

UNIVERSITY OF CALIFORNIA
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**SEARCH FOR WW AND WZ RESONANCES IN $\ell\nu qq$ FINAL
STATES IN pp COLLISIONS AT $\sqrt{s} = 13$ TEV WITH THE ATLAS
DETECTOR**

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requirements for the degree of

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by

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Abstract

Search for WW and WZ Resonances in $\ell\nu qq$ final states in pp collisions at

$\sqrt{s} = 13$ TeV with the ATLAS detector

by

Natasha Woods

This thesis presents a search for WW and WZ resonances using data from pp collisions at $\sqrt{s} = 13$ TeV using the ATLAS detector, corresponding to an integrated luminosity of 139 fb^{-1} . Diboson resonances are predicted in a number of Standard Model (SM) extensions, such as Extended Gauge Models, and Extra dimensional models. This search looks for resonances where one W boson decays leptonically and the other W or Z boson decays hadronically. This search is sensitive to diboson resonance production via vector-boson fusion as well as quark-antiquark annihilation and gluon-gluon fusion mechanisms. No significant excess of events is observed with respect to the Standard Model backgrounds, and constraints on the masses of new W' , Z' , and bulk-RS Gravitons are extended to up to 3.3 TeV, depending on the model. As the dominant backgrounds in this search contain gluons, classifying jets as quark-initiated or gluon-initiated would make this analysis more sensitive to new physics. Towards this end, this thesis provides a calibrated quark-gluon tagger based on the multiplicity of charged particles within a jet.

Loving Dedication

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Acknowledgments

Proper acknowledgments of everyone else who helped you graduate. Write later.

Part I

Introduction

³ Chapter 1

⁴ Introduction

⁵ In general, humanity has continually strived to understand the structure and
⁶ dynamics of reality for widely varying reasons. Each academic field uses a spe-
⁷ cific set of concepts and models to describe nature. Physics is one such field,
⁸ that uses mathematical objects to systematically develop testable models about
⁹ the universe. Currently, the most fundamental types particles are fermions and
¹⁰ bosons. Fermions are the particles that make up the "ordinary" matter of the
¹¹ universe, while bosons are the quanta of the fundamental forces. The Standard
¹² Model (SM) of particle physics describes the quantum behavior of three of the
¹³ four fundamental forces: electromagnetic, strong, and weak forces.

¹⁴ The Standard Model has consistently described much of reality to an extreme
¹⁵ degree of accuracy. It has predicted cross sections for strong and electroweak
¹⁶ processes that span over ten orders of magnitude [see Fig. 3.1] and contains no
¹⁷ known logical inconsistencies. Despite the reality of the Standard Model, it still
¹⁸ fails to describe aspects of reality and suffers from aesthetic issues. The SM fails
¹⁹ to account for dark matter, dark energy, neutrino masses, the hierarchy of the
²⁰ fundamental force strengths, and other issues that may have not been noticed
²¹ yet! This incompleteness may indicate that a more fundamental theory exists. It

22 is hoped that such a theory would address the aforementioned phenomena and
23 the ad-hoc structure and parameter values of the SM. In particular the relative
24 scales of the fundamental forces impose oddly fine-tuned SM parameters, unless
25 there is additional structure at higher energies (e.g. between the electroweak and
26 Planck scales). This and other theoretical arguments motivate the search for new
27 physics at the TeV scale. The set of theories that hope to explain more of reality
28 are known as Beyond the Standard Model theories (BSM). Many of these theories,
29 if true, would revolutionize concepts of symmetry and space-time, which would
30 be intrinsically meaningful.

31 To probe the physics at this high energy frontier, physicists often collide ener-
32 getic particles that combine to produce massive particles, such as the Higgs boson
33 and top quark. The more energetic the colliding particles are the more massive
34 produced particles can be. Currently, the world's highest energy particle collider
35 is the Large Hadron Collider (LHC) at the European Organization for Nuclear
36 Research (CERN).

37 This thesis presents a search for WW and WZ resonances using data from pp
38 collisions at $\sqrt{s} = 13$ TeV using the ATLAS detector at CERN, corresponding
39 to an integrated luminosity of 139 fb^{-1} . Diboson resonances are predicted in a
40 number of BSM theories, such as Extended Gauge Models and Extra dimensional
41 models. This search looks for resonances where one W boson decays leptonically
42 and the other W or Z boson decays hadronically. This search is sensitive to
43 diboson resonance production via vector-boson fusion as well as quark-antiquark
44 annihilation and gluon-gluon fusion mechanisms (which will be collectively called
45 non-VBF modes).

46 To search for these new resonances, Monte-Carlo simulations are used to model
47 SM backgrounds and BSM signals. In these simulations, a series of optimized cuts

48 are used to create signal regions (SR) to identify the leptonic and hadronic decay
49 products of the resonance, maximize signal acceptance, and minimize background
50 contamination. In these regions, the resonance mass is calculated as the combined
51 system mass of the leptonic and hadronic system. The expected resonance mass
52 distribution from the simulated backgrounds and anticipated signal are compared
53 to data to search for the existence of these BSM signals (also known as a "bump
54 hunt"). Control regions enriched in the dominant backgrounds, $t\bar{t}$ and $W+\text{jets}$
55 (TCR and WCR, respectively) are constructed to be orthogonal to SRs and used
56 to determine the normalization of the $t\bar{t}$ and $W+\text{jets}$ backgrounds in SRs.

57 The VBF W' and Z' and ggF W' and Z' resonances studied have unique
58 SR and CR selections to maximize analysis sensitivity. RS Graviton signals are
59 probed using the same selections as the ggF Z' signal. Additionally, more mas-
60 sive resonances are more likely to have boosted W/Z bosons. As the boost of
61 the hadronically decaying boson increases the separation of its hadronic decay
62 products decreases. When the hadronically decaying boson has sufficient boost,
63 the two quarks will overlap and not be identified separately. For this reason, a
64 set of "resolved" selections are used when the hadronic decay products are recon-
65 structed separately, and "merged" selections when the decay products overlap and
66 identified as a single object in the event. A W/Z tagger identifies merged jets as
67 originating from a W/Z bosons based on jet substructure and mass cuts. However,
68 the more boosted the jet is the less likely it is to pass the jet substructure cut, due
69 to track merging. Consequently, the merged selection uses a high purity region
70 (HP), which requires that the jet pass both cuts, and low purity (LP) region where
71 the jet can fail the jet substructure cut.

72 The aforementioned SR definitions veto events with b -jets to minimize $t\bar{t}$ con-
73 tamination. However, b -jets are anticipated from W' resonances from the hadron-

74 ically decaying Z boson. To increase the signal acceptance of these resonances,
75 a $Z \rightarrow bb$ tagger is used to construct additional SR and CRs called the "tagged"
76 regions (and "un>tagged" if the event fails the $Z \rightarrow bb$ tagger).

77 For each signal model, the simulated and measured resonance mass distribu-
78 tions in the relevant SR and CRs are combined to construct a likelihood. This
79 likelihood is parameterized by the signal strength parameter, μ and systematic
80 uncertainties of the resonance mass distribution. This likelihood is used to quan-
81 tify the likelihood of a certain signal model given the anticipated backgrounds and
82 measured data.

83 No significant excess of events is observed with respect to the Standard Model
84 backgrounds, and constraints on the masses of new W' , Z' , and bulk-RS Gravi-
85 tons are extended to up to 3.3 TeV, depending on the model. As the dominant
86 backgrounds in this search contain gluons, classifying jets as quark-initiated or
87 gluon-initiated would improve the sensitivity of this analysis to new physics. To-
88 wards this end, this thesis provides a calibrated quark-gluon tagger based on the
89 multiplicity of charged particles within a jet.

90 Part II reviews the SM, its successes and shortcomings, and the aforemen-
91 tioned BSM theories that address the incompleteness of the SM. The structure
92 and performance of the ATLAS detector used is given in Part III. Part IV summa-
93 rizes the search for the diboson resonances using ATLAS data from pp collisions
94 at $\sqrt{s} = 13$ TeV. Finally, Part V examines the prospects for a quark gluon tagger
95 based on the track multiplicity of jets and details the calibration of this tagger.

Part II

96

Theoretical Motivation

97

₉₈ **Chapter 2**

₉₉ **The Standard Model of Particle**

₁₀₀ **Physics**

₁₀₁ **2.1 Introduction**

₁₀₂ By determining the dynamics of the most elementary degrees of freedom, par-
₁₀₃ ticle physics hopes to uncover the fundamental laws of the universe. The definition
₁₀₄ of elementary has evolved through time and currently refers to matter and force
₁₀₅ mediating particles: fermions and bosons, respectively. The Standard Model of
₁₀₆ Particle Physics (SM) describes the quantum behavior of three of the four funda-
₁₀₇ mental forces: weak, strong, and electromagnetic, via boson and fermion interac-
₁₀₈ tions. Gravity is not included in the SM and still under investigation.

₁₀₉ **2.2 Quantum Field Theory**

₁₁₀ In the SM, forces (and particles) are represented as fields. In this context,
₁₁₁ fields are mathematical objects that define a tensor (e.g. scalar, vector, etc) at
₁₁₂ every point on a manifold, here the manifold is space-time. These fields obey laws

₁₁₃ dictated by Quantum Field Theory (QFT). Particles arise naturally in QFT as
₁₁₄ quantized field excitations localized in spacetime.

₁₁₅ According to Noether's theorem, symmetries of a field give rise to conserved
₁₁₆ quantities (e.g. time-translation invariance leads to energy conservation). Often
₁₁₇ in the history of physics, a conserved quantity of a field is found and then the
₁₁₈ underlying symmetry of the field is inferred. Gauge symmetries are symmetries
₁₁₉ among the internal degrees of freedom of the field (components of the tensor),
₁₂₀ which give rise to quantities associated with fields. By specifying the symmetries
₁₂₁ of a system the dynamics and conserved quantities of the system may be succinctly
₁₂₂ defined.

₁₂₃ 2.3 $U(1)_{EM}$ Local Gauge Invariance

₁₂₄ The Lagrangian of Quantum Electrodynamics (QED) describes the electro-
₁₂₅ magnetic force. QED may be derived by requiring local $U(1)_{EM}$ gauge invariance
₁₂₆ of the free dirac fermion Lagrangian, ψ :

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi - m\bar{\psi}\psi \quad (2.1)$$

₁₂₇ This symmetry may be represented as a complex number with unit modulus,
₁₂₈ $e^{i\theta}$. $U(1)$ gauge invariance requires this gauge transformation of ψ will leave the
₁₂₉ Lagrangian unchanged.

$$\psi(x) \rightarrow \psi'(x) = e^{i\theta(x)}\psi(x) \quad (2.2)$$

₁₃₀ NB: This transformation is a local gauge transformation as θ depends on the
₁₃₁ spacetime coordinate.

₁₃₂ By requiring this symmetry of the free Dirac fermion Lagrangian:

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi - m\bar{\psi}\psi \quad (2.3)$$

¹³³ The mass term is unaffected, but the kinetic term is modified due to $\theta(x)$.

$$\mathcal{L} \rightarrow \mathcal{L}' = i\bar{\psi}e^{-i\theta(x)}\gamma^\mu\partial_\mu\psi e^{i\theta(x)} - m\bar{\psi}e^{-i\theta(x)}\psi e^{i\theta(x)} \quad (2.4)$$

¹³⁴

$$= i\bar{\psi}\gamma^\mu(\partial_\mu\psi + i\psi\partial_\mu\theta) - m\bar{\psi}\psi \quad (2.5)$$

¹³⁵ The $\partial_\mu\theta$ terms breaks the gauge invariance of the Lagrangian. By introducing a
¹³⁶ new field, A_μ we can recover the gauge invariance of the derivative. Now redefining
¹³⁷ the derivative as the covariant derivative:

$$D_\mu\psi \equiv (\partial_\mu - iqA_\mu)\psi \quad (2.6)$$

¹³⁸ And letting A_μ transform under $U(1)$ as:

$$A_\mu \rightarrow A_\mu + \delta A_\mu \quad (2.7)$$

¹³⁹ The transformed covariant derivative becomes:

$$D_\mu\psi \rightarrow D_\mu\psi' = (\partial_\mu - iqA_\mu)\psi' \quad (2.8)$$

¹⁴⁰

$$= (\partial_\mu - iq(A_\mu + \delta A_\mu))\psi e^{i\theta} \quad (2.9)$$

¹⁴¹

$$= e^{i\theta}D_\mu + ie^{i\theta}\psi(\partial_\mu\theta - q\delta A_\mu) \quad (2.10)$$

¹⁴² The covariant derivative can be made gauage invariant by setting the last term
¹⁴³ to zero.

$$\delta A_\mu = \frac{1}{q} \partial_\mu \theta \quad (2.11)$$

¹⁴⁴ So now A_μ transforms as:

$$A_\mu \rightarrow A_\mu + \frac{1}{q} \partial_\mu \theta \quad (2.12)$$

¹⁴⁵ Finally, replacing the derivative with the covariant derivative the Dirac La-
¹⁴⁶ grangian we have:

$$\mathcal{L} = i\bar{\psi}\gamma^\mu D_\mu \psi - m\bar{\psi}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} \quad (2.13)$$

¹⁴⁷

$$= \mathcal{L}_{QED} \quad (2.14)$$

¹⁴⁸ Here $F_{\mu\nu} \equiv \partial_\mu A_\nu - \partial_\nu A_\mu$. This last term in the Lagrangian is the kinetic
¹⁴⁹ energy of the gauge boson field.

¹⁵⁰ So we have derived the QED Lagrangian. By requiring the free Dirac La-
¹⁵¹ grangian to be invariant under local U(1) transformations we have generated a
¹⁵² new gauge boson field, A_μ , which describes the photon. As expected the photon
¹⁵³ interacts with fermions.

¹⁵⁴ Stepping back, a global U(1) gauge symmetry of the free Dirac Lagrangian
¹⁵⁵ implies we cannot measure the absolute phase of a charged particle. A local U(1)
¹⁵⁶ gauge symmetry changes the phase of fields differently across space time. For this
¹⁵⁷ type of transformation to leave the Lagrangian invariant, we had to introduce an
¹⁵⁸ additional field, A_μ , which "communicates" these phase changes across space-time.
¹⁵⁹ In less formal language this effectively means: if the field at one location changes,
¹⁶⁰ this change is conferred to other particles via A_μ .

¹⁶¹ 2.4 Yang-Mills Gauge Theories

¹⁶² Requiring $U(1)_{EM}$ gauge invariance of the free Dirac Lagrangian gave us QED.

¹⁶³ Requiring different gauge symmetries we can derive the structure of other inter-

¹⁶⁴ actions. Any gauge symmetry may be written as:

$$\psi_i \rightarrow \exp(i\theta^a T_{ij}^a) \psi_j \quad (2.15)$$

¹⁶⁵ Here θ is a dimensionless real parameter and T is the generator of the gauge

¹⁶⁶ symmetry group. With this the covariant derivative can be written as:

$$D_\mu \psi_i \equiv \partial_\mu \psi_i + ig A_\mu^a T_{ij}^a \psi_j \quad (2.16)$$

¹⁶⁷ Then the gauge field must transform as:

$$A_\mu^a \rightarrow A_\mu^a - \frac{1}{g} \partial_\mu \theta^a - f^{abc} \theta^b A_\mu^c \quad (2.17)$$

¹⁶⁸ Here f is the structure constant of the gauge group. The field strength tensor

¹⁶⁹ is given by:

$$F_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - g f^{abc} A_\mu^b A_\nu^c \quad (2.18)$$

¹⁷⁰

$$F_{\mu\nu}^a \rightarrow F_{\mu\nu}^a - f^{abc} \theta^b F_{\mu\nu}^c \quad (2.19)$$

¹⁷¹ This gives the Yang-Mills Lagrangian:

$$\mathcal{L}_{YM} = -\frac{1}{4} F_{\mu\nu}^{a\mu\nu} F_{\mu\nu}^a + i \bar{\psi}_i \gamma^\mu D_\mu \psi_i + m \bar{\psi}_i \psi_i \quad (2.20)$$

¹⁷² 2.5 Particles in the Standard Model

¹⁷³ The SM consists of fermions (half-integer spin matter constituents) and bosons
¹⁷⁴ (integer spin force mediators). Fermions are spinor representations of the Poincare
¹⁷⁵ group and can be further separated into leptons and quarks. Bosons are the result
¹⁷⁶ of requiring a particular symmetry among the spinor fields:

$$SU(3)_C \times SU(2)_L \times U(1)_Y \quad (2.21)$$

¹⁷⁷ $SU(3)_C$ is the symmetry group of the strong force and generates eight gluon
¹⁷⁸ fields, G_μ . $SU(2)_L$ is the symmetry group of the Electroweak force and generates
¹⁷⁹ three electroweak boson fields. The mixing of this $SU(2)_L$ and $U(1)_Y$ gives rise
¹⁸⁰ to the photon field, where Y is the weak-hypercharge:

$$Y = 2(Q - T_3) \quad (2.22)$$

¹⁸¹ Q is the electromagnetic charge, and T_3 is the z-component of the weak isospin.
¹⁸² Weak isospin is the charge associated with the $SU(2)_L$ symmetry. The correspond-
¹⁸³ ing covariant derivative is then:

$$D_\mu \phi \equiv (\partial_\mu + ig_1 B_\mu Y_{L/R} + [ig_2 W_\mu^\alpha T^\alpha]_L + [ig_3 G_\mu^\alpha \tau^\alpha]_C) \psi \quad (2.23)$$

¹⁸⁴ It is important to note that the gauge symmetry of the SM yields a particular
¹⁸⁵ structure of the fermion representations. So for a given fermion to interact with
¹⁸⁶ a given gauge field it must have a non-zero corresponding Noether charge for
¹⁸⁷ that gauge symmetry. If the corresponding Noether charge is zero, that fermion
¹⁸⁸ transforms as a singlet and does not participate in that gauge interaction.

¹⁸⁹ Fermions are divided into quarks and leptons based on their transformations
¹⁹⁰ under $SU(3)_C$. Quarks transform as color triplets. Leptons transform as color

singlets and consequently do not interact with gluons. Fermions may be further
 classified by their $SU(2)_L$ interactions. Only the left-chiral part of fermions (denoted by L here) transform as $SU(2)_L$ doublets, the right-chiral part forms singlets under this gauge. Lastly, all these groups of particles come in three generations, each a heavier copy of the previous, but with differing flavor quantum numbers.
 This is summarized in Table 2.1 and shown in Figures 2.1 and 2.2.

SM Fermion Gauge Group	First Generation	Second Generation	Third Generation	$(SU(3)_C, SU(2)_L, U(1)_Y)$ Representations
Left-handed quarks	$\begin{pmatrix} u_L^r & u_L^g & u_L^b \\ d_L^r & d_L^g & d_L^b \end{pmatrix}$	$\begin{pmatrix} c_L^r & c_L^g & c_L^b \\ s_L^r & s_L^g & s_L^b \end{pmatrix}$	$\begin{pmatrix} t_L^r & t_L^g & t_L^b \\ b_L^r & b_L^g & b_L^b \end{pmatrix}$	$(3, 2, \frac{1}{6})$
Right-handed quarks	(u_R^r, u_R^g, u_R^b) (d_R^r, d_R^g, d_R^b)	(c_R^r, c_R^g, c_R^b) (s_R^r, s_R^g, s_R^b)	(t_R^r, t_R^g, t_R^b) (b_R^r, b_R^g, b_R^b)	$(3, 1, \frac{2}{3})$ $(3, 1, -\frac{1}{3})$
Left-handed leptons	$\begin{pmatrix} \nu_e^L \\ e_L \end{pmatrix}$	$\begin{pmatrix} \mu_e^L \\ \mu_L \end{pmatrix}$	$\begin{pmatrix} \tau_e^L \\ \tau_L \end{pmatrix}$	$(1, 2, -\frac{1}{2})$
Right-handed leptons	e_R	μ_R	τ_R	$(1, 1, -1)$

Table 2.1: Representations of the SM fermions under $SU(3)_C \times SU(2)_L \times U(1)_Y$ gauge symmetry group. Rows are correspond to different weak isospin states and columns to different QCD color states.

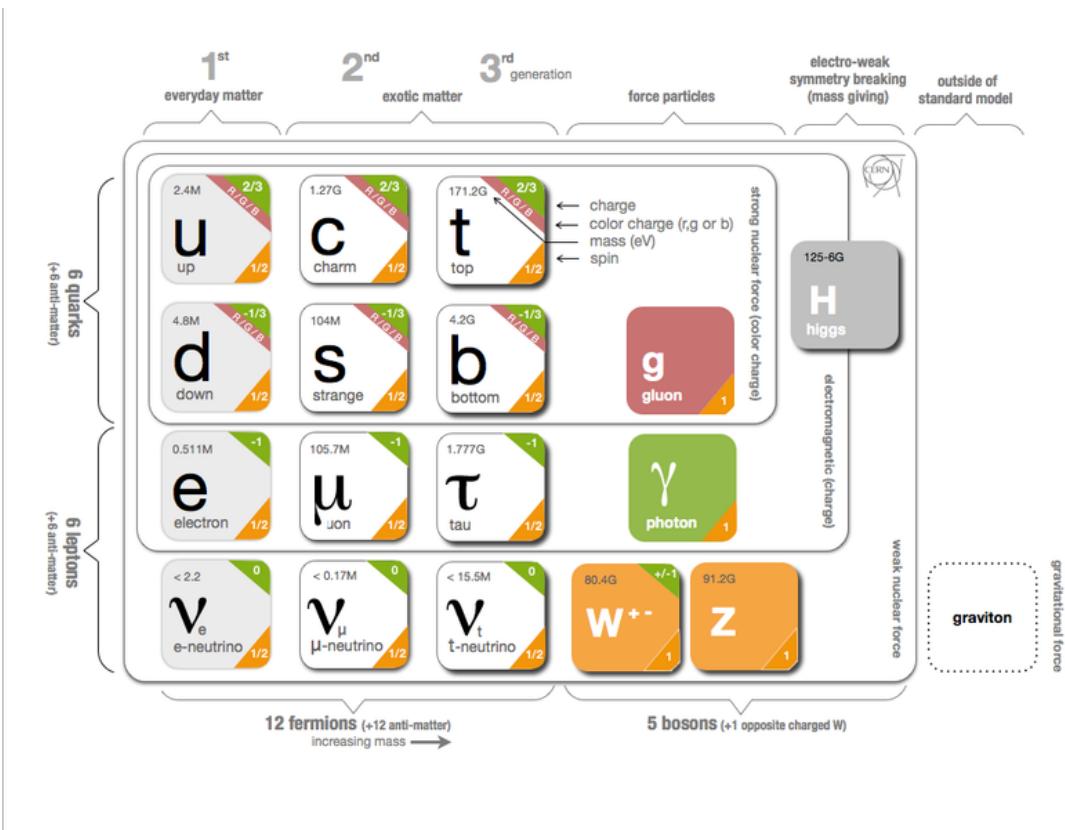


Figure 2.1: The particles of the Standard Model.



Figure 2.2: Summary of how Standard Model particles interact with other Standard Model particles.

197 Now we can understand the SM Lagrangian density as a Yang-Mills theory
 198 with the gauge group: $SU(3)_C \times SU(2)_L \times U(1)_Y$ with an additional $SU(2)$ complex
 199 scalar Higgs field doublet that will be discussed later.

$$\begin{aligned} \mathcal{L}_{SM} = & \underbrace{-\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}W_{\mu\nu}^a W^{a\mu\nu} - \frac{1}{4}G_{\mu\nu}^\alpha G^{\alpha\mu\nu}}_{\text{Kinetic Energies and Self-Interactions of Gauge Bosons}} \\ & + \underbrace{\bar{L}_i \gamma^\mu (i\partial_\mu - \frac{1}{2}g_1 Y_{iL} B_\mu - \frac{1}{2}g_2 \sigma^a W_\mu^a) L_i}_{\text{Kinetic Energies and EW Interactions of Left-handed Fermions}} \\ & + \underbrace{\bar{R}_i \gamma^\mu (i\partial_\mu - \frac{1}{2}g_1 Y_{iR} B_\mu) R_i}_{\text{Kinetic Energies and EW Interactions of Right-Handed Fermions}} \\ & + \underbrace{\frac{ig_3}{2} \bar{Q}_j \gamma^\mu \lambda^\alpha G_\mu^\alpha Q_j}_{\text{Strong Interactions between Quarks and Gluons}} \\ & + \underbrace{\frac{1}{2} |(i\partial_\mu - \frac{1}{2}g_1 B_\mu - \frac{1}{2}g_2 \sigma^a W_\mu^a)\Phi|^2 - V(\Phi)}_{\text{Electroweak Boson Masses and Higgs Couplings}} \\ & - \underbrace{(y_{kl}^d \bar{L}_k \Phi R_l + y_{kl}^u \bar{R}_k \tilde{\Phi} L_l + h.c.)}_{\text{Fermion Mass terms and Higgs Couplings}} \end{aligned}$$

200 Here several abstract spaces are being spanned:

- 201 – a spans the three $SU(2)_L$ gauge fields with generators expanded in Pauli
 202 matrices, $T^\alpha = \frac{1}{2}\sigma^\alpha$
- 203 – α spans the eight $SU(3)_C$ gauge fields, with generators expanded in Gell-
 204 Mann matrices, $\tau^\alpha = \frac{1}{2}\lambda^\alpha$
- 205 – L/R represent left and right projections of Dirac fermion fields. The Strong
 206 interaction is not chiral, so $Q = L+R$

- 207 – μ and ν are four-vector indices
- 208 – i, j, k are summed over the three generations of SM particles.

209 2.6 Higgs Mechanism

210 The SM Lagrangian without the addition of a Higgs field does not allow for
211 gauge boson and fermion mass terms: $\frac{1}{2}m_A^2 A_\mu A_\mu$ and $m(\bar{\psi}\psi)$, as these terms are
212 not gauge invariant. By introducing the Higgs field, mass terms for these particles
213 may be included in a gauge invariant way. This field is a complex doublet with a
214 potential $V(\Phi)$:

$$\Psi = \begin{pmatrix} \Phi^\dagger \\ \Phi^0 \end{pmatrix} \quad (2.24)$$

215

$$V(\Phi) = \mu^2 \Phi^\dagger \Phi + \lambda |\Phi^\dagger \Phi|^2 \quad (2.25)$$

216 The minima of this field occurs for $|\Phi| = \sqrt{\frac{\mu^2}{2\lambda}} \equiv \frac{v}{2}$. This yields degenerate
217 minima, this symmetry is broken by choosing a specific minima (a.k.a. sponta-
218 neous symmetry breaking). By convention $\Phi_{min} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$ is chosen. This means
219 the ground state of the Higgs field (Higgs vacuum) is non-zero, $\sqrt{\frac{-\mu^2}{\lambda}}$. The Higgs
220 Field may now be expanded around this new ground state:

$$\Phi(x) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h(x) \end{pmatrix} \quad (2.26)$$

221 This non-zero Higgs vacuum now generates mass terms for the gauge bosons
222 from the following term in the Lagrangian:

$$|(-\frac{1}{2}g_1B_\mu - \frac{1}{2}g_2\sigma^aW_\mu^a)\Phi|^2 = \frac{1}{2}m_W^2W_\mu^+W^{-\mu} + \frac{1}{2}m_Z^2Z_\mu Z^\mu \quad (2.27)$$

223 where:

$$W_\mu^\pm \equiv \frac{1}{\sqrt{2}}(W_\mu^1 \mp iW_\mu^2) \quad (2.28)$$

$$\begin{aligned} \text{224} \quad Z_\mu &\equiv \frac{1}{\sqrt{g_1^2 + g_2^2}}(g_2W_\mu^2 - g_1B_\mu) \end{aligned} \quad (2.29)$$

$$\begin{aligned} \text{225} \quad m_W &= \frac{vg_2}{\sqrt{2}} \end{aligned} \quad (2.30)$$

$$\begin{aligned} \text{226} \quad m_Z &= \frac{v}{\sqrt{2}}\sqrt{g_1^2 + g_2^2} \end{aligned} \quad (2.31)$$

227 The Higgs field also generates a mass term for the Higgs boson and self-
228 interactions for the Higgs boson.

229 2.7 Electroweak Theory

230 $SU(2)_L$ generates W^\pm, W^0 gauge bosons, which would be massless if $SU(2)_L$
231 was a perfect symmetry. These bosons are massive as this symmetry is broken.

232 The mass eigenstates, Z and γ given by:

$$\begin{pmatrix} Z_\mu \\ A_\mu \end{pmatrix} = \begin{pmatrix} \cos\theta_W & -\sin\theta_W \\ \sin\theta_W & \cos\theta_W \end{pmatrix} \begin{pmatrix} W_\mu^3 \\ B_\mu \end{pmatrix} \quad (2.32)$$

233 Here θ_W is the Weinberg angle given by:

$$\cos\theta_W = \frac{g_2}{\sqrt{g_1^2 + g_2^2}} = \frac{m_W}{m_Z} \quad (2.33)$$

²³⁴ 2.8 Quantum ChromoDynamics

²³⁵ As mentioned earlier the Strong Force, which binds the proton together, is
²³⁶ mediated by gluons. Quantum Chromodynamics is the QFT which describes the
²³⁷ interactions of quarks and gluons via $SU(3)_C$ symmetry. QCD contains features
²³⁸ not present in Electroweak Interactions due to $SU(3)_C$ generators not commuting
²³⁹ (a.k.a. $SU(3)_C$ is a non-abelian group) and the number of quark flavors (n_f).
²⁴⁰ For example, in QCD there is color confinement and asymptotic freedom due to
²⁴¹ the structure constants being non-zero. Requiring $SU(3)_C$ local gauge invariance
²⁴² implies:

$$\psi(x) \rightarrow \psi(x)' = \exp[i g_S \alpha(x) \cdot \hat{T}] \psi(x) \quad (2.34)$$

²⁴³ where $\alpha(x)$ is the local phase function, g_S is the strong coupling constant, and
²⁴⁴ \hat{T} are the eight generators of $SU(3)$ (note $\hat{T}^a = \frac{1}{2}\lambda^a a$, where λ^a are the Gell-Mann
²⁴⁵ matrices). As the Gell-Mann matrices are 3x3, this means ψ has three degrees of
²⁴⁶ freedom under these $SU(3)$ rotations. So we represent ψ under $SU(3)$ rotations
²⁴⁷ as:

$$\psi = \begin{pmatrix} \psi_{red} \\ \psi_{green} \\ \psi_{blue} \end{pmatrix} \quad (2.35)$$

²⁴⁸ Consequently, particle fields transforming under $SU(3)$ rotations have three
²⁴⁹ components which physicists describe as color components (red, green, and blue).
²⁵⁰ A particle's corresponding antiparticle has the corresponding anticolor. This color
²⁵¹ is the "charge" of QCD and is conserved under $SU(3)$ rotations. Combining colors,
²⁵² color neutral states (e.g. red and antired, or red, green and blue) may be created.
²⁵³ For the Free Dirac Lagrangian to remain invariant under $SU(3)$ transformations,

254 we must again postulate a boson field that modifies the derivative. The gluon
255 field tensor is given by ($\alpha = 1, \dots, 8$):

$$G_{\mu\nu}^k = \partial^\mu G_\alpha^\nu - \partial^\nu G_\alpha^\mu - g_S f^{\alpha\beta\gamma} G_\beta^\mu G_\gamma^\nu \quad (2.36)$$

256 Here $f^{\alpha\beta\gamma}$ are the structure constants of $SU(3)$. Combining all this gives the
257 QCD Lagrangian:

$$\mathcal{L}_{QCD} = \bar{\psi}_q i\gamma^\mu (D_\mu)_{ij} \psi^{qj} - m \bar{\psi}_q \psi_q - \frac{1}{4} G_{\mu\nu}^\alpha G^{\alpha\mu\nu} \quad (2.37)$$

258 Here i are the color indices, and q are the quark flavors. It is important to
259 note that quarks transform under the fundamental representation of $SU(3)$, while
260 gluons transform under the adjoint representation. This means quarks carry a
261 single color charge (red, green, blue, antired, antigreen, antiblue) and gluons carry
262 a color and anticolor charge.

263 Figure 2.3 shows the three dominant QCD interactions. Since gluons carry
264 color charge, they interact with one another. This does not occur in QED, as
265 photons do not have electric charge and therefore do not interact with each other.
266 In QED, a bare electron's effective charge is largest closest to the electron and
267 decreases as a function of distance. This is because the QED vacuum fills with
268 particle antiparticle pairs spontaneously, which screen the charge of the bare elec-
269 tron. The larger the distance from the electron, the smaller the effective charge
270 and therefore the weaker the force.

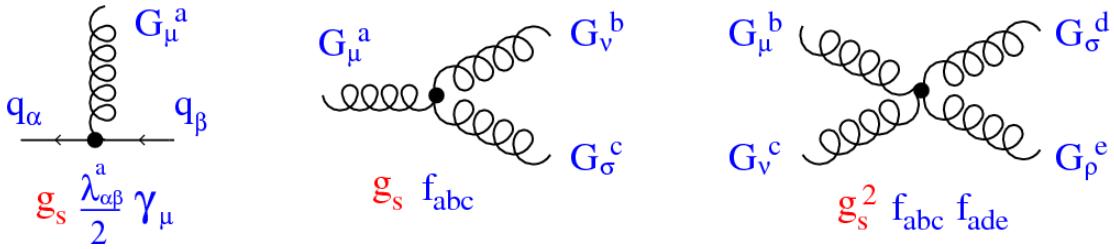


Figure 2.3: This figure shows the three dominant QCD interactions. From Ref. [18]

271 As the distance from a quark increases it's effective color charge increases due
 272 to the vacuum polarization in QCD. Color charge grows as the distance from
 273 the source increases (a.k.a. color is anti-screened in QCD). In this way, strong
 274 interactions become stronger at large distances (low momenta interactions). At
 275 small distances (large momenta interactions) strong interactions are significantly
 276 weaker and considered nearly free. This effect of referred to as asymptotic freedom.
 277 At large distances, a quark's effective charge is large and the strong force is more
 278 significant. This force becomes so strong that quarks form colorless bound states
 279 instead of remaining free particles. This effect is known as color confinement.
 280 This running of all SM fields is shown in Figure 2.4.



Figure 2.4: Strength of the U(1), SU(2), and SU(3) gauge couplings as a function of the energy scale of the interaction (Q). From Ref. [12]

Commonly the change in a particle's effective charge under a given force is quantified with $\beta(r) \equiv -\frac{de(r)}{d\ln r}$, where $e(r)$ is the effective charge of a given particle under a force. In QED this function is positive but in QCD this function is negative leading to confinement and asymptotic freedom. Moreover, one can calculate how the coupling (α) of a force varies with energies. (More deeply this amounts to incorporating renormalization and vacuum polarization in the boson propagators).

For QCD this is:

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu^2)}{1 + \frac{\alpha_s(\mu^2)}{12\pi}(33 - 2n_f)\ln(Q^2/\mu^2)} \quad (2.38)$$

where Q is the momentum of the force is probed at, μ^2 is the renormaliza-

289 tion scale, n_f is the number of quark flavors. There are six quark flavors in SM
290 QCD, making $33 - 2n_f > 0$. This factor being positive and the $\ln(Q^2/\mu^2)$ being
291 in the denominator means that as Q^2 increases α_s decreases. So for large Q^2 , α_s
292 is small and SM QCD is asymptotically free, while for small Q^2 , α_s is large and
293 SM QCD is confined, as mentioned earlier.

294 As stated previously, quarks and gluons have not been observed in isolation.
295 Instead they form bound colorless states. Hadronization is the process by which
296 quarks and gluons form hadrons. The process of hadronization is still an active
297 area of research. One qualitative description is show in Figure 2.5. In this figure,
298 as two quarks separate the color field between them is restricted to a tube with
299 energy density of $1\text{GeV}/\text{fm}$. As they separate further, the energy in the color field
300 increases, until there is enough energy to produce $q\bar{q}$ pairs, which breaks the color
301 field. This process repeats until quarks and antiquarks have low enough energy
302 to form colorless hadrons. The resulting spray of hadrons is called a jet.

303 Since quarks and gluons carry different color charges, their respective jets have
304 different properties. As quarks carry only a single color charge (vs. gluons which
305 have color and anticolor charge), so their jets have less constituent particles. More
306 precisely, the Altarelli-Parisi splitting functions [3] contain a factor C_A for gluon
307 radiation off a gluon and C_F for gluon radiation off a quark ($C_A/C_F = 9/4$). These
308 color factors are the prefactor in the Feynman diagrams for these processes [1],
309 which leads to gluon jets having more constituents and therefore more tracks than
310 quark jets. Gluon jets also tend to have a larger radius with lower momentum
311 constituents than quarks. There are many novel techniques to distinguish quarks
312 from gluons. For this study the number of charged particles will be focused on.

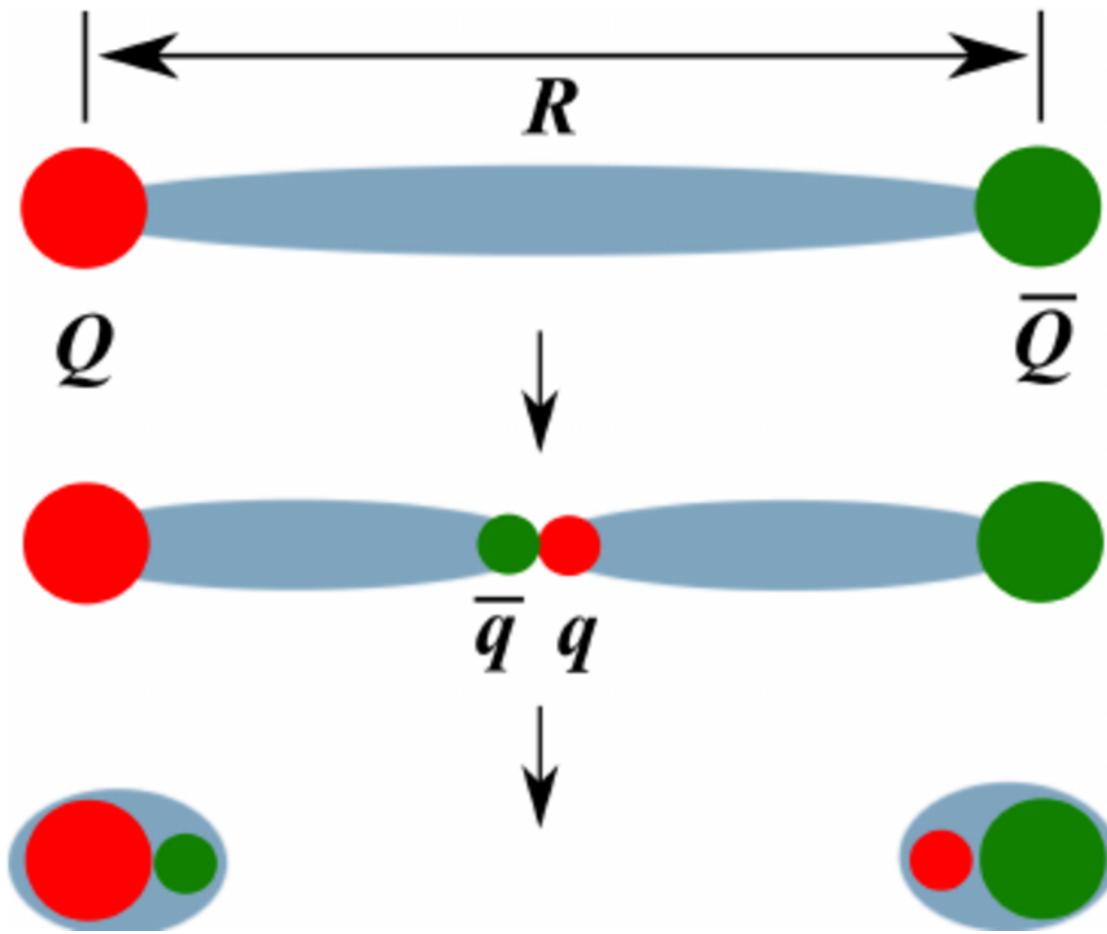


Figure 2.5: A cartoon of string breaking: the QCD string spanned between quark Q and antiquark \bar{Q} breaks due to $q\bar{q}$ creation [4]

³¹³ **Chapter 3**

³¹⁴ **Standard Model Successes and
Limitations**

³¹⁶ The Standard Model has accurately described most of the underlying principles
³¹⁷ of nature. It has predicted cross sections for strong and electroweak processes that
³¹⁸ span over ten orders of magnitude correctly [see Fig. 3.1] and contains no known
³¹⁹ logical inconsistencies. Despite the strength and reality of the Standard Model, it
³²⁰ still fails to describe some important aspects of reality and suffers from aesthetic
³²¹ issues. To date, dark matter and energy comprise 95% of the universe, but are
³²² not accounted for in the SM. Additionally, neutrinos are known to have mass but
³²³ are massless in the SM. There are mechanisms for introducing massive neutrinos
³²⁴ in the SM, but these mechanisms create hierarchy problems.

³²⁵ Possibly the most significant aesthetic issue is the hierarchy between the elec-
³²⁶ troweak and Planck scales. The electroweak scale is the scale of electroweak
³²⁷ symmetry breaking. The Planck scale is the scale where the gravitational force
³²⁸ is comparable in strength to the other forces. The Planck scale is where the SM
³²⁹ breaks down, as there is not an experimentally verified theory of quantum gravity,
³³⁰ and at this scale gravity cannot be ignored (like it can at the electro-weak scale).

331 These scales differ by ~ 30 orders of magnitude. Understanding the difference
332 in these energy scales may help explain the weakness of gravity at electroweak
333 scales, and possibly a QFT for gravity. (NB: This hierarchy can also be framed in
334 terms of the corrections to the Higgs mass, which depend on the UV cutoff scale -
335 where the SM is suppose to break, which is taken at the Planck scale. This leads
336 the quantum corrections to the Higgs mass that would force the Higgs mass to
337 $\sim 10^{18}$ TeV.)

338 These stark contrasts in scales may indicate that a more fundamental theory
339 exists. It is hoped that such a theory would explain and motivate some of the
340 ad-hoc features of the SM. In particular, the values of the 19 SM parameters (6
341 quark masses, 3 charged lepton masses, 3 gauge couplings, Higgs parameters (μ^2 ,
342 λ)), the structure of the fermion representations, etc.

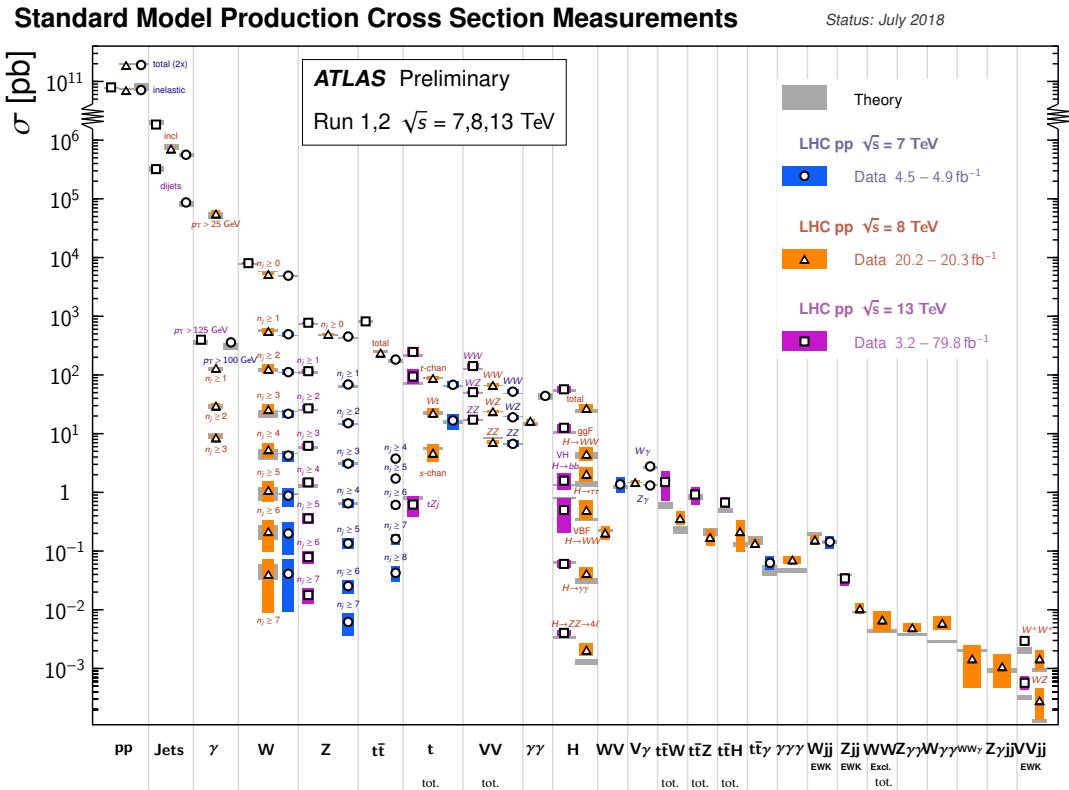


Figure 3.1: A comparison of cross section measurements at $\sqrt{s} = 7,8,13$ TeV from ATLAS compared to theoretical measurements. From Ref. [6]

³⁴³ **Chapter 4**

³⁴⁴ **New Physics Models with**

³⁴⁵ **Diboson Resonances**

³⁴⁶ **4.1 Randall Sundrum Bulk Model**

³⁴⁷ The electroweak-planck hierarchy may be explained by the existence of extra
³⁴⁸ dimensions, like the 5D Randall Sundrum Bulk Model ([19], [2]). In this model,
³⁴⁹ there is one extra warped spatial dimension, y , with a metric:

$$ds^2 = e^{-2k|y|} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 \quad (4.1)$$

³⁵⁰ where $e^{-k|y|}$ is the warp factor of the extra dimension, which is compactified on
³⁵¹ a S^1/Z_2 orbifold (a.k.a. a circle where $y \rightarrow -y$). This can be visualized as every
³⁵² point in space time having a line extending from it a distance L , representing
³⁵³ this fifth dimension. At the end of this line is the Planck brane. This fourth
³⁵⁴ spatial dimension separates two 4-D branes: Planck brane and TeV brane. We
³⁵⁵ live on the TeV brane, as shown in Figure 4.1. The Higgs field (and to a lesser
³⁵⁶ degree the top quark and graviton fields) is localized near the TeV Brane, while

357 the light fermion fields are localized more near the Planck brane. Fundamental
 358 parameters are set on the Planck brane. The warp factor may be scaled away from
 359 all dimensionless SM terms by field redefinitions. However, the only dimensionful
 360 parameter, $m_H^2 = v^2$ is rescaled by $\tilde{v} \sim e^{-kL} M_{Pl} \sim 1\text{TeV}$ for $kL \sim 35$, explaining
 361 why gravity is so weak on the TeV brane. Also, by localizing the light fermion
 362 fields near the Planck brane and top and graviton fields near the TeV brane, the
 363 light quarks will have smaller masses.

364 The two free parameters of this theory are M_{Pl} and k . Based on this RS Bulk
 365 model, all SM particles should have Kaluza-Klein (KK) excitations. In particular,
 366 the graviton would have KK excitations that prefer to decay to WW or ZZ, which
 367 is why this analysis searches for RS Gravitons.

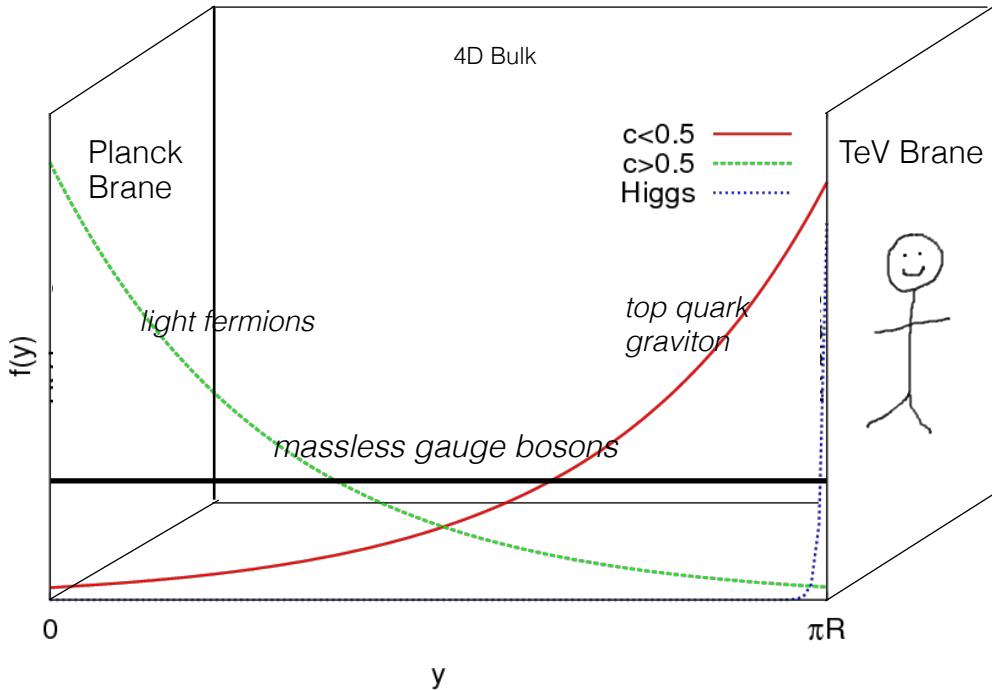


Figure 4.1: Cartoon of RS Bulk Model

368 4.2 Simple Standard Model Extensions

369 The RS Bulk model is motivated by resolving SM hierarchies, but it does not
370 address all of the other SM issues. There are many other interesting and well
371 motivated new physics frameworks that address these issues, but there is a lack
372 of completely predictive models, due to model flexibility (free parameters). It is
373 difficult for experimentalists to know which theories to search for in data. There-
374 fore, developing a model-independent resonance search that can be reinterpted
375 in the context of a given BSM theory is ideal.

376 This search is sensitive to the resonance mass and its interactions, but not
377 all of a given BSM model's parameters. Therefore, the BSM Lagrangian may be
378 reduced to only retain this information (mass parameters and couplings) following
379 the procedure in [17]. In this simplified approach, the new resonance searched for
380 is represented as an additional heavy vector triplet (HVT), which is a real vector
381 field in the adjoint representation of $SU(2)_L$ with vanishing hypercharge. This
382 results in one neutral and two charged bosons, defined as:

$$V^\pm = \frac{V_\mu^1 \mp i V_\mu^2}{\sqrt{2}} \quad (4.2)$$

$$V_\mu^0 = V_\mu^3 \quad (4.3)$$

384 The SM Lagrangian is then augmented with the additional terms:

$$\mathcal{L} \supset -\frac{1}{4} D_{[\mu} V_{\nu]}^a D^{[\mu} V^{\nu]}_a + \frac{m_V^2}{2} V_\mu^a V^{a\mu} + ig_V c_H V_\mu^a H^\dagger \tau^a \overset{\leftrightarrow}{D}{}^\mu H + \frac{g^2}{g_V} c_F V_\mu^a J_F^{\mu a} \quad (4.4)$$

385 In order the terms represent: the kinetic, V mass, Higgs- V interaction, and
386 V -left-handed fermion interaction terms. Phenomenologically the three physical
387 particles this predicts are degenerate, where V couples most strongly to VV , via

388 the g_V coupling factor. The dominant production modes are DY and VBF.

389 Two versions of HVT are considered, Model A and B. Model A is a weakly
390 coupled model where $g_V \sim 1$, like the extended gauge symmetry discussed in Ref
391 . [16]. Model B is a strongly coupled model, where $1 < g_V < 4\pi$. The width
392 of the resonance grows with g_V so for this narrow resonance search only g_V is
393 chosen to be less than 6 (so $\Gamma/M < 10\%$). More precisely, the coupling of these
394 resonances to fermions scales as $g_f = g^2 c_F/g_V$, where g is the SM $SU(2)_L$ gauge
395 coupling and c_F is the free parameter (expected to be of order 1 for Model A and
396 B). This then means that for Model B the coupling to fermions is more suppressed
397 than for Model A, leading to a smaller DY production rate and BR to fermionic
398 final states. The coupling of V to SM bosons scales as $g_H = g_V c_H$, where c_H is
399 a free parameter on the order of 1 for Model A and B. So for small values of g_V
400 (i.e. Model A - weakly coupled theories) the BR to gauge bosons is smaller than
401 for Model B. So weakly coupled vectors have large production cross sections and
402 decay prominently to leptons or jets, while strongly coupled vectors are produced
403 less and decay predominantly to gauge bosons.

404 Vectors in Model A and B are generally produced via quark-anti-quark annihi-
405 lation. The more rare production via vector-boson-fusion is considered by setting
406 $g_H = 1$ and $g_F = 0$. In Model B diboson final states are enhanced as stated
407 previously due to g_H and moreover the BR to WZ , WH , WW , and ZH are the
408 same.

409 In summary, V couples most strongly to left-handed fermions and VV depen-
410 dent on g_V .

Part III

411

Experimental Setup

412

⁴¹³ **Chapter 5**

⁴¹⁴ **LHC**

⁴¹⁵ The Large Hadron Collider (LHC) is the highest-energy particle collider in the
⁴¹⁶ world. It was designed to expand the frontier of high energy particle collisions in
⁴¹⁷ energy and luminosity. This enables LHC experiments to test the Standard Model
⁴¹⁸ and search for new physics at higher energies than tested with previous colliders.
⁴¹⁹ Collisions at higher energies not only produce more massive particles but also
⁴²⁰ more weakly interacting particles. Fig 5.1 shows production cross sections for
⁴²¹ various processes at hadron colliders. The rate for electroweak physics processes
⁴²² including W and Z scale with the center-of-momentum energy, \sqrt{s} .

⁴²³ The LHC consists of a 26.7 km (17 miles) ring, approximately 100 m un-
⁴²⁴ derground, outside Geneva, Switzerland. Counter-circulating proton (and occa-
⁴²⁵ sionally heavy ions) beams collide inside four experiments along the beam line:
⁴²⁶ ATLAS, CMS, LHCb, ALICE. ATLAS and CMS are general purpose detectors de-
⁴²⁷ signed to explore the high energy frontier. LHCb is designed to study the physics
⁴²⁸ of b -quarks. ALICE specializes in studying heavy ion collisions.

proton - (anti)proton cross sections

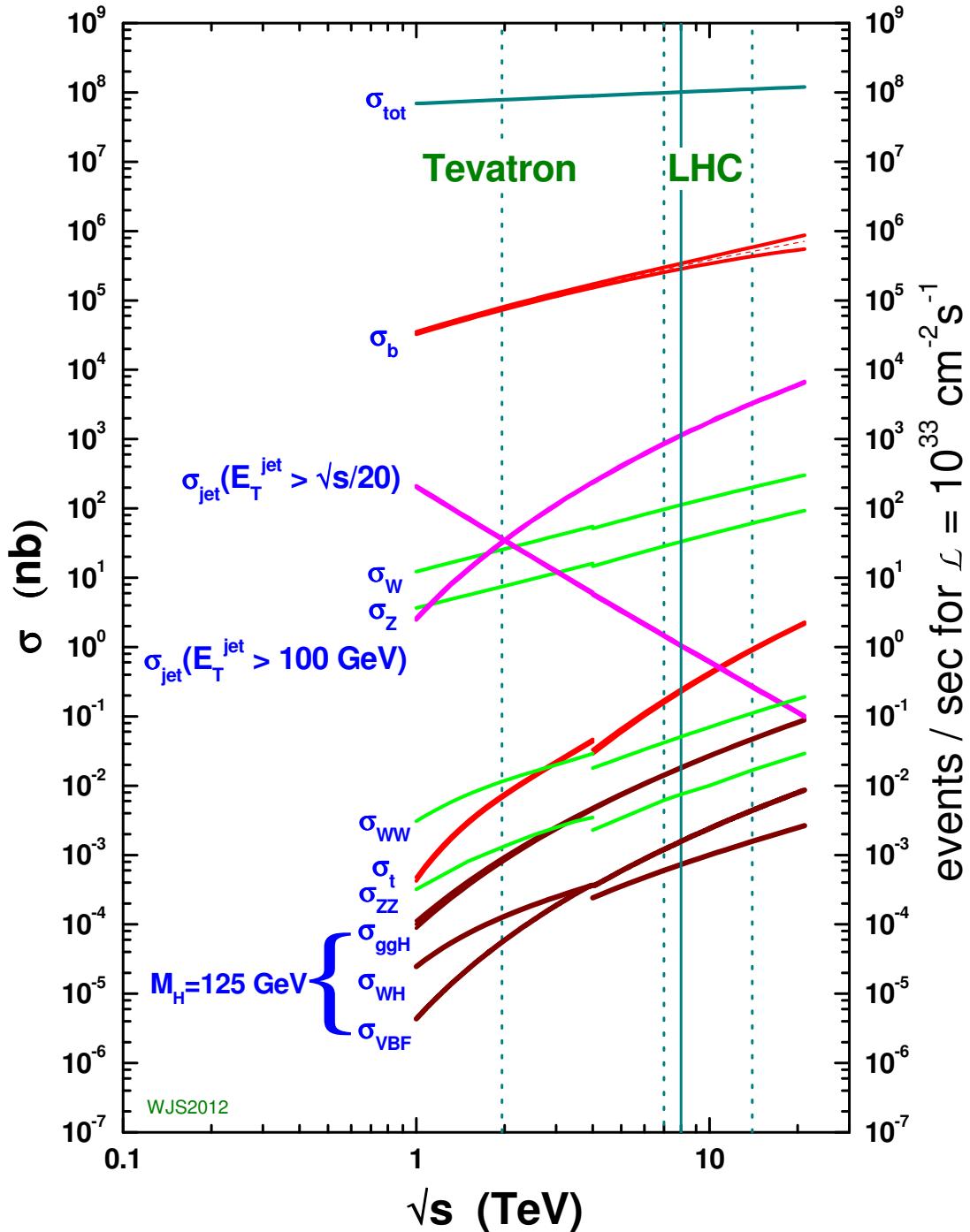


Figure 5.1: Scaling of cross sections with \sqrt{s} . Natasha: write more here

429 The first proton beams circulated in September, 2008. Nine days later an elec-
430 trical fault lead to mechanical damage and liquid helium leaks in the collider. This
431 incident delayed further operation until November 2009, when the LHC became
432 the world's highest energy particle collider, at 1.18TeV per beam. This first oper-
433 ational run continued until 2013, reaching 7 and 8 TeV collision energies. During
434 this run a particle with properties consistent with the Standard Model Higgs bo-
435 son was discovered. The next run began after a two year shutdown after upgrades
436 to the LHC and experiments. This run lasted from 2013 to 2018 reaching 13 TeV
437 collision energies. This analysis uses data from the second operational run.

438 5.1 LHC Layout and Design

439 The layout of the LHC is shown in Figure 5.2. The red and blue lines in the
440 figure represent the counter-circulating proton beams. The LHC is divided into
441 eight octants. Octant 4 contains the RF cavities that accelerate the protons and
442 octant 6 contains the beam dump system. Octants 3 and 7 house the collimation
443 systems for beam cleaning. The beams collide inside the four aforementioned
444 experiments. Each octant contains a curved and straight section. The LHC
445 magnets are built with NbTi superconductors cooled with super-fluid Helium to
446 2K, creating a 8.3T magnetic field to bend the proton beams.

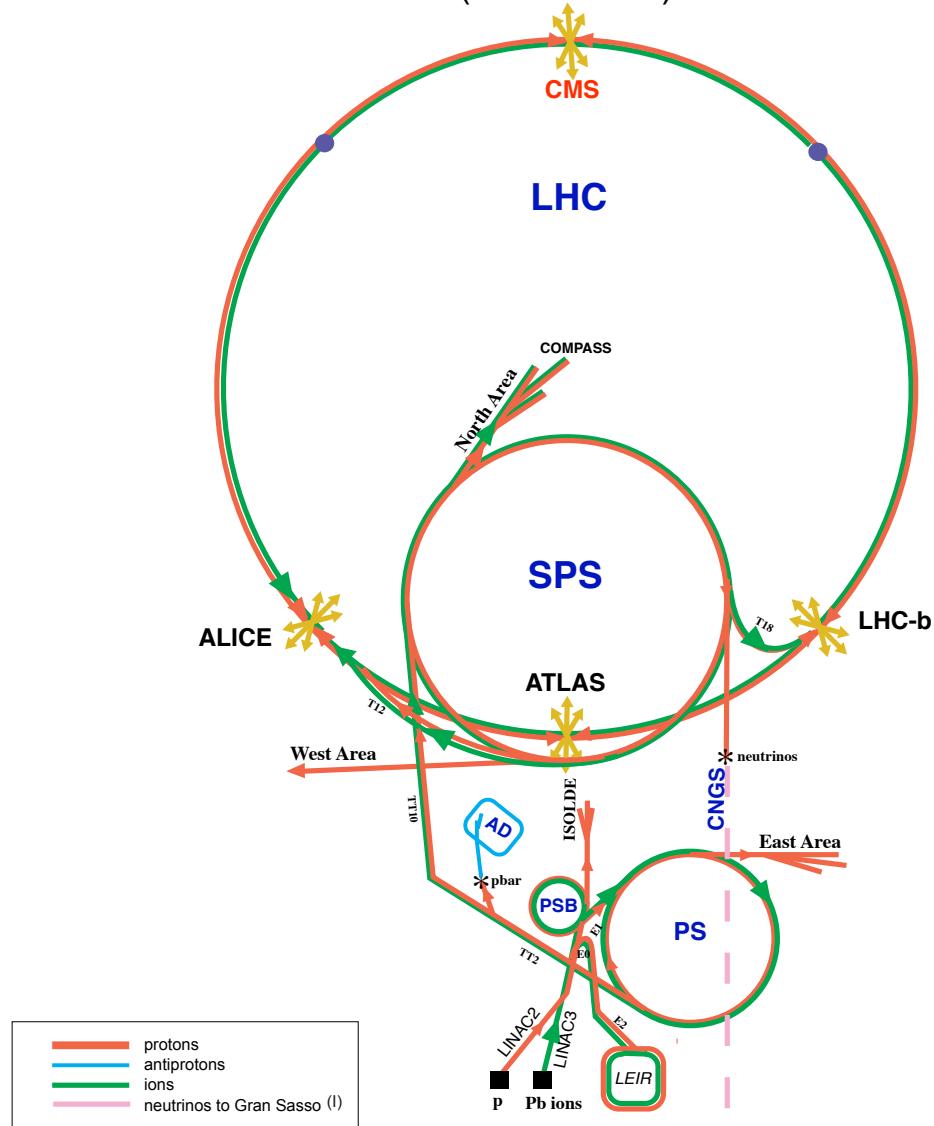


Figure 5.2: LHC Layout. Natasha write more

447 Four sequential particle accelerators are used to accelerate protons from rest
 448 as shown in Figure 5.3. First, Hydrogen gas is ionized to produce protons which
 449 are then accelerated to 50 MeV using Linac 2, a linear accelerator. The result-
 450 ing proton beam is then passed to three circular particle accelerators: Proton
 451 Synchrotron Booster, Proton Synchrotron, and Super Proton Synchrotron (SPS),

452 accelerating protons to 1.4, 25, and 450 GeV, respectively. Once the protons exit
453 the SPS, they are injected into the LHC at octant 2 and 8. Each proton bunch
454 contains $\sim 10^{11}$ protons. The spacing between bunches is 25 ns, which means
455 each beam contains 3564 bunches. However, some bunches are left empty due
456 to injection and safety requirements, yielding 2808 bunches per beam. Once the
457 proton beams are injected they are accelerated to 13 TeV.

CERN Accelerators (not to scale)



LHC: Large Hadron Collider

SPS: Super Proton Synchrotron

AD: Antiproton Decelerator

ISOLDE: Isotope Separator OnLine DEvice

PSB: Proton Synchrotron Booster

PS: Proton Synchrotron

LINAC: LINear ACcelerator

LEIR: Low Energy Ion Ring

CNGS: Cern Neutrinos to Gran Sasso

Rudolf LEY, PS Division, CERN, 02.09.96
Revised and adapted by Antonella Del Rosso, ETT Div.,
in collaboration with B. Desforges, SL Div., and
D. Manglunki, PS Div. CERN, 23.05.01

Figure 5.3: LHC Accelerator. Natasha write more

458 As many new physics models predict cross-sections below the weak scale it was
459 important to design the LHC to be capable of collecting enough data, by running
460 in high luminosity conditions. The machine luminosity depends only on beam
461 parameters:

$$L = \frac{N_p^2 f}{4\epsilon\beta^*} F \quad (5.1)$$

462 where N_p is the number of protons per bunch, f is the bunch crossing frequency,
463 ϵ is the transverse beam emittance, β^* is the amplitude function at the collision
464 point, and F is the geometric luminosity reduction factor due to the beams crossing
465 at an angle (rather than head-on).

466 **Chapter 6**

467 **The ATLAS Detector**

468 The ATLAS detector measures the position, momentum and energy of parti-
469 cles produced in the proton collisions by using magnetic fields, silicon detectors,
470 sampling calorimeters, and gaseous wire detectors. It is located approximately
471 100 m underground at Point-1 around the LHC beam line and weighs 7000 metric
472 tons. The detector is 46 m long, 25 m high, 25 m wide as shown in Figure 6.2.
473 The detector can be divided into three subsystems: the Inner Detector (ID), the
474 Calorimeters, and the Muon Spectrometer (MS). Figure 6.3 shows an overview of
475 how different particles interact in the detector.

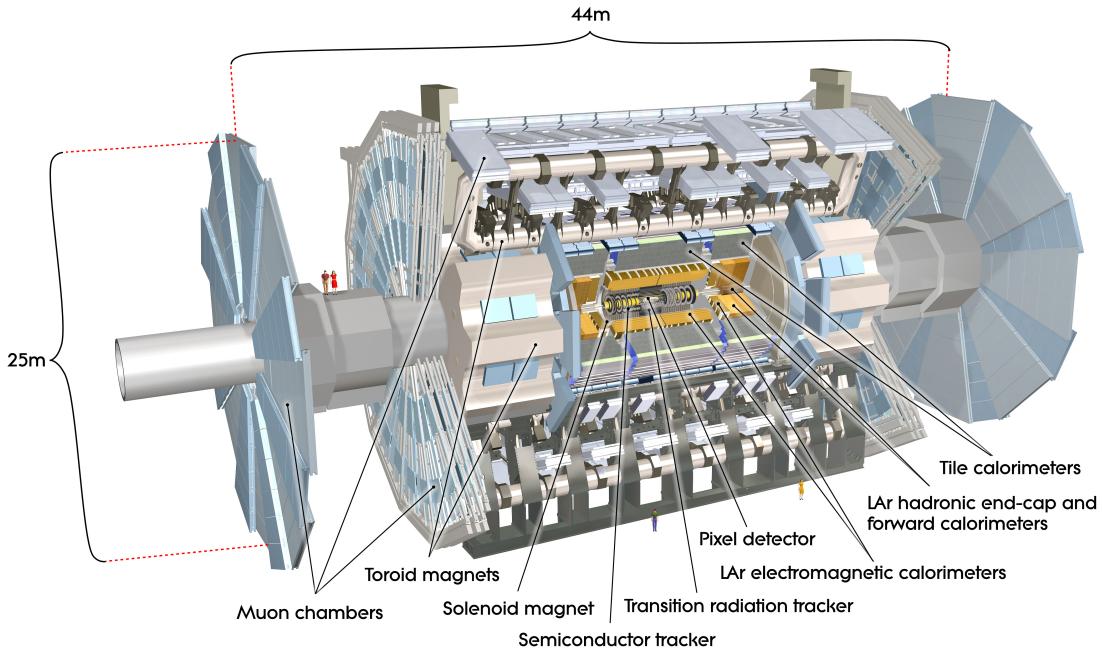


Figure 6.1: Big picture layout of ATLAS detector. Natasha: write more

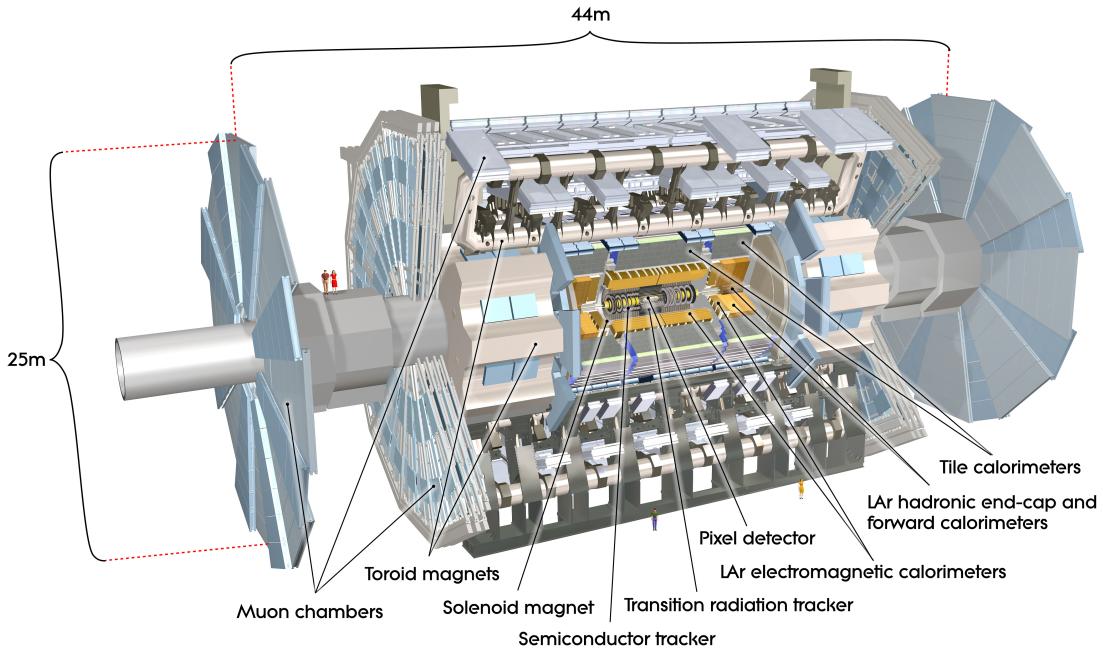


Figure 6.2: Big picture layout of ATLAS detector. Natasha: write more

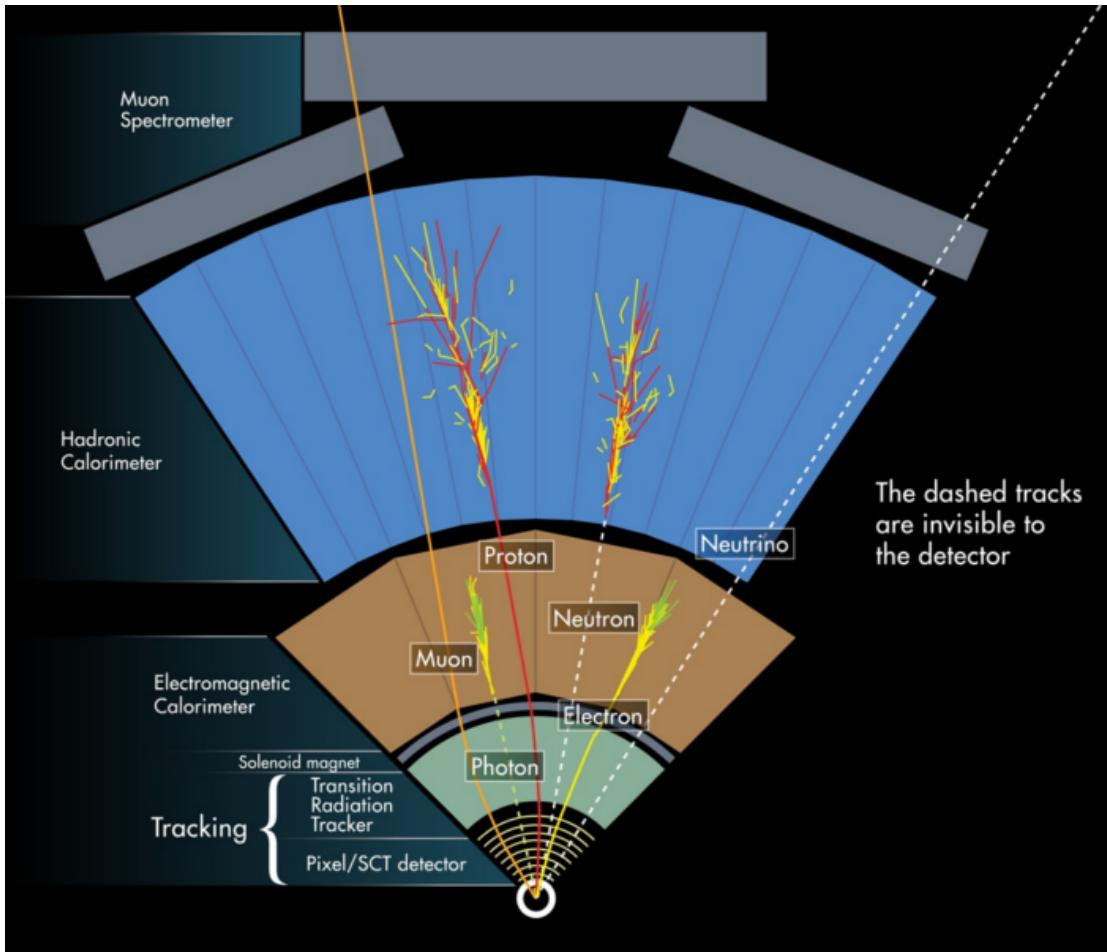


Figure 6.3: A simplified schematic of how different particles interact and are detected within ATLAS.

476 6.1 Coordinate System

477 The trajectory of particles within ATLAS is measured relative to the nominal
 478 interaction point. The z -axis points along the beam line, such that when the
 479 LHC is viewed from above, the counter-clockwise circulating beam points along
 480 the positive- z direction. The $x - y$ plane is transverse to the beam line, with the
 481 positive x -axis pointing towards the center of the LHC ring. The positive y -axis
 482 points vertically upward. The azimuthal angle, ϕ , is the angular distance about

483 the z -axis, with $\phi = 0$ along the x -axis. The polar angle from the z -axis is denoted
484 as θ . However, this quantity is not Lorentz invariant, like rapidity, $y = \frac{1}{2} \ln \frac{E+p_z}{E-p_z}$,
485 where E is the energy of the particle considered, and p_z , is it's momentum along
486 the z -axis. Pseudo-rapidity is preferred as $\Delta\eta$ is invariant under boosts along z
487 and particle production is approximately invariant under η . For massless particles,
488 rapidity and a related quantity, pseudorapidity, are the identical. The pseudora-
489 pidity is defined as: $\eta = -\ln \tan(\frac{\theta}{2})$. This quantity is preferred as it is purely a
490 geometric quantity, independent of particle energy. Angular separation between
491 particles in ATLAS are given by $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2}$. The distance from the
492 beamline is given by $r = \sqrt{x^2 + y^2}$

493 6.2 Inner Detector

494 The Inner Detector (ID) was designed to identify and reconstruct vertices,
495 distinguish pions from electrons, and measure the momentum of charged particles.
496 The ID uses three different technologies for particle reconstruction: the Pixel
497 Detector, Semiconductor Tracker (SCT), and the Transition Radiation Tracker
498 (TRT), shown in Figure 6.4 and 6.5. The entire ID is immersed in a 2T solenoidal
499 magnetic field parallel to the $+z$ -axis, causing charged particles to bend in the
500 transverse-plane, allowing particle momentum measurements.

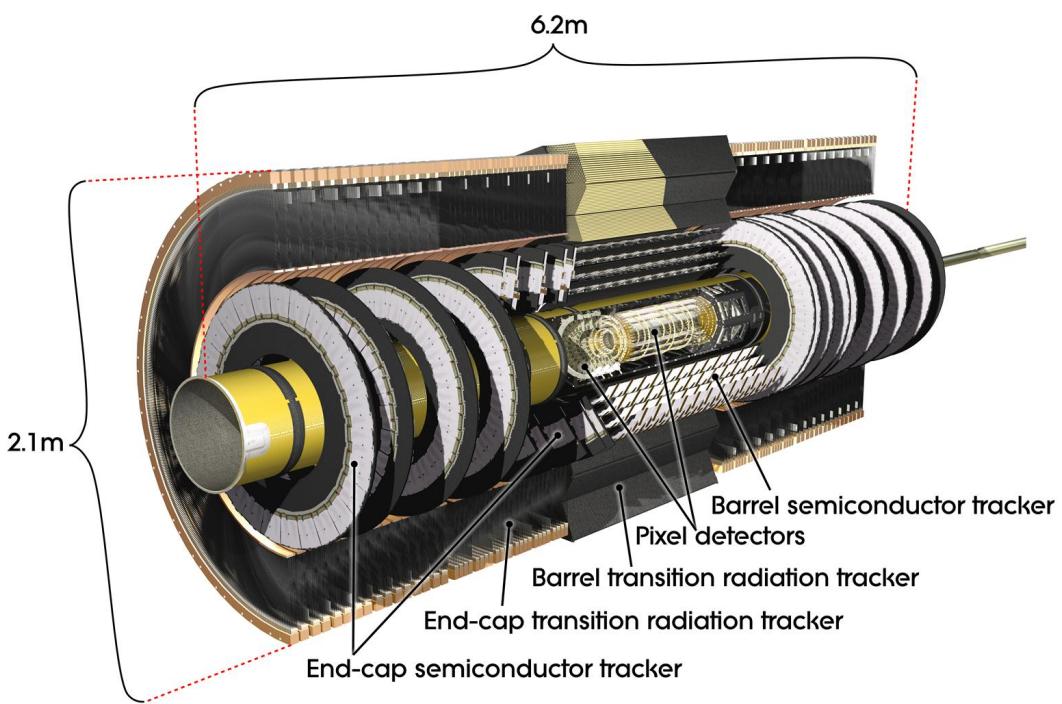


Figure 6.4: Layout of ATLAS Inner Detector

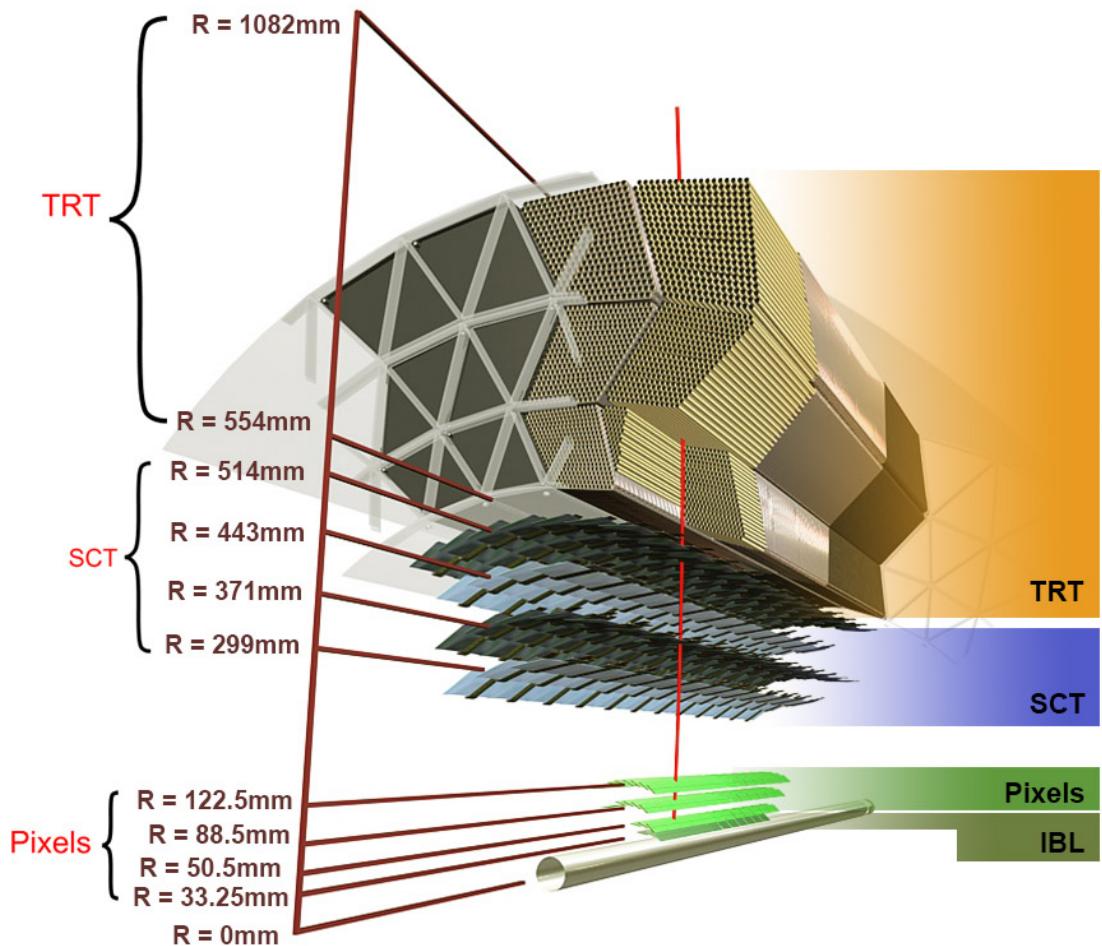


Figure 6.5: Layout of ATLAS ID Barrel System.

501 **6.2.1 Pixel Detector**

502 The pixel detector consists of four barrel layers between $r = 32.7$ and 122.5
503 mm, extending to $|z| = 400.5$ mm. The remaining detectors are arranged in bar-
504 rels and forward and backward rings. The innermost pixel barrel, the Insertable
505 b-Layer (IBL), only extends to $|z| = 332$ mm. The pixel detectors closer to the
506 beam line (larger η values) consists of six parallel cylindrical rings of pixel de-
507 tectors transverse to the beam line. The entire pixel detector consists of 1744
508 identical pixel sensors each with 46080 readout channels, totaling about 80 mil-
509 lion individual pixels. Most of the pixel sensors are $50 \times 400 \mu\text{m}^2$. Each pixel has
510 a position resolution of $14 \mu\text{m}$ in ϕ and $115 \mu\text{m}$ in the z direction.

511 **6.2.2 Semiconductor Tracker**

512 The SCT is located outside the pixel detector and has the same barrel and
513 endcap geometry as the pixel detector. SCT sensors are $80 \mu\text{m} \times 12$ cm with
514 a $80 \mu\text{m}$ strip pitch. In the barrel the strips are parallel to the z -axis and are
515 segmented in ϕ . In the endcaps, the strips extend radially. Sensors are grouped in
516 modules containing two layers of strips rotated 40 mrad with respect to each other.
517 This offset allows for the two-dimensional position of a track to be determined by
518 identifying the crossing point of the strips that registered a hit. SCT modules
519 measure tracks with an accuracy of $17 \mu\text{m}$ in $r - \phi$ and $580 \mu\text{m}$ in $z(r)$ in the
520 barrel (end-cap) region.

521 **6.2.3 Transition Radiation Tracker**

522 The transition radiation tracker (TRT), enveloping the SCT, is a gaseous
523 straw-tube tracker mainly used for electron/pion track separation. Each straw
524 is 4 mm in diameter and filled with a Xe- CO_2 - O_2 gas mixture. An anode wire at

525 the center of the straw is held at ground potential, while the walls of the straw
526 are kept at -1.4kV. When a charged particle passing through the TRT ionizes the
527 gaseous mixture, the resulting ions form an avalanche on the anode wire with a
528 gain of $\sim 10^4$. The signal from the anode wire is then digitized and discriminated.
529 Signals passing a low threshold cutoff are used to distinguish noise from tracks.
530 Signals passing a high threshold cutoff are sensitive to transition radiation (TR).
531 TR photons are emitted when charged particles pass between materials with dif-
532 ferent dielectric constants. The probability that a charged particle with energy E
533 and mass m passing between two materials emits a TR photon in the keV range
534 is proportional to $\gamma = E/m$. In the TRT straws these often then convert via the
535 photoelectric effect, causing a large avalanche triggering the high-threshold. Since
536 electrons have a smaller mass than pions, electron tracks are more likely to trig-
537 ger the high threshold. This then provides discrimination between electrons and
538 charged hadrons.

539 The barrel region of the TRT extends from $r = 563\text{-}1066$ mm and $|z| < 712$
540 mm. Barrel Straws are 144 cm long (divided $\sim \eta \approx 0$) and orientated parallel to
541 the beam direction. End-cap straws extend radially and are 37 cm long. There
542 are 53,544 straws in the barrel and 160,000 straws in the end-caps. Radiator mats
543 of polypropylene/polyethylene fibers in the barrel are aligned perpendicular to the
544 barrel straws (with holes for the straws to pass through). In the end-cap region,
545 radiator foils are layered between the radial TRT straws.

546 The arrival time of the signal pulse is sensitive to the distance between the
547 charged particle track and the anode wire and allows for a hit resolution of $130\mu\text{m}$.
548 The TRT extends to $|\eta| = 2.0$ and provides about 36 hits per track.

549 6.3 Calorimeters

550 The ATLAS electromagnetic and hadronic calorimeters (EMC and HCAL,
551 respectively) absorb and measure the energy of high energy hadrons, photons,
552 and electrons with $|\eta| < 4.9$. Both systems use sampling calorimeters which
553 consist of alternating layers of dense absorbing and active layers. In the absorbing
554 layer particles interact and lose energy, creating showers. These showers are then
555 detected and measured in the active layer. The amount of charge measured in the
556 active material scales with the energy of the incident particle, and thus provides a
557 measurement of the particle's energy. An overview of the layout of the calorimeter
558 system is shown in Figure 6.6.

559 The EMC measures and contains the energy of electromagnetically interacting
560 particles. It consists of layered accordion-shaped Lead absorber plates and elec-
561 trodes immersed in liquid Argon with 170k channels.. Using accordion-shaped
562 electrode and absorbers ensures ϕ symmetry and coverage. The EMC is com-
563 posed of a barrel part ($|\eta| < 1.475$), two end-caps ($1.375 < |\eta| < 3.2$), and a
564 presampler ($|\eta| < 1.8$). The presampler, containing only liquid Argon, corrects
565 for upstream energy losses of electrons and photons. The EMC barrel is segmented
566 into three layers. The first layer has finest segmentation with readout cells ex-
567 tending $\Delta\eta \times \Delta\phi = 0.025/8 \times 0.1$. This provides a precise shower measurements
568 used to separate prompt photons from $\pi^0 \rightarrow \gamma\gamma$ decays. The second layer has
569 coarser segmentation and is approximately 16 radiation lengths long. A radiation
570 length is the average distance an electron travels before losing all but $1/e$ of its
571 energy to bremsstrahlung. The last layer is the most coarse and measures the tail
572 of the electromagnetic shower. A schematic of the ECAL is shown in Figure 6.7.

573 The hadronic calorimeter located outside the EMC and is used to contain
574 and measure the energy of hadronically interacting particles. It consists of a tile

575 calorimeter (TileCal), hadronic end-cap calorimeter (HEC), and liquid Argon for-
 576 ward calorimeter (FCAL). TileCal is located behind the LAr EMC and uses steel
 577 absorbers and liquid Argon as the active material. TileCal consists of three barrel
 578 layers in the central and forward regions, extending up to $|\eta| < 1.7$. Photons
 579 generated from hadronic interactions are collected via wavelength-shifting fibers
 580 connected to photomultiplier tubes, as shown in Figure 6.8. The HEC lies behind
 581 the EMC endcap wheels. It uses copper absorbers and liquid Argon as the active
 582 material and covers $1.5 < |\eta| < 3.2$. Finally, the FCAL covers $3.1 < |\eta| < 4.9$
 583 and consists of three modules all using liquid Argon as the active material. The
 584 first module uses copper absorber and was designed for electromagnetic measure-
 585 ments. The second and third modules consist of tungsten absorber and are used
 586 to measure the kinematics of hadronically interacting particles. A schematic of
 587 the HCAL is shown in Figure 6.8.

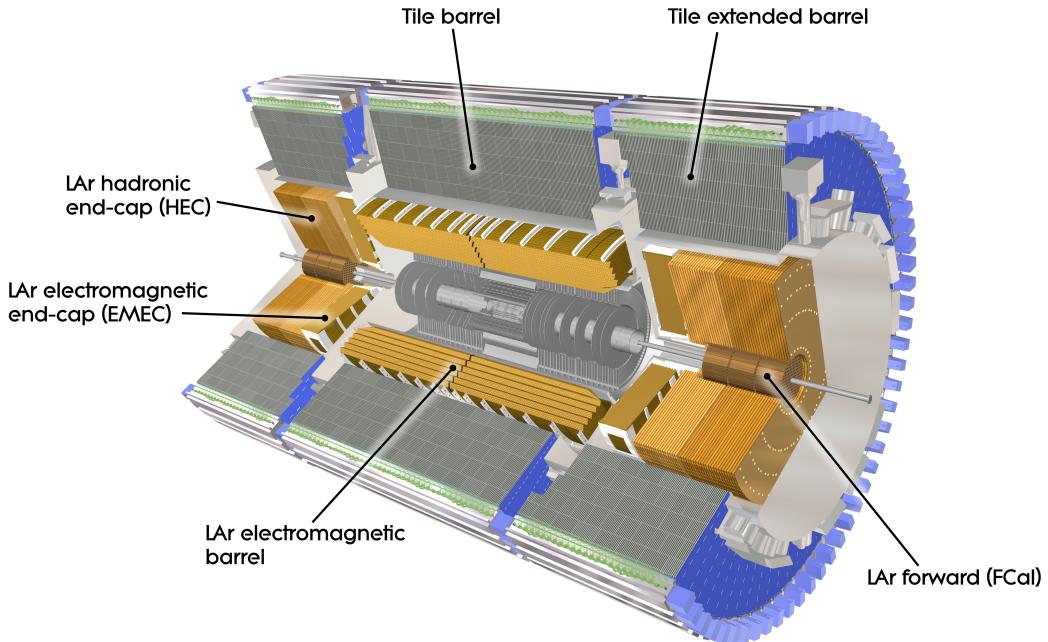


Figure 6.6: Overview of ATLAS electromagnetic and hadronic calorimeters.

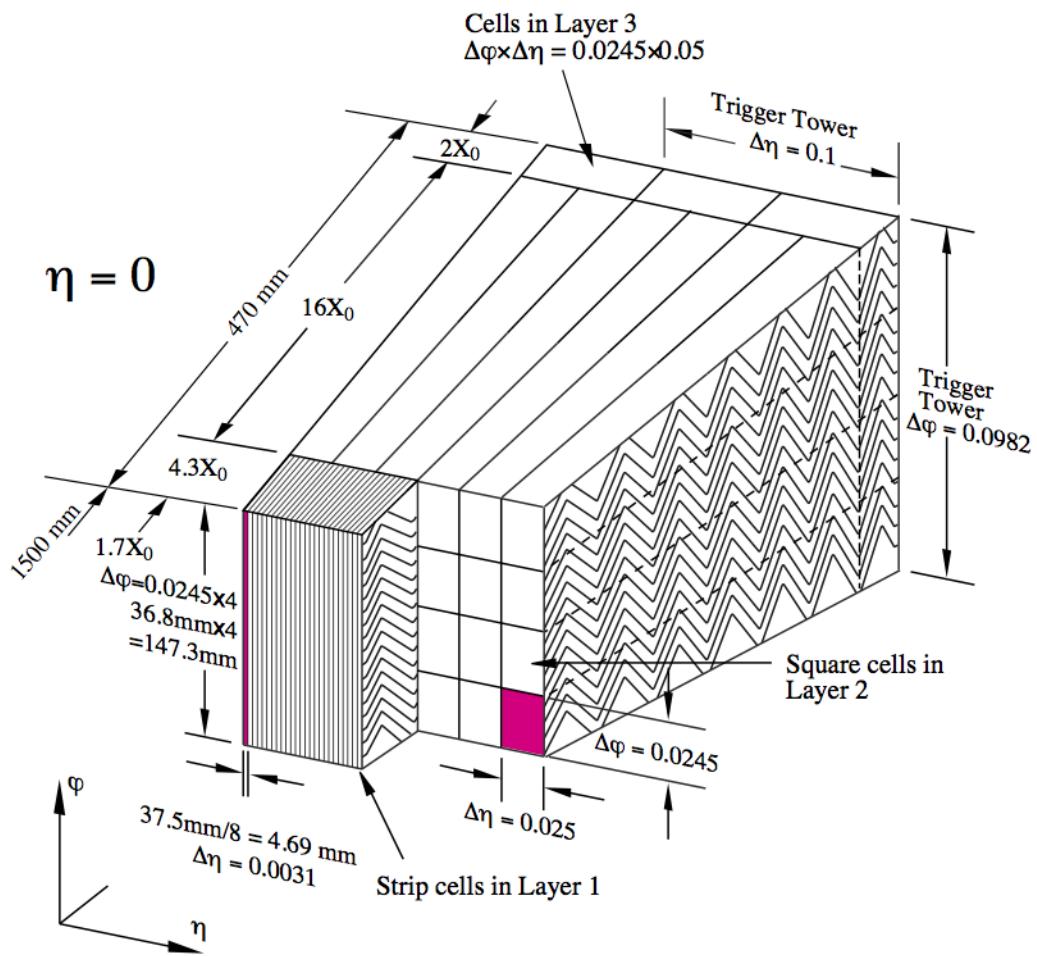


Figure 6.7: Schematic of ECAL.



Figure 6.8: Schematic of HCAL.

588 The energy resolution of the calorimeter subsystems are:

589 $\frac{\sigma_E}{E} = \frac{10\%}{\sqrt{E}} \oplus 0.7\% \text{ EMC}$

590 $\frac{\sigma_E}{E} = \frac{50\%}{\sqrt{E}} \oplus 3\% \text{ hadronic barrel}$

591 $\frac{\sigma_E}{E} = \frac{100\%}{\sqrt{E}} \oplus 10\% \text{ hadronic end-cap}$

592

6.4 Muon Spectrometer

593 The muon spectrometer (MS) is the outermost detector system in ATLAS.

594 Muons with a $p_T > 4$ GeV are energetic enough to reach the MS. To measure the
595 momentum of these muons barrel and end-cap toroid magnets are used covering
596 $|\eta| < 1.4$ and $1.6 < |\eta| < 2.7$. For $1.4 < |\eta| < 1.6$, a combination of the barrel
597 and end-cap toroidal magnetic fields bend muon trajectories. The detector in the
598 barrel region form three concentric rings at $R = 5, 7.5, 10$ m and are segmented
599 in ϕ to accommodate the magnets. The end-cap region consists of three circular
600 planes perpendicular to z and located at $|z| = 7.4, 14, 21.5$ m from the interaction
601 region. An additional detector at $|z| = 10.8$ m covers the transition region between
602 the barrel and end-cap.

603 The MS readout consists of four subsystems: Monitored Drift Tubes (MDT),
604 Cathode Strip Chambers (CSC), Resistive Plate Chambers (RPC), and Thin Gap
605 Chambers (TGC). The first two subsystems are used primarily for measuring
606 muon track parameters, while the RPC and TGC subsystems are used for muon
607 triggering. A schematic of this system is shown in Figure 6.9.

608 The MDT subsystem consists of precision tracking chambers for $|\eta| < 2.7$,
609 except for the inner most end-cap layer ($2.0 < |\eta| < 2.7$), where CSCs are used.
610 The basic unit of MDT chambers are thin walled Aluminum tubes with a diameter
611 of 3 cm and length of 0.9-6.2 m. These tubes are filled with a mixture of Ar-CO₂
612 gas with a 50μm W-Rn wire running down the center of the tube, which is kept at

613 3080 V. Since the maximum drift time of these chambers is ~ 700 ns, they are not
614 used for triggering. MDT chambers consist of 3-4 layers of tubes mounted on a
615 rectangular support system, as seen in Figure 6.10, orientated along ϕ to measure
616 the coordinate in the bending plane of the magnetic field with a resolution of 35
617 μm .

618 The MDT subsystem can only handle hit rates below $150\text{Hz}/\text{cm}^2$. For this
619 reason, CSCs are used in the innermost end-cap layer where hit rates are larger.
620 CSCs can handle hit rates up to $1000\text{Hz}/\text{cm}^2$. CSC are multiwire proportional
621 chambers. These chambers are filled with a Ar- CO_2 gas mixture and evenly spaced
622 wires kept at 1900 V. These wires are orientated in the radial direction but not
623 read out. Instead on one side of the cathode are copper strips parallel to the wires,
624 measuring η , while on the other side of the cathode are strips parallel to the wires
625 measuring ϕ . The width between strips is approximately 1.5 mm providing a
626 resolution of 60 μm in the bending-plane and 5 mm in the non-bending plane.

627 Since the CSC and MDT systems do not have prompt timing signals, the RPC
628 and TGC systems are used for triggering. The RPC system is used in the barrel
629 region ($|\eta| < 1.05$). RPC consist of two parallel resistive plates separated by a 2
630 mm insulated spacer with 100 mm spacing kept at 9.8 kV, as shown in Figure 6.11.
631 A gaseous mixture of $\text{C}_2\text{H}_2\text{F}_4$, C_4H_{10} , and SF_6 fills the space between the two
632 plates. Metallic strips on the outer faces of the plates are used to read out signals
633 produced by the gas ionizing. The middle barrel layer consists of two layers of
634 RPCs on either side of the MDT layer and one layer on the outermost MDT
635 layer. Each layer contains two orthogonal sets of metallic strips providing η and
636 ϕ measurements. The timing resolution of RPCs is 1.5 ns, and therefore may be
637 used to identify bunch crossings.

638 Finally, the TGCs are used in the end-cap regions and are primarily used to

639 provide L1 trigger decisions and ϕ measurements. TGCs are multi-wire propor-
 640 tional chambers consisting of arrays of gold-coated tungsten wires placed between
 641 two cathode planes. These wires are separated by 1.8 mm and cathodes are 1.4 mm
 642 from the wires. Orthogonal to the wires, on the opposite side of the cathode plane
 643 are copper strips held at 2900 V. The chambers are filled with a mixture of CO_2
 644 and n-pentane gas, the latter acts as a quenching gas to prevent avalanches initi-
 645 ated by secondary γ -rays from the primary avalanche. Figure 6.12 is a schematic
 646 of a TGC. The timing resolution of TGCs is less than 25 ns and therefore they
 647 are used for bunch crossing measurements.

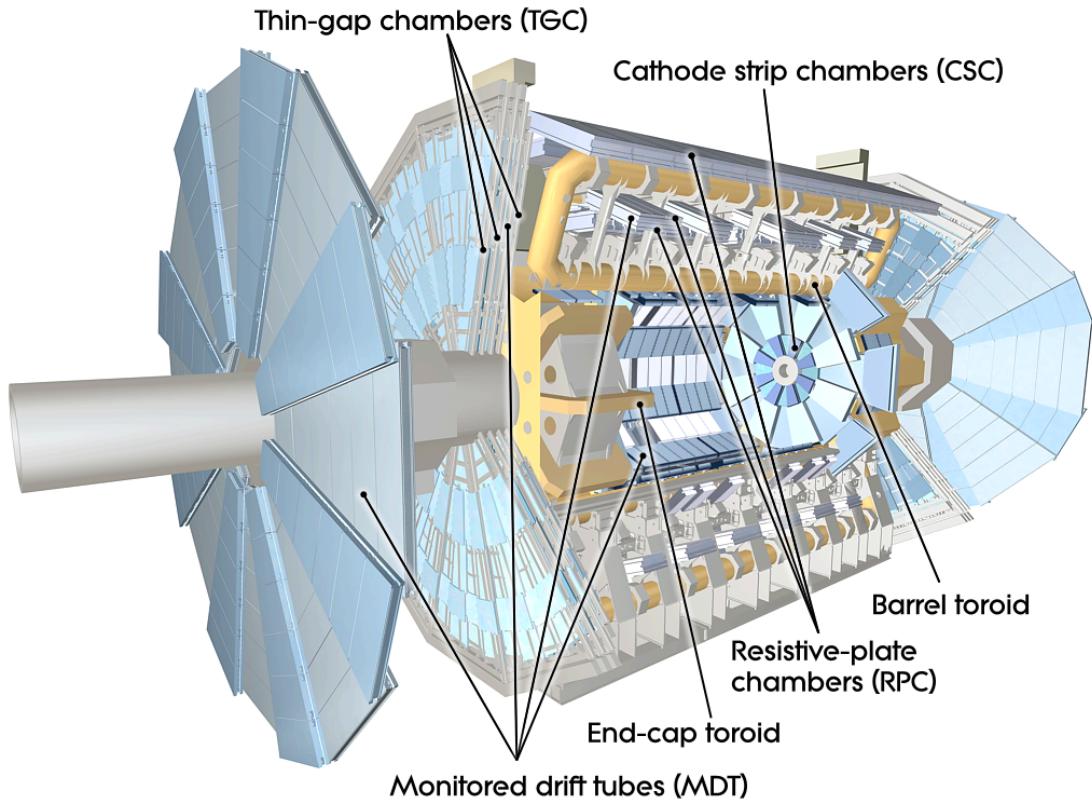


Figure 6.9: Schematic of Muon Spectrometer [cite G35]

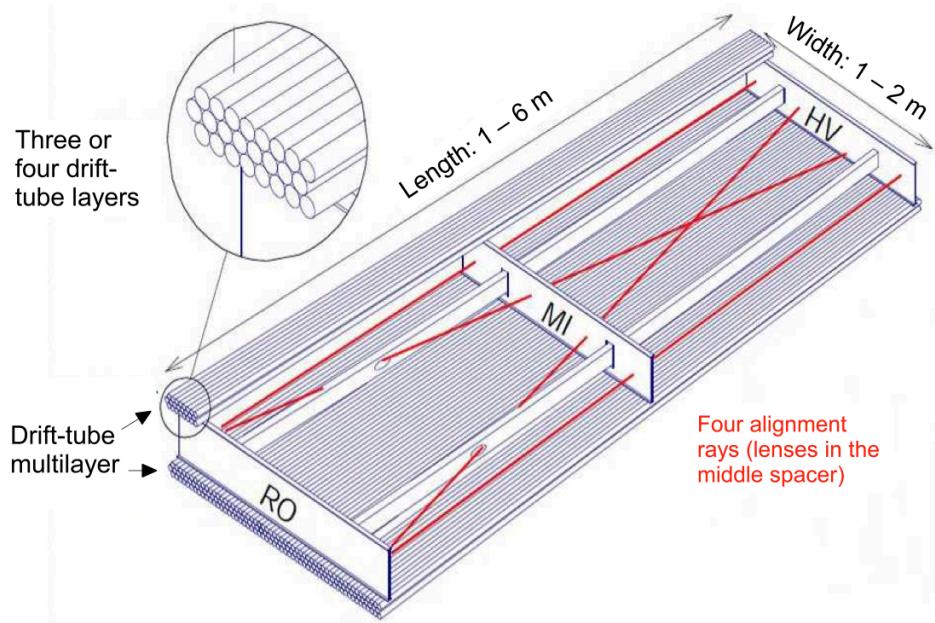


Figure 6.10: Schematic of MDT chamber. [cite G35]

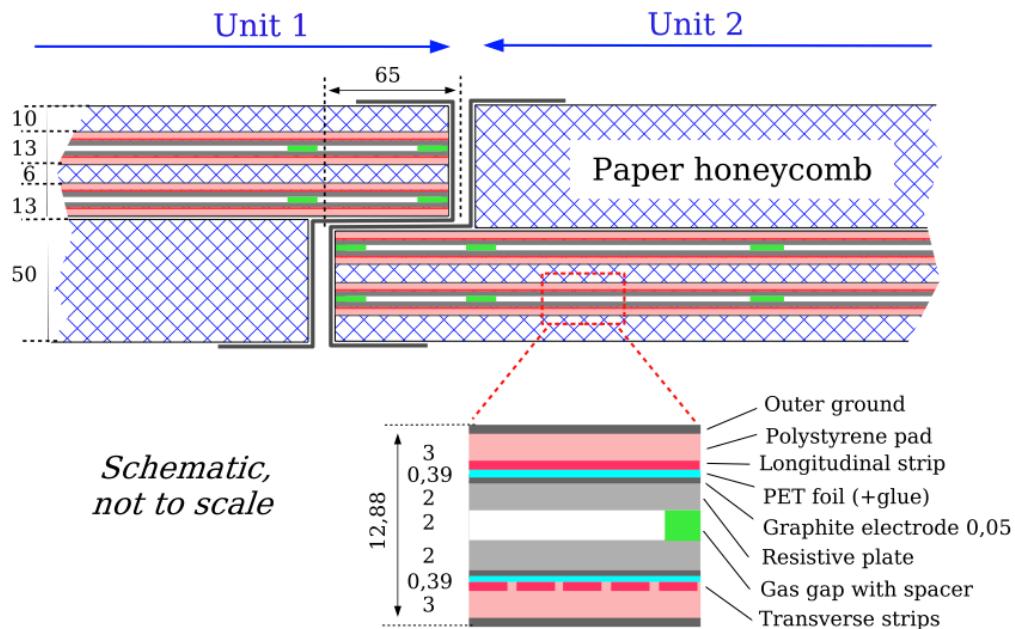


Figure 6.11: Schematic of RPC chamber, which is used for triggering in the central region of the detector [cite G35].

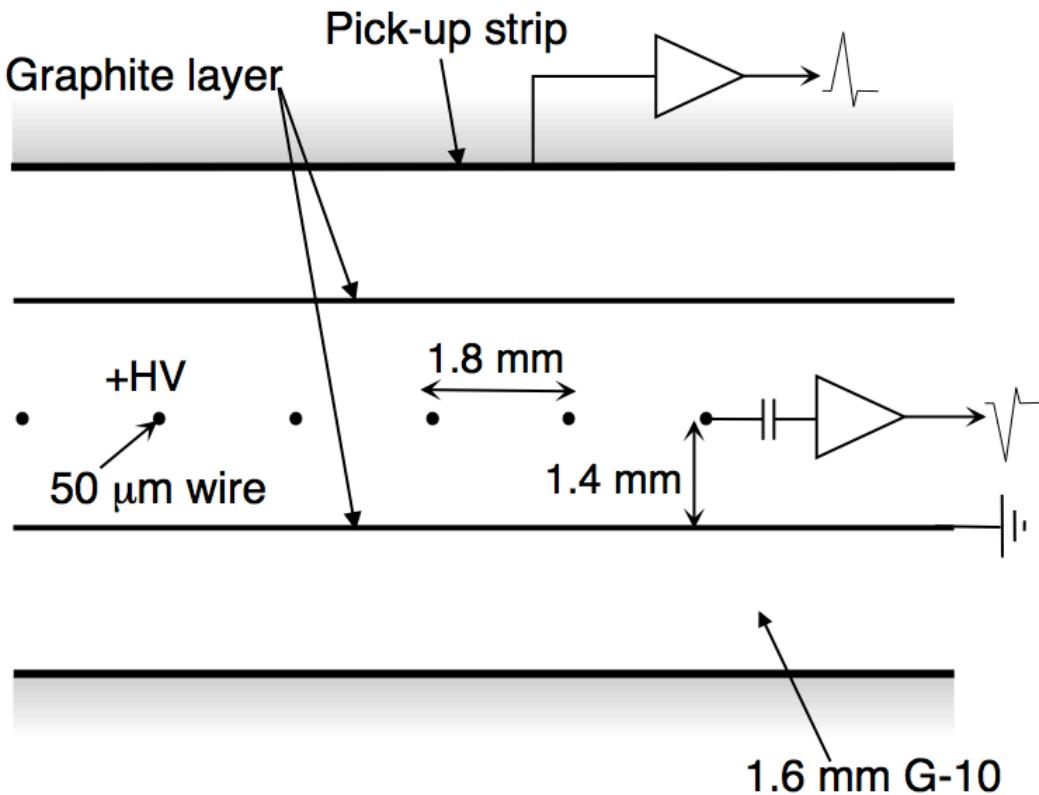


Figure 6.12: Schematic of TGC chamber, which is used for triggering in the muon end-cap region. [cite G35]

648 6.5 Magnet System

649 A particles with charge, q , and velocity v , moving in magnetic field, B , ex-
 650 periences a force, $F = qv \times B$. This force can cause charged particles to have a
 651 curved trajectory in magnetic fields, which the ID and MS use to determine the
 652 particles p_T . The central solenoid provides the magnetic field for the ID and the
 653 toroidal magnets provide the magnetic field for the MS.

654 The layout of the magnet system is shown in Figure 6.13. The central solenoid
 655 consists of a single-layer Al-stabilized NbTi conductor coil wound inside an Al

support cylinder. The solenoid is 5.8 m long, 50 cm thick and has an inner radius of 1.23 m. It is cooled to 4.5 K to reach superconducting temperatures and shares the liquid argon calorimeter vacuum vessel to minimize material in the detector. A current of 7.730kA produces a 1.998 T solenoidal magnetic field, pointing in the $+z$ direction.

The toroidal magnet system consists of a barrel and two end-cap toroidal magnets used to create a magnetic field outside the calorimeters that is orientated along ϕ . Each barrel toroid is 25.3 m long with an inner and outer diameter of 9.4 and 20.1 m and weighs 830 tonnes. Endcap toroids are 5 m long with an inner and outer radius of 1.65 and 10.7 m. Both toroid systems use Al-stabilized Nb/Ti/Cu conductors. The magnetic field strength in the barrel and endcap regions are 0.5 and 1 T, respectively.

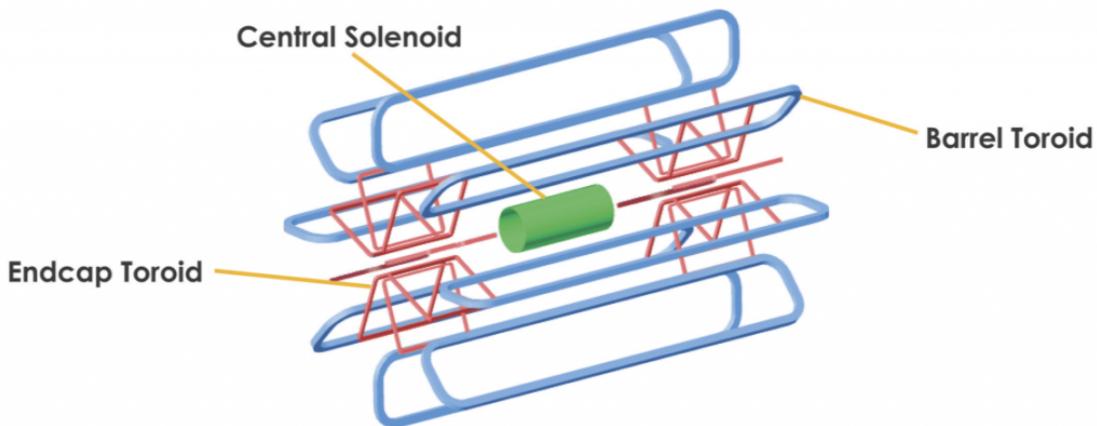


Figure 6.13: Layout of ATLAS magnet systems.

6.6 Trigger System

Since collisions occur every 25 ns and reading out all detector channels and storing that information is not currently feasible (would require saving 60 million megabytes per second), the majority of events are not kept for analysis. ATLAS

672 uses a multi-stage trigger system to select approximately 1,000 of the 1.7 billion
673 collisions that occur each second (corresponding to a rate of 1 kHz from the 40
674 MHz proton collision rate). The first stage of the trigger system is the hardware
675 level (L1) trigger. This trigger reduces the event rate to \sim 100 kHz by identifying
676 Regions-of-Interest (ROIs) containing high p_T leptons, photons, jets, or E_T^{miss} by
677 using information from RPCs, TGCs, and calorimeters to make a $2.5 \mu\text{s}$ decision.
678 This information is then passed to a high-level trigger (HLT) which further de-
679 creases event rates to \sim 1 kHz. The HLT uses finer granularity measurements
680 from the MS and ID to perform simplified offline reconstruction to decide which
681 events to keep.

Part IV

682

Method

683

684 **Chapter 7**

685 **Dataset and Simulated Samples**

686 **7.1 Dataset**

687 This analysis uses pp collision data collected from 2015 to 2018 at $\sqrt{s} = 13$
688 TeV, corresponding to 139/fb of data as shown in Figure 7.1 and 7.2. From this
689 dataset, only those events in which the tracker, calorimeters, and muon spectrom-
690 eter have good data quality are used. For a given event, the solenoid and toroidal
691 magnets must also be operating at their nominal field strengths. In addition to
692 this, events must pass further quality checks to reject events where detector sub-
693 systems may have failed. These selections reject events that containing LAr noise
694 bursts, saturation in the electromagnetic calorimeter, TileCal errors, and failures
695 in event recovery due to tracker failures. Events with information missing from
696 subsystems (usually due to busy detector conditions) are rejected. Events must
697 also contain a primary vertex with at least two associated tracks, where the pri-
698 mary vertex is selected as the vertex with the largest $\sum p_T^2$ over tracks associated
699 with the vertex and $p_T > 0.5$ GeV.

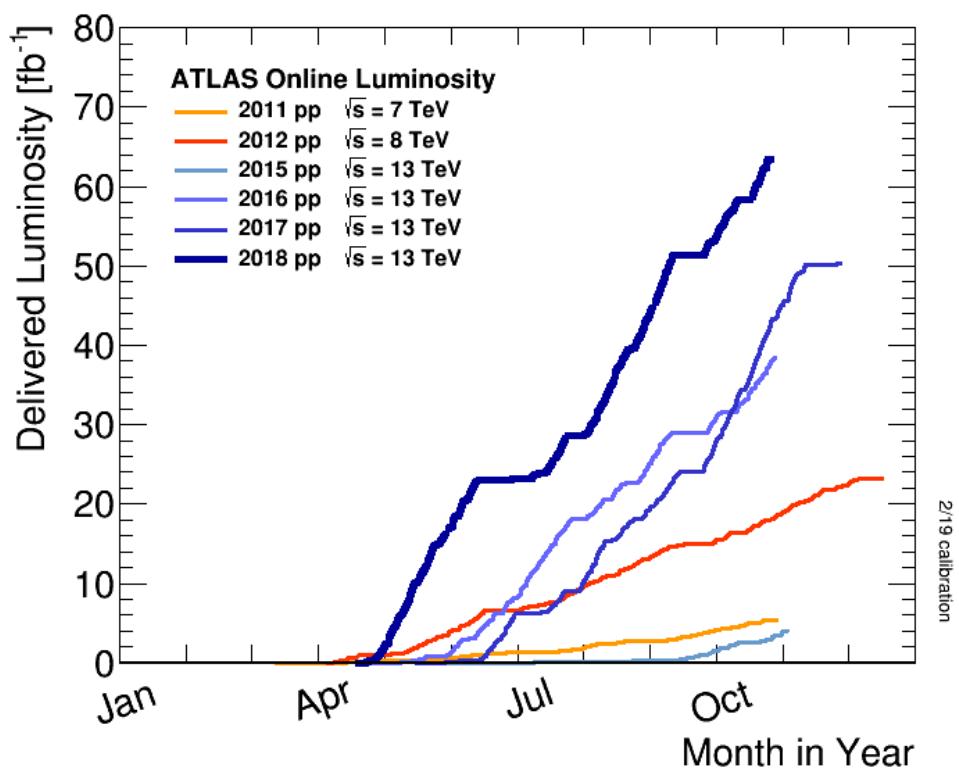


Figure 7.1: Integrated luminosity for data collected from ATLAS from 2011 - 2018



Figure 7.2: Mean number of interactions per crossing for data collected from ATLAS from 2011 - 2018

700 **7.2 Simulated Samples**

701 Samples are simulated in order to model backgrounds, evaluate signal ac-
702 ceptance, optimize event selection and estimate systematic and statistical uncer-
703 tainties. The dominant backgrounds for this analysis are $W/Z + \text{jets}$, diboson
704 (WZ/WW), $t\bar{t}$, single top and multijet production.

705 $W/Z+\text{jet}$ events are simulated using Sherpa 2.2.1 at NLO [cite [29]] and merged
706 with the Sherpa parton shower using the ME+PS@NLO prescription [13]. These
707 events are then normalized to NNLO cross sections. The $t\bar{t}$ and single-top back-
708 grounds are generated with Powheg-Box with NNPDF3.0NLO PDF sets in the
709 matrix element calculation [cite[35]]. For all processes, the parton shower, frag-
710 mentation, and underlying event are simulated using Pythia 8.320 with the A14
711 tune set[cite[ATL-PHYS-PUB-2014-02]]. Diboson processes are generated using
712 Sherpa 2.2.1.

713 Signal samples are simulated using MadGraph 5-2.2.2 [cite 42] and Pythia
714 8.186 with NNPDF230LO. RS Graviton samples are generated with $k/M_{PL}=1$.
715 HVT Model A (B) samples are simulated with $g_V = 1(3)$, as the difference in the
716 width of the samples is smaller than detector resolution. Model C is generated by
717 setting $g_H = 1$ and $g_f = 0$ to model VBF production of HVT bosons. Signals are
718 generated for masses between 300 GeV and 5 TeV.

719 **Chapter 8**

720 **Objects**

721 **8.1 Electrons**

722 Electrons are reconstructed from electromagnetic showers in the LAr EM
723 calorimeter. During reconstruction cells of $\Delta\eta \times \Delta\phi = 0.025 \times 0.025$ are grouped
724 into 3×5 clusters. These clusters are then scanned for local maxima that seed
725 electron clusters. These clusters must then be matched to ID track from the PV.
726 This requirement minimizes non-prompt electron and fake electron backgrounds.
727 Electrons must pass identification and isolation requirements. Electron identifica-
728 tion (loose, medium, tight) classification is based on a multivariate discriminant
729 that identifies electrons using a likelihood based method. For this analysis, events
730 are required to have one tight electron and no additional loose electrons. Elec-
731 trons are also required to be isolated. The electrons are considered isolated if the
732 quotient of the sum of the transverse momentum (of calorimeter energy deposits)
733 in a cone around the electron of size $\Delta R = 0.2$ and the transverse momentum
734 of the electron to be less than $0.015 * p_T$ or 3.5 GeV, whichever is smaller. This
735 requirement rejects non-prompt photons and other fake leptons. Electrons in this
736 analysis are also required to have $p_T > 30$ GeV and $|\eta| < 2.47$. Electrons are also

737 required to have $p_T > 30$ GeV.

738 Electrons are calibrated to determine data-driven scale factors using $J/\Psi \rightarrow$
739 ee , $Z \rightarrow ee$, $Z \rightarrow \ell\ell\gamma$ processes. These corrections account for the non-uniform
740 response of the detector which introduces modeling and reconstruction uncertain-
741 ties.

742 8.2 Muons

743 As muons traverse the entire detector, they are reconstructed from ID and
744 MS tracks. For this analysis the muon identification and isolation working points
745 are chosen to minimize the contributions from non-prompt muons. Towards this
746 end, each selected event must contain exactly one muon that passes the medium
747 identification working point, and no additional muons (that pass the loose working
748 point). For the medium working point, two types of reconstructed muons are
749 used: combined and extrapolated muons (CB and ME, respectively). For CB
750 muons, ID and MS tracks are reconstructed independently and a combined track
751 fit is performed by adding or removing MS tracks to improve the fit quality.
752 ME muons are reconstructed from only MS tracks with hits in at least two layers,
753 which ensures the track originates from the PV. ME muons extend the acceptance
754 for muon reconstruction outside the ID from $2.5 < |\eta| < 2.7$. The medium
755 identification working point uses CB and ME tracks. CB tracks must have at
756 least 3 hits in two MDT layers. ME tracks are required to have at least three
757 MDT/CSC hits. To further minimize contributions from fake muons, the selected
758 muons are required to be isolated from other tracks, as muons from W, Z decays are
759 often isolated from other particles. To insure the selected muons are isolated, the
760 scalar sum of the transverse momentum of tracks in a cone of $\Delta R = 0.3$ compared
761 to the transverse momentum of the muon must be less than 0.06. Muons are also

762 required to have $p_T > 30$ GeV.

763 Muons are calibrated using well-studied resonances $J/\Psi \rightarrow \mu\mu$ (low- p_T), $Z \rightarrow$
764 $\mu\mu$ (high- p_T). Figure 8.1 shows the combined muon p_T uncertainty from this
765 calibration. The total systematic uncertainty is less than 1% for all p_T ranges
766 considered in this analysis.

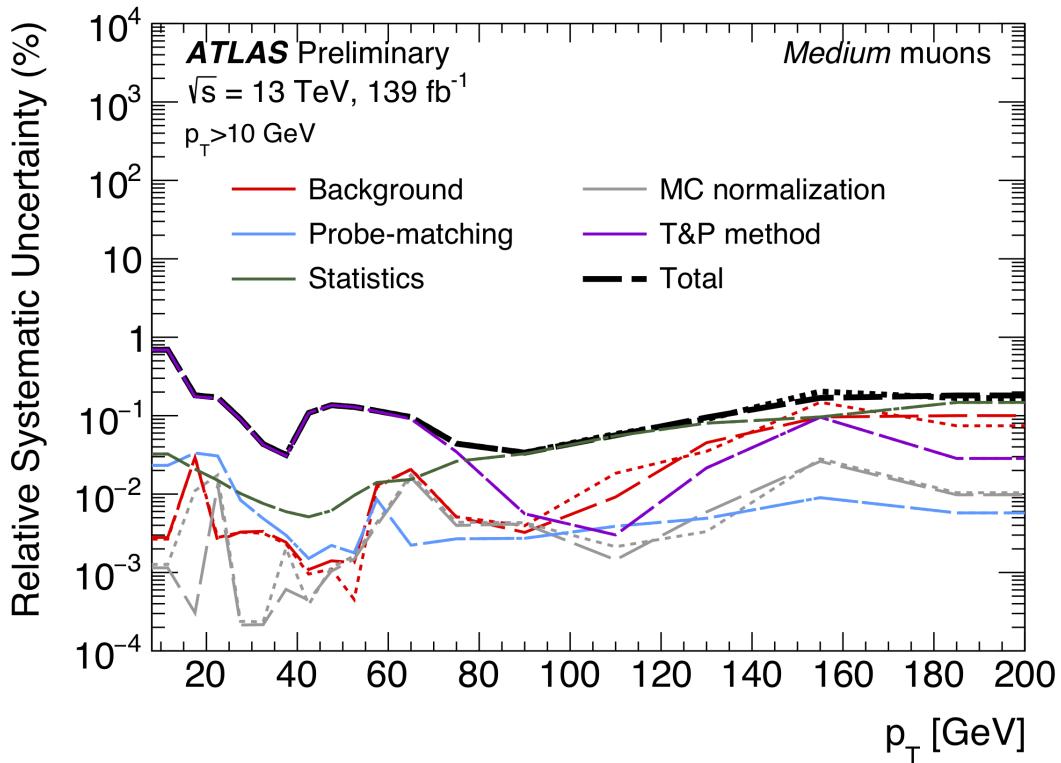


Figure 8.1: This figure shows the breakdown of the muon reconstruction efficiency scale factor measured in $Z \rightarrow \mu\mu$ as a function of p_T [5].

767 8.3 Jets

768 Three types of jets are used in this analysis: variable radius, small-R and
769 large-R jets. Variable radius jets are used to reconstruct Z bosons decaying to
770 two b -jets in the jet catchment area of large-R jet in the Merged regime. Small-R

771 jets are used to reconstruct the hadronically decaying W/Z candidates in the
772 resolved analysis and the forward jets from resonances produced through vector
773 boson fusion. Large-R jets are used to reconstruct the hadronically decaying boson
774 in the merged regime.

775 For all of these jet collections, the jet energy is calibrated sequentially as shown
776 in Figure 8.2. After the jet direction is corrected to point to the PV, the energy
777 of the jet is corrected. First, the jet energy is corrected to account for pileup
778 contributions based on the p_T and area of the jet (these corrections are extracted
779 from a $pp \rightarrow jj$ sample). Following this, another pileup correction is applied that
780 scales with μ and N_{PV} .

781 MC-based corrections are then applied that are meant to transform the jet
782 energy and η back to truth level. Therefore, these corrections account for the
783 non-compensating nature of the ATLAS calorimeters and inhomogeneity of the
784 detector. Following this, the Global Sequential Calibration is applied that re-
785 duces flavor dependence of jet calibrations and accounts for energy leakage of jets
786 outside the calorimeters. Finally, in-situ corrections are applied that account for
787 differences in jet responses between data and simulation ($\gamma/Z + \text{jet}$ and fake lep-
788 ton samples are used). These differences can be due to mismodelling of the hard
789 scatter event, pile-up, jet formation, etc.

790 To further reject fake jets, jets must pass quality requirements based on the
791 following variables ([cite P42]):

- 792 - f_Q^{LAr} : fraction of energy of jet's LAr cells with poor signal shape
- 793 - f_Q^{HEC} : fraction of energy of jet's HEC cells with poor signal shape
- 794 - E_{neg} : sum of cells with negative energy
- 795 - f_{EM} : fraction of jet's energy deposited in EM calorimeter

- 796 - f_{HEC} : fraction of jet's energy deposited in HEC calorimeter
- 797 - f_{max} : maximum energy fraction in any single calorimeter layer
- 798 - f_{ch} : ratio of the scalar sum of the p_T of a jet's charged tracks to the jet's p_T
- 799 Jets selected for the resolved analysis must pass one of the following criteria:
- 800 - $f_{HEC} > 0.5$ and $|f_Q^{HEC}| > 0.5$ and $\langle Q \rangle > 0.8$
- 801 - $|E_{neg}| > 60$ GeV
- 802 - $f_{EM} > 0.95$ and $f_Q^{LAr} > 0.8$ and $\langle Q \rangle > 0.8$ and $|\eta| < 2.8$
- 803 - $f_{max} > 0.99$ and $|\eta| < 2$
- 804 - $f_{EM} < 0.05$ and $f_{ch} < 0.05$ and $|\eta| < 2$
- 805 - $f_{EM} < 0.05$ and $|\eta| > 2$

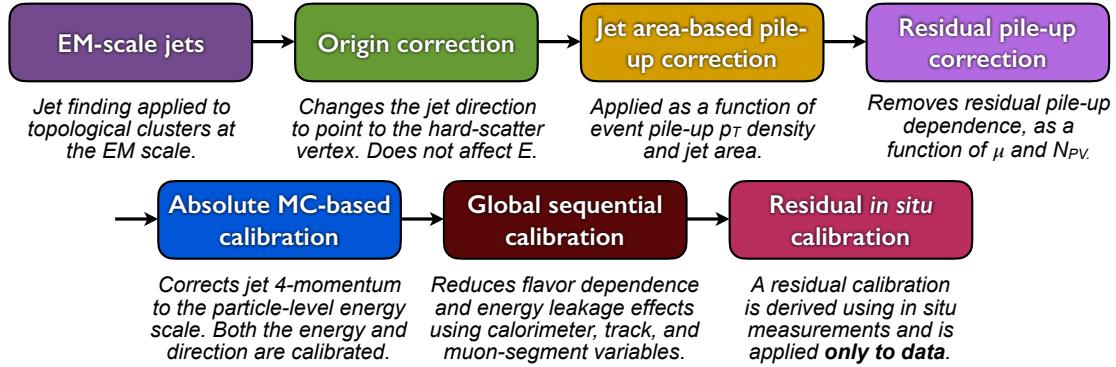


Figure 8.2: [8] This diagram shows the calibration stages for EM jets.

806 **8.3.1 Small-R jets**

807 Small-R jets are used to reconstruct the hadronically decaying W/Z candi-
808 date when the two resulting jets are well-separated in $\eta\text{-}\phi$ space. Small-R jets
809 are also used to identify forward jets from resonances produced through vector
810 boson fusion. Small-R jets are constructed from topologically connected clusters
811 of calorimeter cells (topoclusters), seeded from calorimeter cells with energy de-
812 posits significantly above the noise threshold. These cells are then used as inputs
813 to the $anti-k_t$ algorithm [16] with a radius parameter, $R = 0.4$.

814 Jets used in this analysis must have $p_T > 30$ GeV and $|\eta| < 2.5$. To further
815 reduce fake jets the jet-vertex-tagger (JVT) is used to reject pile-up jets [7]. The
816 JVT uses two track-based variables, corrJVF and R_{p_T} to calculate the likelihood
817 that the jet originated from the PV. The corrJVF compares the scalar sum of the
818 p_T of tracks associated with the jet and PV to the scalar sum of the p_T of tracks
819 associated with the jet. This variable also includes a correction that reduces the
820 dependency of corrJVF with the number of reconstructed vertices in the event.
821 The other discriminant, R_{p_T} , is given by the ratio of the scalar sum of the p_T of
822 tracks associated with the jet and PV to the p_T of the jet. Both of these variables
823 peak around zero for pileup jets, as these jets are unlikely to have tracks associated
824 with the PV. JVT cuts are applied to all jets with $p_T > 120$ GeV. Central jets
825 ($|\eta| < 2.4$) are required to have a $JVT > 0.59$ and forward jets ($2.4 < |\eta| < 2.5$)
826 are required to have $JVT > 0.11$.

827 **8.3.2 Large-R jets**

828 Large-R ($R = 1.0$) jets are used to reconstruct the hadronically decaying W/Z
829 candidate when the resulting jets are not well-separated in $\eta\text{-}\phi$ space, and overlap
830 forming one large-R jet. Track-Calorimeter Clusters (TCCs) are used to reconstruct these

jets [cite ANA 50]. These jets are constructed using a pseudo particle flow method
 using ID tracks matched to calorimeter clusters. The angular resolution of the
 calorimeter degrades sharply with jet p_T , but the jet energy resolution improves.
 The tracker has excellent angular resolution which improves with p_T . Therefore,
 by matching tracks to jets, TCCs have more precise energy and angular resolution
 than jets constructed from calorimeter information only. These jets are required
 to have $p_T > 200$ GeV, $|\eta| < 2.0$ and $m_J > 50$ GeV.

TCC jets are trimmed as detailed in [cite ANA 45], which suppresses pileup
 and soft radiation in the jet, the jet mass is calculated as the four-vector sum
 of the jet's constituents (assuming massless constituents). The jet mass peaks
 around the W/Z boson mass for $W/Z \rightarrow qq$ jets, and more broadly for quark and
 gluon induced jets.

These jets are then tagged as W jet if it passes optimized jet mass and D_2
 cuts for W bosons, and a Z jet if it passes the optimized cuts for the Z boson.
 The jet substructure variable D_2 is given by the ratio of energy correlation func-
 tions. These fuctions are derived from the energies and pair-wise angles of a jet's
 constituents [cite ANA 46, 47]:

$$D_2^{\beta=1} = E_{CF3} \left(\frac{E_{CF1}}{E_{CF2}} \right)^3 \quad (8.1)$$

Where the energy correlation functions are defined as:

$$E_{CF1} = \sum_i p_{T,i} \quad (8.2)$$

$$E_{CF2} = \sum_{ij} p_{T,i} p_{T,j} \Delta R_{ij} \quad (8.3)$$

$$E_{CF3} = \sum_{ijk} p_{T,i} p_{T,j} p_{T,k} \Delta R_{ij} \Delta R_{jk} \Delta R_{ki} \quad (8.4)$$

851 A two-dimensional optimization of the jet mass and D_2 thresholds was per-
 852 formed to provide maximum sensitivity for this analysis. This optimization was
 853 done by maximizing the signal sensitivity (using HVT W' and G_{KK} samples)
 854 against the single quark and gluon jet backgrounds in bins of jet p_T . Figure 8.3
 855 shows the optimized thresholds on D_2 and jet mass as a function of jet p_T . Figure
 856 8.4 shows the efficiency of the optimized W/Z taggers as a function of jet p_T .

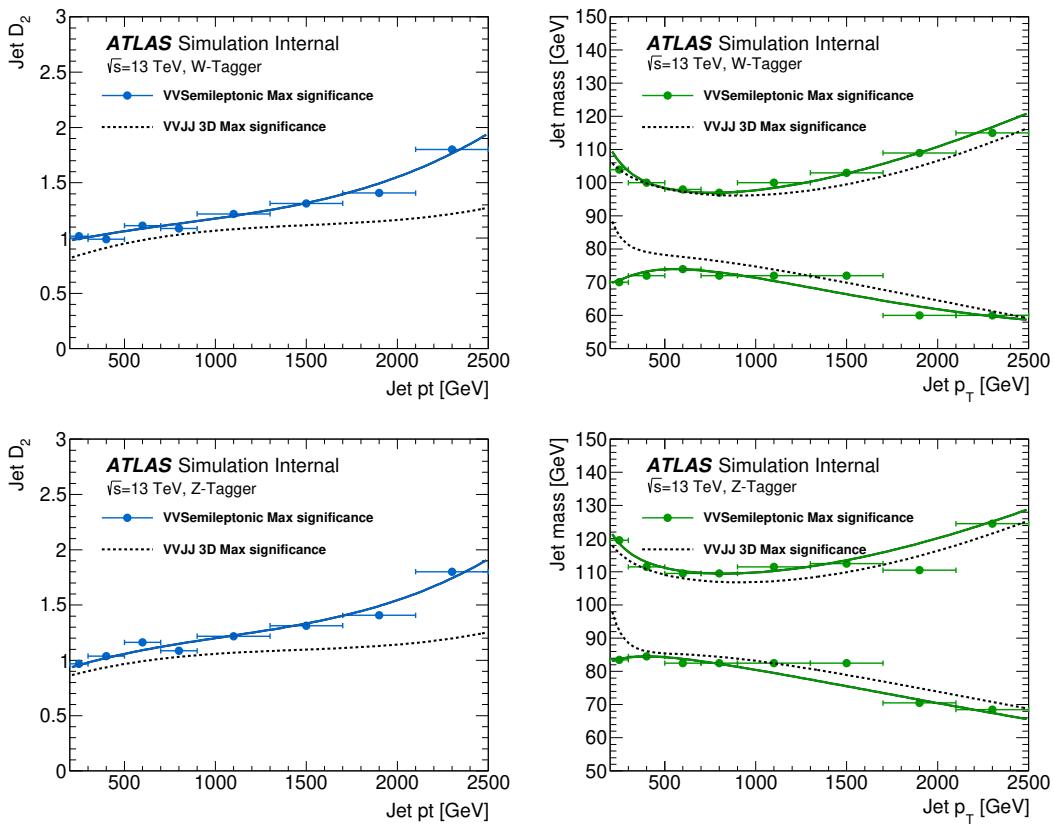


Figure 8.3: The upper cut on D_2 (a) and jet mass window cut i.e. the upper and lower boundary of the mass (b) of the W -tagger as a function of jet p_T . Corresponding values for Z -tagger are shown in (c) and (d). The optimal cut values for maximum significance are shown as solid markers and the fitted function as solid lines. Working points from $VV \rightarrow JJ$ [ATLAS-HDBS-2018-31-002] is also shown as dashed lines as a reference. Natasha reword?

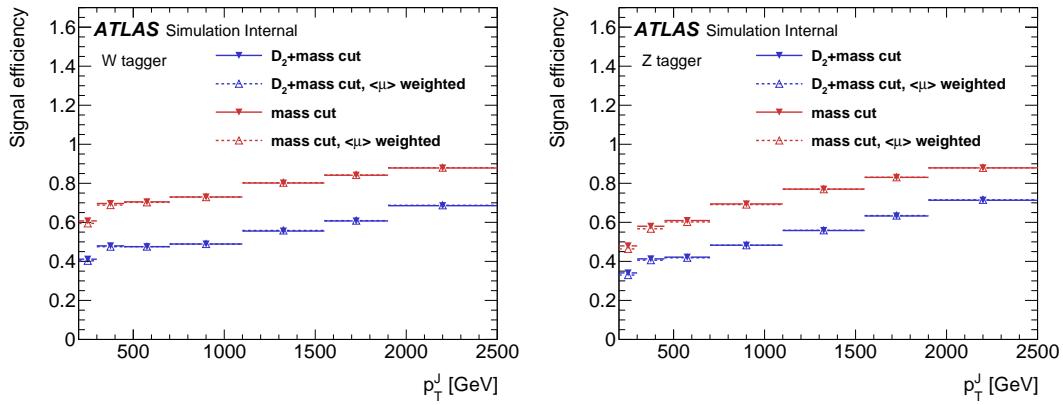


Figure 8.4: Natasha write caption

857 8.3.3 Variable Radius jets

858 To more accept more boosted Z bosons decaying to $b\bar{b}$ that would normally be
 859 rejected due to topological cuts discussed ?? variable radius (VR) track jets are
 860 used to identify b -jets within the catchment area of large-R jets [15]. VR jets are
 861 constructed from ID tracks using the anti- k_t algorithm with a radius parameter
 862 that depends on the p_T of the track, shown in Equation 8.5.

$$R_{eff}(p_{T,i}) = \frac{\rho}{p_{T,i}} \quad (8.5)$$

863 For this search $\rho = 30$ GeV and an upper and lower limit on cone size are set
 864 to 0.02 and 0.4, respectively, to prevent unphysical asymptotic behavior of ρ .
 865 Collinear VR jets are possible, so track jets that are not separated by the the
 866 smaller jet's cone size are not used. Additionally, VR jets are required to have
 867 $p_T > 10$ GeV and $|\eta| < 2.5$.

868 **8.3.4 Jet Flavor Tagging**

869 To further classify events, the small-R and VR jets originating from a b-quark
870 are classified using a multivariate b -tagging algorithm (BDT), MV2c10 [cite G 210
871 199]. This algorithm uses the impact parameters of the jet's ID tracks, secondary
872 vertices (if they exist), and reconstructed flight paths of b and c hadrons in the
873 jet to determine if the jet was induced by a b -quark. For this analysis the 85%
874 efficient working point of this algorithm is used giving c , τ , and light-flavor jet
875 rejection of 3, 8, and 34 respectively in simulated $t\bar{t}$ samples.

876 **8.4 MET/Neutrinos**

877 As neutrinos are uncharged and colorless they do not leave tracks or jets in the
878 detector. For this reason, neutrinos are reconstructed calculated the E_T^{miss} , the
879 negative vector sum of p_T all the physics objects and an extra "soft" term. The
880 "soft" term accounts for energy deposits not associated with any of the objects in
881 the event. For this analysis the soft term is given by the sum p_T of all ID tracks
882 not associated with objects in the event. The selected tracks must be matched to
883 the primary vertex, which decreases pile-up contamination [cite G 217 218]. The
884 tight working point is used [Natasha look up what this means].

885 **8.5 Overlap Removal**

886 Reconstructed jets and leptons in this analysis can arise from the same energy
887 deposits. For instance, a cluster of energy from an electron can also be a valid
888 calorimeter seed for a jet. To mitigate this confusion of multiple objects originating
889 from a single jet or lepton overlapping objects are removed via a procedure referred
890 to a overlap removal. In this procedure the separation of the two objects, $\Delta(R) =$

891 $\sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ determines which object is removed from the event.

892 The overlap selections used in this analysis are:

- 893 - when an electron shares a track with another electron with the lower p_T
894 electron is rejected, as it is more likely to be a fake electron
- 895 - when a muon and electron share a track the muon is rejected if it is a
896 calo-muon, otherwise the electron is rejected
- 897 - when $\Delta R < 0.2$ for an electron and jet, the jet is rejected to maximize signal
898 acceptance
- 899 - when $\Delta R > 0.2$ for an electron and jet, the electron is rejected as likely
900 originated from decays within the jet
- 901 - when $\Delta R < \min(0.4, 0.04 + 10\text{GeV}/p_T^\mu)$ the muon is rejected, again maximiz-
902 ing signal acceptance, otherwise the jet is rejected
- 903 - when $\Delta R < 1.0$ for the a large-R jet and electron, the jet is rejected

904 8.6 Reconstructed Resonance Mass (m_{WV})

905 The WV system mass, m_{WV} is reconstructed from the lepton, neutrino, and
906 hadronically-decaying boson candidate. The momentum of the neutrino along the
907 z -direction is obtained by constraining the W boson mass of the lepton neutrino
908 system to be $80.3 \text{ GeV}/c^2$. For complex solutions to this constraint, p_Z is taken
909 as the real component of the solution. For real solutions, the one with the smaller
910 absolute value is used. For the resolved analysis, m_{WV} is reconstructed by con-
911 straining the $W(Z)$ dijet system in the SRs, which improves the mass resolution:

$$p_{T,jj}^{corr} = p_{T,jj} \times \frac{m_{W/Z}}{m_{jj}} \quad (8.6)$$

$$m_{jj}^{corr} = m_{W/Z} \quad (8.7)$$

912 where m_{jj} and $m_{W/Z}$ are the reconstructed invariant mass of the hadronically-
913 decaying W/Z boson and the PDG values of the W/Z boson masses, respectively.
914 The reconstructed resonance mass is the final discriminating variable in this anal-
915 ysis. The distribution of this variable in the CR and SRs are used in the final
916 likelihood fit to search for evidence of an excess of events due to BSM resonances.
917 The distribution of m_{WV} are shown in Figures 12.1-12.9.

918 **Chapter 9**

919 **Event Selection and**
920 **Categorization**

921 **9.1 Pre-selection**

922 Before applying topological cuts to suppress backgrounds and reduce data
923 size in this search, preselection cuts are applied which include trigger and event
924 requirements. Events must contain exactly one tight lepton (no additional loose
925 leptons), the $p_T^{\ell\nu} > 75$ GeV, and there must be at least two small-R jets or one
926 large-R jet.

927 **9.2 Trigger**

928 The data were collected using the lowest unprescaled single-lepton or E_T^{miss}
929 triggers, as summarized in Table 9.1. Since the muon term is not considered in the
930 trigger E_T^{miss} calculation, the E_T^{miss} trigger is fully efficient to events with high- p_T
931 muons. For this reason, the E_T^{miss} trigger is used for events where $p_T^\mu > 150$ GeV, to
932 compensate for the poor efficiency of the single muon trigger above $p_T^\mu > 150$ GeV.

Table 9.1: The list of triggers used in the analysis.

Data-taking period	$e\nu qq$ channel	$\mu\nu qq$ ($p_T(\mu\nu) < 150$ GeV) channel	$\mu\nu qq$ ($p_T(\mu\nu) > 150$ GeV) channel
2015	HLT_e24_lhmedium_L1EM20 OR HLT_e60_lhmedium OR HLT_e120_lhloose	HLT_mu20_iloose_L1MU15 OR HLT_mu50	HLT_xe70
2016a (run < 302919) $(L < 1.0 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1})$	HLT_e26_lhtight_nod0_ivarloose OR HLT_e60_lhmedium_nod0 OR HLT_e140_lhloose_nod0 HLT_e300_etcut	HLT_mu26_ivarmedium OR HLT_mu50	HLT_xe90_mht_L1XE50
2016b (run ≥ 302919) $(L < 1.7 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1})$	same as above	same as above	HLT_xe110_mht_L1XE50
2017	same as above	same as above	HLT_xe110_pufit_L1XE55
2018	same as above	same as above	HLT_xe110_pufit_xe70_L1XE50

933 **9.3 non-VBF/VBF RNN**

934 To classify events as originating from non-VBF or VBF production a recursive
935 neural network (RNN [22]) is used. This approach is more powerful than a cut-
936 based classification as it improves signal efficiency and analysis sensitivity by
937 exploiting correlations between variables that the RNN learns. In particular, a
938 RNN architecture is ideal as it can handle variable numbers of jets in the events.

939 The RNN uses the four-momentum of candidate VBF jets to classify events
940 as VBF or non-VBF topologies. Sometimes jets are incorrectly reconstructed,
941 so the number of jets in the event is expected to vary across the input samples.
942 VBF candidate jets are identified by removing jets from the event that are likely
943 from $W/Z \rightarrow qq$. For the resolved regime this means removing the two leading
944 small-R jets from the VBF candidate jet list. For the merged regime this means
945 removing small-R jets separated by less than 1.0 in dR from the large-R jet. VBF
946 candidate jets are also required to be within $|\eta| < 4.5$. From the list of remaining
947 VBF candidate jets, the two highest- p_T jets are chosen.

948 The architecture of the RNN is shown in Figure 9.2. The RNN is composed
949 of Long Short Term Memory Cells (LSTM) that extract meaningful information
950 and retains³ it. The logic embedded in the LSTM is shown in Figure ???. LSTMs
951 are useful for VBF event classification for events with two jets, where using the
952 kinematic properties of both jets (and their correlations) will lead to more efficient
953 event classification.

954 In this RNN architecture, the VBF candidates are first passed to a masking
955 layer which checks the number of jets in the event. If there is only one jet, only one
956 vertical LSTM layer is used. The output of masking is then passed to a LSTM,
957 with a tanh activation function. The output of the LSTM is then passed to a
958 second horizontal LSTM layer (and vertical LSTM layer if there are two jets in

959 the event). Finally the output of the last LSTM cell is passed to a dense layer
960 and then to a sigmoid activation layer, leading to an overall RNN score.

961 The weights and other parameters of the network are learned by training the
962 network with VBF and non-VBF signals over 200 epochs with an Adam Optimizer
963 [14]. To prevent overfitting during training, dropout is applied to RNN weights
964 and training is truncated if the network parameters are unchanged after ten it-
965 erations [24]. Figure 9.4 shows the ROC curve for the RNN using k-fold cross
966 validation [20].

967 Figure 9.3 shows the RNN discriminant for backgrounds, non-VBF signals,
968 and VBF signals. The RNN score is ~ 0 for non-VBF signals and background
969 processes and ~ 1 for VBF processes. Figure 9.5 shows the limits for various signal
970 processes based on the RNN cut applied. Requiring the RNN score to be > 0.8
971 was chosen as it provided the best significance (and signal efficiency) for this final
972 state and the $\nu\nu qq$ and $\ell\ell qq$ channels, which this channel will be combined with
973 for future publications.

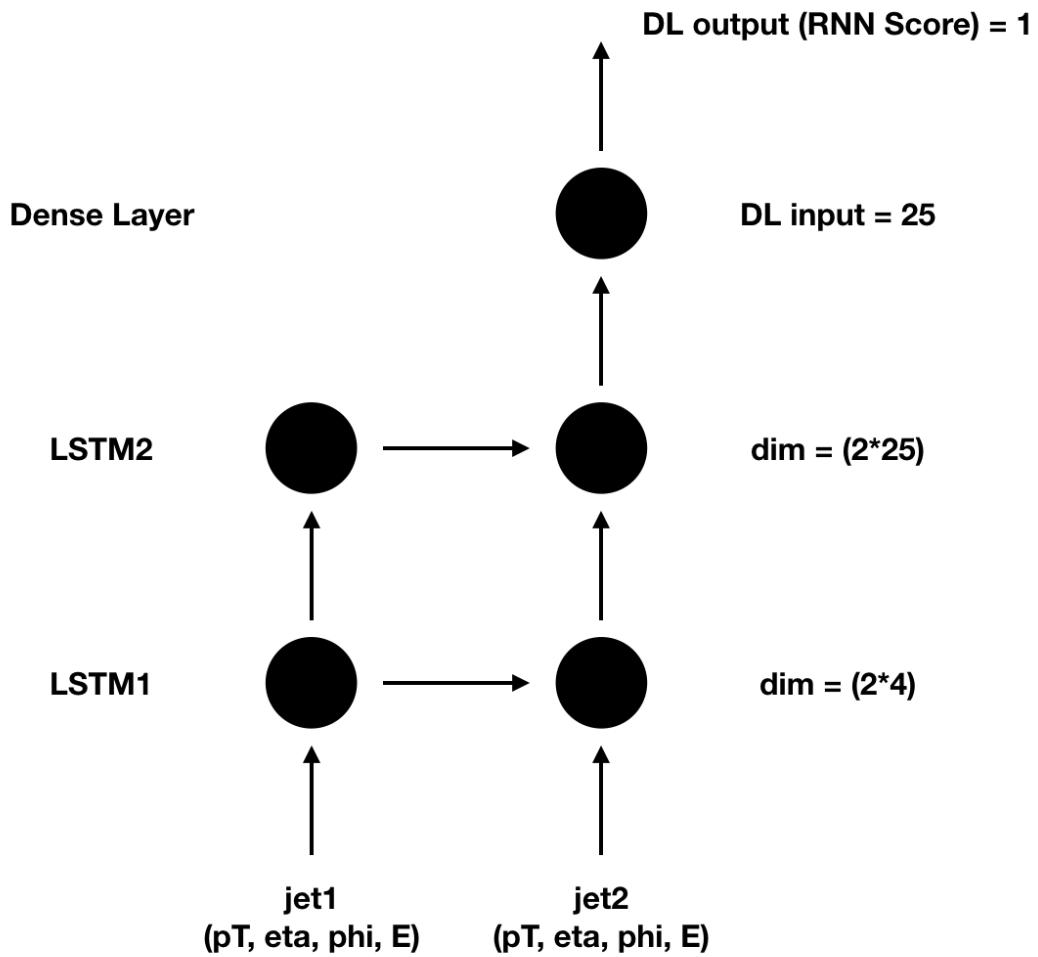


Figure 9.1: RNN architecture. Natasha add caption

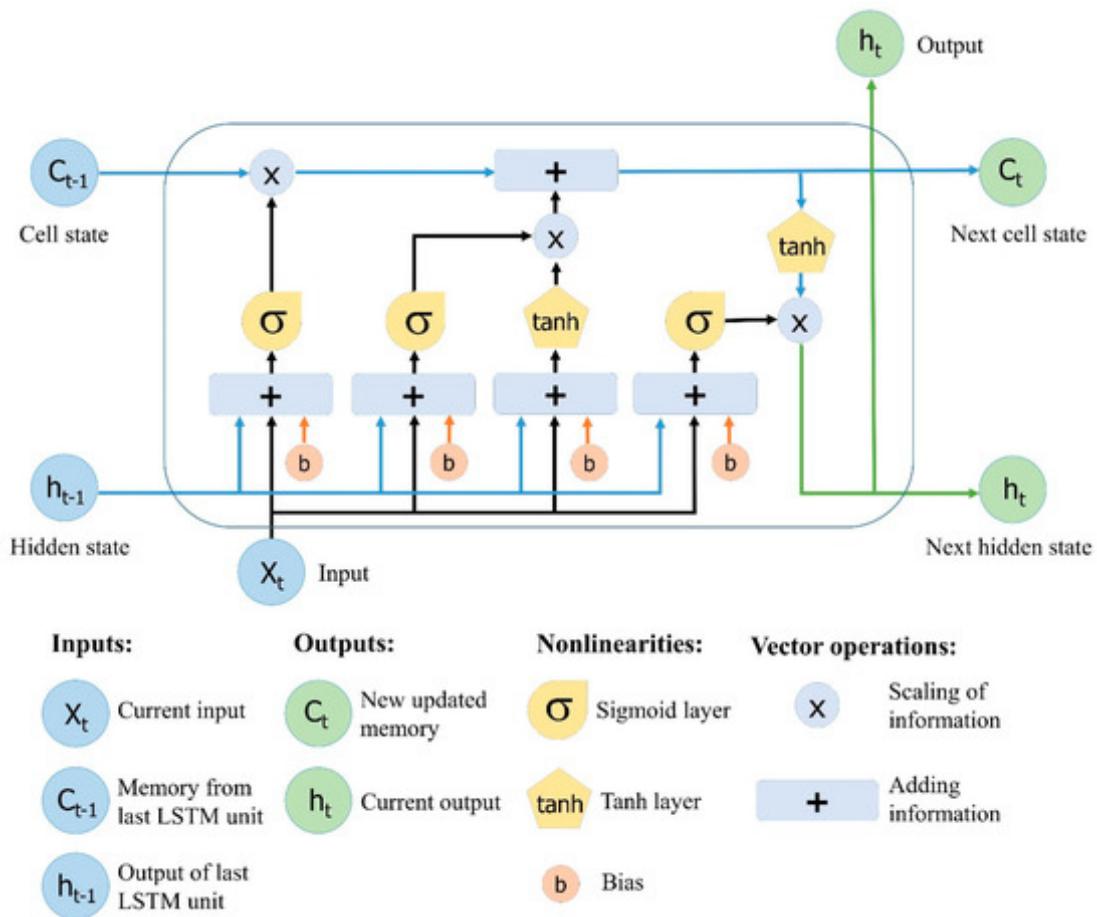


Figure 9.2: This figure shows the embedded logic in LSTM cells. This image was taken from [23], where a more in depth discussion about LSTMs may be found.

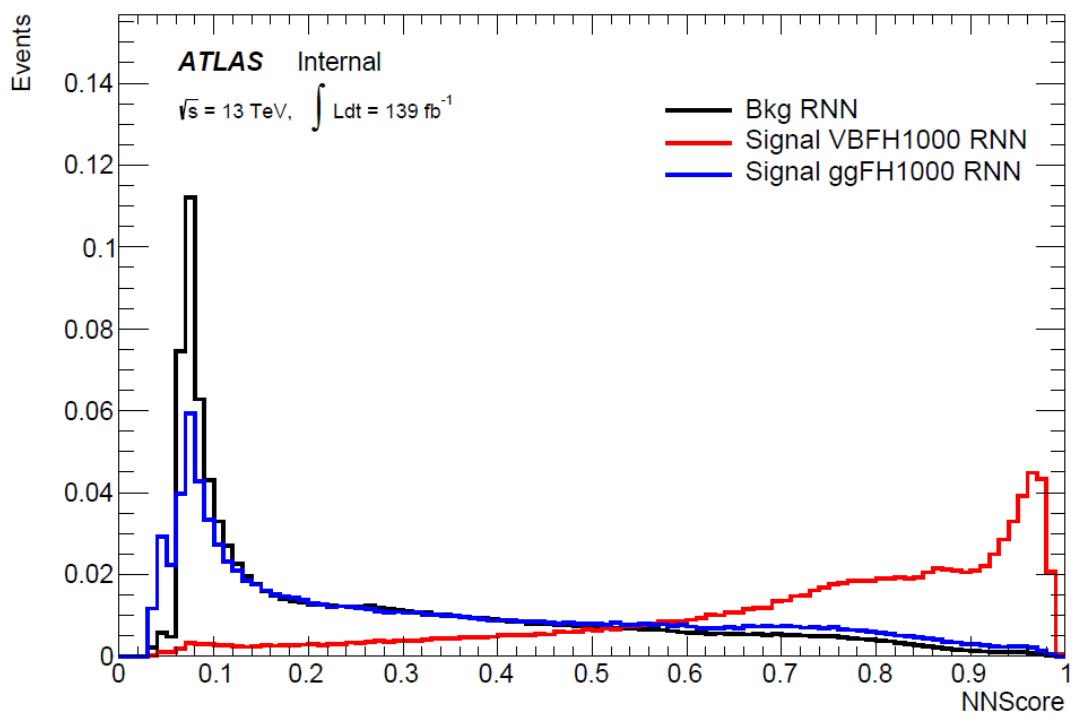


Figure 9.3: RNN Score distribution for ggF and VBF signals and backgrounds.

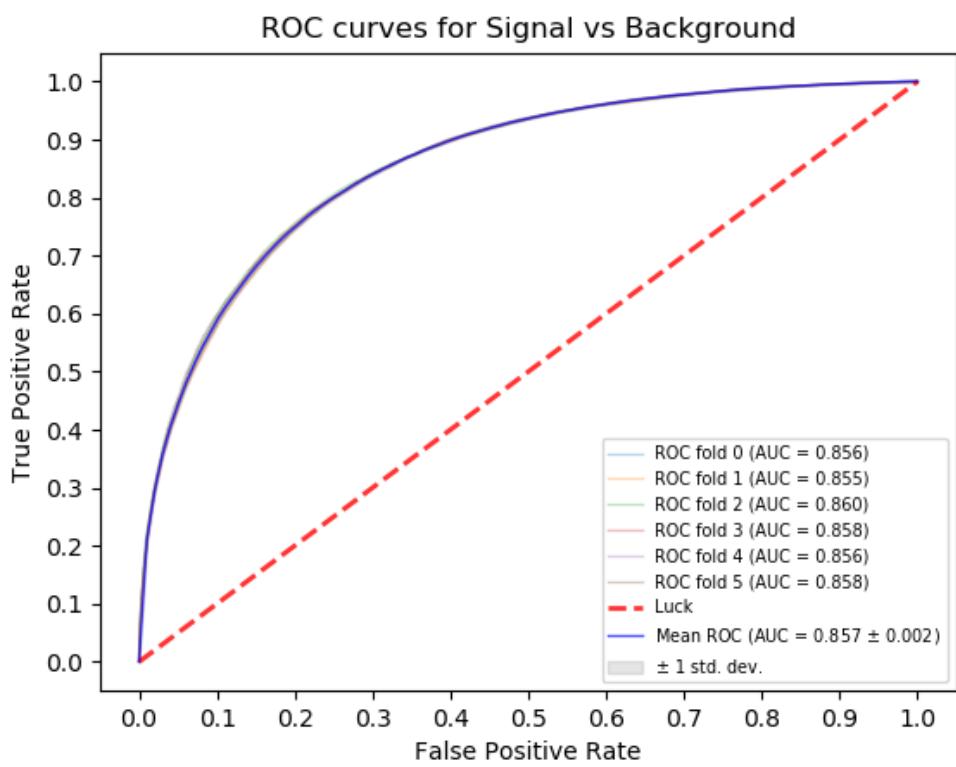


Figure 9.4: ROC curve using k-fold validation for RNN.

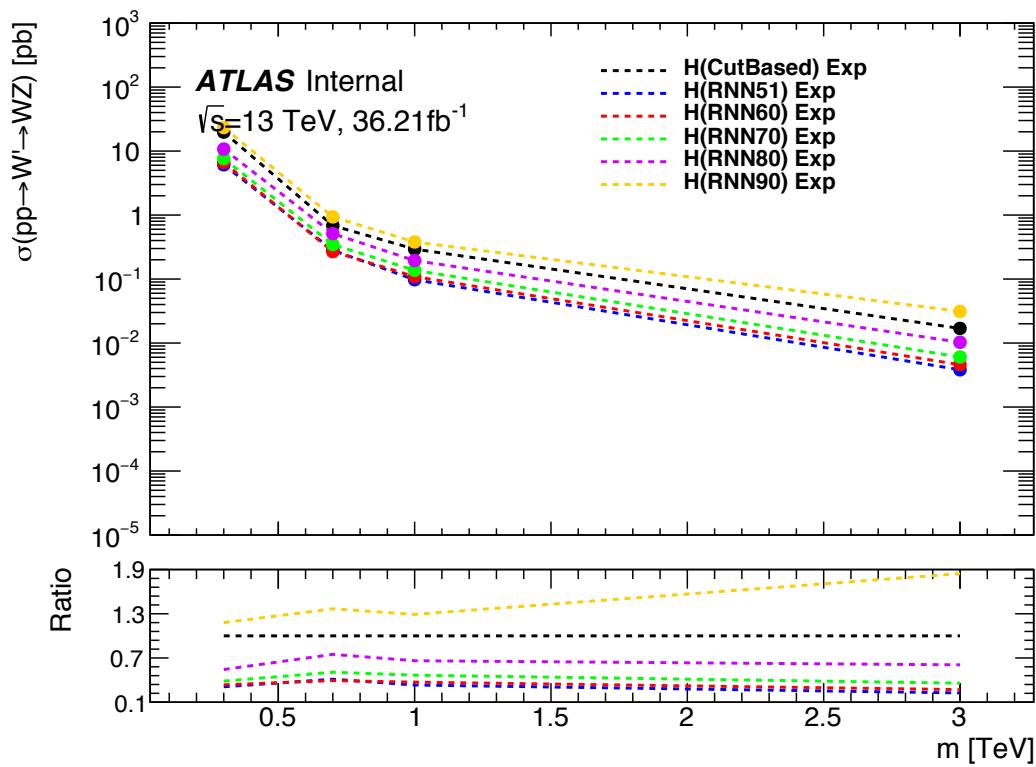


Figure 9.5: Comparison of GGF Z' limits for different RNN score selections. The bottom panel shows the ratio of the upper limits set for different RNN cuts to the cut-based analysis. In this panel smaller numbers, indicate that the expected upper limit is smaller than the cut-based analysis, which is desired.

974 9.4 Signal Region Definitions

975 Signal regions are constructed to be dominated by signal and used in the final
976 likelihood fit to look for a bump in the reconstructed resonance mass distribution.
977 Once an event is classified by the RNN, it must pass topological cuts that maximize
978 S/\sqrt{B} . To efficiently select events with a $W \rightarrow \ell\nu$ candidate exactly one tight
979 lepton is required and $E_T^{miss} > 100(60)$ GeV and $p_{T,\ell\nu} > 200(75)$ GeV in the
980 merged (resolved) analysis to suppress the fake lepton backgrounds.

981 For the merged analysis, in addition to the $W \rightarrow \ell\nu$ and $W/Z \rightarrow J$ selections
982 above, the $\min(p_{T,\ell\nu}, p_{T,J})/m_{WV} > 0.35(0.25)$ for the non-VBF (VBF) category.
983 To reduce $t\bar{t}$ contamination, events with at least one b jet with $\Delta R > 1.0$ from
984 the large-R jet are excluded. High purity signal regions require the D_2 and W/Z
985 mass window cut to be passed, whereas the low purity region only requires the
986 W/Z mass window cut to be passed.

987 For the HVT W' resonance search tagged and untagged regions are used to
988 increase signal acceptance. For events to be classified as tagged the large-R jet
989 must contain exactly two b-tagged jets. Untagged events must have no more than
990 one b-tagged jet matched to the large-R jet. These selections are shown in Table
991 9.2.

992 Events failing the merged selection are then re-analyzed in the resolved cat-
993 egory. To enhance resolved signals, the event should contain two high- p_T boson
994 candidates that are back-to-back in the ϕ as shown by the cuts in Table 9.3.
995 Again, to suppress the $t\bar{t}$ backgrounds, events are required to have no additional
996 b-jets. A summary of the resolved selections is shown in Table 9.3.

997 The analysis cutflow in Figure 9.6 which shows how the different categories
998 are prioritized. Events classified as VBF events are classified as Merged high
999 purity, low purity or resolved signal region selections sequentially. If the event

Table 9.2: Summary of selection criteria used to define the signal region (SR), W +jets control region (W CR) and $t\bar{t}$ control region ($t\bar{t}$ CR) for merged 1-lepton channel.

Selection		SR		W CR (WR)		$t\bar{t}$ CR (TR1)	
		HP	LP	HP	LP	HP	LP
$W \rightarrow \ell\nu$	Num of Tight leptons	1					
	Num of Loose leptons	0					
	E_T^{miss}	$> 100 \text{ GeV}$					
	$p_T(\ell\nu)$	$> 200 \text{ GeV}$					
$W/Z \rightarrow J$	Num of large- R jets	≥ 1					
	D_2 cut	pass	fail	pass	fail	pass	fail
	W/Z mass window cut	pass	pass	fail	fail	pass	pass
	Numb. of associated VR track jets b -tagged	For $Z \rightarrow J$: ≤ 1 ($= 2$) for untagged (tagged) category					
	Topology cut	$\min(p_{T,\ell\nu}, p_{T,J}) / m_{WV} > 0.35(0.25)$ for DY/ggF (VBF) category					
	Top-quark veto	Num of b -tagged jets outside of large- R jet	0	≥ 1			
Pass VBF selection			no (yes) for DY/ggF (VBF) category				

1000 does not pass any of these selections but passes a VBF control region selection it
 1001 is classified as a VBF CR event. If the event fails all VBF categories, it is then
 1002 checked if it passes the Merged high purity, low purity or resolved signal region
 1003 selections (NB: for the WZ decay modes all the regions have tagged and untagged
 1004 categories). If the event fails all the non-VBF signal region selections, it is then
 1005 kept for non-VBF control region selections, if it passes those selections.

Table 9.3: The list of selection cuts in the resolved analysis for the WW and WZ signal regions (SR), $W+jets$ control region (WR) and $t\bar{t}$ control region (TR).

cuts	SR	W CR (WR)	$t\bar{t}$ CR (TR1)		
$W \rightarrow \ell\nu$	Number of Tight leptons	1			
	Number of Loose leptons	0			
	E_T^{miss}	> 60 GeV			
	$\cancel{p}_T(\ell\nu)$	> 75 GeV			
$W/Z \rightarrow jj$	Number of small-R jets	≥ 2			
	Leading jet p_T	> 60 GeV			
	Subleading jet p_T	> 45 GeV			
	$Z \rightarrow q\bar{q}$	$78 < m_{jj} < 105$ GeV	$50 < m_{jj} < 68$ GeV or $68 < m_{jj} < 98$ GeV	$105 < m_{jj} < 150$ GeV	$50 < m_{jj} < 150$ GeV
	Num. of b -tagged jets	For $Z \rightarrow jj$: ≤ 1 ($= 2$) for untagged (tagged) category			
Topology cuts	$\Delta\phi(j, \ell)$	> 1.0			
	$\Delta\phi(j, E_T^{\text{miss}})$	> 1.0			
	$\Delta\phi(j, j)$	< 1.5			
	$\Delta\phi(\ell, E_T^{\text{miss}})$	< 1.5			
	$\min(p_{T,\ell\nu}, p_{T,jj}) / m_{WW}$	> 0.35(0.25) for DY/ggF (VBF) category			
Top veto	Number of additional b -tagged jets	0	≥ 1		
	Pass VBF selection	no (yes) for DY/ggF (VBF) category			

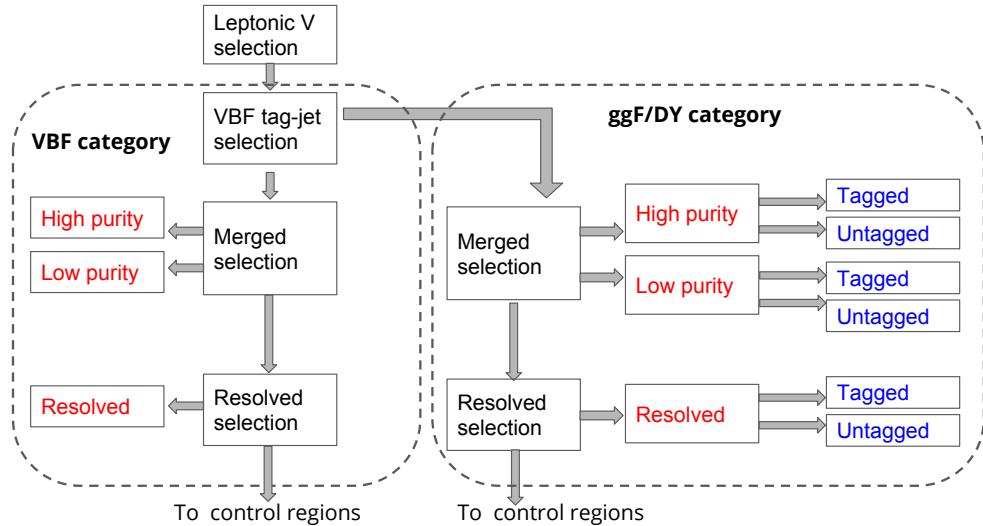


Figure 9.6: Event Categorization. Natasha write more.

1006 9.5 Selection Acceptance and Efficiency

1007 The acceptance times efficiency for the signal region selection is defined as:

$$A \cdot \epsilon = \frac{N_{\text{events selected}}^{\text{truth}}}{N_{\text{events generated}}^{\text{truth}}} \cdot \frac{N_{\text{events selected}}^{\text{reco}}}{N_{\text{events selected}}^{\text{truth}}} = \frac{N_{\text{events selected}}^{\text{reco}}}{N_{\text{events generated}}^{\text{truth}}} \quad (9.1)$$

1008 The distributions of $A \cdot \epsilon$ as a function of the resonance mass for the different spin
1009 models are shown in Figures 9.8 - ??.

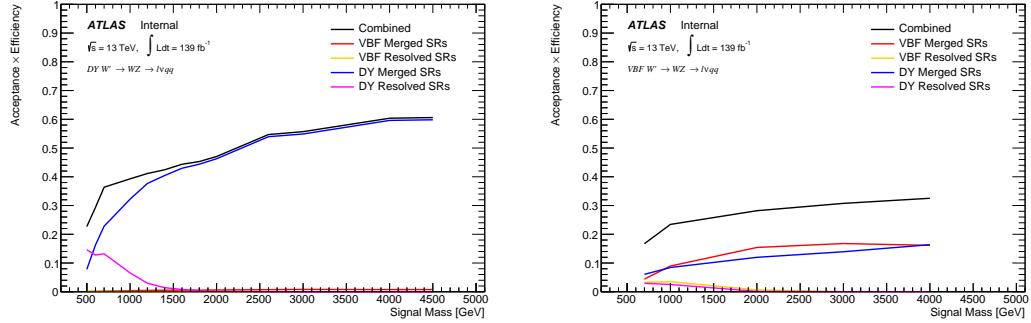


Figure 9.7: Selection acceptance times efficiency for the $W' \rightarrow WZ \rightarrow \ell\nu qq$ events from MC simulations as a function of the W' mass for (a) Drell-Yan and (b) VBF production, combining the merged HP and LP signal regions of the $WW \rightarrow \ell\nu J$ selection and the resolved regions of the $WW \rightarrow \ell\nu jj$ selection.

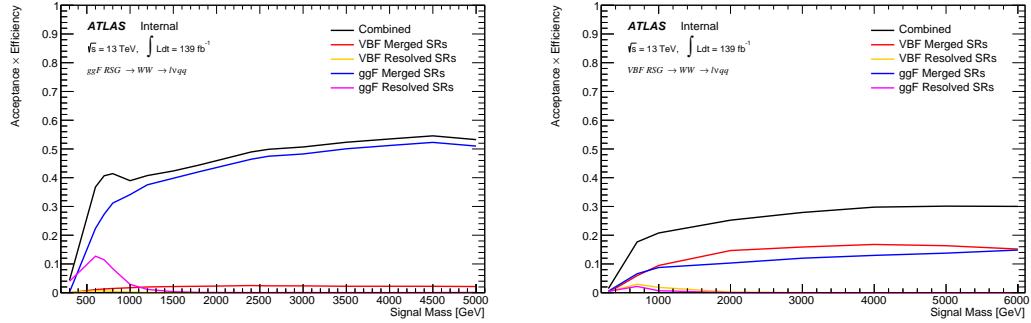


Figure 9.8: Selection acceptance times efficiency for the $G \rightarrow WW \rightarrow \ell\nu qq$ events from MC simulations as a function of the G mass for (a) Drell-Yan and (b) VBF production, combining the merged HP and LP signal regions of the $WW \rightarrow \ell\nu J$ selection and the resolved regions of the $WW \rightarrow \ell\nu jj$ selection.

1010 **9.6 Background Estimate**

1011 Backgrounds from VV , $t\bar{t}$, single-top, $W+\text{jets}$, $Z+\text{jets}$ are simulated as de-
1012 scribed in 7.2. The dominant backgrounds for this search are from $W+\text{jet}$ and $t\bar{t}$
1013 processes. To more accurately model the distribution of m_{VV} from these back-
1014 grounds in the SRs, control regions are constructed for each as described in 9.6.1.
1015 The $t\bar{t}$ and $W+\text{jets}$ control regions are called TCR and WCR, respectively. There
1016 are separate control regions for VBF and non-VBF regions as well as for each
1017 region (merged HP, merged LP, resolved). For the HVT W' search there are also
1018 tagged and untagged control regions (where tagged refers to events with two b -jets
1019 inside the large-R jet).

1020 The aforementioned backgrounds containing real leptons and are well-modeled
1021 with simulated samples. Backgrounds with fake leptons (also referred to as the
1022 multijet background) are not well-modeled with simulation. For this reason, the
1023 multijet background is extracted from data as described in 9.6.2.

1024 **9.6.1 Control Regions**

1025 The distributions for the variables used in merged analysis for top control
1026 regions are shown in Figure 9.9-9.12.

1027 The distributions for the variables used in the resolved analysis in the TCR
1028 are shown in Figure 9.13, 9.14.

1029 To more accurately model the two dominant backgrounds in this analysis,
1030 $W+\text{jets}$ and $t\bar{t}$, control regions are constructed for each. These control regions
1031 are dominated by these processes and used to extract normalization factors in
1032 the final likelihood fit that are then used in the signal region estimates. For the
1033 $t\bar{t}$ control region the event must contain at least one such b jet. The WCR is
1034 constructed using the $m_{jj/J}$ mass window sidebands. All other backgrounds are

₁₀₃₅ estimated using simulation, except fake lepton backgrounds, which are derived
₁₀₃₆ using a data-driven method.

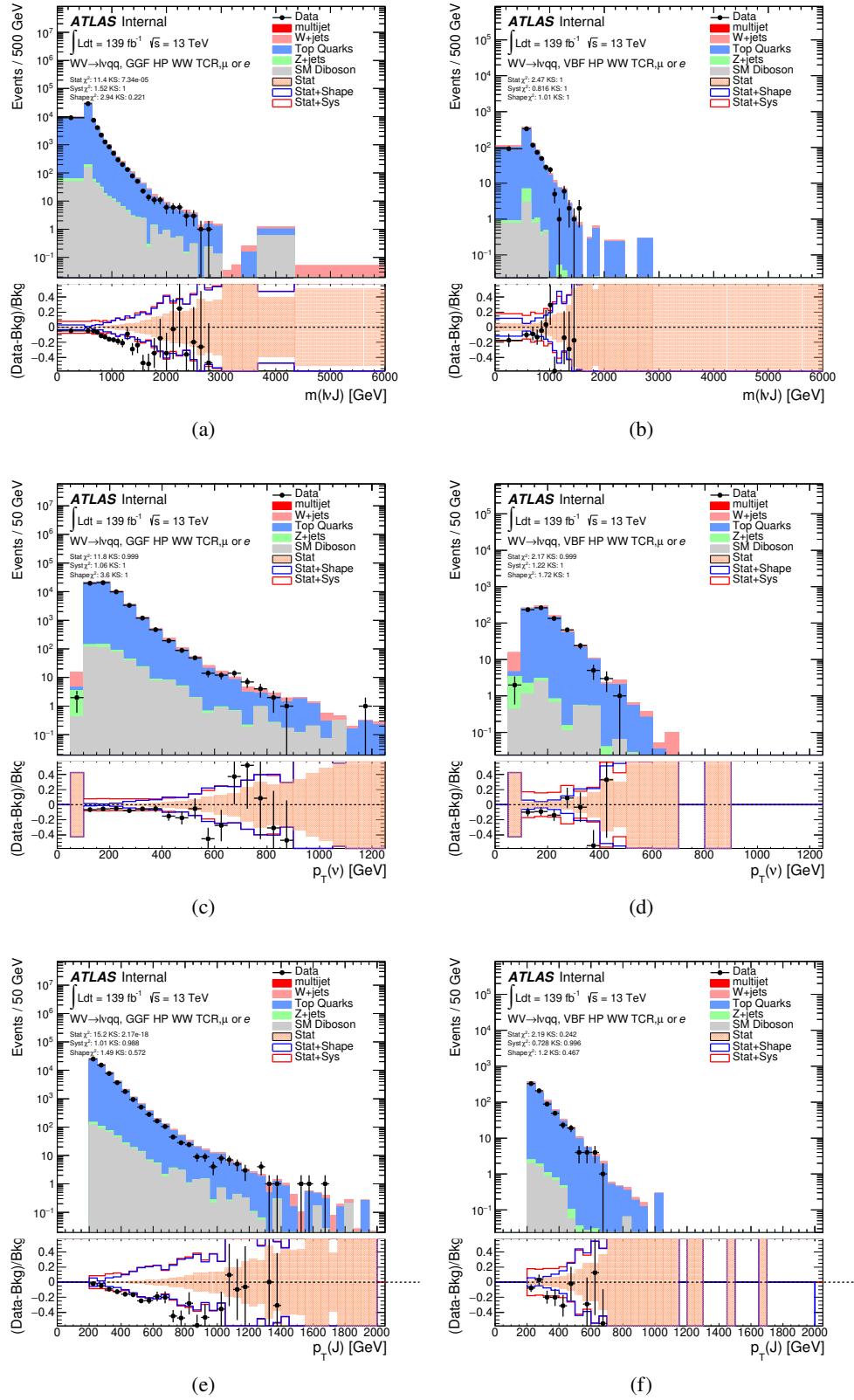


Figure 9.9: Data MC comparison for the merged WW HP TCR. The bottom panel shows the ratio of the difference between data and simulation to simulation. The red bands include the all systematic and statistical uncertainties on the background.

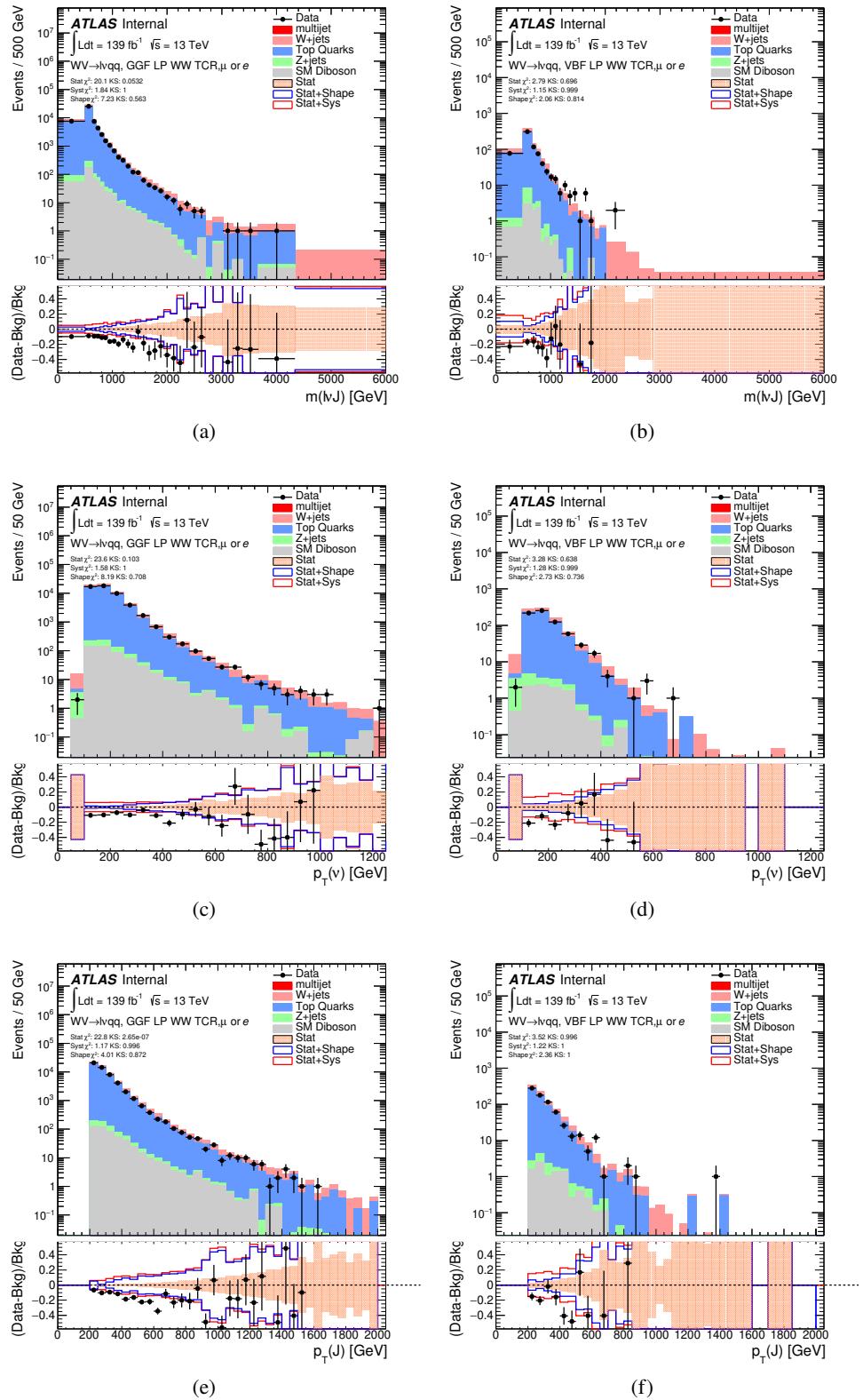


Figure 9.10: Data MC comparison for the merged WW LP TCR. The bottom panel shows the ratio of the difference between data and simulation to simulation. The red bands include the all systematic and statistical uncertainties on the background.

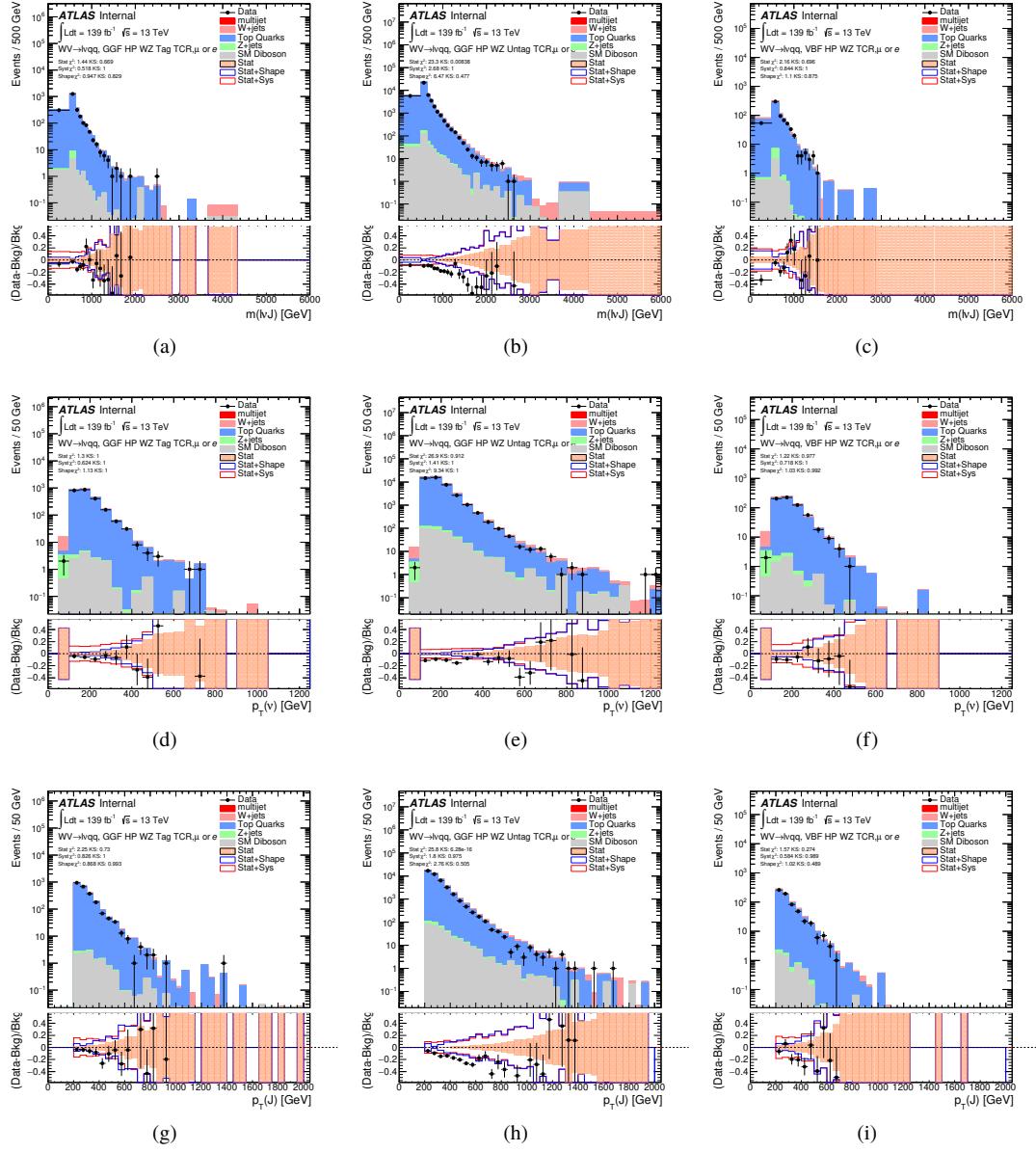


Figure 9.11: Data MC comparison for the merged WZ HP TCR. The bottom panel shows the ratio of the difference between data and simulation to simulation. The red bands include the all systematic and statistical uncertainties on the background.

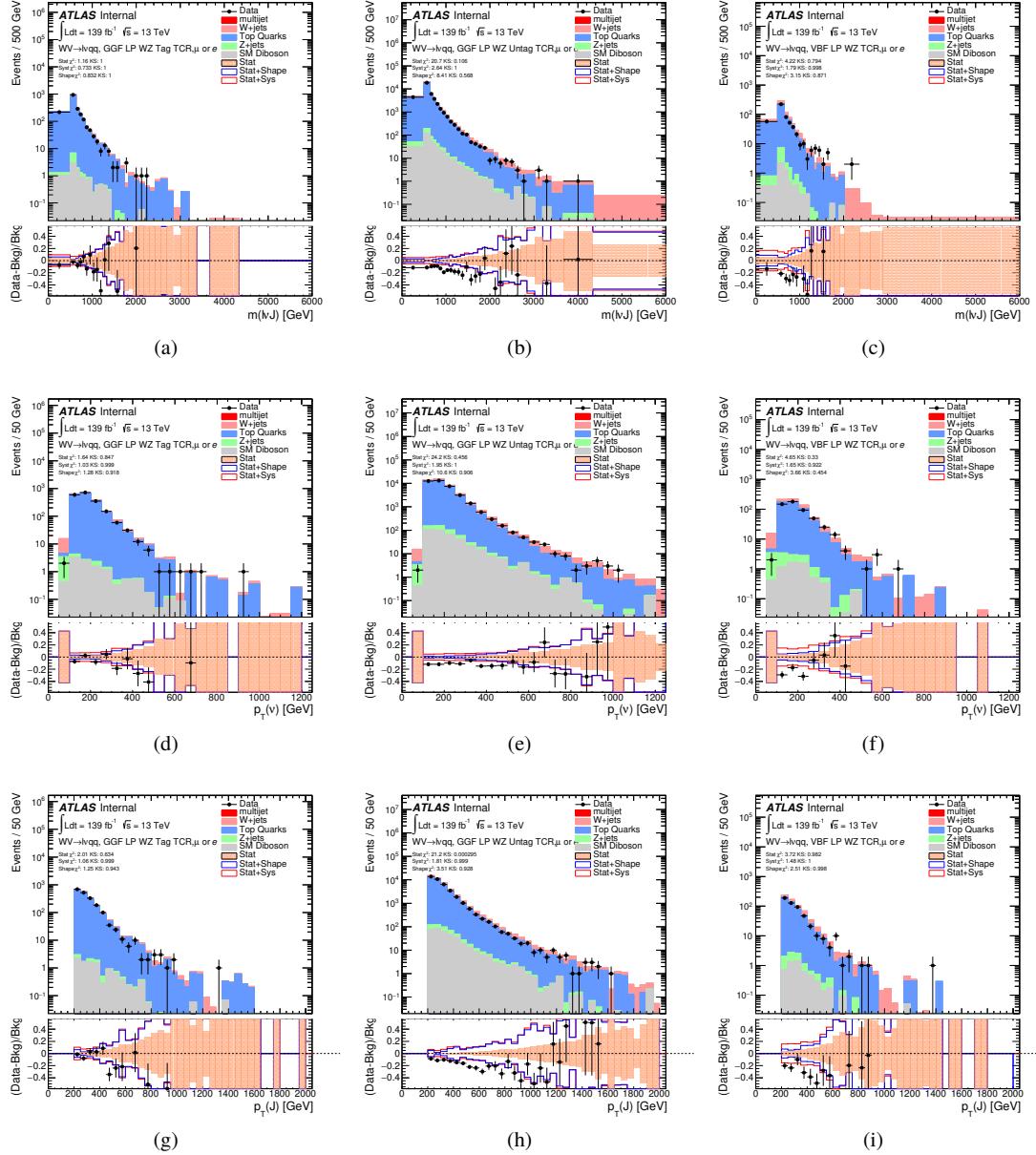


Figure 9.12: Data MC comparison for the merged WZ LP TCR. The bottom panel shows the ratio of the difference between data and simulation to simulation. The red bands include the all systematic and statistical uncertainties on the background.

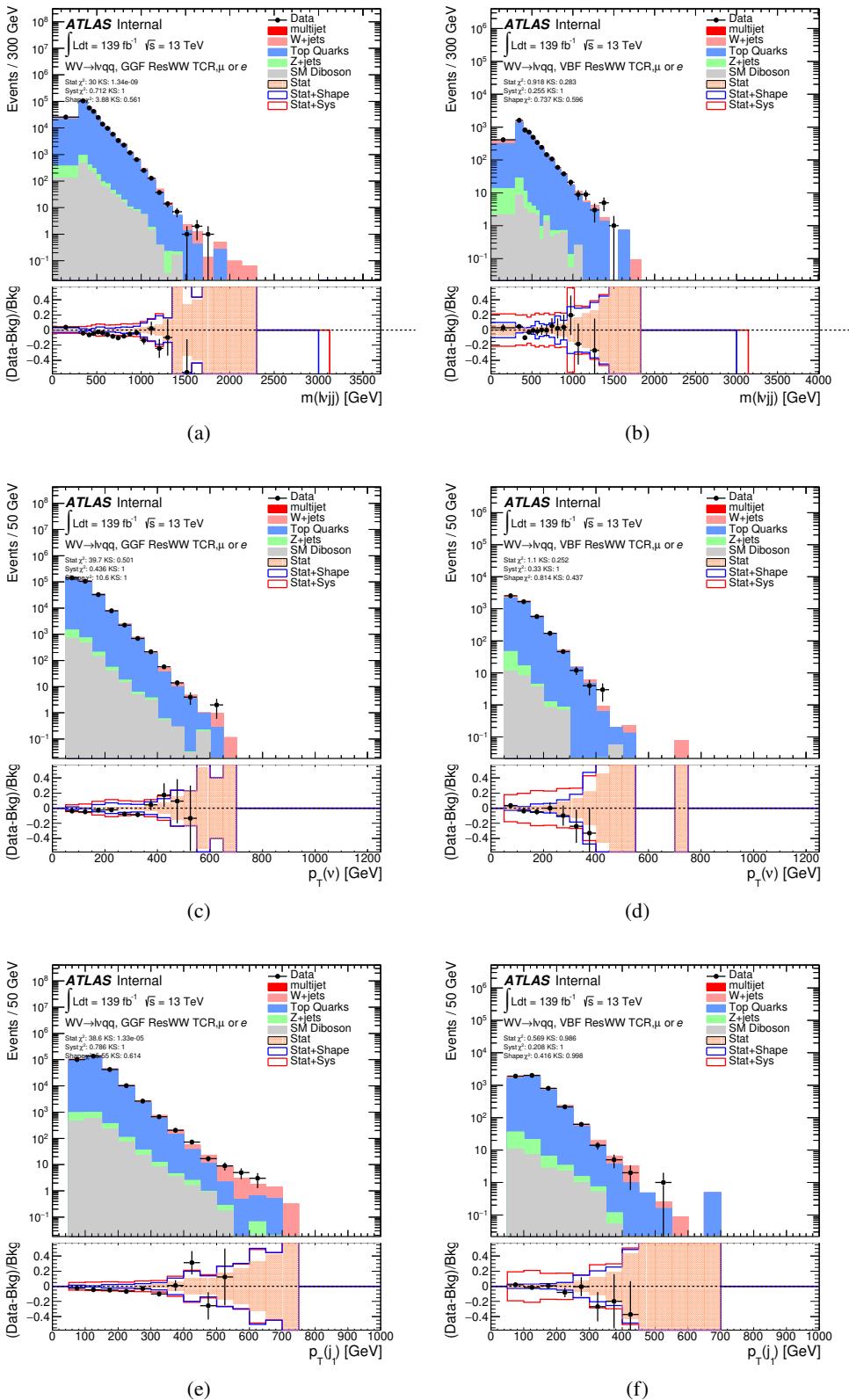


Figure 9.13: Data MC comparison for the resolved WW TCR. The bottom panel shows the ratio of the difference between data and simulation to simulation. The red bands include the all systematic and statistical uncertainties on the background.⁹⁵

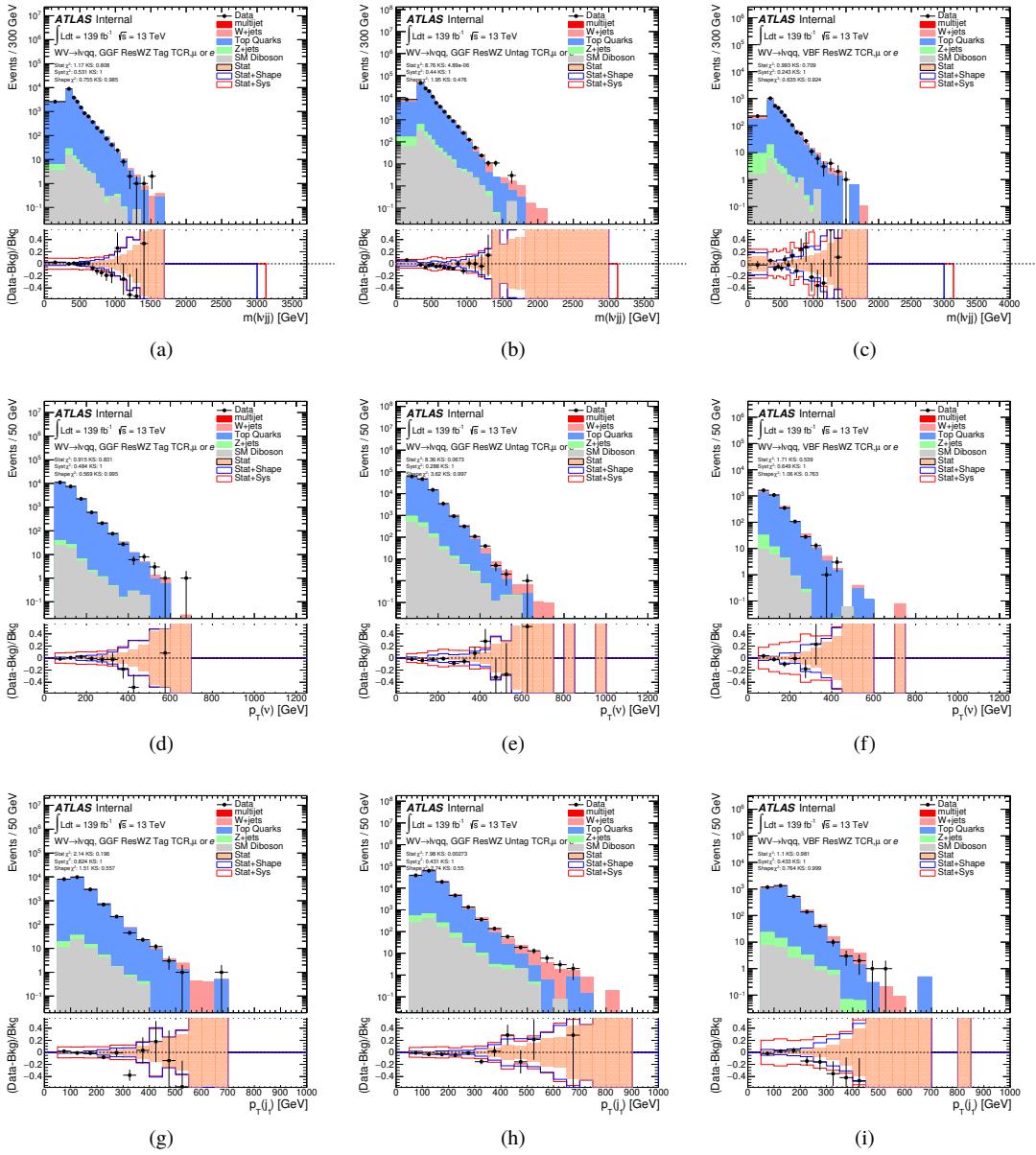


Figure 9.14: Data MC comparison for the resolved WZ TCR. The bottom panel shows the ratio of the difference between data and simulation to simulation. The red bands include the all systematic and statistical uncertainties on the background.

1037 **9.6.2 Fake Lepton Backgrounds**

1038 The fake lepton backgrounds for this search are not well-modeled with simu-
1039 lation. For this reason, this background is extracted from data. Fake electrons
1040 often arise from fake jets and converted photons while non-prompt muons usually
1041 arise from heavy flavor decay products.

1042 Fake electrons generally fail the electron ID criteria and fake muons fail the
1043 muon isolation requirement. Therefore, separate multijet samples are derived for
1044 the fake electron and muon samples. For each sample the m_{WV} template shape
1045 is derived for the SR and WCR selections using the same SR and WCR cuts but
1046 with inverted lepton requirements as seen in Table 9.4. NB: By inverting the
1047 lepton isolation/identification criteria the SRs and CRs are orthogonal.

1048 To derive the multijet template in a given SR, first the multijet template in
1049 the WCR is derived, called the MJCR template. This template is calculated using
1050 events that pass the WCR selection but with the inverted lepton criteria. The
1051 E_T^{miss} distribution for the MJCR is given by the difference between data and the
1052 simulated samples in the MJCR. The E_T^{miss} distribution of those events is then
1053 added to the simulated backgrounds in the WCR. The floating background and
1054 multijet normalizations of the MJCR in this region are then fit to the data. The
1055 fitted MJCR is then used as the multijet sample in the WCR.

1056 The fitted normalizations from the MJCR template are then used to construct
1057 the multijet template in the SR (MJSR). The MJSR is constructed from events
1058 that pass the SR selections but with the inverted lepton criteria. Again, the
1059 difference between the data and simulated backgrounds in this region gives MJSR
1060 template shape in m_{WV} . This shape is then scaled by the fitted normalizations
1061 from the MJCR. These fitted electron and muon multijet templates are then
1062 used as the multijet samples in the SRs. The normalizations of the electron and

1063 muon multijet samples are parameters in the final likelihood fit.

1064 This template method was validated using WCR and full Run 2 data. The
1065 results of the fit are shown in Table 9.5. The multijet contribution in the muon
1066 channel for $p_T^W > 150$ GeV is consistent with zero, and therefore neglected in
1067 the final fit. Applying the extracted normalization factor to MJCR in WCRs for
1068 various kinematic variables such as E_T^{miss} , W transverse mass, lepton p_T , and the
1069 invariant mass as show in Figures 9.16 -9.25. These figures show good agreement
1070 between the data and background estimate.

Table 9.4: Definitions of “inverted” leptons used in multijet control region

	Criterion	signal lepton	inverted lepton
Electron	ID	TightLH	MediumLH !TightLH
	Calo Isolation	FixedCutHighPtCaloOnlyIso	FixedCutHighPtCaloOnlyIso
Muon	ID	WHSignalMuon	WHSignalMuon
	Track Isolation	FixedCutTightTrackOnlyIso	!FixedCutTightTrackOnlyIso $ptvarcone30/pt < 0.07^*$

*Only applied to events with $pTW < 150\text{GeV}$

1071

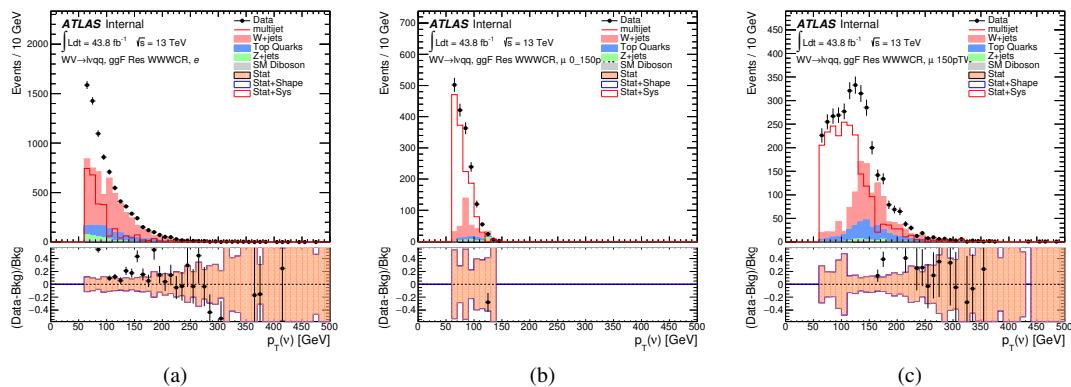


Figure 9.15: The E_T^{miss} distribution in MJCR for 2017 data in the electron channel(left), muon channel with W-boson pT < 150 GeV (center) and > 150 GeV (right). Multi-jet templates are calculated as remaining data components after excluding known MC

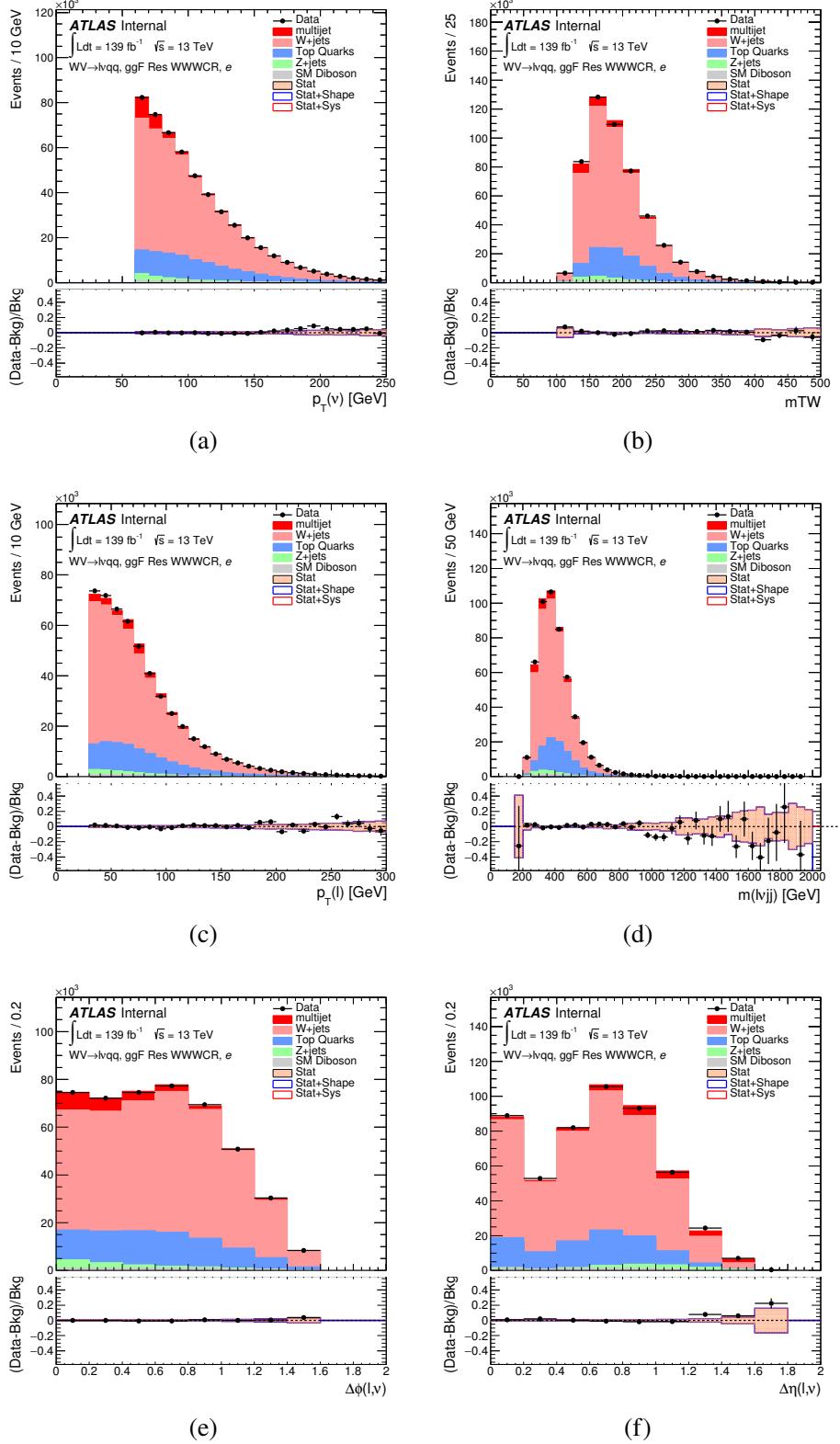


Figure 9.16: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the WW electron channel. The MJ template is obtained from the pre-MJ-fit.

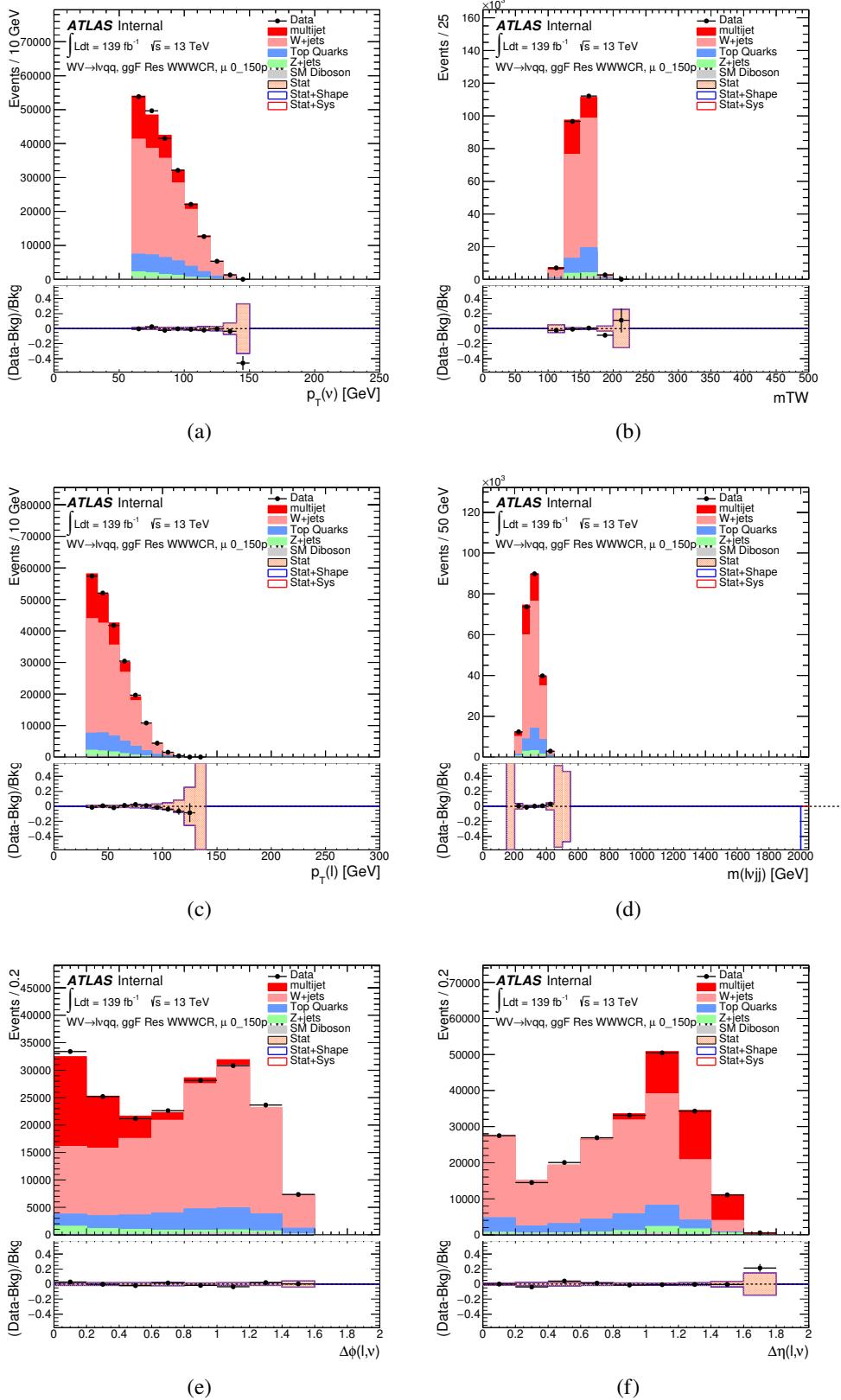


Figure 9.17: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the WW muon channel. The MJ template is obtained from the pre-MJ-fit.

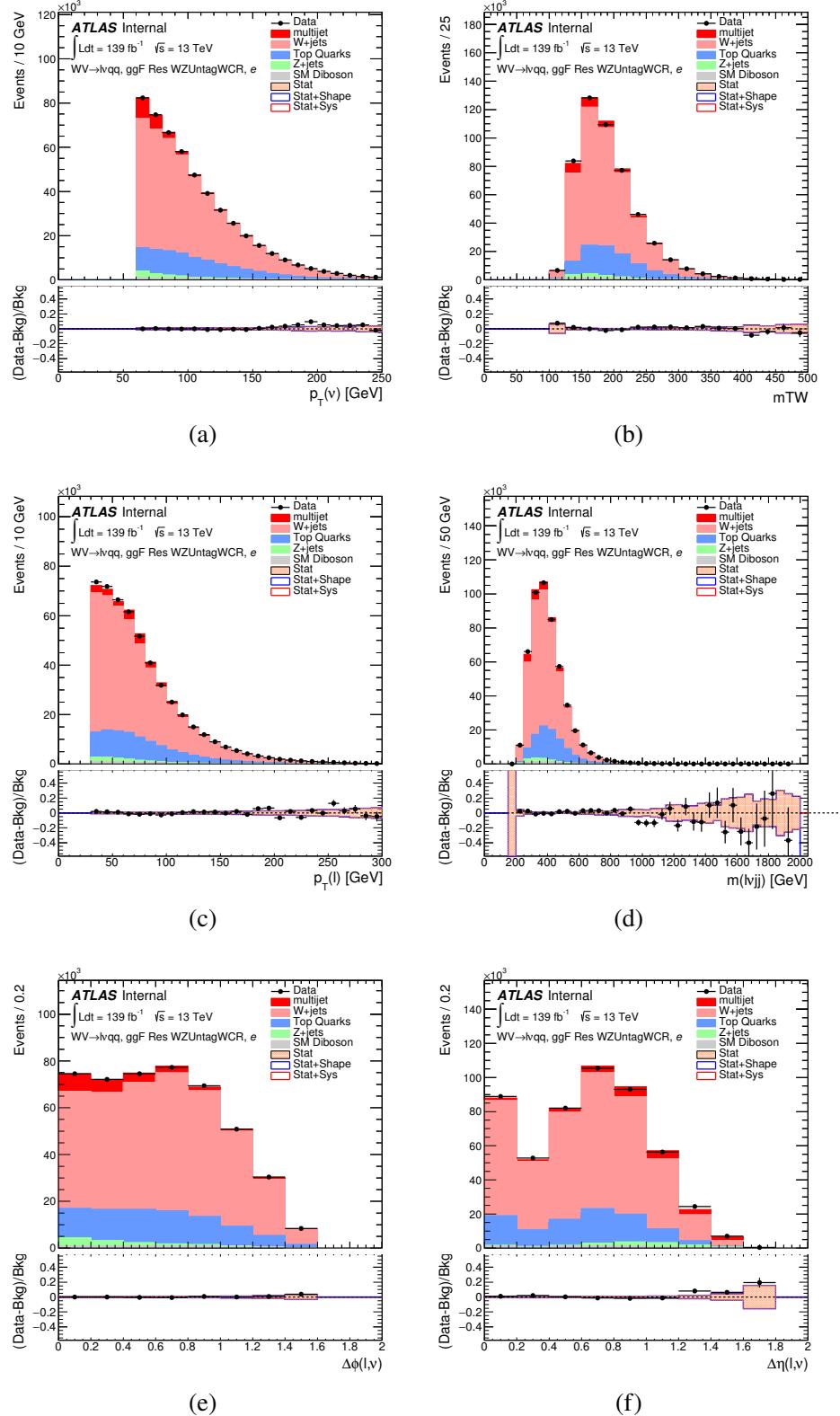


Figure 9.18: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the WZ untag electron channel. The MJ template is obtained from the pre-MJ-fit.

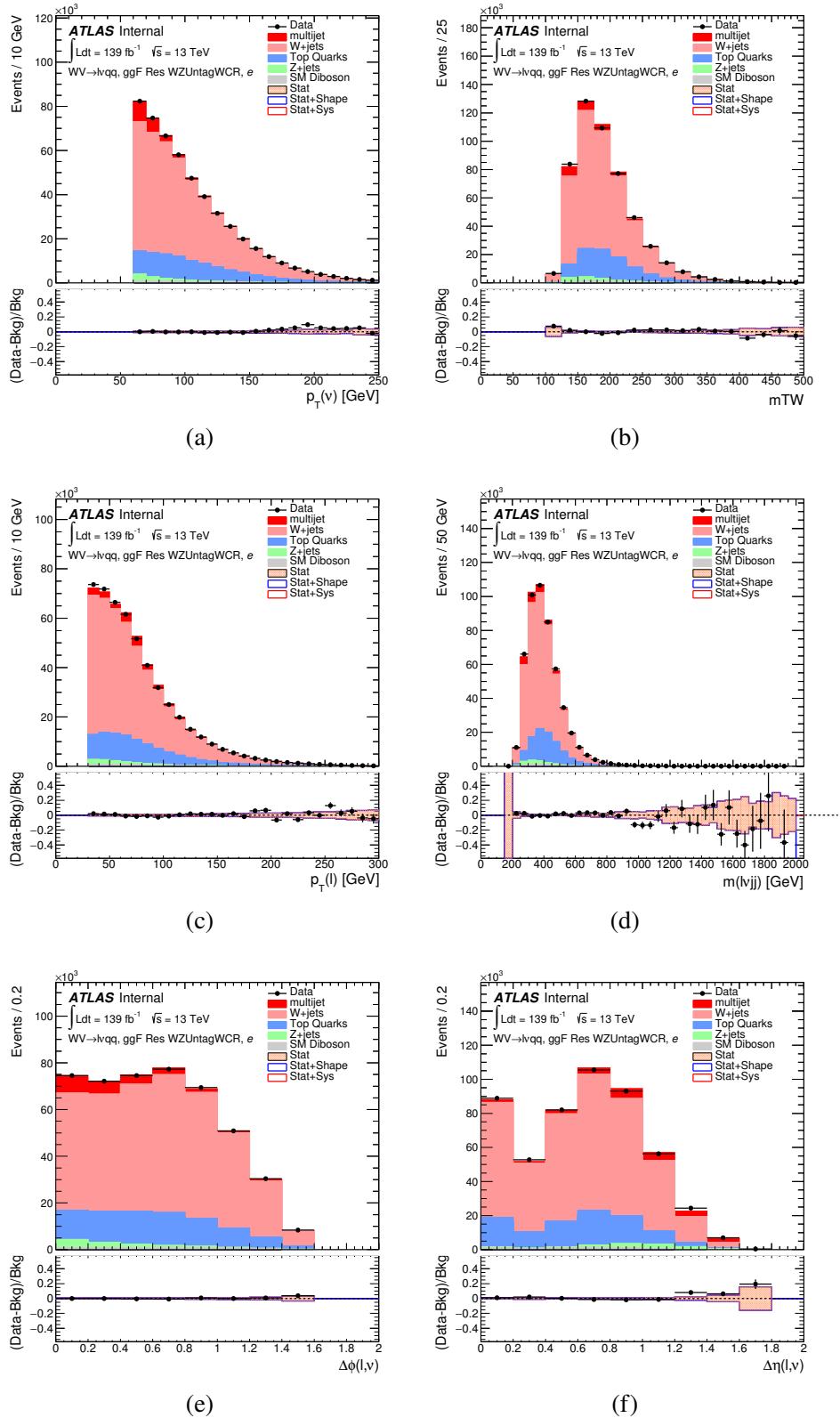


Figure 9.19: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the WZ untag muon channel. The MJ template is obtained from the pre-MJ-fit.

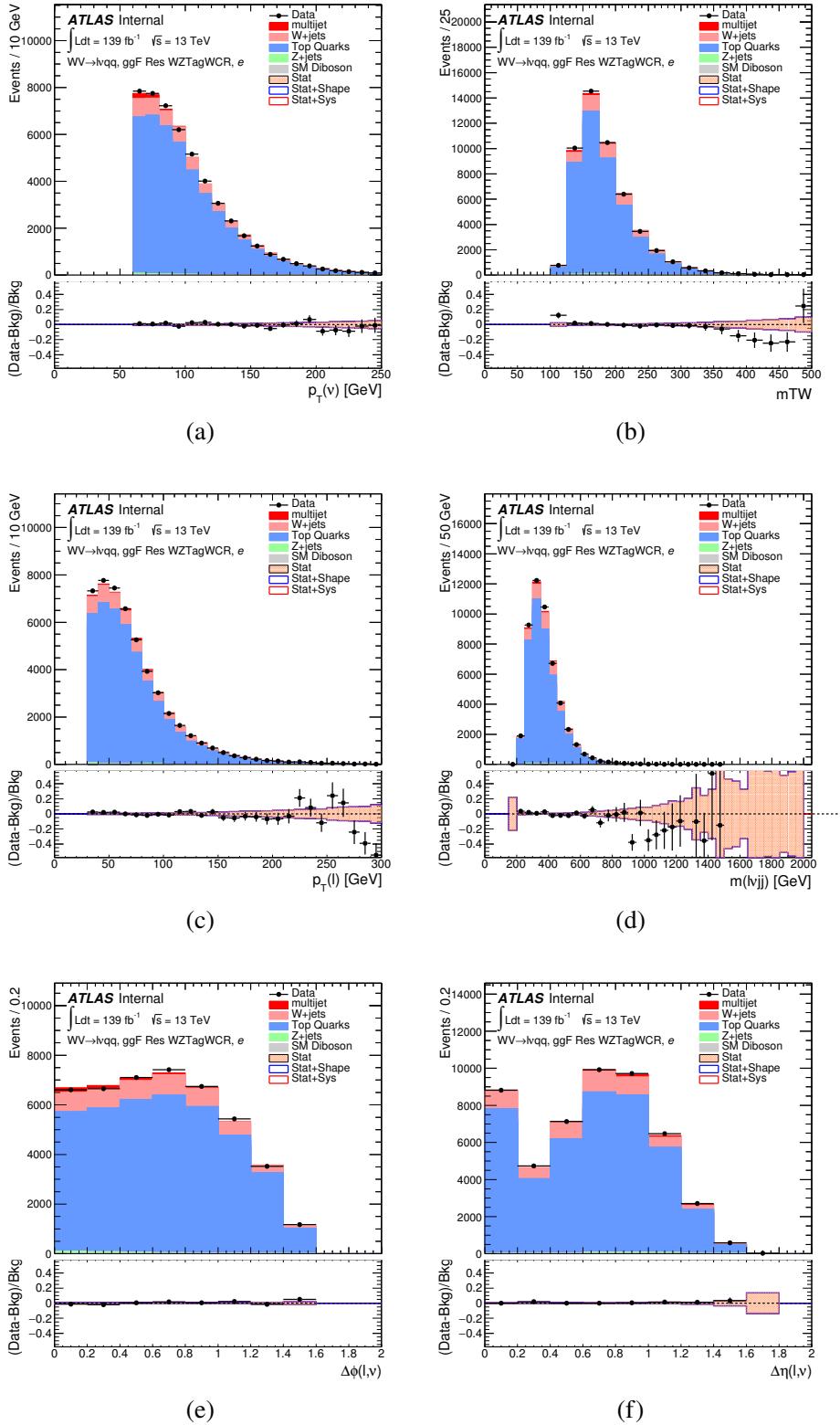


Figure 9.20: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the WZ untag electron channel. The MJ template is obtained from the pre-MJ-fit.

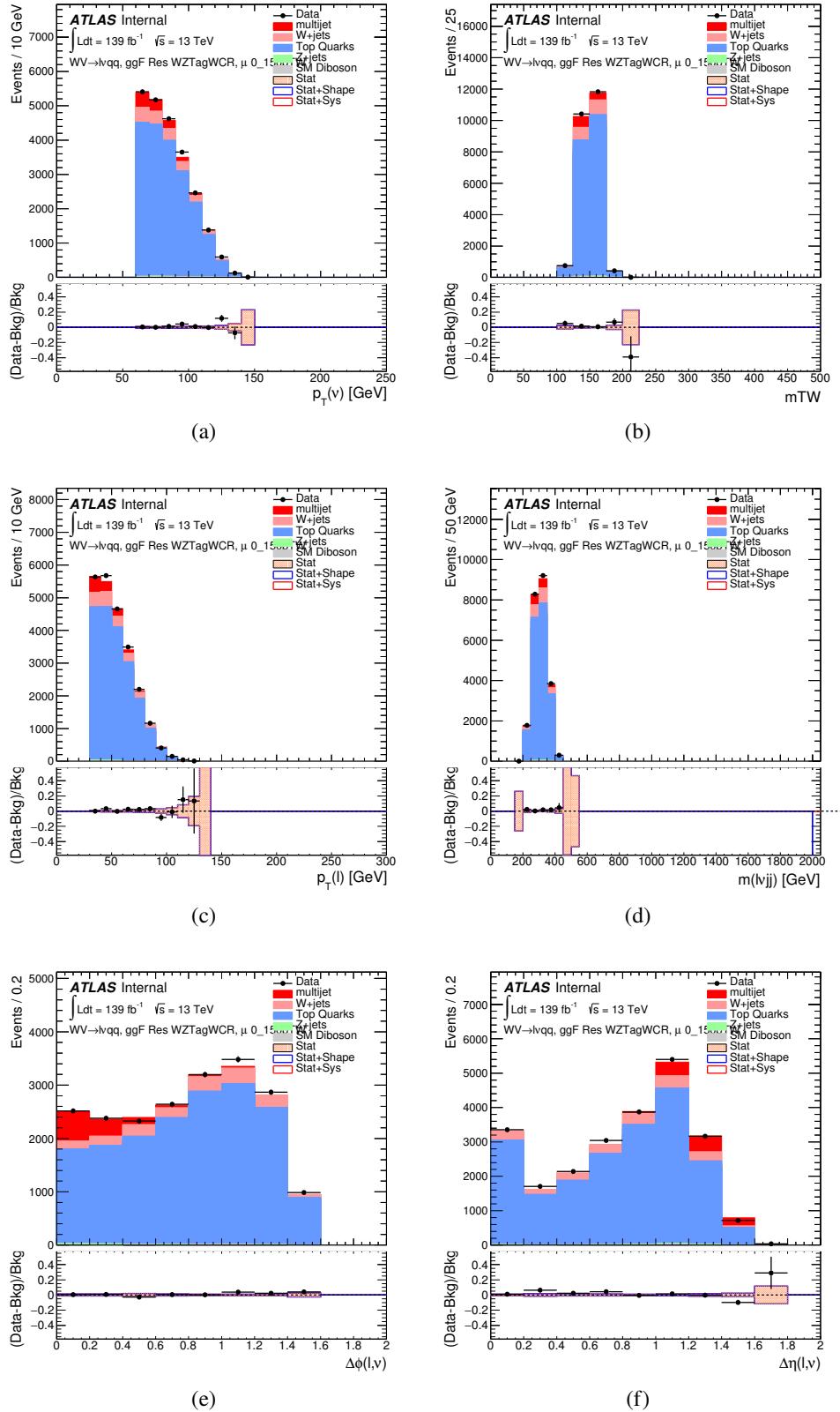


Figure 9.21: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the WZ untag muon channel. The MJ template is obtained from the pre-MJ-fit.

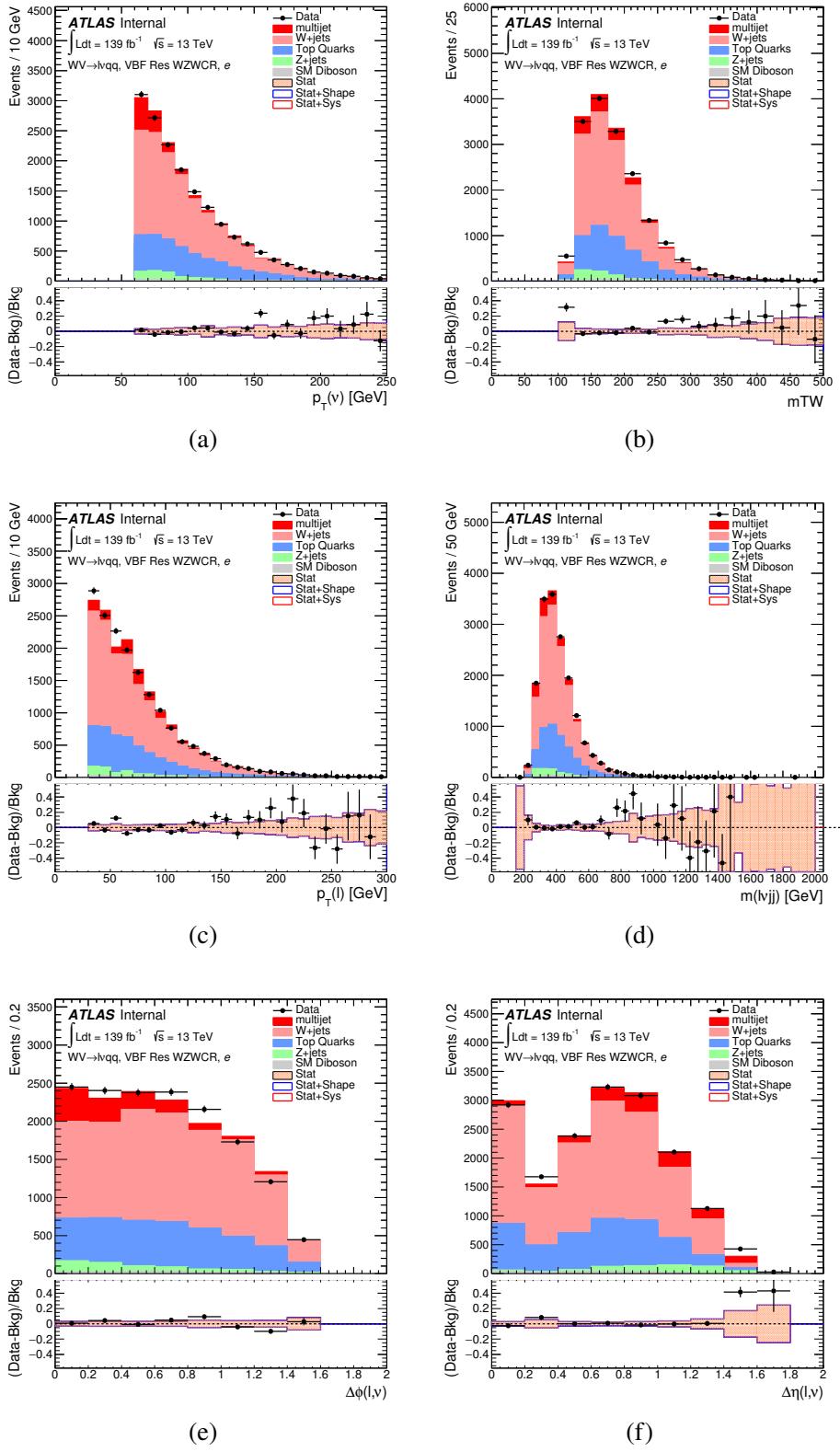


Figure 9.22: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the VBF WW electron channel. The MJ template is obtained from the pre-MJ-fit.

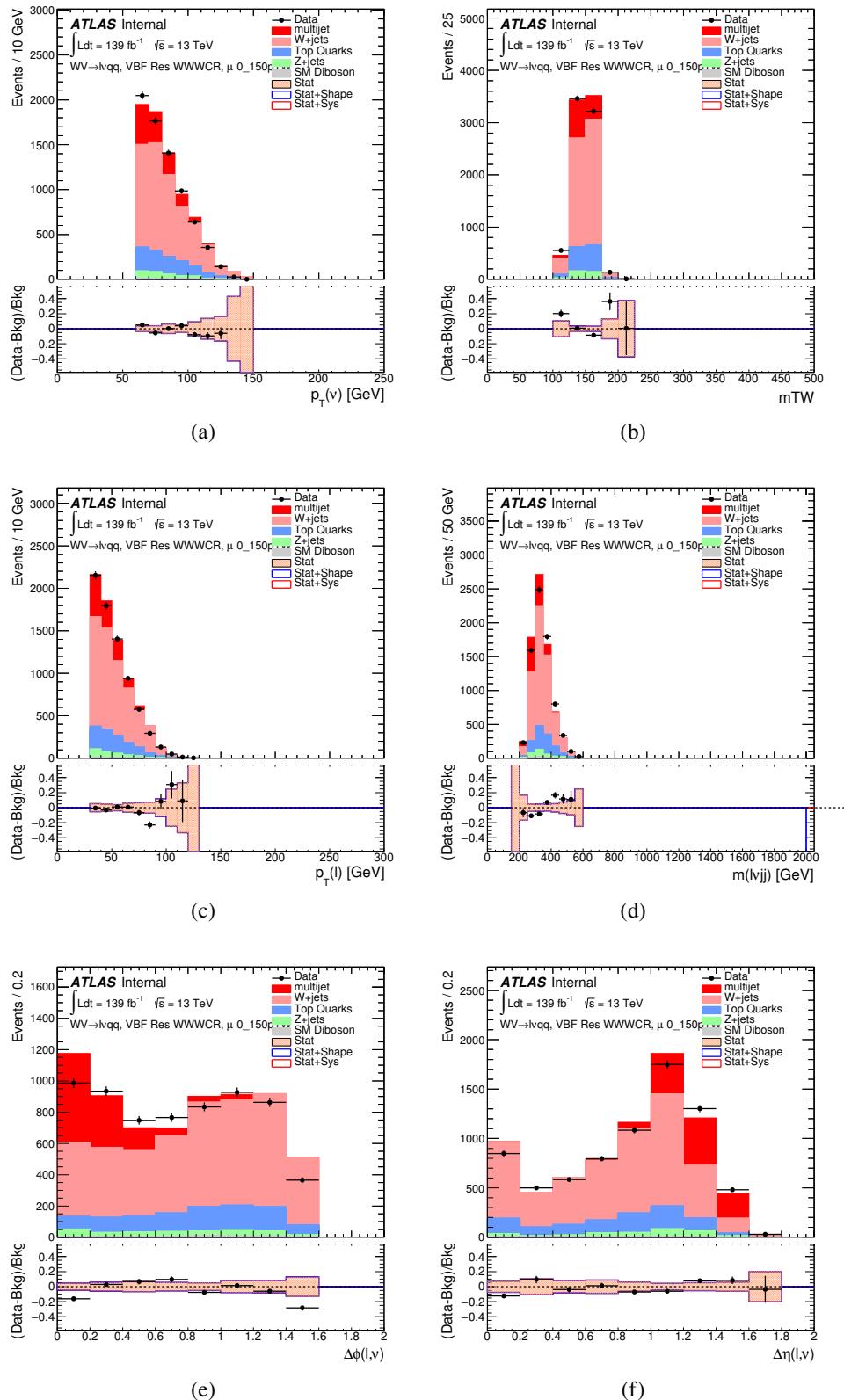


Figure 9.23: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{\ell\nu jj}$, lepton- ν angular distance in the VBF WW muon channel. The MJ template is obtained from the pre-MJ-fit.

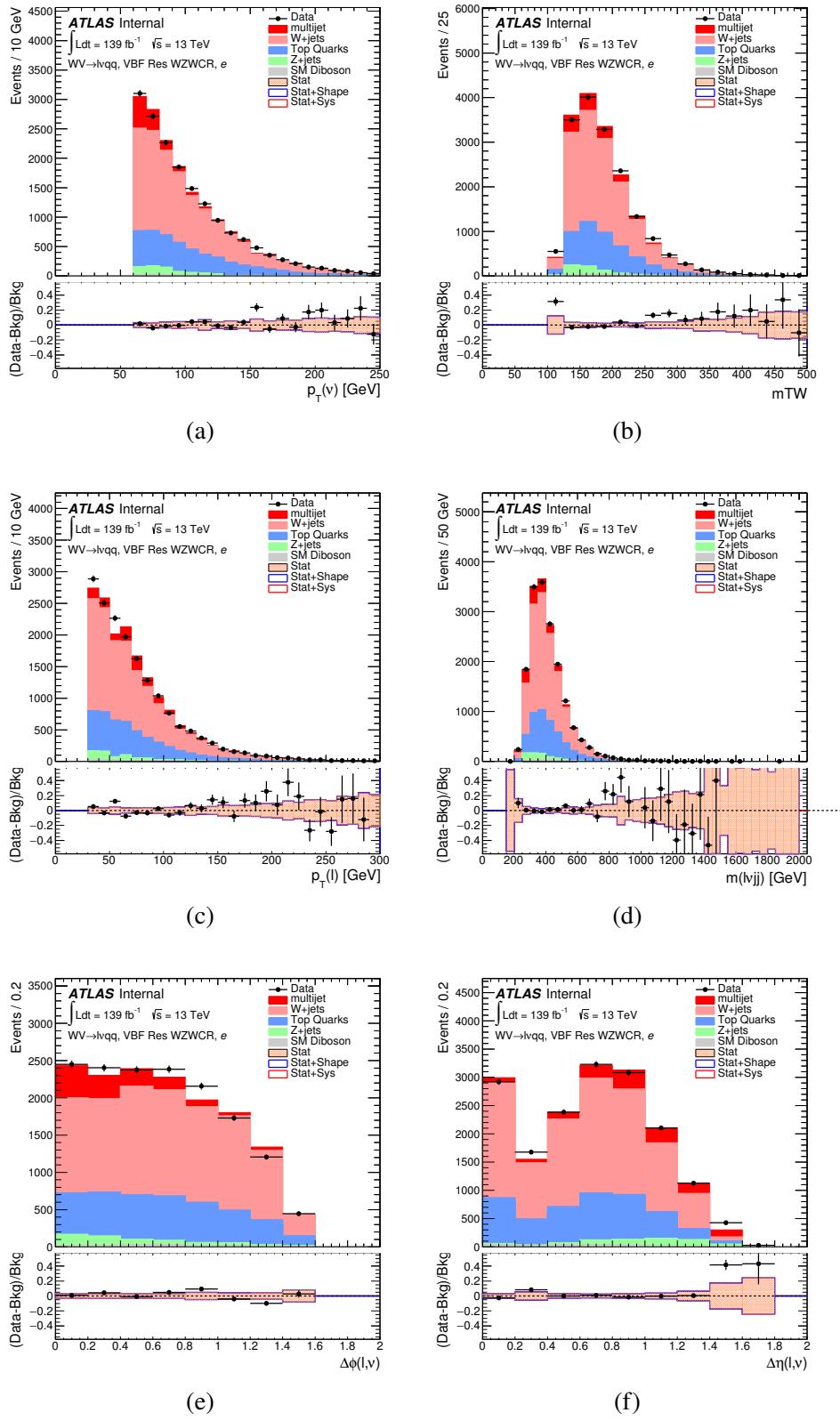


Figure 9.24: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{l\nu jj}$, lepton- ν angular distance in the VBF WZ electron channel. The MJ template is obtained from the pre-MJ-fit.

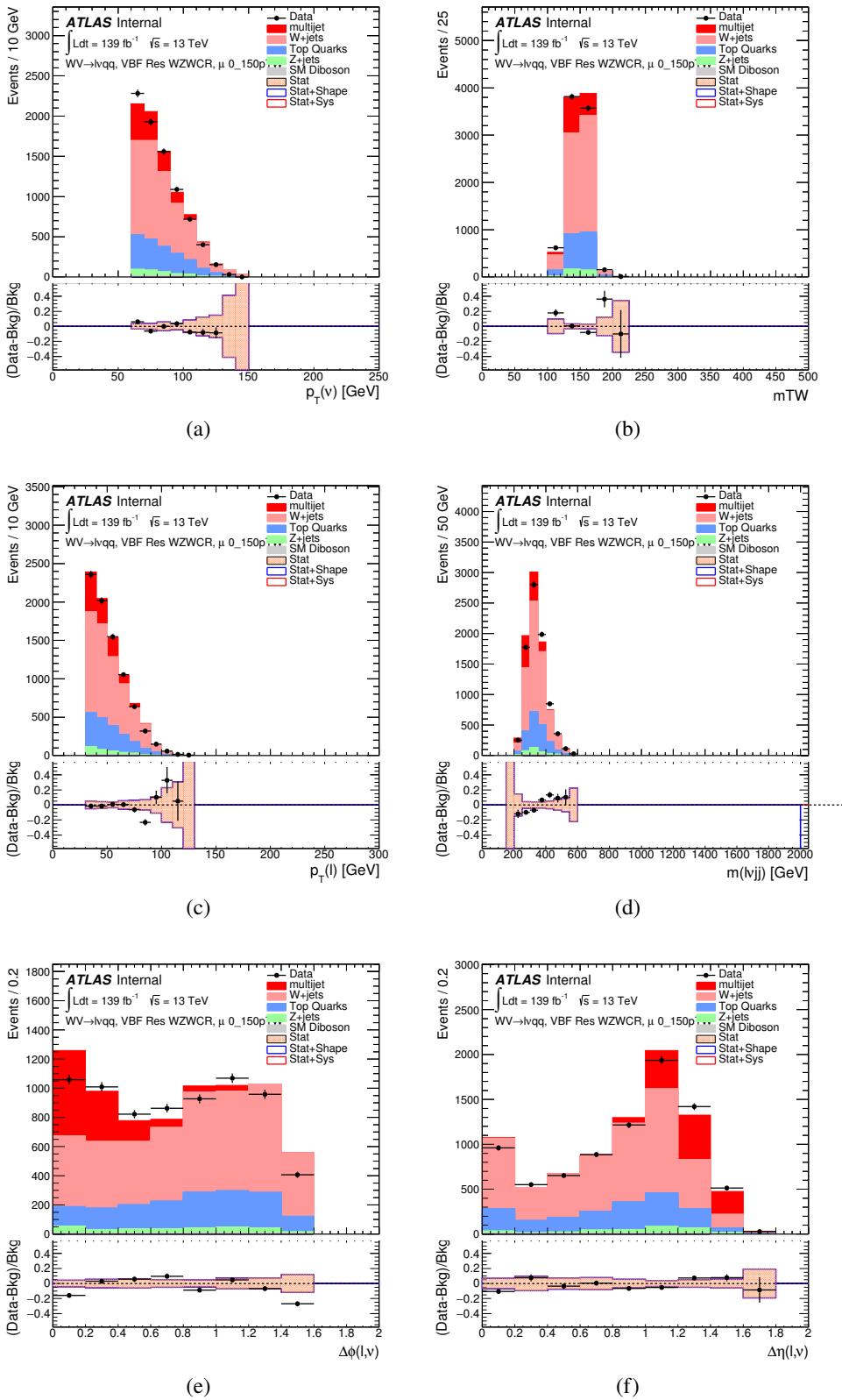


Figure 9.25: Postfit Data/MC comparison of distributions of E_T^{miss} , m_T^W , lepton and neutrino p_T , $m_{\ell\nu jj}$, lepton- ν angular distance in the VBF WZ muon channel. The MJ template is obtained from the pre-MJ-fit.

Full Run 2
ggF Res WWWCR

Sample	Yield	R.U.	SF
Top&W	645040 ± 1971.68	0.31%	0.998
Z&VV	24075.9		fixed
MJ_el	24156.3 ± 1224.62	5.06%	3.973
MJ_mu	35528.5 ± 923.94	2.60%	9.019

ggF Res WZ01bWCR

Sample	Yield	R.U.	SF
Top&W	644690 ± 1981.4	0.31%	0.997
Z&VV	24075.9		fixed
MJ_el	24366.5 ± 1232.69	5.05%	3.874
MJ_mu	35528.5 ± 921.27	2.58%	8.746

ggF Res WZ2bWCR

Sample	Yield	R.U.	SF
Top&W	71236.5 ± 688.74	0.97%	1.031
Z&VV	518.5		fixed
MJ_el	595.63 ± 449.34	75.44%	0.094
MJ_mu	1196.9 ± 222.13	18.56%	0.294

VBF Res WWWCR

Sample	Yield	R.U.	SF
Top&W	19032.3 ± 364.43	1.91%	0.928
Z&VV	1091.63		fixed
MJ_el	1425.73 ± 214.42	15.03%	0.235
MJ_mu	1281.36 ± 157.21	11.83%	0.314

VBF Res WZWCR

Sample	Yield	R.U.	SF
Top&W	21341.8 ± 392.21	1.84%	0.942
Z&VV	1111.75		fixed
MJ_el	1413.76 ± 230.36	16.29%	0.225
MJ_mu	1281.36 ± 157.21	12.27%	0.314

Table 9.5: Fit validation result in WCRs for 2015+16 data. The fit is done in various WCRs, in order to obtain the corresponding scale factors for MJ templates: ggF resolved WCR for the $WW \rightarrow lvqq$ selection, ggF resolved untagged WCR for the $WZ \rightarrow lvqq$ selection, ggF resolved tagged WCR for the $WZ \rightarrow lvqq$ selection, VBF resolved WCR for the $WW \rightarrow lvqq$ selection, and VBF resolved WCR for the $WZ \rightarrow lvqq$ selection. Post-fit event yields for electroweak processes and MJ contributions are shown. The SF column shows the corresponding normalization scale factors for electroweak processes from the fit. R.U. stands for relative uncertainty.

1072 **Chapter 10**

1073 **Systematic Uncertainties**

1074 This section describes the sources of systematic uncertainties of the m_{WV} dis-
1075 tribution. These uncertainties are divided into experimental and modeling uncer-
1076 tainties. Each systematic uncertainty is used as a nuisance parameter in the final
1077 likelihood fit.

1078 **10.1 Experimental Systematics**

1079 The uncertainty on the integrated luminosity of the dataset used is 1.7% and
1080 a systematic in the final fit. This uncertainty was calculated using $x - y$ beam
1081 separation scans [natasha ref P55].

1082 An additional source of systematic uncertainty is assigned to the pileup mod-
1083eling in MC samples. This ensures simulated detector response and particle re-
1084construction conditions are as similar as possible. The distribution of the average
1085 number of interactions per bunch crossing applied to simulation is called the μ pro-
1086file. The pileup modeling uncertainty is accounted for by re-weighting simulated
1087events so the average number of interactions per bunch crossing varies within its
1088uncertainty due to systematics from vertex reconstruction [ref ATL-COM-SOFT-

1089 2015-119]. The associated re-weighting factors are propagated through the entire
1090 analysis chain to construct a systematic uncertainty on m_{VV} .

1091 The single-lepton and E_T^{miss} triggers used are not fully efficient, so scale factors
1092 are applied to simulation to more accurately model the data. These scale factors
1093 are given by the ratio of the distribution of offline objects before trigger selection
1094 and after trigger selection. The associated uncertainty on these scale factors are
1095 used in the final fit.

1096 Uncertainties on small-R jet energy scale and resolution are measured in-situ
1097 by calculating the response between data and simulation. This analysis uses a
1098 reduced set of JES and JER uncertainties (totaling 30 and 8 systematics, re-
1099 spectively). These reduced sets of systematics are calculated using a principal
1100 component analysis, yield largely uncorrelated independent systematics. These
1101 uncertainties on JES and JER account for the dependence on p_T , η , μ , flavor re-
1102 sponse and global sequential corrections. Systematic uncertainties associated with
1103 b -tagging are also considered. These systematics are evaluated as uncertainties on
1104 the scale factor which account for the difference in b -tagging efficiencies in data
1105 and MC, and the flavor dependence (between b, c, and light jets).

1106 The uncertainty on the p_T scale of the large-R jets is determined by comparing
1107 the jet's p_T^{calo} to p_T^{track} in di-jet simulation and data. In addition to this uncertain-
1108 ties from tracking, modeling (Pythia vs Herwig), and statistical constraints are
1109 also calculated. The large-R jet p_T resolution is given by smearing the jet p_T with
1110 a Gaussian with a 2% width.

1111 The W/Z tagging efficiency SF is estimated by comparing the tagging effi-
1112 ciency in simulation with that in data for four regions of the W/Z tagger (D_2 fail,
1113 m_J fail; D_2 pass, m_J fail; D_2 fail, m_J pass; D_2 pass, m_J pass). (Additionally,
1114 separate scale factors are determined for events with large-R jets from W bosons

and top backgrounds.) A simultaneous template fit is used to fit the signal jets (jets initiated by W/Z bosons or top quarks) and background jets (all other jets from the simulated backgrounds) to the data in the four regions. using the m_J distributions. The SF for a given region is then given by:

$$SF = \frac{\epsilon_{data} = \frac{N_{fitted-signals}^{region}}{N_{all-regions}^{fitted-signals}}}{\epsilon_{MC} = \frac{N_{signal}^{region}}{N_{signal}^{all-regions}}} \quad (10.1)$$

The effects of experimental and theoretical uncertainties on the efficiency scale factor are determined by taking the ratio of efficiencies in data and simulation. By taking this ratio the uncertainties not arising for jet mass and D_2 cancel.

Lepton identification, reconstruction, isolation systematic uncertainties are determined by reconstructing the Z mass peak with a tag and probe method. The lepton energy and momentum scales are also measured with the Z mass peak.

As E_T^{miss} is calculated using all the physics objects in the event, all those objects associated errors result in an uncertainty on E_T^{miss} . Additionally, the unassociated tracks used to construct E_T^{miss} contribute to the uncertainty on E_T^{miss} .

10.2 Theory Systematics

Theoretical uncertainties for signal and background processes arise from uncertainties in the parameters used in Monte Carlo simulation. In particular for the $t\bar{t}$, $W/Z+jets$, diboson backgrounds and signal samples, the QCD scale, PDF, generator and hadronization uncertainties were evaluated. To assess the QCD scale uncertainty the renormalization and factorization scales were scaled up and down by a factor of two at the event generation stage of sample production. Uncertainties due to the choice of the parton distribution functions were evaluated by

1136 re-weighting samples from the nominal PDF to a set of error PDFs which account
1137 for the uncertainty of the fits used to produce the PDF set. In addition to this,
1138 samples are re-weighted to different PDF sets to account for the arbitrariness of
1139 the PDF choice. The difference between the m_{WV} distributions using different
1140 event generators is assessed by comparing samples generated with different gen-
1141 erators. Similarly, the uncertainty in hadronization models is accounted for by
1142 comparing samples created using different hadronization models (e.g. $t\bar{t}$ Powheg
1143 is compared to AMC@NLO, $W + jets$ compares Sherpa and MadGraph+Pythia
1144 samples). Figures 10.2 - 10.8 show the impact of these uncertainties on the $t\bar{t}$ and
1145 $W/Z + jets$ backgrounds. Additionally, contributions to the diboson background
1146 for the VBF analysis were found to be small and were accounted for by including
1147 a 5(10)% systematic in the diboson normalization in the final fit.

1148 The normalization of the $t\bar{t}$ and $W+jets$ processes impact the multijet tem-
1149 plate shape. The impact of these normalizations was assessed by including a
1150 shape systematic on the multijet background from varying the $t\bar{t}$ and $W+jets$
1151 normalization factors.

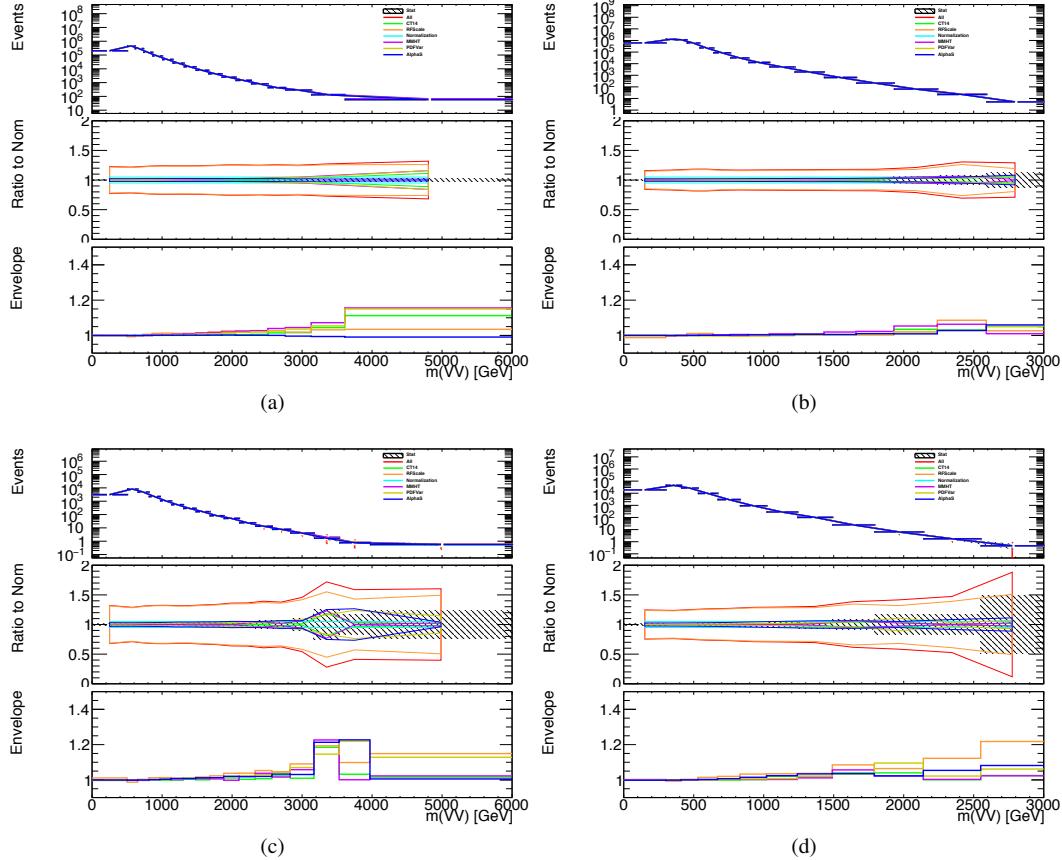


Figure 10.1: The $W/Z + \text{jet}$ systematics for the a) Merged ggF, b) Resolved ggF, c) Merged VBF, and d) Resolved VBF regions. The top subplot shows the nominal and variation distributions/bands, the middle shows the ratio of the two, and the final shows just the shape of the envelope (the final uncertainty).

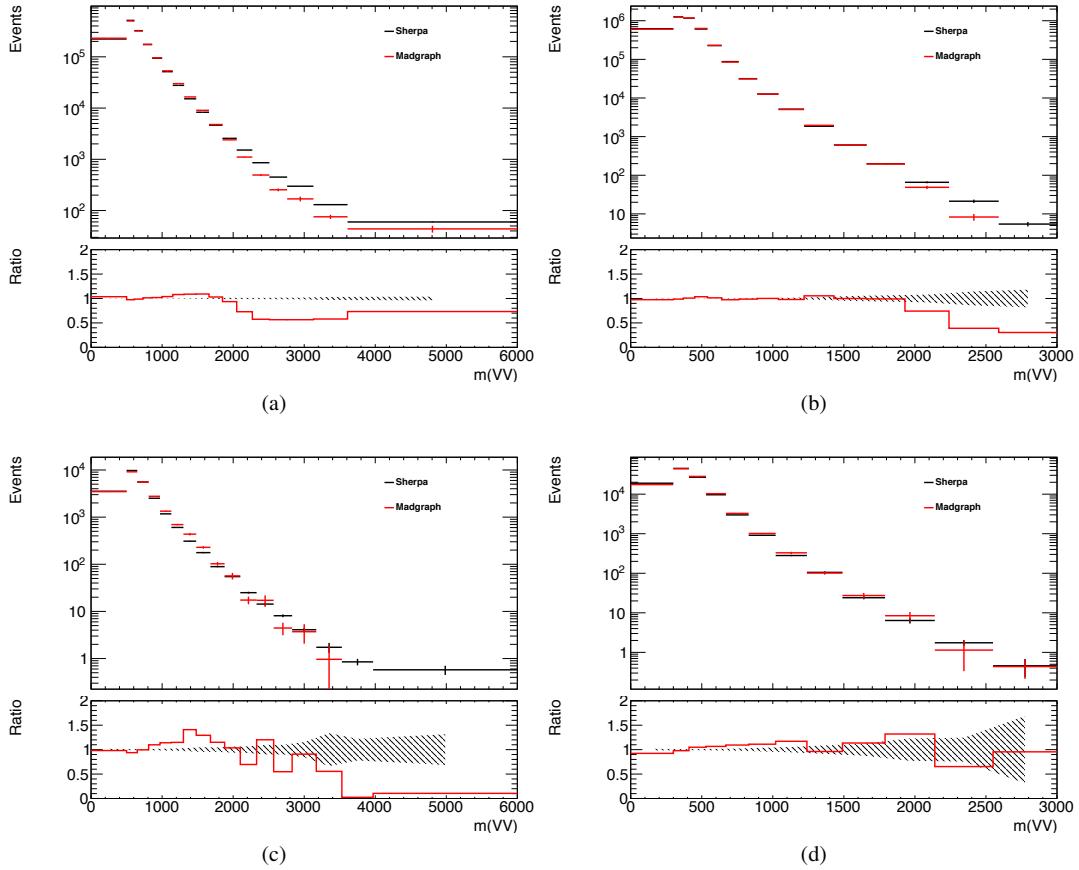


Figure 10.2: The two-point generator comparison between Sherpa and MadGraph for the $W/Z + \text{jet}$ samples in the a) Merged ggF, b) Resolved ggF, c) Merged VBF, and d) Resolved VBF regions. The normalization of the Madgraph sample is set to the Sherpa value to consider only shape effects. The bottom inset shows the ratio of the two.

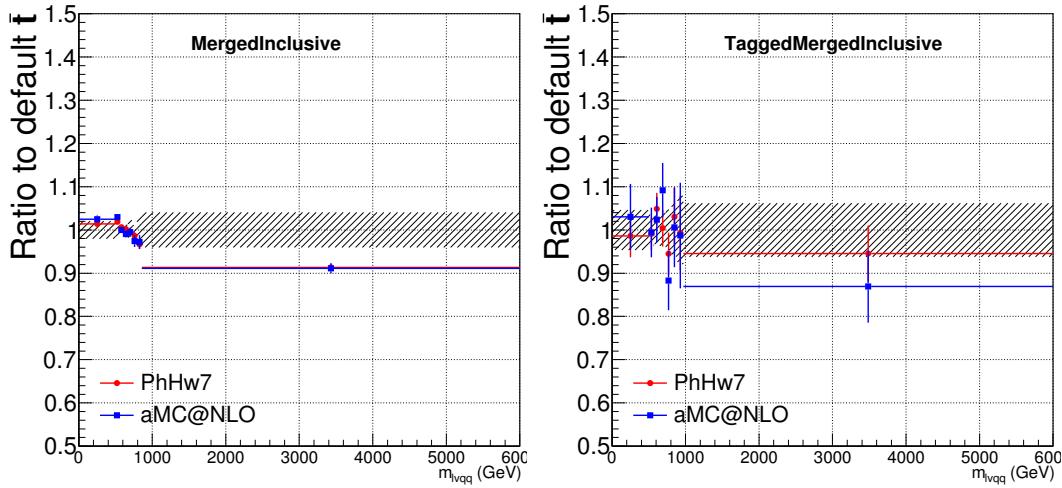


Figure 10.3: Ratio between the variations of generator (red) and hadronization (blue) variations for the Merged regime for $t\bar{t}$ sample.

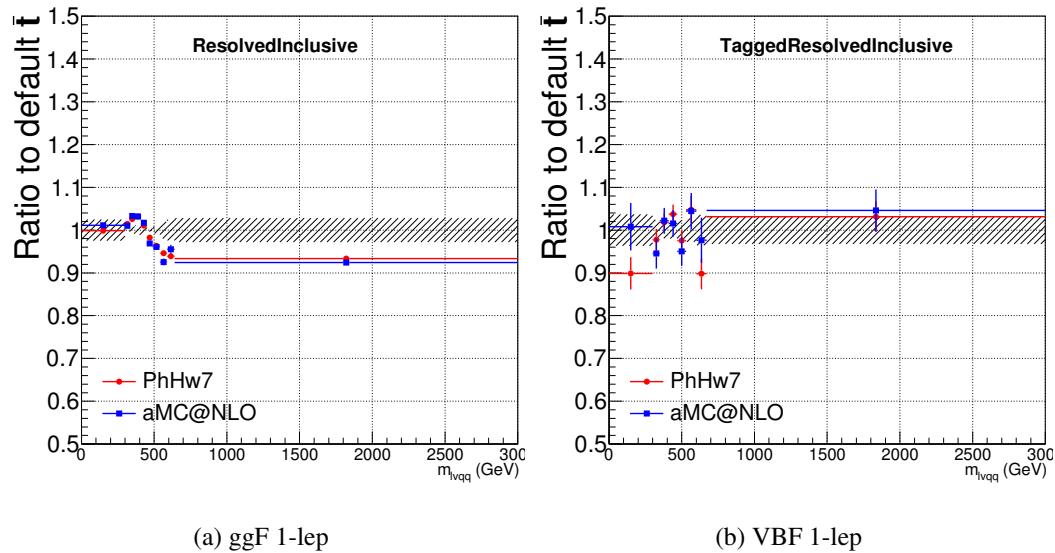


Figure 10.4: Ratio between the variations of generator (red) and hadronization (blue) variations for the Resolved regime for $t\bar{t}$ sample.

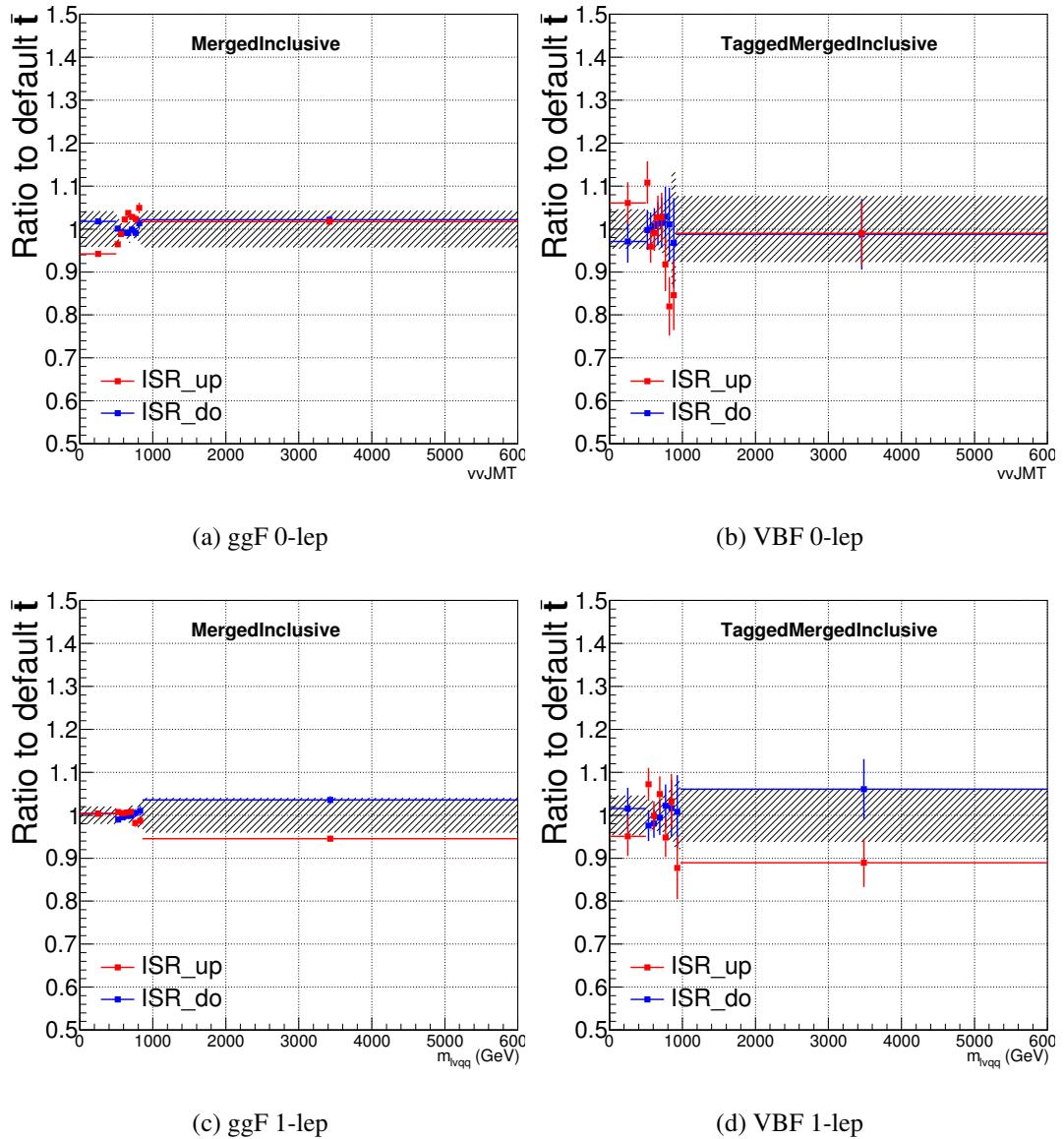


Figure 10.5: Ratio between the variations of ISR up (red) and down (blue) variations for the Merged regime for $t\bar{t}$ sample.

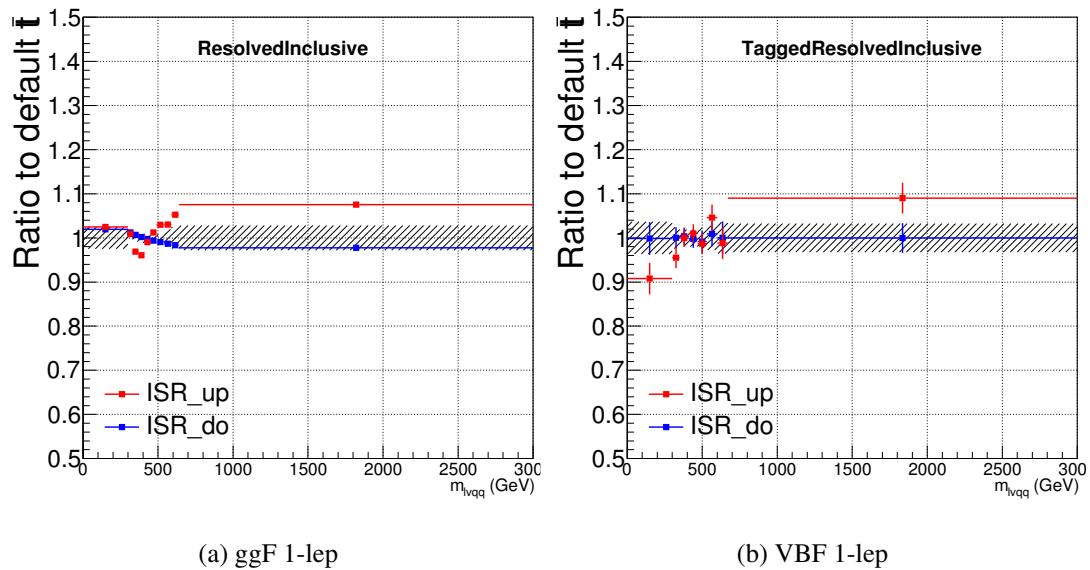


Figure 10.6: Ratio between the variations of ISR up (red) and down (blue) variations for the Resolved regime for $t\bar{t}$ sample.

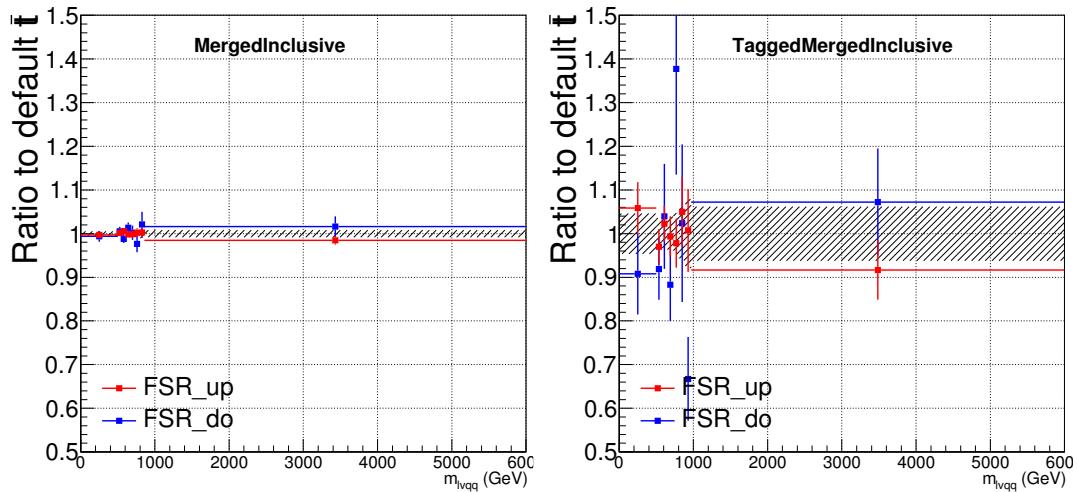


Figure 10.7: Ratio between the variations of FSR up (red) and down (blue) variations for the Merged regime for $t\bar{t}$ sample.

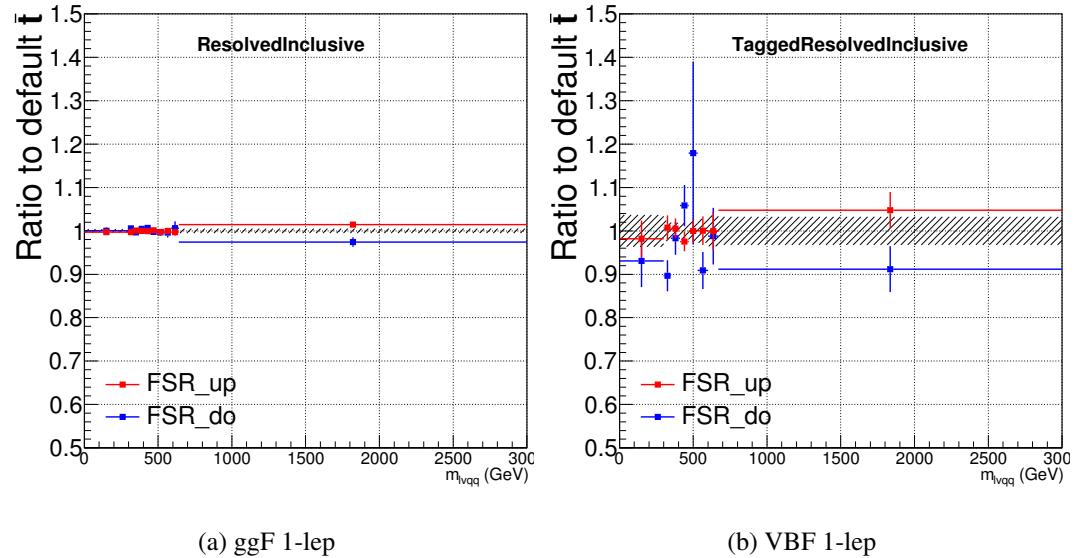


Figure 10.8: Ratio between the variations of FSR up (red) and down (blue) variations for the Resolved regime for $t\bar{t}$ sample.

₁₁₅₂ **Chapter 11**

₁₁₅₃ **Statistical Analysis**

₁₁₅₄ To determine the compatibility of the data collected with the proposed reso-
₁₁₅₅ nances a statistical procedure based on a likelihood function is used. A discovery
₁₁₅₆ test is used to measure the compatibility of the observed data with the back-
₁₁₅₇ ground only hypothesis. If the observed data is sufficiently incompatible with the
₁₁₅₈ background only hypothesis, this could indicate a discovery. In the absence of
₁₁₅₉ discovery, upper limits on the signal strength parameter, μ , are assessed using the
₁₁₆₀ CLs method.

₁₁₆₁ **11.1 Likelihood Function Definition**

₁₁₆₂ The likelihood function is product of Poisson probabilities for all analysis bins
₁₁₆₃ and systematic constraints:

$$\mathcal{L}(\mu, \boldsymbol{\theta}) = \prod_c \prod_i \frac{(\mu s_{ci}(\boldsymbol{\theta}) + b_{ci}(\boldsymbol{\theta}))^{n_{ci}}}{n_{ci}!} e^{-(\mu s_{ci}(\boldsymbol{\theta}) + b_{ci}(\boldsymbol{\theta}))} \prod_k (\theta'_k | \theta_k) \quad (11.1)$$

₁₁₆₄ Here c are the analysis channels considered and i runs over all the $m_{\ell\nu qq}$ bins

1165 used in the fit. The signal strength parameter, μ , multiplies the expected signal
1166 yield in each analysis bin, s_{ci} . The background content for channel c and bin i is
1167 given by b_{ci} . The dependence of signal and background predictions on system-
1168 atic uncertainties is described by the aforementioned set of nuisance parameters
1169 $\boldsymbol{\theta}$, which are parameterized by Gaussian or log-normal priors denoted here as
1170 θ_k . Statistical uncertainties of the simulated bin contents are also included as
1171 systematic uncertainties. Most systematics are correlated among all the analysis
1172 regions and considered to be independent from each other. The validity of this
1173 assumption is checked by evaluating the covariance of nuisance parameters.

1174 11.2 Fit Configuration

1175 The binning of $m_{\ell\nu qq}$ in signal regions for likelihood fit is determined by the
1176 statistical uncertainty of signal mass width. For each signal mass point, the signal
1177 mass resolution is given by the fitted Gaussian width of the $m_{\ell\nu qq}$. The fitted
1178 signal widths are then fit to a line to give a parameterized signal mass width, as
1179 shown in Figures 11.1 and 11.2. Bin widths are set first to this parameterized
1180 signal mass resolution. Then if the statistical uncertainty of the data or simulated
1181 background is more than 50%, bins are merged until the statistical uncertainty is
1182 less than 50%. All control regions contain only a single bin.

1183 For this analysis, each signal model is fit in the Merged and Resolved chan-
1184 nels for the relevant signal production mode simultaneously. The $W + \text{jets}$ and $t\bar{t}$
1185 normalizations are given by the best fit values in the overall fit and these fitted
1186 normalizations are then applied to those backgrounds in the SRs.

1187 Systematics may be affected by low statistics, leading to unsmooth m_{VV} dis-
1188 tributions with unphysically large fluctuations. This can lead to artificial pulls
1189 and constraints in the fit. To remove such issues a multi-step smoothing pro-

1190 cedure is applied to all systematic variation distributions in all regions. First,
1191 distributions are rebinned until the statistical error per bin is at least 5%. Next
1192 all local extrema are identified. The bins around smallest extrema are iteratively
1193 merged until only four local extrema remain. Then distributions are rebinned so
1194 that statistical uncertainties in each bin are $< 5\%$.

1195 For some systematics, up and down variations may be in the same direction
1196 with respect to the nominal distributions. This causes the variations to not cover
1197 the nominal choice, and the interpretation of the confidence interval is skewed as
1198 the nominal distribution should be bracketed by the up and down variations. This
1199 asymmetry may also lead to unconstrained systematics in the fit. To handle such
1200 asymmetric systematics, if the up and down variation for a given systematic are in
1201 the same direction for at least three m_{VV} bins the variation is averaged for those
1202 bins. The averaging procedure replaces bin-by-bin the up and down variation bins
1203 by $b_{\pm}^{new} = b_{nom} \pm \frac{|b_+ - b_-|}{2}$, where b_{nom} is the nominal bin content and b_{\pm} are the
1204 original up and down variation bin content. The same procedure is also applied to
1205 any variations where the integral of the difference between the up/down variation
1206 and the nominal distribution is twice that of the other down/up variation, further
1207 ensuring variations are symmetric around the nominal distribution.

1208 Finally, systematics that have a negligible effect on the m_{VV} distribution are
1209 not considered in the fit. Shape systematics where no bin in the variational dis-
1210 tribution deviates more than 1% from the nominal distribution (after normalizing
1211 all histograms to the nominal) are not included in the fit. Also, statistical bin
1212 uncertainties $< 1\%$ are ignored.

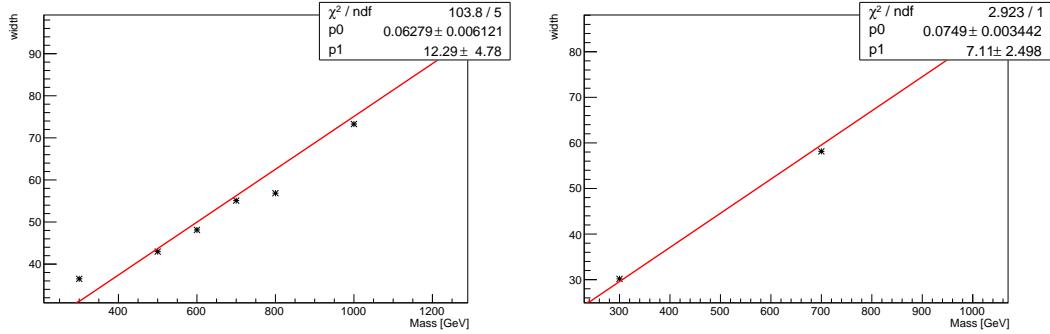


Figure 11.1: The HVT signal mass resolution as a function of mass fit with a straight line in the Resolved ggF region (left) and VBF (right) region.

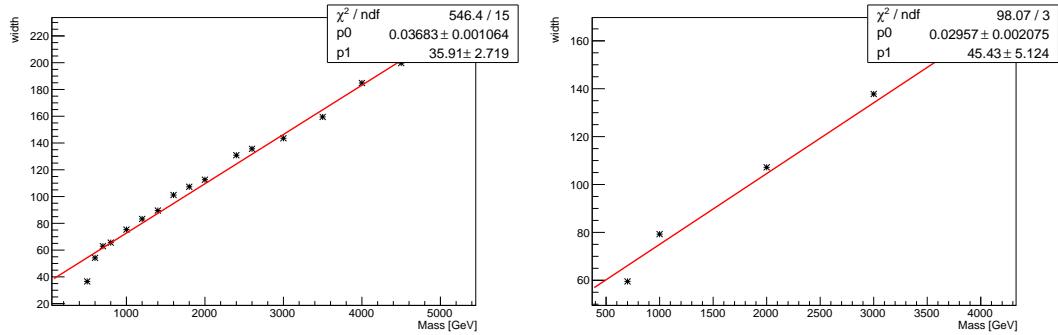


Figure 11.2: The HVT signal mass resolution as a function of mass fit with a straight line in the Merged ggF region (left) and VBF (right) region.

1213 11.3 Best Fit μ

1214 The best fit signal strength parameter is denoted by $\hat{\mu}$ and calculated by
 1215 maximizing the likelihood function with respect to all systematics and μ . The
 1216 corresponding set of systematics that maximize the likelihood are given by $\hat{\mu}$.
 1217 The first term in the likelihood is maximized when the expected number of signal
 1218 and background events is equal to the number of events in data ($n_{ci} = \mu s_{ci} +$
 1219 b_{ci}). Thus, by maximizing the likelihood, the fit determines values of μ and θ
 1220 that give the best agreement between expected and measured event yields. The

1221 second term in the likelihood is a penalty term which decreases the likelihood
1222 when systematics are shifted from their nominal values. This prevents the fit
1223 from profiling systematics in unphysical ways to maximize the likelihood. The
1224 uncertainty on μ is calculated by varying μ up and down until the natural log of
1225 the likelihood function shifts by one-half.

1226 11.4 Discovery Test

1227 To determine if the observed dataset is consistent with tested signal model a
1228 likelihood ratio is constructed:

$$\lambda(\mu) = \frac{\mathcal{L}(\mu, \hat{\theta}_\mu)}{\mathcal{L}(\hat{\mu}, \hat{\theta})} \quad (11.2)$$

1229 The denominator in this equation is the maximized value of \mathcal{L} over all system-
1230 atics and μ . The numerator is the maximized likelihood over all systematics for
1231 a given μ value, where the maximized systematics are given by $\hat{\mu}_\mu$. To test for
1232 the existence of signal the observed dataset the null hypothesis (H_0) is defined as
1233 the background only hypothesis and the alternate hypothesis includes signal and
1234 background (H_1). This test quantifies the compatibility of observed data with
1235 H_0 by calculating a p-value representing the probability of observing data as dis-
1236 crepant or more than the observed data under the H_0 . The test statistic used to
1237 calculate this p-value is given by (r_0):

$$r_0 = \begin{cases} -2 \ln \lambda(0), \hat{\mu} > 0 \\ +2 \ln \lambda(0), \hat{\mu} < 0 \end{cases} \quad (11.3)$$

1238 The expected distribution of the the test statistic under H_0 ($f(r_0|0)$) is used to
1239 calculate the p-value:

$$p_0 = \int_{r_0, obs}^{\infty} f(r_0 | 0) dr_0 \quad (11.4)$$

1240 Small p-values indicate the observed data is poorly described by H_0 . This
 1241 equivalent Z-score of a given p-value is usually used to further quantify the agree-
 1242 ment between the observed data and H_0 . The Z-score is given by the number of
 1243 standard deviations away from the mean of a Gaussian distribution, the integral
 1244 of the upper tail of the distribution would equal the p-value. Mathematically:

$$Z = \Phi^{-1}(1 - p_0) \quad (11.5)$$

1245 where Φ is the Gaussian cumulative distribution function. The statistical
 1246 significance of these tests are expressed as the Z -score. In particle physics, 3σ is
 1247 considered evidence for new phenomena and 5σ is the threshold for discovery.

1248 11.5 Exclusion Limits

1249 In the absence of discovery, upper limits on the signal strength, μ are set using
 1250 the CLs method [cite P60]. The test statistic for this test, q_μ , is constructed as:

$$\tilde{\lambda}_\mu = \begin{cases} \frac{\mathcal{L}(\mu, \hat{\theta}_\mu)}{\mathcal{L}(\hat{\mu}, \hat{\theta})}, \hat{\mu} > 0 \\ \frac{\mathcal{L}(\mu, \hat{\theta}_\mu)}{\mathcal{L}(0, \hat{\theta}_0)}, \hat{\mu} < 0 \end{cases} \quad (11.6)$$

$$\tilde{q}_\mu = \begin{cases} -2 \ln \tilde{\lambda}(\mu), \hat{\mu} < \mu \\ +2 \ln \tilde{\lambda}(\mu), \hat{\mu} > \mu \end{cases} \quad (11.7)$$

1251 As defined, larger values of q_μ correspond to increasing incompatibility between
 1252 the observed data and the background + signal hypothesis. The observed value
 1253 of the test statistic, $q_{\mu, obs}$, is then compared to its expected distribution, f , to

1254 calculate p-values to assess the likelihood of the background+signal hypothesis.

1255 Using these distributions, CL_s values are computed as:

$$1256 \quad CL_{s+b} = \int_{q_{\mu,obs}}^{\infty} f(q_{\mu}|\mu) dq_{\mu} \quad (11.8)$$

$$1257 \quad CL_b = \int_{q_0^{obs}}^{\infty} f(q_{\mu}|\mu = 0) dq_{\mu} \quad (11.9)$$

$$1258 \quad CL_s = \frac{CL_{s+b}}{CL_b} \quad (11.10)$$

1259 CL_{s+b} is the p-value for the signal + background hypothesis and CL_b is the
1260 p-value for the background only hypothesis. The CL_s value is interpreted as
1261 the probability to observe the background + signal hypothesis normalized to the
1262 probability of background-only hypothesis. Normalizing by CL_b prevents setting
1263 artificially strong exclusion limits due to downward fluctuations in data.

1264 For a given signal hypothesis, μ values are scanned simultaneously over all
1265 m_{WV} bins to find the μ value that yields $CL_s=0.05$, meaning the likelihood of
1266 finding data more incompatible with the signal+background hypothesis (relative
1267 to the background only hypothesis) is 5%. The 95% upper limit on the cross
1268 section is then calculated as the product of the μ value found, branching ratio,
1269 and theory cross section.

₁₂₆₉ **Chapter 12**

₁₂₇₀ **Results**

₁₂₇₁ **12.1 Expected and Measured Yields**

₁₂₇₂ The yield tables for the four analysis regions are shown in Tables ?? - ??.

₁₂₇₃ The fitted background normalizations are shown in Tables ??-??. The control

₁₂₇₄ region $m_{\ell\nu qq}$ distributions are shown in Figures 12.1 - 12.4. The signal region

₁₂₇₅ $m_{\ell\nu qq}$ distributions are shown in Figures 12.5 - ??.

	HP WCR	LP WCR	Resolved WCR
Electron Multi-jet	-	-	16507.83 ± 2314.87
Muon Multi-jet	-	-	19977.12 ± 2816.06
Diboson	1833.41 ± 177.78	3323.93 ± 320.92	9147.67 ± 961.63
Single-top	2160.62 ± 402.34	3551.09 ± 660.00	20058.36 ± 3817.26
$t\bar{t}$	15518.86 ± 338.22	24069.54 ± 453.15	138866.23 ± 1989.71
$W+jets$	40141.57 ± 357.79	88113.06 ± 487.87	673200.38 ± 4120.53
$Z+jets$	778.83 ± 78.93	1765.54 ± 179.10	16570.50 ± 1672.71
Total	60433.29 ± 664.92	120823.16 ± 1006.99	894328.12 ± 7247.12
Data	60264.00	120852.00	895362.00
	HP TCR	LP TCR	Resolved TCR
Electron Multi-jet	-	-	-
Muon Multi-jet	-	-	-
Diboson	421.11 ± 37.98	550.44 ± 53.10	996.87 ± 119.63
Single-top	4691.44 ± 846.11	3466.26 ± 631.03	16848.71 ± 3258.26
$t\bar{t}$	38945.18 ± 848.77	33836.95 ± 637.04	224226.14 ± 3212.76
$W+jets$	2258.34 ± 20.13	6564.78 ± 36.35	23466.41 ± 143.63
$Z+jets$	66.35 ± 6.72	213.26 ± 21.63	846.66 ± 85.47
Total	46382.43 ± 1199.25	44631.70 ± 899.23	266384.78 ± 4580.43
Data	46354.00	44629.00	266443.00
	WW SR	LP SR	Resolved 1-lepton SR
Electron Multi-jet	-	-	10788.40 ± 1512.85
Muon Multi-jet	-	-	15759.50 ± 2221.53
Diboson	4990.30 ± 376.50	3901.07 ± 313.22	16971.29 ± 1523.77
Single-top	3117.71 ± 565.07	2176.46 ± 400.52	20422.85 ± 3731.94
$t\bar{t}$	13785.77 ± 302.14	11005.12 ± 207.41	126965.25 ± 1819.66
$W+jets$	24718.56 ± 223.72	60080.66 ± 333.12	444133.56 ± 2719.02
$Z+jets$	478.18 ± 48.46	1226.69 ± 124.44	11686.32 ± 1179.69
Total	47090.52 ± 777.65	78389.98 ± 654.22	646727.19 ± 5963.98
Data	47330.00	78380.00	645610.00

Table 12.1: Expected and Measured for DY WW $W+jets$, $t\bar{t}$ control regions and signal regions.

	HP Untagged WCR	LP Untagged WCR	Resolved Untagged WCR
Electron Multi-jet	-	-	15080.03 ± 2277.99
Muon Multi-jet	-	-	27347.10 ± 2950.07
Diboson	1508.48 ± 154.20	2758.24 ± 284.50	9038.55 ± 728.69
Single-top	1756.59 ± 306.69	2913.18 ± 515.93	20511.74 ± 3523.47
$t\bar{t}$	13134.00 ± 238.30	21815.37 ± 334.98	140157.77 ± 2636.96
$W+jets$	40654.84 ± 333.65	87657.76 ± 501.96	665909.12 ± 4420.62
$Z+jets$	768.72 ± 77.97	1759.87 ± 178.96	16512.46 ± 1673.23
Total	57822.63 ± 540.40	116904.42 ± 862.16	894556.75 ± 7492.20
Data	57699.00	117306.00	895362.00
	HP Tagged WCR	LP Tagged WCR	Resolved Tagged WCR
Electron Multi-jet	-	-	384.58 ± 57.11
Muon Multi-jet	-	-	602.93 ± 190.12
Diboson	30.22 ± 4.69	48.95 ± 7.16	264.64 ± 28.24
Single-top	308.44 ± 56.19	371.59 ± 69.43	5752.39 ± 1029.97
$t\bar{t}$	1683.82 ± 48.73	2041.48 ± 70.00	58431.49 ± 614.30
$W+jets$	583.55 ± 75.37	1109.45 ± 85.78	11891.68 ± 903.01
$Z+jets$	13.19 ± 1.34	23.06 ± 2.34	324.74 ± 32.85
Total	2619.22 ± 106.00	3594.53 ± 130.90	77652.45 ± 1514.89
Data	2565.00	3546.00	77973.00
	HP Untagged TCR	LP Untagged TCR	Resolved Untagged TCR
Electron Multi-jet	-	-	-
Muon Multi-jet	-	-	-
Diboson	289.45 ± 28.45	346.78 ± 35.85	650.85 ± 65.56
Single-top	3107.99 ± 538.03	2250.64 ± 385.41	9606.87 ± 1698.22
$t\bar{t}$	30992.40 ± 562.33	26954.21 ± 413.89	91893.59 ± 1728.91
$W+jets$	2236.29 ± 18.35	4874.03 ± 27.91	16122.97 ± 107.03
$Z+jets$	71.54 ± 7.26	155.50 ± 15.81	577.71 ± 58.54
Total	36697.66 ± 779.03	34581.16 ± 567.59	118851.98 ± 2427.40
Data	36677.00	34573.00	118928.00
	HP Tagged TCR	LP Tagged TCR	Resolved Tagged TCR
Electron Multi-jet	-	-	-
Muon Multi-jet	-	-	-
Diboson	9.72 ± 1.13	8.75 ± 1.16	34.06 ± 4.98
Single-top	105.87 ± 20.65	119.66 ± 22.68	656.89 ± 132.96
$t\bar{t}$	1904.75 ± 50.61	1483.86 ± 47.05	17965.33 ± 188.87
$W+jets$	32.36 ± 4.28	85.74 ± 6.96	489.01 ± 37.13
$Z+jets$	1.27 ± 0.13	1.93 ± 0.20	19.14 ± 1.94
Total	2053.98 ± 54.84	1699.93 ± 52.70	19164.43 ± 234.01
Data	2047.00	1708.00	19143.00

Table 12.2: Expected and Measured for DY WZ $W+jets$, $t\bar{t}$ tag and untag control regions.

	HP Untagged SR	LP Untagged SR	Resolved Untagged SR
Electron Multi-jet	-	-	7782.17 ± 1175.56
Muon Multi-jet	-	-	17004.81 ± 1834.40
Diboson	3041.17 ± 273.77	2266.35 ± 212.79	14724.12 ± 1224.31
Single-top	2123.28 ± 373.83	1379.35 ± 240.92	18336.88 ± 3082.47
$t\bar{t}$	11678.86 ± 213.63	8906.34 ± 136.88	112669.24 ± 2122.46
$W+jets$	22741.32 ± 191.47	41726.76 ± 240.56	342934.00 ± 2280.21
$Z+jets$	442.03 ± 44.84	849.79 ± 86.42	9271.83 ± 939.52
Total	40026.65 ± 546.81	55128.59 ± 432.90	522723.03 ± 5131.71
Data	40193.00	54735.00	521813.00
	HP Tagged SR	LP Tagged SR	Resolved Tagged SR
Electron Multi-jet	-	-	199.22 ± 29.58
Muon Multi-jet	-	-	393.43 ± 124.06
Diboson	102.58 ± 11.59	65.44 ± 8.05	624.07 ± 58.10
Single-top	178.21 ± 33.62	155.53 ± 28.95	3470.39 ± 617.48
$t\bar{t}$	1017.93 ± 31.95	706.76 ± 26.20	38189.30 ± 401.91
$W+jets$	325.58 ± 41.62	575.36 ± 43.29	6161.96 ± 467.71
$Z+jets$	7.81 ± 0.80	11.62 ± 1.19	183.36 ± 18.55
Total	1632.11 ± 63.39	1514.70 ± 58.86	49221.74 ± 884.06
Data	1699.00	1559.00	48919.00

Table 12.3: Expected and Measured for DY WZ $W+jets$, $t\bar{t}$ tag and untag signal regions.

	HP WCR	LP WCR	Resolved WCR
Electron Multi-jet	-	-	898.48 ± 137.82
Muon Multi-jet	-	-	601.46 ± 182.74
Diboson	107.45 ± 45.20	166.87 ± 68.11	292.10 ± 235.29
Single-top	78.19 ± 18.22	132.71 ± 31.93	879.82 ± 216.89
$t\bar{t}$	400.71 ± 28.35	569.70 ± 48.88	5067.51 ± 155.69
$W+jets$	864.49 ± 63.44	1940.80 ± 89.41	18563.70 ± 408.99
$Z+jets$	19.51 ± 2.00	46.63 ± 4.77	795.20 ± 80.89
Total	1470.35 ± 84.89	2856.71 ± 126.74	27098.28 ± 594.01
Data	1495.00	2898.00	27120.00
	HP TCR	LP TCR	Resolved TCR
Electron Multi-jet	-	-	-
Muon Multi-jet	-	-	-
Diboson	14.95 ± 6.61	27.57 ± 14.12	24.33 ± 20.32
Single-top	68.31 ± 16.17	58.93 ± 13.56	278.60 ± 73.04
$t\bar{t}$	496.60 ± 31.72	401.23 ± 32.13	3834.49 ± 104.60
$W+jets$	50.68 ± 4.19	144.02 ± 7.86	450.01 ± 11.87
$Z+jets$	1.32 ± 0.14	5.35 ± 0.55	29.96 ± 3.07
Total	631.87 ± 36.45	637.10 ± 38.44	4617.39 ± 129.77
Data	636.00	634.00	4615.00
	HP SR	LP SR	Resolved SR
Electron Multi-jet	-	-	596.34 ± 91.52
Muon Multi-jet	-	-	481.01 ± 144.48
Diboson	148.84 ± 48.64	181.42 ± 67.30	395.52 ± 318.06
Single-top	79.49 ± 19.80	56.82 ± 14.89	782.07 ± 190.79
$t\bar{t}$	338.42 ± 24.14	236.80 ± 20.88	4261.70 ± 138.98
$W+jets$	501.13 ± 39.36	1347.76 ± 64.50	11445.73 ± 291.49
$Z+jets$	9.25 ± 0.95	28.77 ± 2.95	567.66 ± 57.94
Total	1077.13 ± 69.93	1851.57 ± 96.73	18530.03 ± 523.88
Data	1096.00	1846.00	18530.00

Table 12.4: Expected and Measured for VBF WW $W+jets$, $t\bar{t}$ control regions and signal regions.

	HP WCR	LP WCR	Resolved WCR
Electron Multi-jet	-	-	870.00 ± 132.75
Muon Multi-jet	-	-	618.45 ± 196.90
Diboson	92.92 ± 41.77	145.90 ± 64.26	228.62 ± 114.62
Single-top	71.13 ± 16.29	118.82 ± 27.98	1209.87 ± 281.64
$t\bar{t}$	427.80 ± 29.72	509.19 ± 46.57	6860.87 ± 254.83
$W+jets$	871.68 ± 64.22	2020.67 ± 93.54	19088.50 ± 442.10
$Z+jets$	19.58 ± 2.01	47.39 ± 4.85	800.19 ± 82.02
Total	1483.11 ± 83.79	2841.97 ± 125.92	29676.50 ± 644.96
Data	1495.00	2898.00	29755.00
	HP TCR	LP TCR	Resolved TCR
Electron Multi-jet	-	-	-
Muon Multi-jet	-	-	-
Diboson	10.12 ± 4.51	12.73 ± 6.55	14.23 ± 7.49
Single-top	51.57 ± 12.31	35.07 ± 8.17	169.21 ± 44.54
$t\bar{t}$	470.06 ± 28.97	298.99 ± 25.28	2414.75 ± 75.42
$W+jets$	49.64 ± 4.17	109.69 ± 6.16	378.22 ± 12.05
$Z+jets$	1.28 ± 0.13	4.81 ± 0.50	17.62 ± 1.83
Total	582.67 ± 32.07	461.30 ± 28.05	2994.03 ± 88.75
Data	584.00	459.00	3001.00
	HP SR	LP SR	Resolved SR
Electron Multi-jet	-	-	444.65 ± 67.99
Muon Multi-jet	-	-	397.29 ± 125.59
Diboson	109.66 ± 44.13	112.28 ± 46.45	265.75 ± 139.43
Single-top	63.16 ± 15.20	48.02 ± 11.56	872.16 ± 205.00
$t\bar{t}$	348.95 ± 24.34	190.68 ± 17.75	5134.25 ± 193.57
$W+jets$	467.21 ± 37.12	973.73 ± 47.91	10226.83 ± 254.67
$Z+jets$	8.15 ± 0.84	23.62 ± 2.43	558.48 ± 57.25
Total	997.13 ± 64.42	1348.33 ± 70.06	17899.41 ± 432.98
Data	1018.00	1313.00	17826.00

Table 12.5: Expected and Measured for VBF WZ $W+jets$, $t\bar{t}$ control regions and signal regions.

Background	Fitted Normalization
XS_Top_LP_lvqq_Merg_binned	$0.905^{+0.0166}_{-0.0166}$
XS_Top_Merg	$0.936^{+0.0199}_{-0.0199}$
XS_Top_Res	$0.957^{+0.0134}_{-0.0134}$
XS_Wjets_LP_lvqq_Merg_binned	$0.884^{+0.00489}_{-0.00489}$
XS_Wjets_Merg	$0.931^{+0.00831}_{-0.00831}$
XS_Wjets_Res	$1.03^{+0.00628}_{-0.00628}$

Table 12.6: Fitted background normalizations for $t\bar{t}$ and $W+jets$ backgrounds for the DY WW analysis region.

Background	Fitted Normalization
XS_Top_LP_Tag_lvqq_Merg_binned	$0.973^{+0.0333}_{-0.0333}$
XS_Top_LP_lvqq_Merg_binned	$0.894^{+0.0135}_{-0.0135}$
XS_Top_Merg	$0.893^{+0.016}_{-0.016}$
XS_Top_Res	$0.965^{+0.0179}_{-0.0179}$
XS_Top_Tag_lvqq_Merg_binned	$0.954^{+0.0276}_{-0.0276}$
XS_Top_Tag_lvqq_Res_binned	$0.999^{+0.0105}_{-0.0105}$
XS_Wjets_LP_Tag_lvqq_Merg_binned	$0.912^{+0.0703}_{-0.0703}$
XS_Wjets_LP_lvqq_Merg_binned	$0.876^{+0.00502}_{-0.00502}$
XS_Wjets_Merg	$0.948^{+0.00779}_{-0.00779}$
XS_Wjets_Res	$1.01^{+0.00673}_{-0.00673}$
XS_Wjets_Tag_lvqq_Merg_binned	$0.906^{+0.117}_{-0.117}$
XS_Wjets_Tag_lvqq_Res_binned	$1.2^{+0.0904}_{-0.0904}$

Table 12.7: Fitted background normalizations for $t\bar{t}$ and $W+jets$ backgrounds for the DY WZ analysis region.

Background	Fitted Normalization
XS_Top_LP_lvqq_Merg_binned	$0.79^{+0.0673}_{-0.0673}$
XS_Top_Merg	$0.888^{+0.061}_{-0.061}$
XS_Top_Res	$1.01^{+0.0311}_{-0.0311}$
XS_Wjets_LP_lvqq_Merg_binned	$0.88^{+0.0423}_{-0.0423}$
XS_Wjets_Merg	$0.881^{+0.0677}_{-0.0677}$
XS_Wjets_Res	$0.932^{+0.0202}_{-0.0202}$

Table 12.8: Fitted background normalizations for $t\bar{t}$ and $W+jets$ backgrounds for the VBF WW analysis region.

Background	Fitted Normalization
XS_Top_LP_lvqq_Merg_binned	$0.708^{+0.064}_{-0.064}$
XS_Top_Merg	$0.958^{+0.0644}_{-0.0644}$
XS_Top_Res	$1.02^{+0.038}_{-0.038}$
XS_Wjets_LP_lvqq_Merg_binned	$0.9^{+0.0438}_{-0.0438}$
XS_Wjets_Merg	$0.883^{+0.0685}_{-0.0685}$
XS_Wjets_Res	$0.945^{+0.0219}_{-0.0219}$

Table 12.9: Fitted background normalizations for $t\bar{t}$ and $W+jets$ backgrounds for the VBF WZ analysis region.

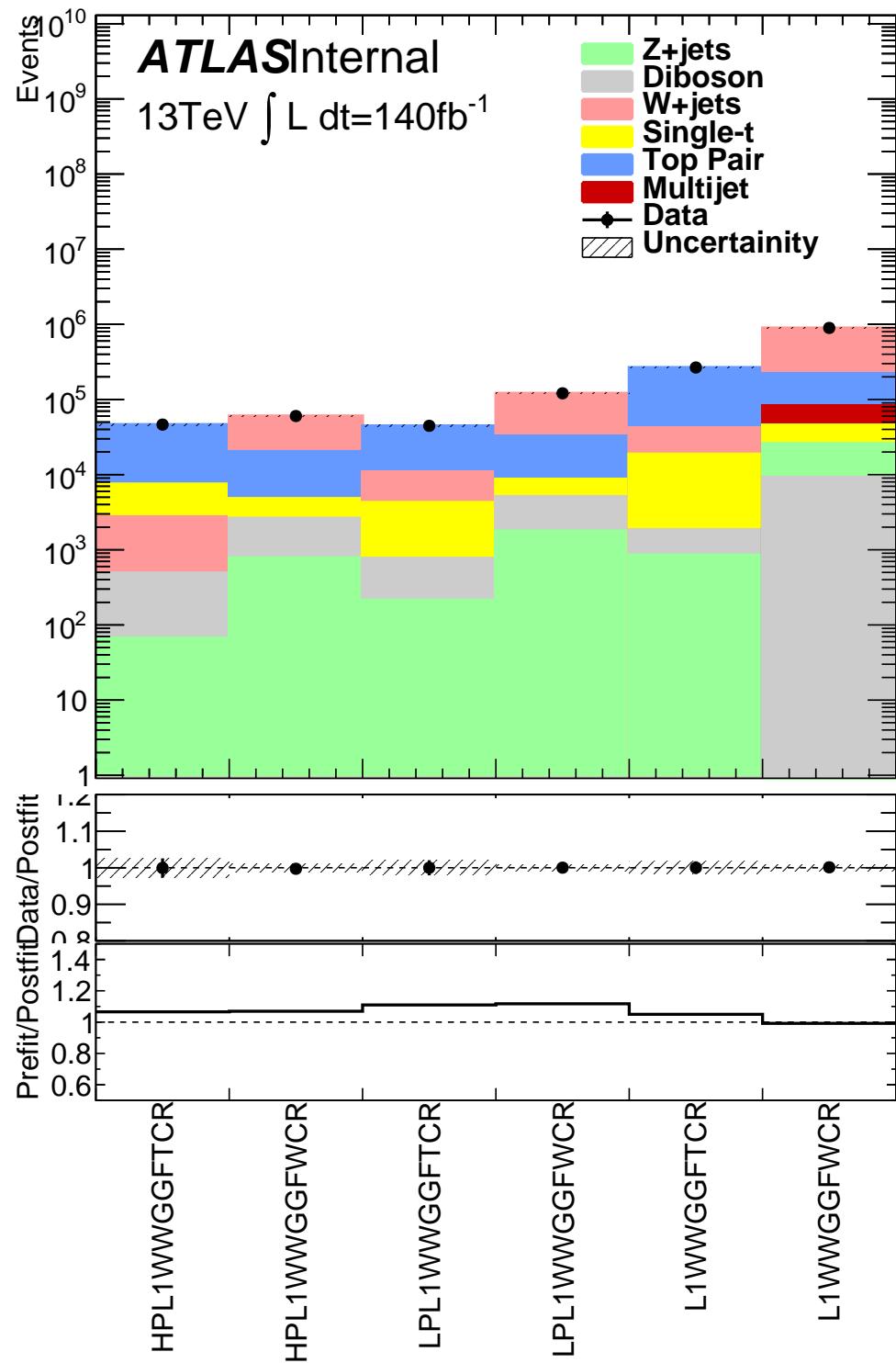


Figure 12.1: This figure shows the distribution of $m_{\ell\nu qq}$ in the DY WW control regions.

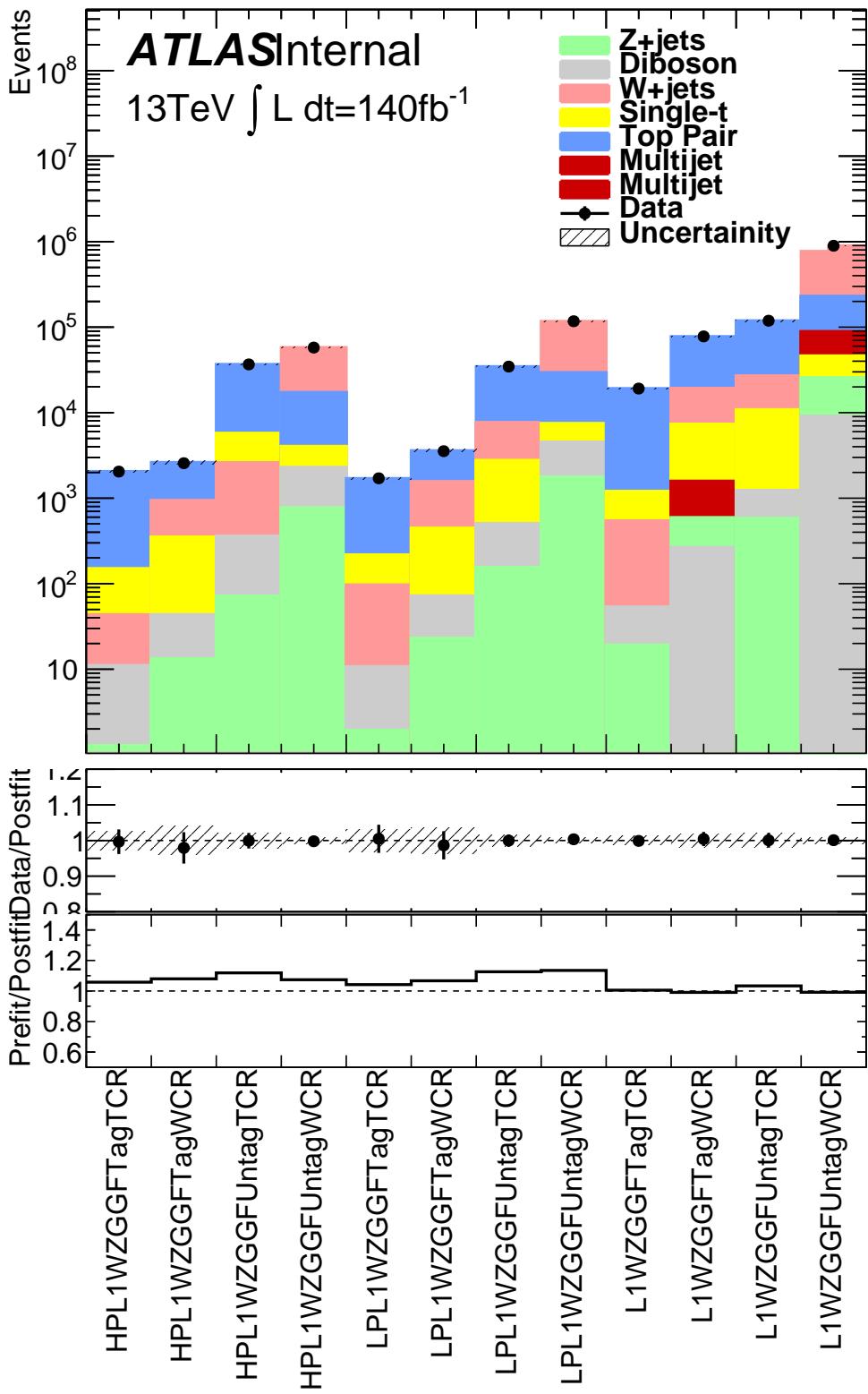


Figure 12.2: This figure shows the distribution of $m_{\ell\nu qq}$ in the DY WZ control regions.

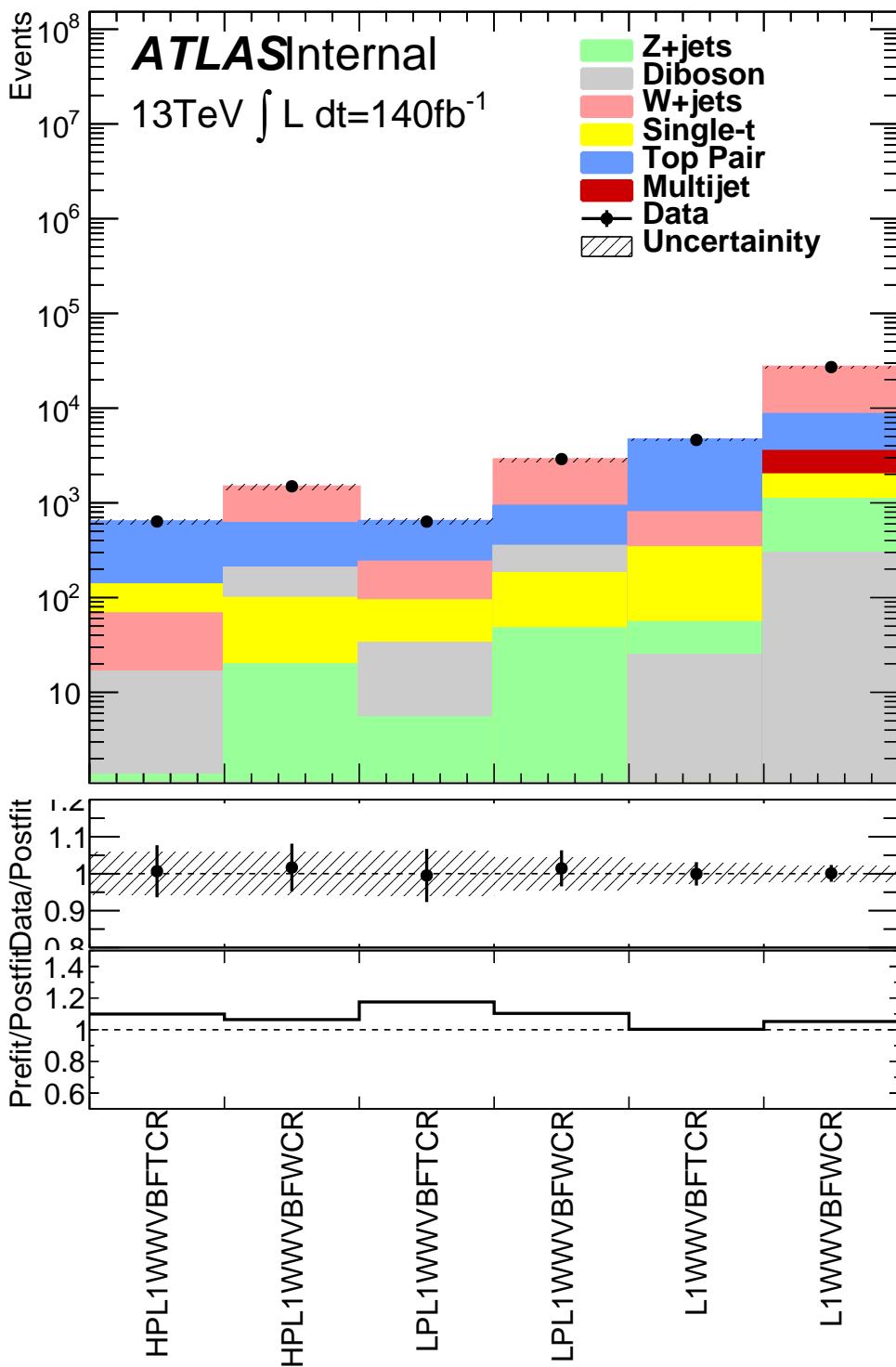


Figure 12.3: This figure shows the distribution of $m_{\ell\nu qq}$ in the VBF WW control regions.

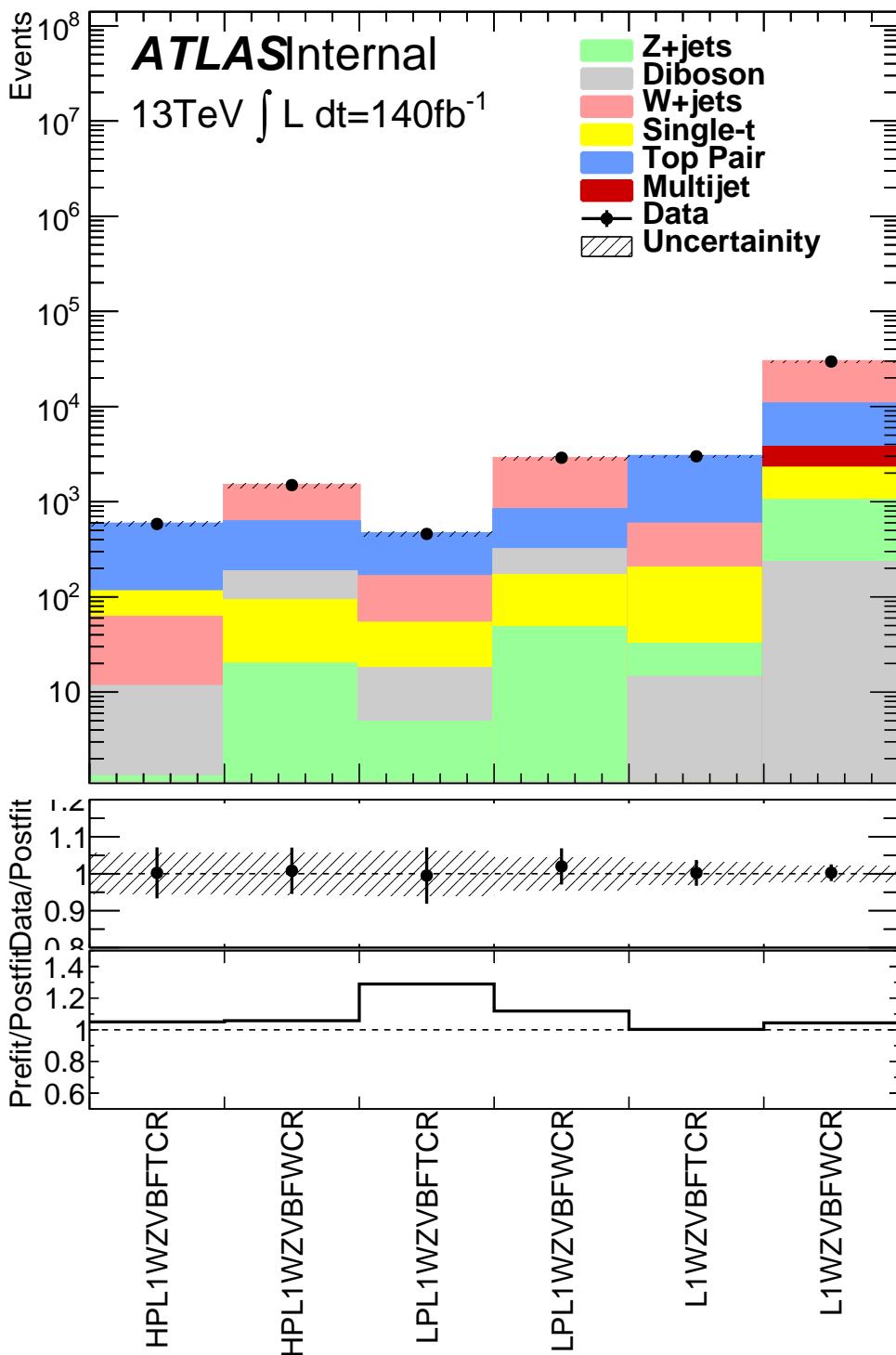


Figure 12.4: This figure shows the distribution of $m_{\ell\nu qq}$ in the VBF WZ control regions.

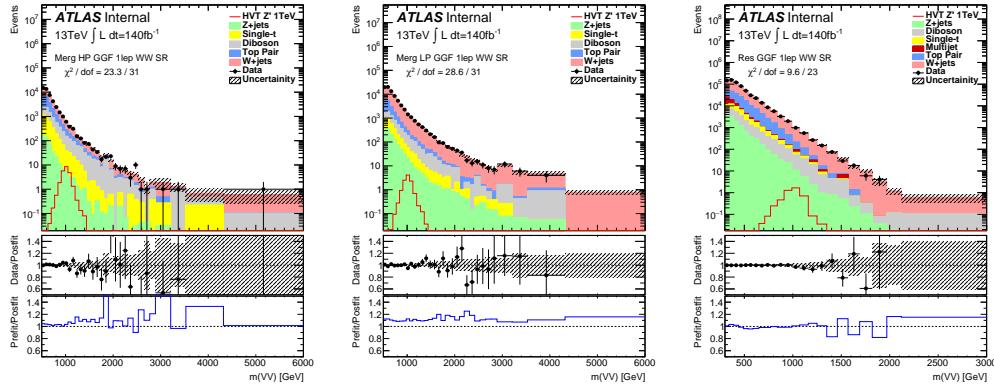


Figure 12.5: This figure shows the distribution of $m_{\ell\nu qq}$ in the GGF WW signal regions.

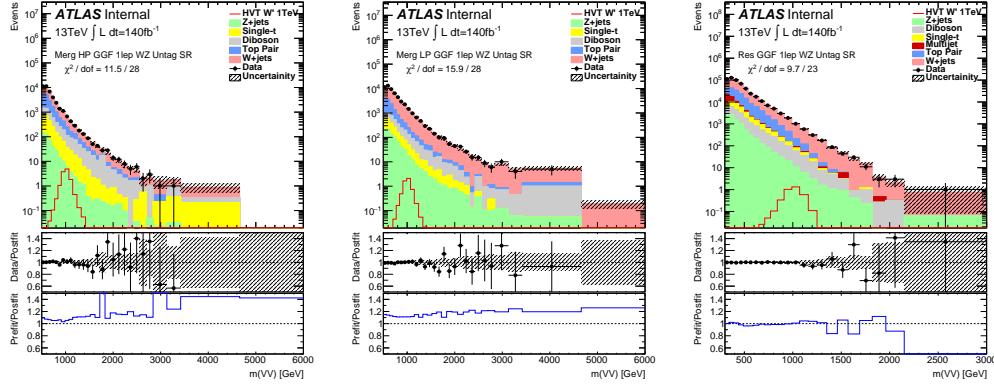


Figure 12.6: This figure shows the distribution of $m_{\ell\nu qq}$ in the GGF WZ Untag signal regions.

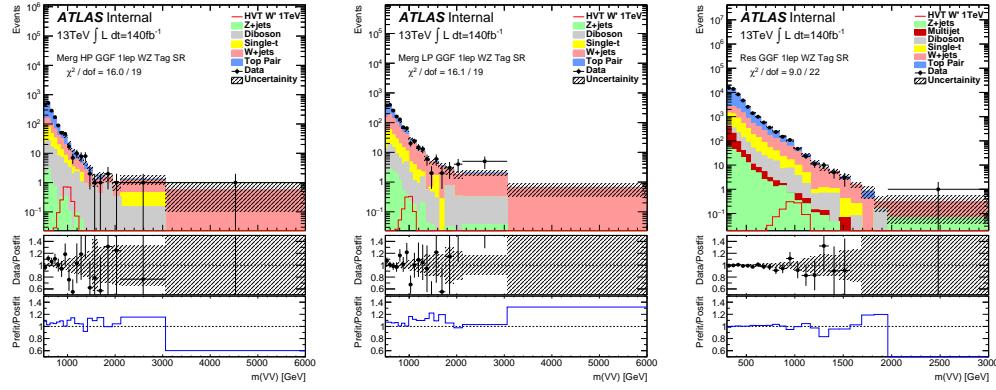


Figure 12.7: This figure shows the distribution of $m_{\ell\nu qq}$ in the GGF WZ Tag signal regions.

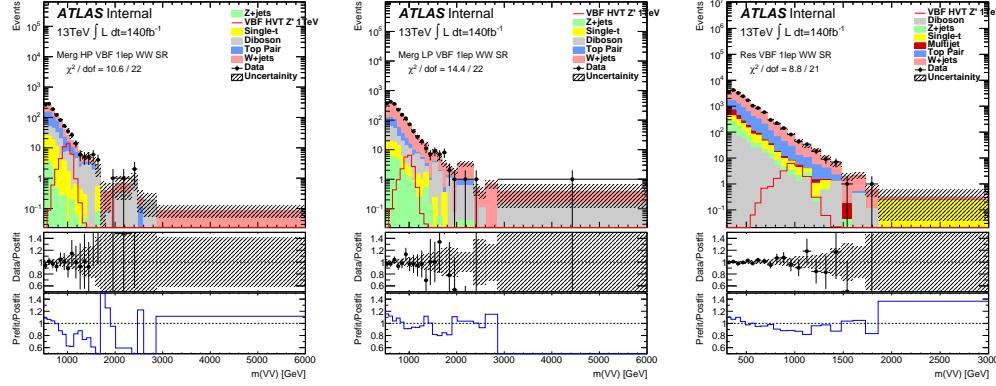


Figure 12.8: This figure shows the distribution of $m_{\ell\nu qq}$ in the VBF WZ Tag signal regions.

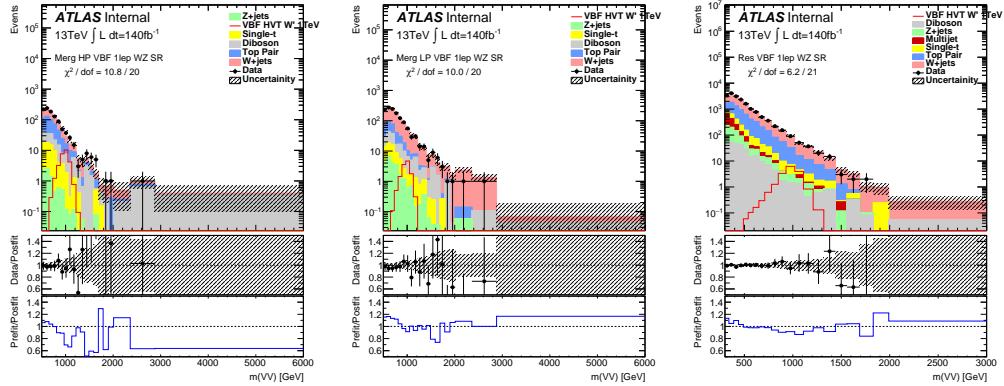


Figure 12.9: This figure shows the distribution of $m_{\ell\nu qq}$ in the VBF WZ Tag signal regions.

1276 12.2 Systematic Profiling and Correlations

1277 The ranked systematics and their fitted values are shown for the different
 1278 analysis regions in Figure 12.10 - 12.12. Note that background normalizations
 1279 for $W+jets$ and $t\bar{t}$ are left free to float in the fit. This means the nominal nor-
 1280 malization values are at one and the uncertainties are not plotted in the ranked
 1281 plots. Overall, systematics are not pulled outside their uncertainties, especially
 1282 for highly ranked nuisance parameters.

1283 The correlation between systematics are shown in Figures ???. Correlations
 1284 between background normalization are expected. The remaining systematic cor-
 1285 relations are not very strong or unexpected.

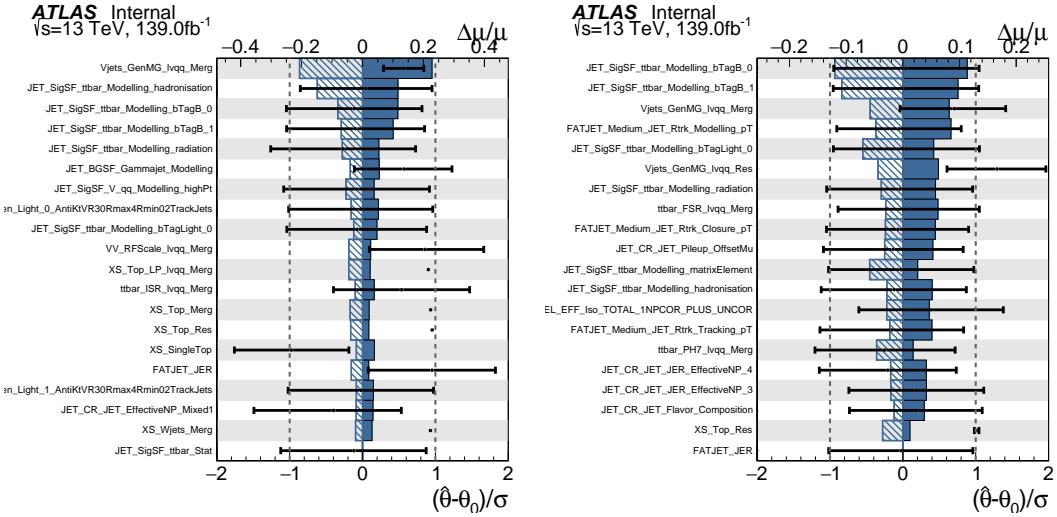


Figure 12.10: Ranked systematics and their fitted values for WW DY (right) and VBF (left) selections.

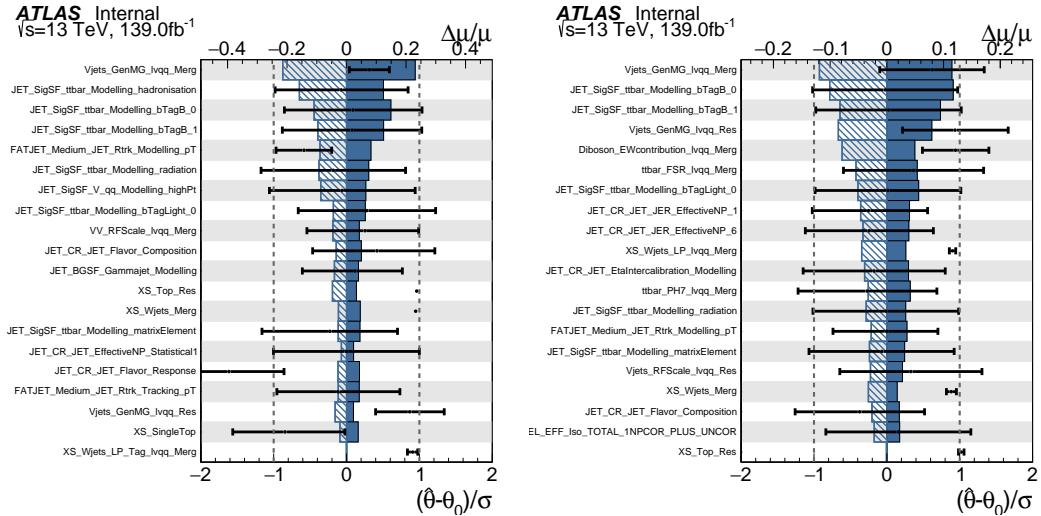


Figure 12.11: Ranked systematics and their fitted values for WZ DY (right) and VBF (left) selections.

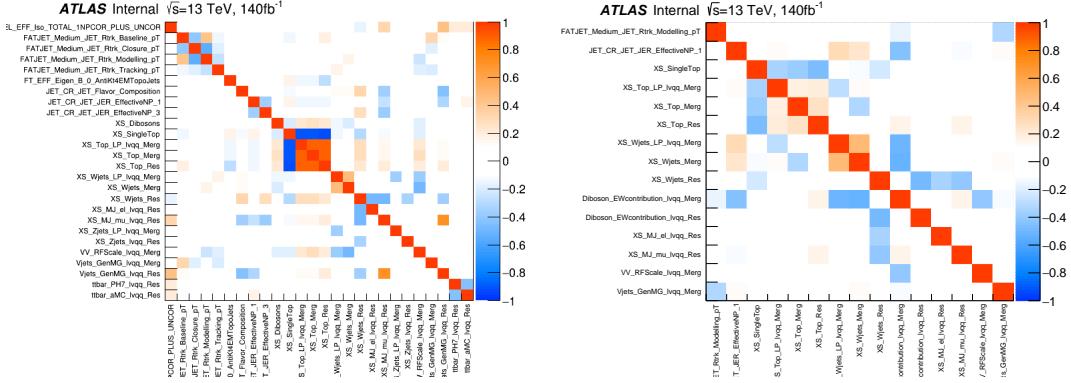


Figure 12.12: Correlations between systematics for WW DY (right) and VBF (left) selections.

12.3 Discovery Tests

To test for the existence of signal in the observed dataset, the discovery tests discussed earlier are used to calculate p-values as a function of resonance mass. The results of these tests are shown in Figures 12.13 - 12.17. Across the different DY signals the largest excesses are $\sim 2.2\sigma$ at 600 GeV and 1.8σ at 2 TeV. The largest excesses for VBF signals are $< 2.5\sigma$ at for 1 TeV resonances. As these deviations do not constitute discoveries, upper limits on μ are calculated.

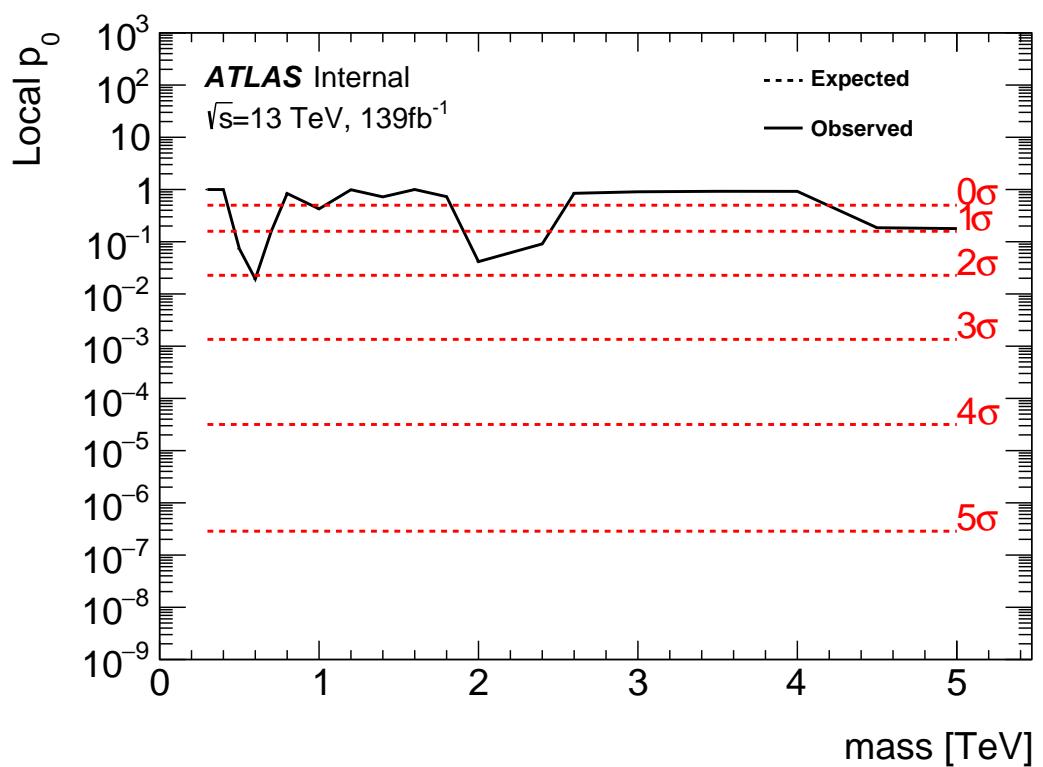


Figure 12.13: These plots show the measured p_0 value as a function of resonance mass for HVT Z' DY production.

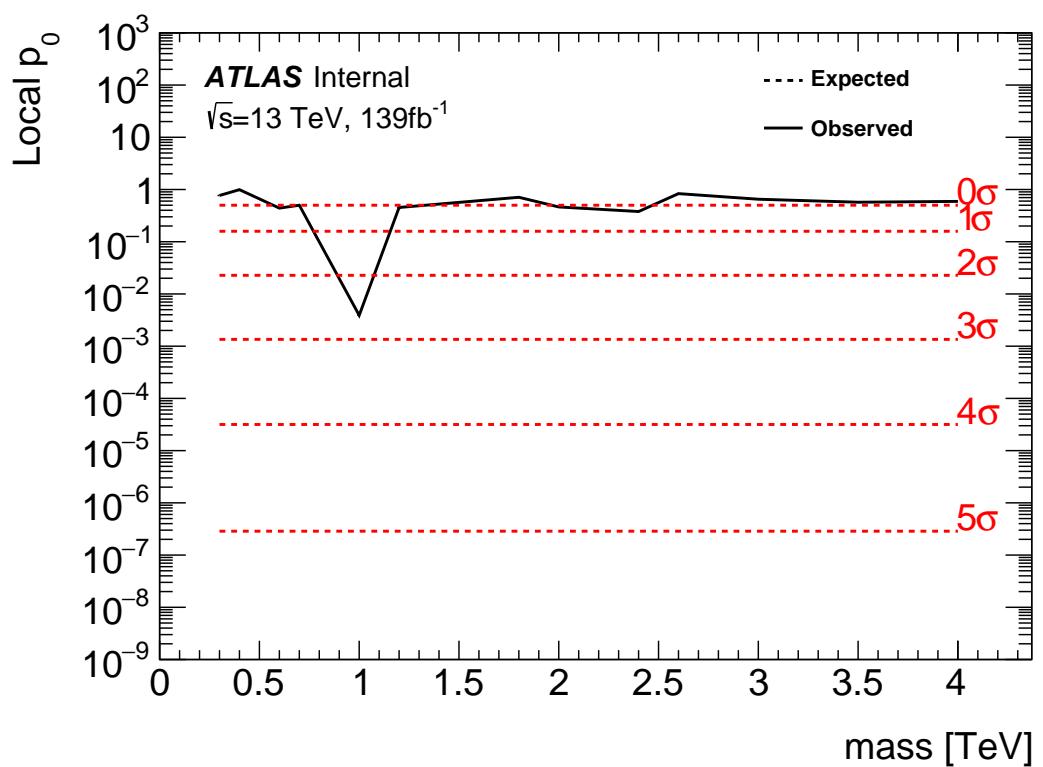


Figure 12.14: These plots show the measured p_0 value as a function of resonance mass for HVT Z' VBF production.

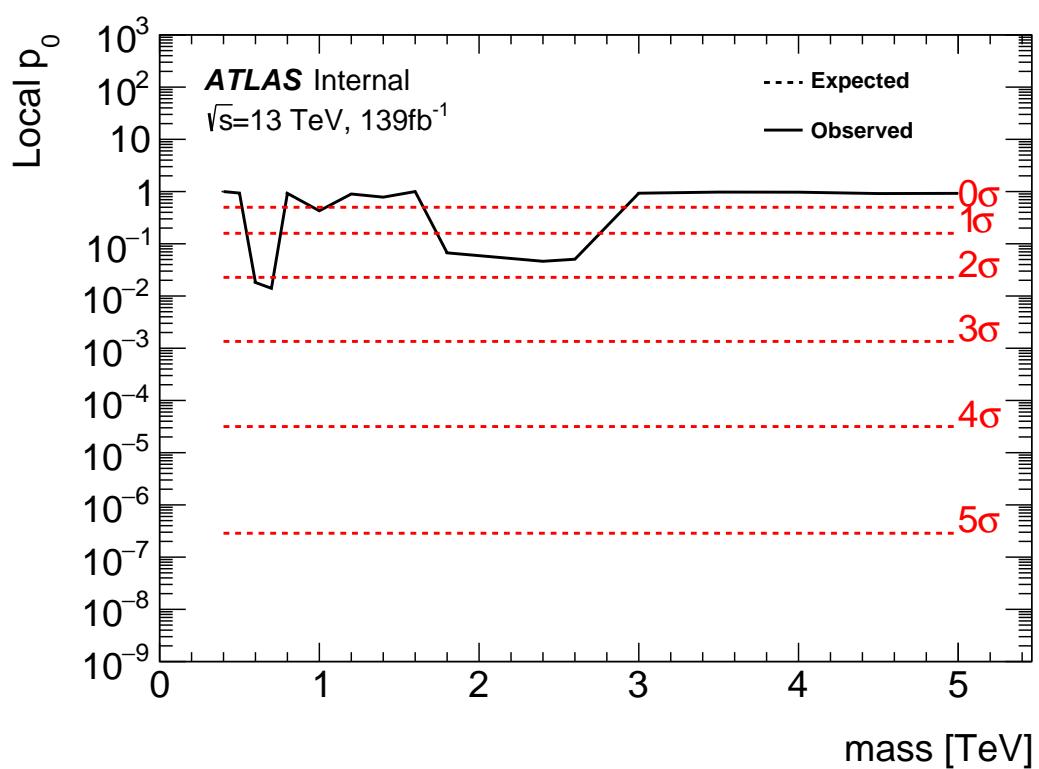


Figure 12.15: These plots show the measured p_0 value as a function of resonance mass for HVT W' DY production.

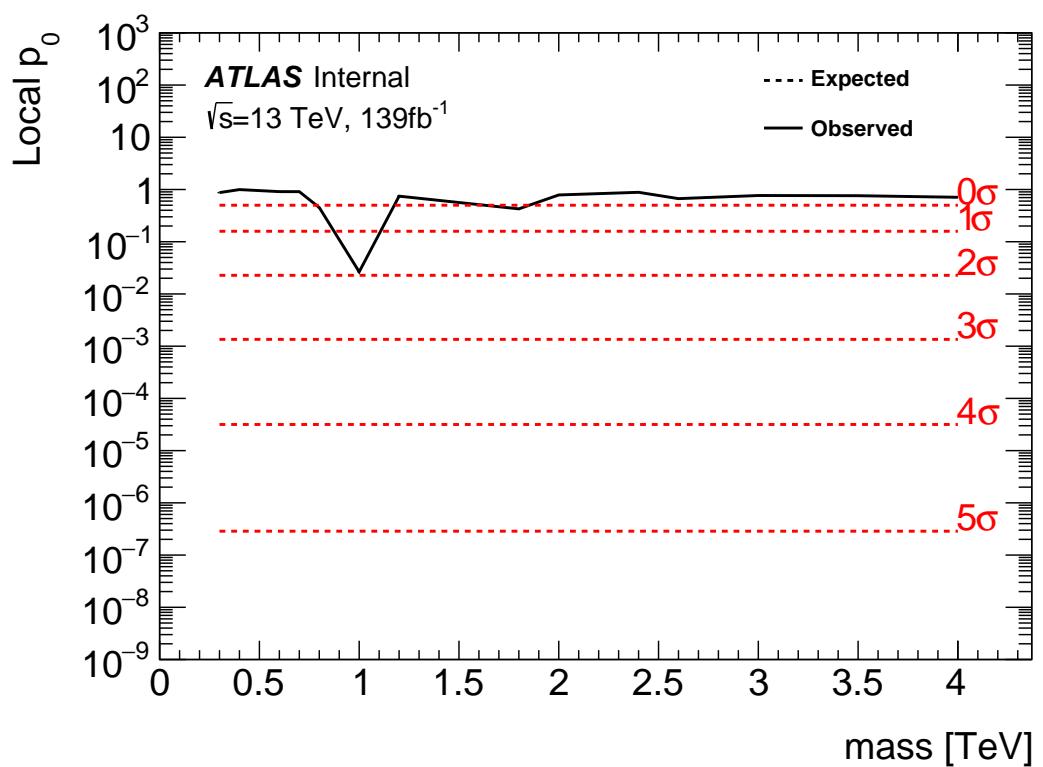


Figure 12.16: These plots show the measured p_0 value as a function of resonance mass for HVT W' VBF production.

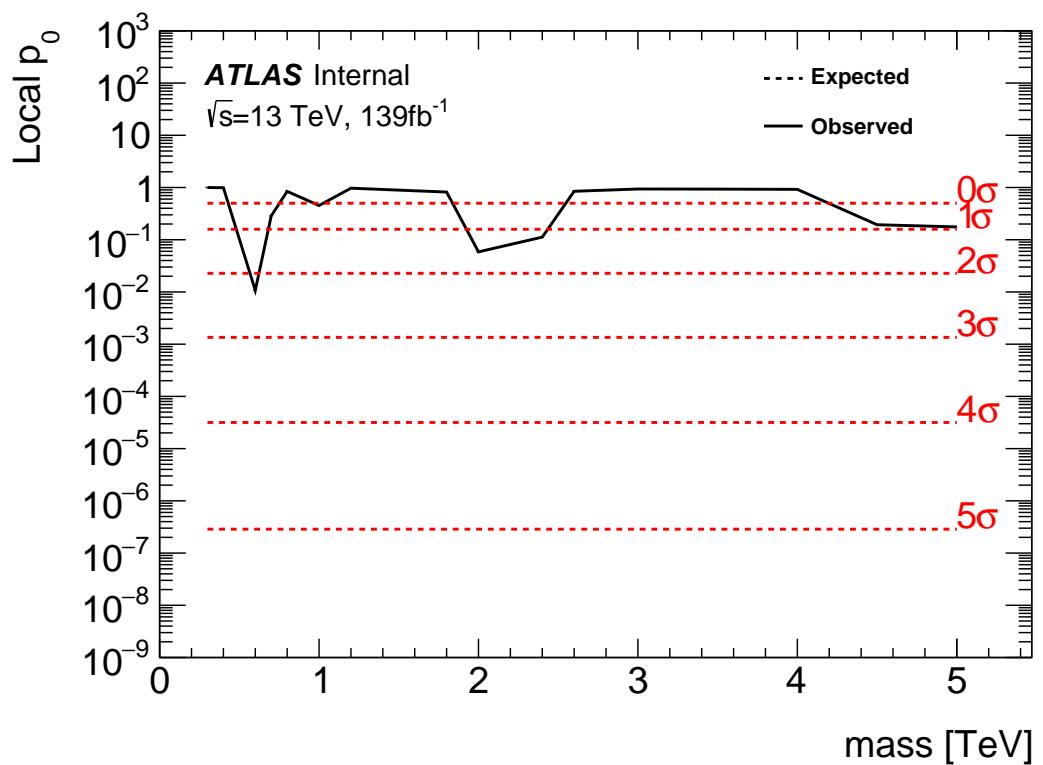


Figure 12.17: These plots show the measured p_0 value as a function of resonance mass for the RS Graviton DY production.

12.4 Limits

Using the exclusion limits tests discussed previously, exclusion limits are set on μ and consequently cross-sections for different signal models. Exclusion limits for the models considered are shown in Figure 12.18 - 12.20. These limits exclude HVT Model A $W' < 3.4\text{TeV}$ and $Z' < 3.3 \text{ TeV}$ and Model B $W' < 3.7 \text{ TeV}$ and $Z' < 3.7 \text{ TeV}$. Randall Sundrum Gravitons are excluded for masses below 1.6TeV .

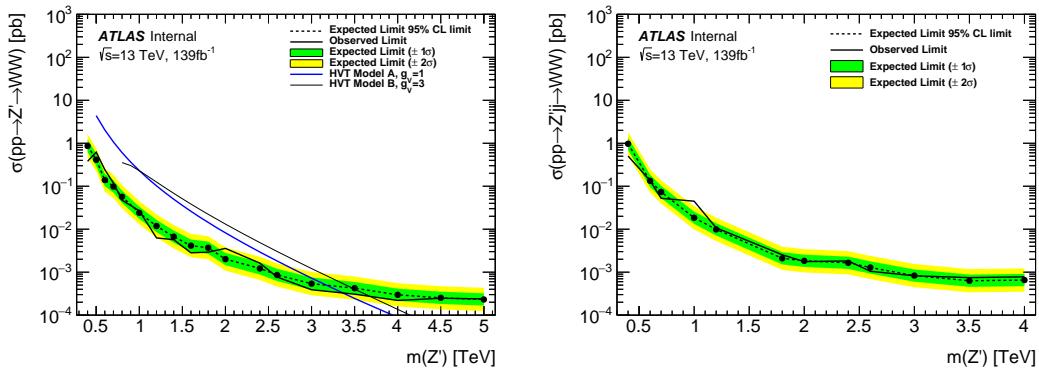


Figure 12.18: This figure shows theory, expected and observed limits for HVT W' DY (left) and VBF (right) production.

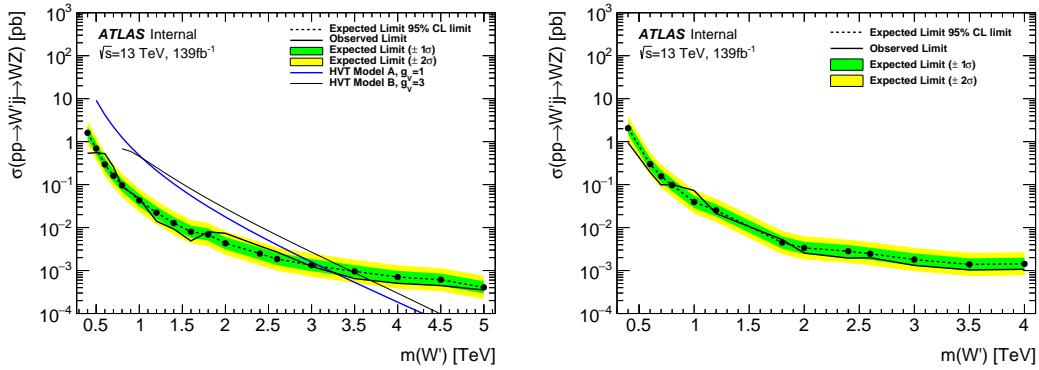


Figure 12.19: This figure shows theory, expected and observed limits for HVT Z' DY (left) and VBF (right) production.

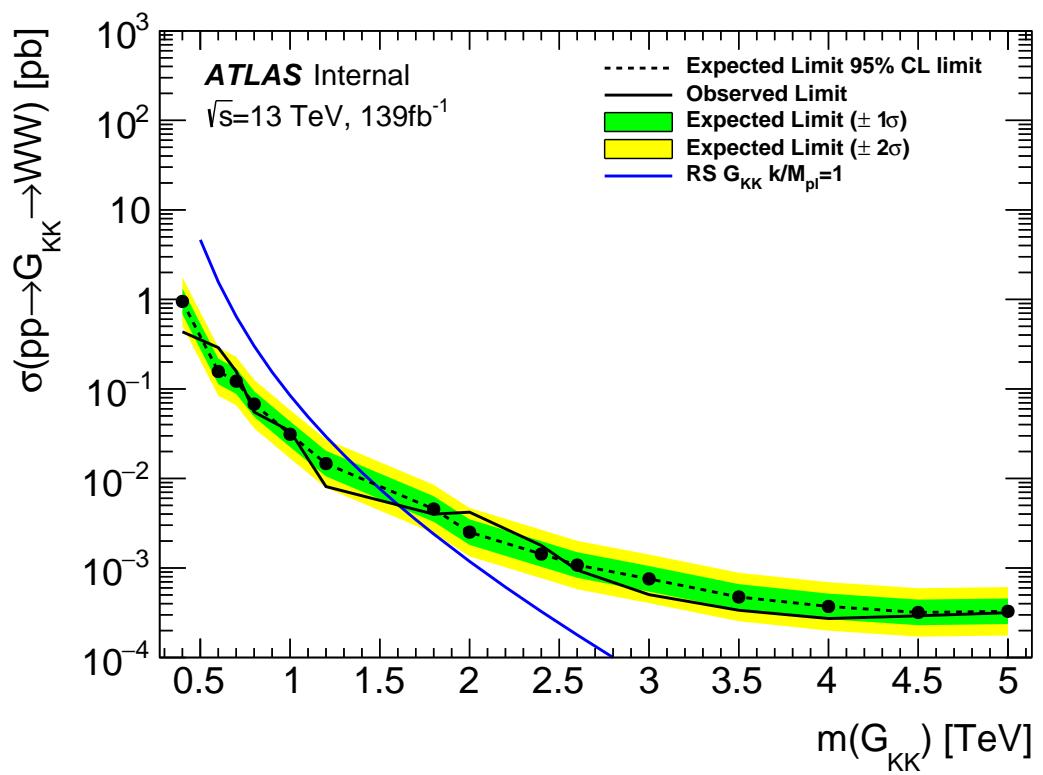


Figure 12.20: This figure shows theory, expected and observed limits for RS Gravitons via DY production.

Part V

1299

1300

Quark and Gluon Tagging

₁₃₀₁ **Chapter 13**

₁₃₀₂ **Prospects**

₁₃₀₃ For the resolved analysis, signal jets are quark enriched and background jets are
₁₃₀₄ gluon dominated. By classifying jets in the event as quark or gluon initiated, less
₁₃₀₅ background would contaminate the signal region. Figure 13.1 shows the PDGID
₁₃₀₆ for the truth parton matched to the jet (meaning the highest energy parton in
₁₃₀₇ the jet catchment area) in events passing the resolved signal region selections.
₁₃₀₈ PDGID = -1 corresponds to pileup jets, $0 < \text{PDGID} < 6$ correspond to quarks
₁₃₀₉ and $\text{PDGID} = 21$ corresponds to gluons. From this Figure, it is evident that a
₁₃₁₀ notable fraction of the background that contaminates the signal region contains
₁₃₁₁ gluon jets, especially for the sub-leading jet.

₁₃₁₂ As gluons jets have more constituents and therefore more tracks (n_{trk}), the
₁₃₁₃ background jets have more tracks than the signal jets. This is shown in Fig-
₁₃₁₄ ure 13.2. Therefore, by cutting on the number of tracks in a jet, quark and gluon
₁₃₁₅ jets may be distinguished (i.e. jets with less than a given number of tracks are
₁₃₁₆ classified as a quark, otherwise the jet is classified as a gluon.) Moreover, as the
₁₃₁₇ momentum of the jet increases the number of tracks also increases logarithmically
₁₃₁₈ [Cite nachman thesis Natasha]. Therefore by applying a cut on the number of
₁₃₁₉ tracks that scales with the $\ln(p_T)$ is more powerful than a threshold cut on the

1320 number of tracks. Figure 13.3-Figure 13.6 show normalized heat maps of $\ln(p_T)$
1321 vs the number of reconstructed tracks for the background and a 300 GeV Z' signal.
1322 In these plots it is evident that the number of tracks in the background jets grows
1323 more quickly with $\ln(p_T)$ than for the signal jets. This is expected given that the
1324 signal is quark dominated and the background is gluon dominated.

1325 In Figure 13.8 is the ROC Curve for quark gluon tagging with cut on the
1326 number of tracks in a jet that depends on $\ln(p_T)$. The sum of the backgrounds in
1327 the signal region were used for this curve. Here the quark tagging efficiency is the
1328 ratio of quarks tagged as quarks to the total number of quarks in the signal region.
1329 The gluon rejection is calculated as the reciprocal of the gluon tagging efficiency.
1330 Choosing a 90% efficient working point with a rejection of 1.4 corresponds to a
1331 slope of 4 and intercept of -5. Tagging both jets in this analysis would yield an
1332 efficiency of $90\%^{n_{jets}}$. Focusing on the background in Figure 13.9, this cut helps
1333 minimize gluon contamination in the signal region. Also, from these heat maps it
1334 is obvious that the number of tracks in gluon jets grows more quickly than those
1335 in quark jets.

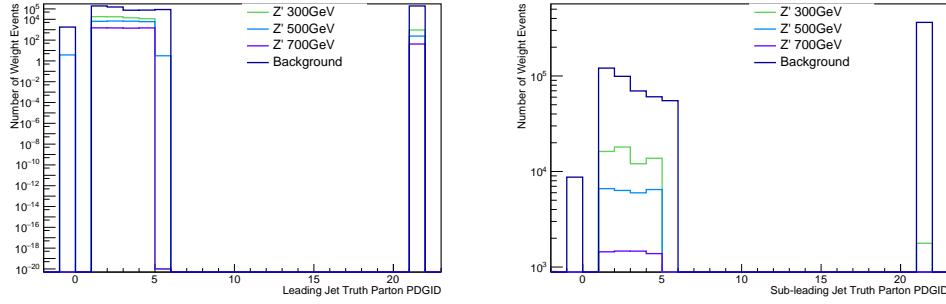


Figure 13.1: PDGID of the truth-level parton matched to the small-R jets passing the Resolved GGF WW Signal Region selections for the (a) Leading (b) Sub-Leading jets . These distributions are shown for 300, 500, and 700GeV Z' signals and the background.

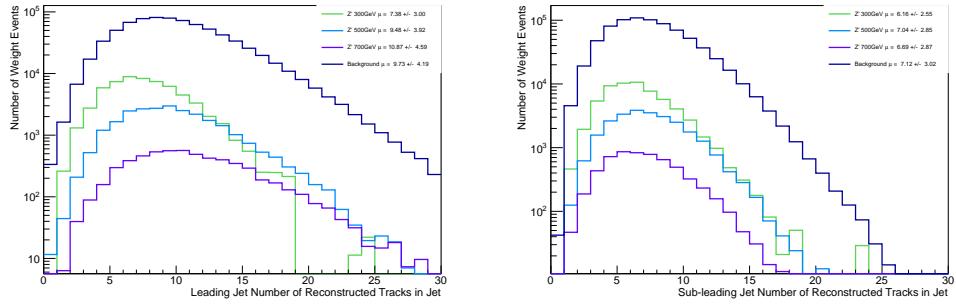


Figure 13.2: The number of tracks in small-R jets in events passing the Resolved GGF WW Signal Region selections for the (a) Leading (b) Sub-Leading jets. These distributions are shown for 300, 500, and 700GeV Z' signals and the background.

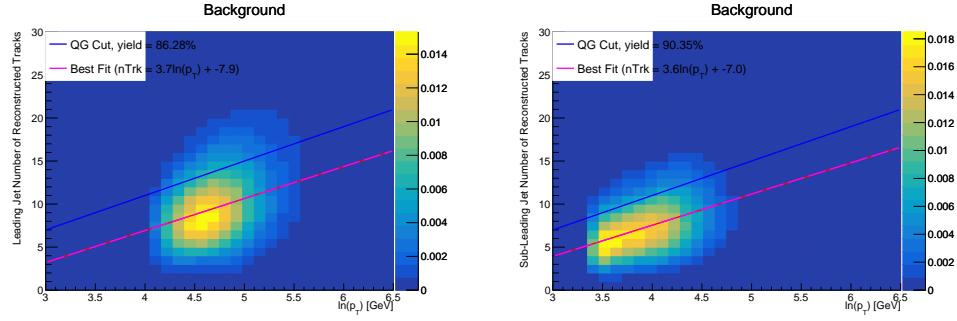


Figure 13.3: The number of tracks in background small-R jets in events passing the Resolved GGF WW Signal region selection vs. $\ln(p_T)$ for (a)Leading (b) Sub-Leading jets. The best fit line for the distribution is also shown, as well as the percentage of jets that pass a cut of number of tracks $< 4 \times \ln(p_T) - 5$. Note the number of total entries in these plots has been normalized to one.

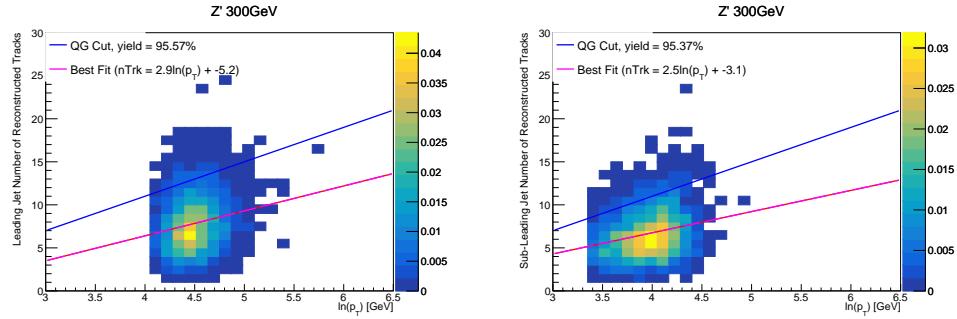


Figure 13.4: The number of tracks in small-R jets in 300GeV Z' events passing the Resolved GGF WW Signal region selection vs. $\ln(p_T)$ for (a)Leading (b) Sub-Leading jets. The best fit line for the distribution is also shown, as well as the percentage of jets that pass a cut of number of tracks $< 4 \times \ln(p_T) - 5$. Note the number of total entries in these plots has been normalized to one.

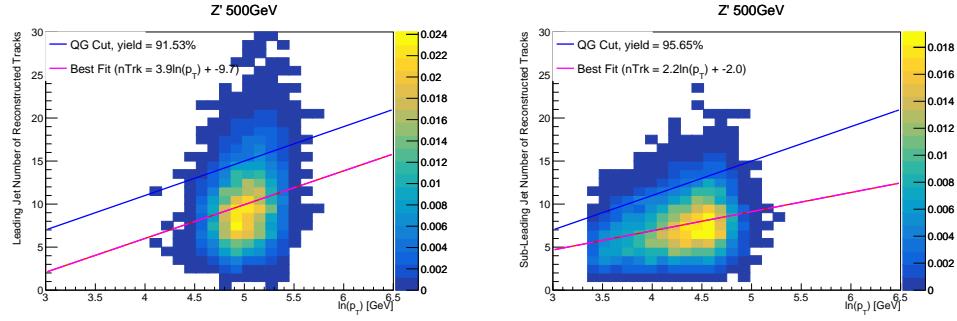


Figure 13.5: The number of tracks in small-R jets in 500GeV Z' events passing the Resolved GGF WW Signal region selection vs. $\ln(p_T)$ for (a)Leading (b) Sub-Leading jets. The best fit line for the distribution is also shown, as well as the percentage of jets that pass a cut of number of tracks $< 4 \times \ln(p_T) - 5$. Note the number of total entries in these plots has been normalized to one.

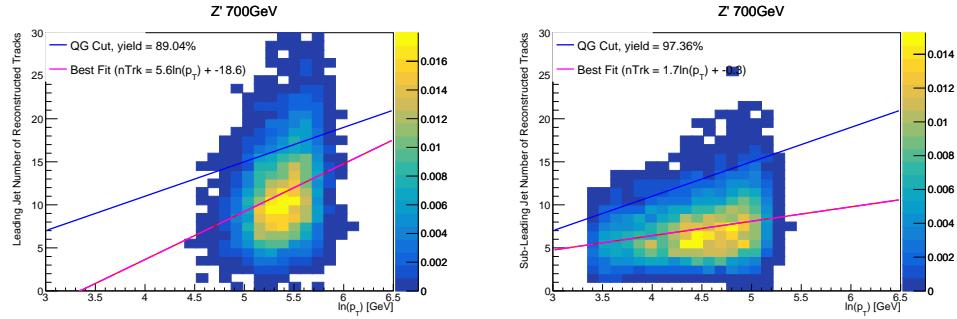


Figure 13.6: The number of tracks in small-R jets in 700GeV Z' events passing the Resolved GGF WW Signal region selection vs. $\ln(p_T)$ for (a)Leading (b) Sub-Leading jets. The best fit line for the distribution is also shown, as well as the percentage of jets that pass a cut of number of tracks $< 4 \times \ln(p_T) - 5$. Note the number of total entries in these plots has been normalized to one.

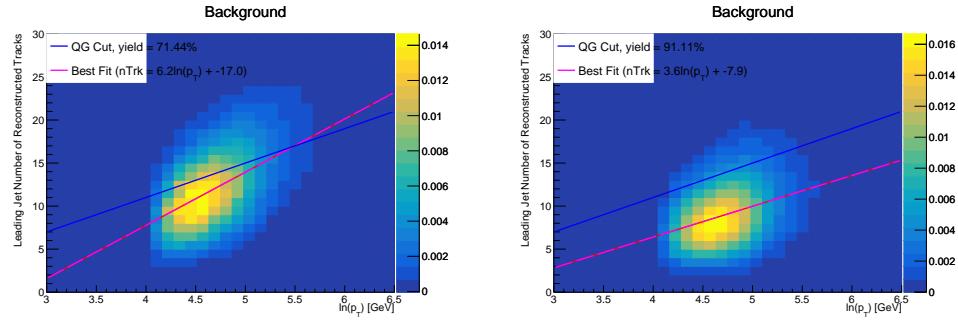


Figure 13.7: The number of tracks in leading small-R jets in background events passing the Resolved GGF WW Signal region selection vs. $\ln(p_T)$ for (a)Gluons (b) Quarks jets. The best fit line for the distribution is also shown, as well as the percentage of jets that pass a cut of number of tracks $< 4 \times \ln(p_T) - 5$.Note the number of total entries in these plots has been normalized to one.

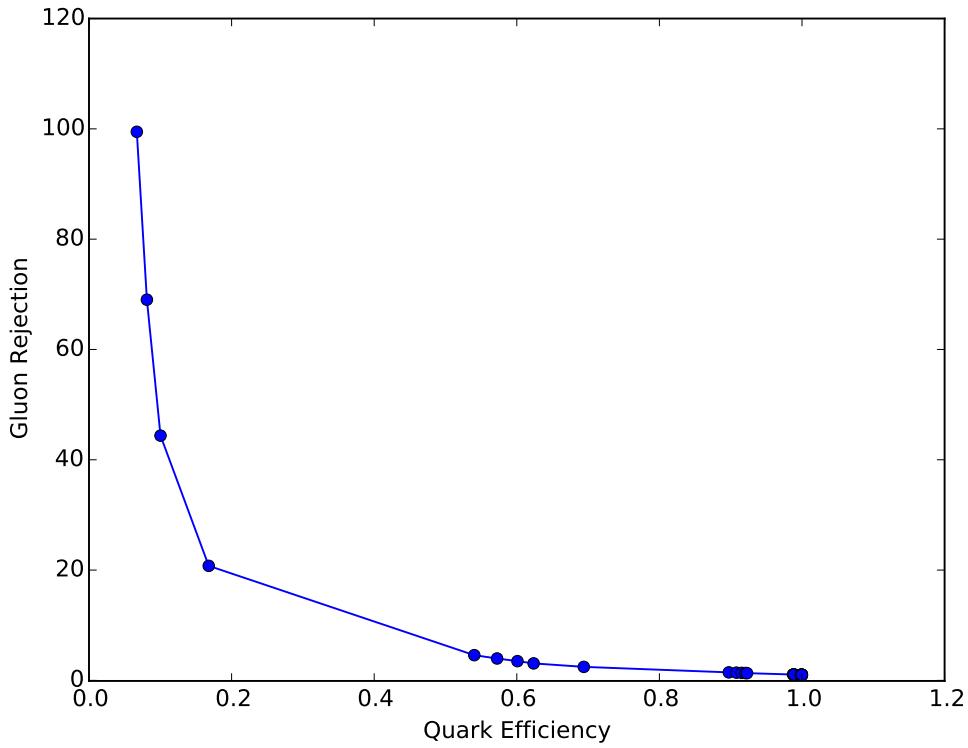


Figure 13.8: ROC Curve for Quark and Gluon Tagging with a cut on the number of tracks that depends on the $\ln(p_T)$.

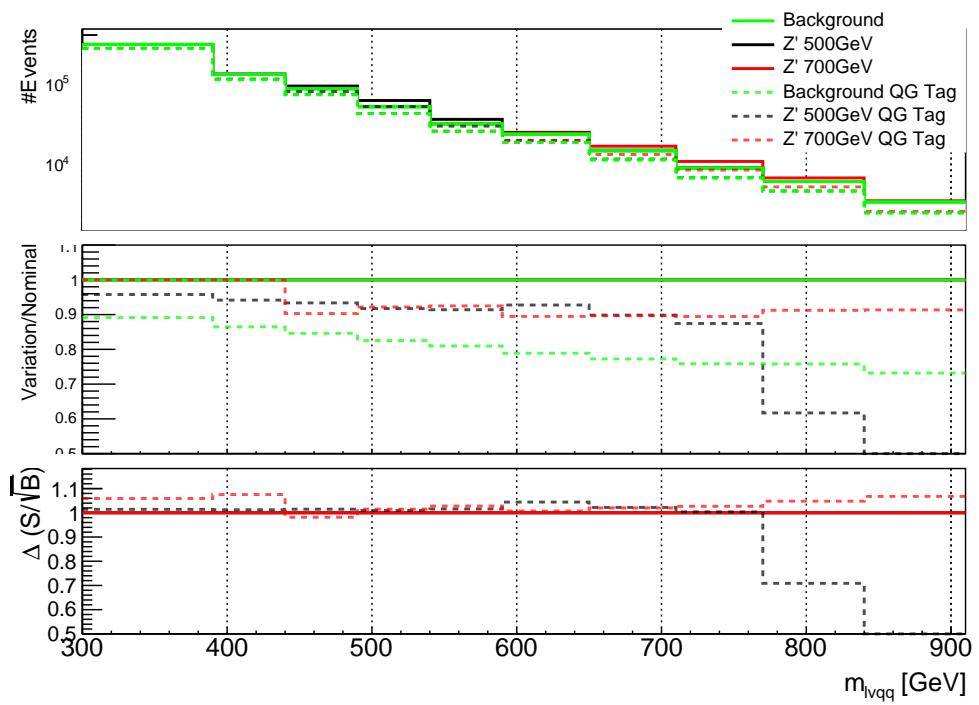


Figure 13.9: The top panel shows the distribution of m_{lvqq} with and without quark gluon tagging. The middle panel shows the ratio of the signals and backgrounds with and without quark gluon tagging. The bottom panel shows the change in S/\sqrt{B} with quark gluon tagging.

1336 Chapter 14

1337 n_{trk} Calibration

1338 As tagger based on nTrk cuts on the number of tracks in jets, a quantity that
1339 is not known with infinite precision, means that relevant systematic uncertainties
1340 must be evaluated. The sources of uncertainty in n_{trk} may be split into modeling
1341 and experimental uncertainties.

1342 Modeling uncertainties are obtained by assessing PDF and ME uncertainties
1343 on the number of charged particles in particle-level jets in dijet events. The
1344 number of charged particles as a function of jet p_T is calculated using an Iterative
1345 Bayesian (IB) technique [cite paper].

1346 This measurement ([9]) uses the ATLAS 2012 pp collision dataset, correspond-
1347 ing to $20.3/\text{fb}$ at center-of-mass energy $\sqrt{s} = 8\text{TeV}$. The number of charged con-
1348 stituents depends on fragmentation modeling and matrix elements, which do not
1349 depend on s . For this reason, it is safe to use these uncertainties for $\text{sqrt}(s)=13\text{TeV}$.
1350 Monte Carlo (MC) samples are used to determine the response matrix. The MC
1351 sample is a dijet sample generated with Pythia 8.175 using CT10 PDF and AU2
1352 tune. The anti- k_T algorithm is used to cluster jets with a radius parameter R
1353 $= 0.4$. Jets are required to have $|\eta| < 2.1$. Tracks in jets are required to have
1354 $p_T > 500\text{MeV}$, $|\eta| < 2.5$, track-fit $\chi^2 < 3.0$ and originate from the primary ver-

tex. Matching tracks to jets is accomplished using ghost-association [cite]. In this technique, jets are re-clustered with the track collection augmented with "ghost" versions of tracks. These "ghosts" tracks have the same direction as their parent track, but infinitesimal track p_T . This insures meta-jet properties (e.g. η , p_T , etc) are unchanged. A track is matched to a jet if it's ghost version remains in the jet after re-clustering. Further details of the data, object, and event selection may be found in [cite 35].

To select dijet topologies events are required to have at least two jets with $p_T > 50GeV$ that are relatively well-balanced ($p_T^{lead}/p_T^{sub-lead} < 1.5$).

In the IB technique, the prior distribution and number of iterations are the inputs [cite Bayesian paper]. The IB response matrix connects number of charged particles to the number of tracks in jets determined using the simulated samples. This response matrix is used to unfold data to extract the n_c . Before applying the response matrix a fake factor is applied. This accounts for jets that pass detector level selections, but not particle level selections. Following this, the IB method iteratively applies the response matrix using the nominal Pythia 8.175 sample as a prior. The number of IB iterations is chosen to minimize unfolding bias and statistical fluctuations. For this measurement four iterations was found to be optimal by minimizing the unfolding bias from pseudodata simulated with Herwig++ with a prior from Pythia 8 AU2. Finally, the inefficiency factor is applied to account for events passing particle level selection but not detector level, yielding the unfolded nCharged distribution.

This process is prone to three main sources of bias: response matrix, correction factor, and unfolding procedure uncertainties. The response matrix is sensitive to experimental uncertainties impacting jet track reconstruction and calorimeter jet p_T . Correction factors are also sensitive to experimental uncertainties (e.g. JES)

1381 as such uncertainties modify detector level acceptance. Sensitivity to particle
 1382 level acceptance is calculated by comparing Pythia and Herwig. Finally, the bias
 1383 from the IB prior choice is determined by reweighting the particle-level spectrum,
 1384 so the simulated detector level spectrum more closely matches the uncorrected
 1385 data. Unfolding this modified detector-level simulation and comparing it to the re-
 1386 weighted particle-level spectrum indicates bias from the prior distribution choice.
 1387

A summary of all the systematic uncertainties associated with this unfolding
 1388 may be found in [ref paper]. Total uncertainties are < 7% for the number of
 1389 charged particles in jets. The unfolded distribution of the nCharged in jets from
 1390 data are further analyzed to extract the quark and gluon nCharged distributions.
 1391 In dijet events, the jet with a larger η is more energetic and therefore more likely
 1392 to be a quark. This is due to the quarks in protons generally having a larger
 1393 fraction of the total momentum of the proton constituents. The more central jet
 1394 is more likely to be a gluon-initiated jet. This correlation between jet η and flavor
 1395 may then be used to extract nCharged in p_T bins using:
 1396

$$\langle n_c^f \rangle = f_q^f \langle n_c^q \rangle + f_g^f \langle n_c^g \rangle \quad (14.1)$$

$$\langle n_c^c \rangle = f_q^c \langle n_c^q \rangle + f_g^c \langle n_c^g \rangle \quad (14.2)$$

1397 In this equation the f and c subscripts denote the more forward and central
 1398 jets, respectively. The q and g subscripts denote quark and gluon. The fraction
 1399 of more forward jets that are say gluons is denoted by f_g^f . The other relevant jet
 1400 fractions are denoted with the same naming scheme. Finally, $\langle n_c \rangle$ is the average
 1401 number of charged particles in a jet in a given p_T bin. To show that Eq. (??) may
 1402 be used to extract quark and gluon n_c distributions the extracted distributions
 1403 are compared to n_c distributions determined using the jet flavor in simulation.
 1404 Figure [add figure natasha] shows that the extracted and true distributions differ

1405 by < 1% over the p_T range probed for this study. Moreover, this implies that n_c
1406 depends only on the flavor of the initiating parton and jet p_T .

1407 These extracted distributions are prone to PDF and ME biases. The bias from
1408 the choice of the CT10 PDF for the Pythia sample is accounted for by comparing
1409 quark/gluon fractions for the nominal CT10 sample with its eigenvector variations.
1410 Comparing the quark/gluon fractions from Pythia 8 and Herwig++ quantify the
1411 uncertainty from the ME calculation. These uncertainties are added in quadra-
1412 ture with the unfolding uncertainty to give the total modelling uncertainty on
1413 the extracted n_c distribution. This is shown in Figure 15.2.

1414 To apply these uncertainties in n_c distributions in data, per-jet event weights
1415 are associated with each uncertainty according to:

$$w_i(n_c) = \frac{P(n_c | n_c > \pm \sigma_{n_c}^i)}{P(n_c | n_c >)} \quad (14.3)$$

1416 In Eq. (??), i denotes the uncertainty considered, P is the Poisson probability,
1417 and $\sigma_{n_c}^i$ represents the average impact of the uncertainty on n_c .

1418 The previous uncertainties described accounted for modeling uncertainty as-
1419 sociated with the number of charged particles in a jet. However, n_c is not a
1420 measurable quantity. Instead the number of tracks in a jet is measured, which is
1421 a proxy for n_c . Therefore the uncertainties associated with the measurement of
1422 nTracks must also be considered ([11]). These uncertainties were calculated using
1423 a Pythia 8 dijet sample with NNPDF 23 and Run 2 data. Track reconstruction
1424 efficiency and fake rates are the dominant sources of nTrack uncertainties.

1425 The track reconstruction efficiency is affected by the uncertainty of the de-
1426 scription of the ID material in simulation and the modeling of charged-particle
1427 interactions with this material. These uncertainties are accounted for by varying
1428 the ID material by 5-25% (dependent on the region of the detector considered).

1429 The difference in the tracking efficiency between the nominal and varied simula-
1430 tion give the uncertainty on the track reconstruction efficiency. Another important
1431 source of track reconstruction inefficiency arises in the core of jets. The high den-
1432 sity of tracks in the jet cores can cause ID clusters to merge. The fraction of lost
1433 tracks due to merging is given by the fraction of tracks that have a charge of two
1434 minimum ionizing particles. This quantity is compared between data and simu-
1435 lation resulting in an uncertainty of 0.4% on tracks with $\Delta R < 0.1$. Combining
1436 these effects gives a total uncertainty as a function of p_T and η that is generally
1437 $< 2\%$ [references figure 44 from [11]].

1438 Fake tracks are the other dominant source of nTrk uncertainty. Fake tracks
1439 are tracks that cannot be associated to a single particle. Often these tracks are a
1440 result of random combinations of hits from charged particles that overlap in space.
1441 In dense environments, such as the core of jets or high-pileup environments, fake
1442 tracks are more likely. Fake tracks are estimated with a 'control region method'
1443 which is briefly summarized here [[10]]. By applying a series of track selections
1444 to enrich the fraction of fake tracks (e.g. $|d_0| > 0.1$, track $\chi^2 > 1.4$, etc) in
1445 simulation, templates for fake track parameters are calculated. These templates
1446 are then fit to data to determine the fraction of fake tracks. On average the fake
1447 rate is found to be 30% (independent of p_T and η).

1448 To assess the impact of these two detector level uncertainties, tracks are ran-
1449 domly dropped according to the rates described above. Reconstruction and fake
1450 uncertainties both lower the number of tracks, hence these uncertainties are one-
1451 sided. By dropping tracks in this way a varied nTrk distribution is calculated for
1452 both uncertainties. The associated per-jet event weights are then calculated in
1453 the same way as the modeling weights as:

$$w_i(n_c) = \frac{P(n_{trk} | < n_{trk} > \pm \sigma_{n_{trk}}^i)}{P(n_{trk} | < n_{trk} >)} \quad (14.4)$$

1454 Adding the modeling and detector level uncertainties in quadrature gives the
 1455 overall nTrack uncertainty. The effects of the individual uncertainties on the nTrk
 1456 distributions can be seen in Fig 15.4. Fig 15.3 shows the m_{lvqq} and nTrk distri-
 1457 butions for the W and Top control regions before likelihood fitting. In these plots
 1458 the nTrk uncertainties improve agreement between data and MC. The remaining
 1459 differences are likely covered by likelihood fitting and improving the analysis itself.

1460 **Chapter 15**

1461 **Application**

1462 Using the 90% WP of the n_{trk} tagger improves S/\sqrt{B} is $\sim 3\%$ as shown in
1463 Figure 13.9. Although, n_{trk} is the single most powerful discriminating variable
1464 for quark and gluon jets, the addition of other jet variables would improve the
1465 classification efficiency. Figure 15.1 shows the possible improvement of 10%
1466 in jet classification using the truth label of the jets to classify jets. This type of
1467 improvement is possible by using variables such as jet width, and energy correlata-
1468 tors. Figure [add BDT figure/use 1612.01551.pdf] shows for a 90% quark tagging
1469 efficiency for a 100 GeV jet, a BDT improve the gluon rejection by 0.4. Once this
1470 tagger is calibrated it would improve the analysis sensitivity of this channel.

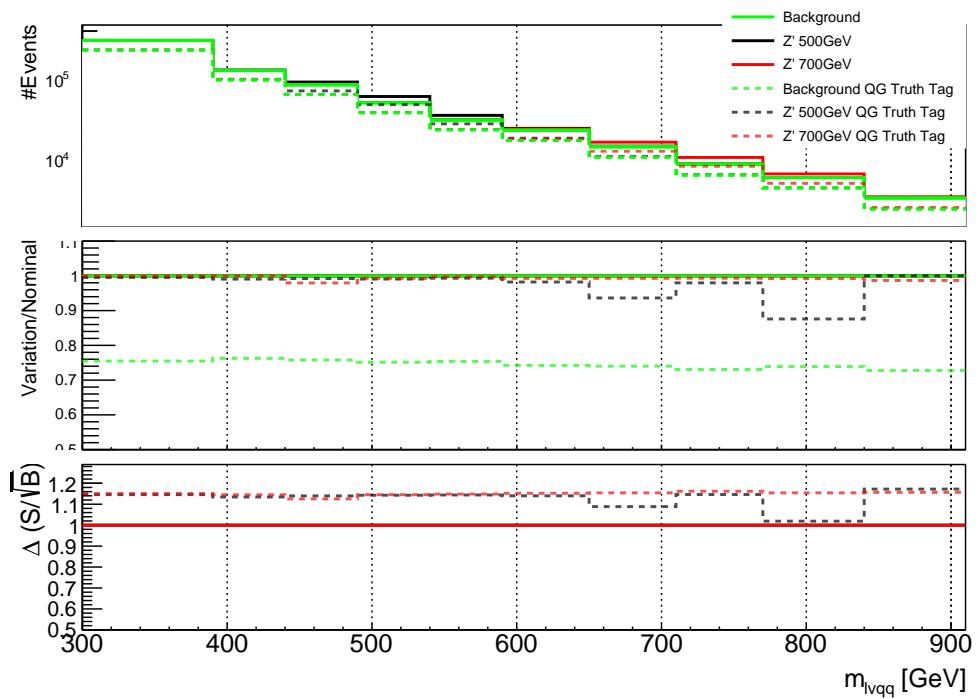


Figure 15.1: The top panel shows the distribution of m_{lvqq} with and without requiring jets to be true quarks. The middle panel shows the ratio of the signals and backgrounds with and without requiring jets to be true quarks. The bottom panel shows the change in S/\sqrt{B} when requiring jets to be true quarks..

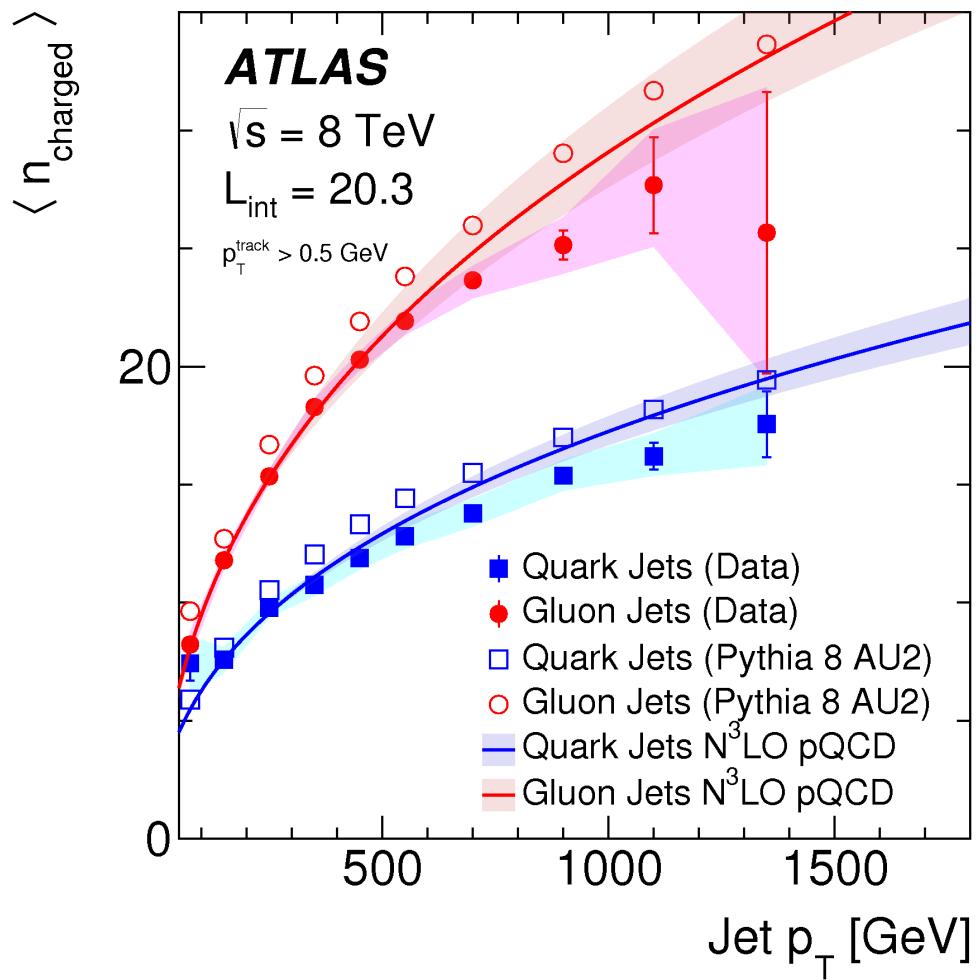


Figure 15.2: Unfolded and extracted n_C qg dstbs..

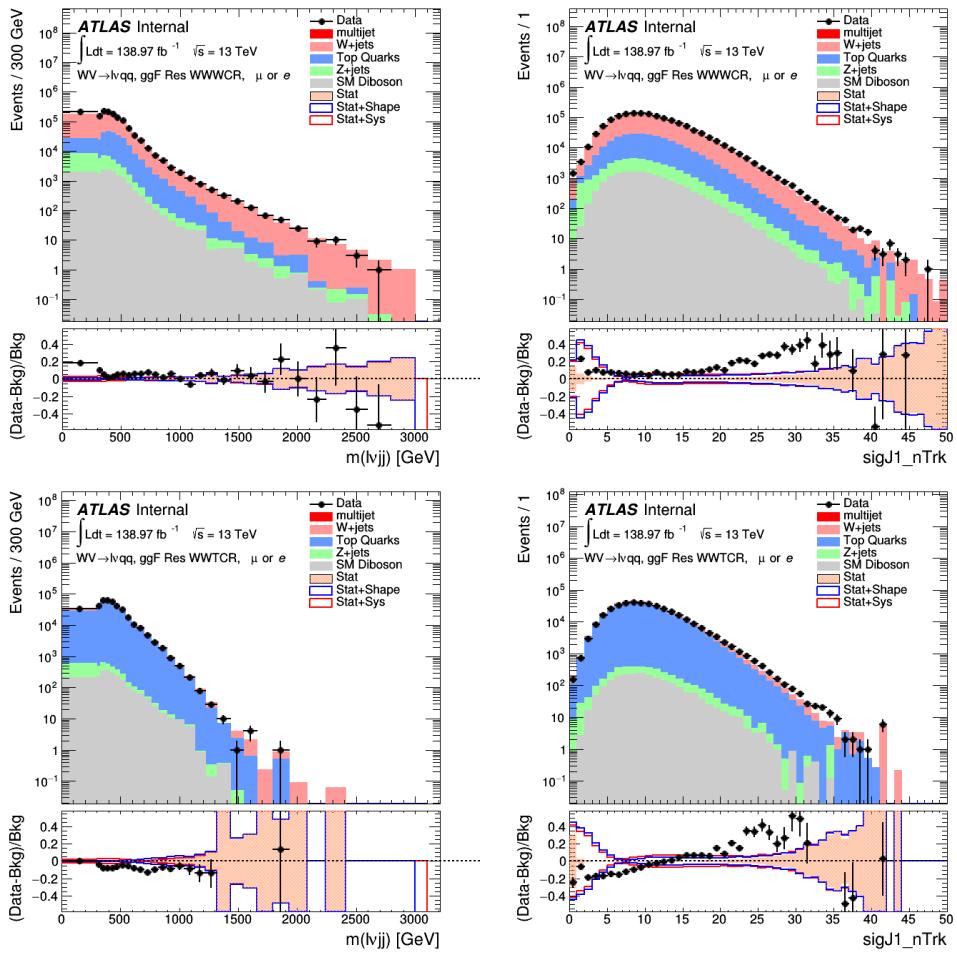


Figure 15.3: PDGID of the truth-level parton matched to the small-R jets passing the Resolved GGF WW Signal Region selections for the (a) Leading (b) Sub-Leading jets . These distributions are shown for 300, 500, and 700GeV Z' signals and the background.

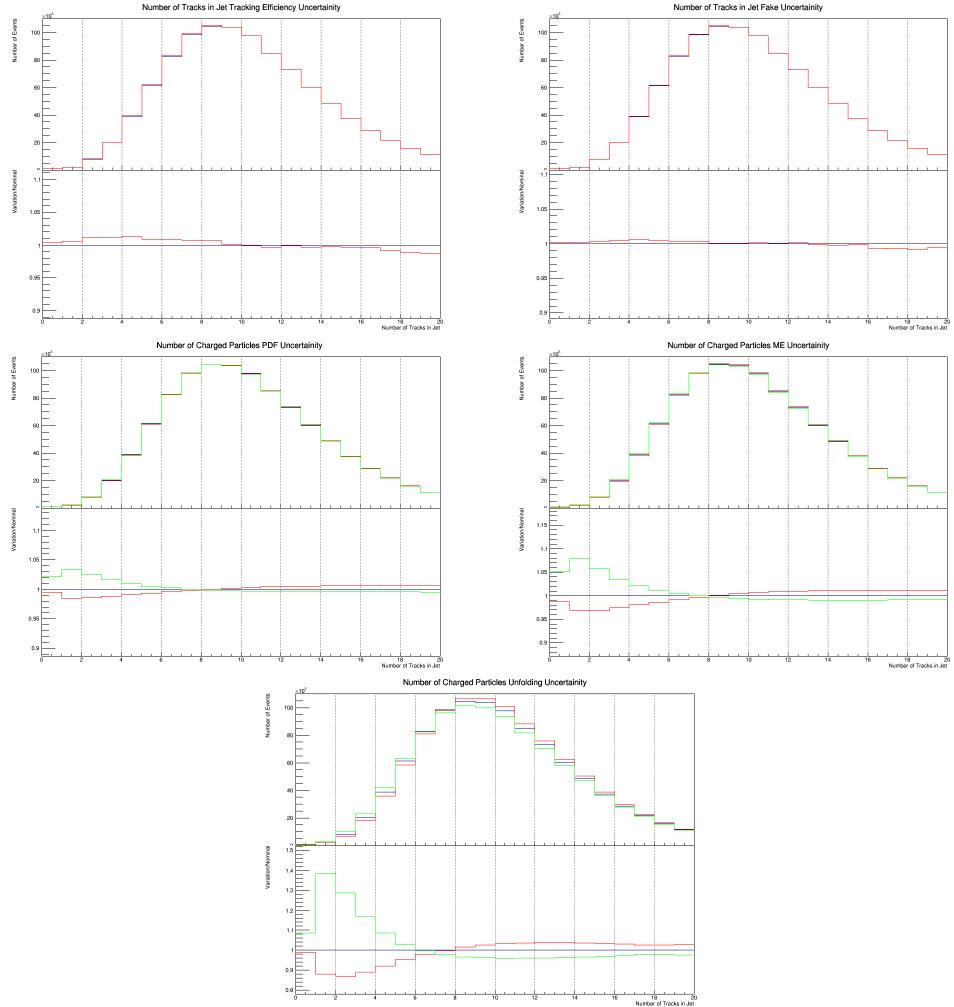


Figure 15.4: These figures show the impact of the uncertainties on the number of tracks in the leading jet in the sum of the background sample in the Resolved GGF WW SR (a) tracking efficiency (b) fake (c) PDF (d) ME (e) unfolding uncertainties.

Part VI

1471

Conclusion

1472

₁₄₇₃ **Chapter 16**

₁₄₇₄ **Conclusions**

₁₄₇₅ A search for WW and WZ diboson resonance production in $\ell\nu qq$ final states
₁₄₇₆ was performed using 139fb^{-1} of pp collision data collected at a center-of-mass
₁₄₇₇ energy of $\sqrt{s} = 13\text{TeV}$ by that ATLAS detector at the LHC between 2015 and
₁₄₇₈ 2018. No excess of events above the background-only expectation was observed.
₁₄₇₉ The largest local excess is approximately 2.7σ , which is not significant. Limits
₁₄₈₀ on the production cross section are obtained for the HVT W' and Z' and RS
₁₄₈₁ Gravitons. Signal masses below 3.4 (3.7) TeV are excluded for HVT W' Model
₁₄₈₂ A(B). Signal masses below 3.3 (3.7) TeV are excluded for HVT Z' Model A(B).
₁₄₈₃ Randall Sundrum Gravitons are excluded for masses below 1.6 TeV. Going forward,
₁₄₈₄ improving the classification of jets in events would improve analysis sensitivity.
₁₄₈₅ To distinguish quark from gluon jets a jet tagger based on the number of tracks in
₁₄₈₆ jets is studied in the context of this search. Finally, the calibration of the number
₁₄₈₇ of tracks in jets is discussed.

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