sigma_{\tilde{z}})^2 \\
&\quad + R^2 \cos^2 \varphi \cos \theta \sin \theta (\sigma_{\tilde{x}} \sigma_{\tilde{y}} + \sigma_{\tilde{y}} \sigma_{\tilde{x}}) \\
&\quad + R^2 \cos \varphi \sin \varphi \cos \theta (\sigma_{\tilde{x}} \sigma_{\tilde{z}} + \sigma_{\tilde{z}} \sigma_{\tilde{x}}) \\
&\quad + R^2 \cos \varphi \sin \varphi \sin \theta (\sigma_{\tilde{y}} \sigma_{\tilde{z}} + \sigma_{\tilde{z}} \sigma_{\tilde{y}}) \\
&= R^2 \cos^2 \varphi \cos^2 \theta + R^2 \cos^2 \varphi \sin^2 \theta + R^2 \sin^2 \varphi \\
&= R^2 \cos^2 \varphi (\cos^2 \theta + \sin^2 \theta) + R^2 \sin^2 \varphi \\
&= R^2 (\cos^2 \varphi + \sin^2 \varphi) \\
&= R^2 \\
&= \left| \Xi(R, \varphi, \theta) \right|^2
\end{aligned}$$

$$(\because \left| \Xi(R, \varphi, \theta) \right| = \sqrt{(R \cos \varphi \cos \theta)^2 + (R \cos \varphi \sin \theta)^2 + (R \sin \varphi)^2})$$

Lastly, we should derive the relative equation $-k_x^2 - k_y^2 - k_z^2 + \frac{\omega^2}{c^2} = \frac{m^2 c^2}{\hbar^2}$ from the dice function. For this purpose, let us introduce a new wave equation.

$$\frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial x^2} + \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial y^2} + \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial t^2} = \frac{m^2 c^2}{\hbar^2} \Xi(R, \varphi, \theta) \quad (4)$$

$$\left[m: \text{the mass of a particle, } c: \text{the velocity of light, } h: \text{the Planck constant } (\hbar = \frac{h}{2\pi}) \right]$$

If we substitute equation (3) for $\Xi(R, \varphi, \theta)$ in equation (4), we can confirm that the wave equation was defined quantum-mechanically.

$$\begin{aligned}
\frac{\partial \varphi}{\partial x} \frac{\partial \varphi}{\partial x} \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} + \frac{\partial \varphi}{\partial y} \frac{\partial \varphi}{\partial y} \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} + \frac{\partial \varphi}{\partial z} \frac{\partial \varphi}{\partial z} \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} - \frac{1}{c^2} \frac{\partial \varphi}{\partial t} \frac{\partial \varphi}{\partial t} \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} &= \frac{m^2 c^2}{\hbar^2} \Xi(R, \varphi, \theta) \\
k_x^2 \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} + k_y^2 \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} + k_z^2 \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} - \frac{\omega^2}{c^2} \frac{\partial^2 \Xi(R, \varphi, \theta)}{\partial \varphi^2} &= \frac{m^2 c^2}{\hbar^2} \Xi(R, \varphi, \theta) \\
(-k_x^2 - k_y^2 - k_z^2 + \frac{\omega^2}{c^2}) \Xi(R, \varphi, \theta) &= \frac{m^2 c^2}{\hbar^2} \Xi(R, \varphi, \theta)
\end{aligned}$$

$$\therefore -k_x^2 - k_y^2 - k_z^2 + \frac{\omega^2}{c^2} = \frac{m^2 c^2}{\hbar^2}$$