Madanalysis 5, a user-friendly framework for collider phenomenology

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Abstract

We present Madanalysis 5, a new framework for phenomenological investigations at particle colliders. Based on a C++ kernel, this program allows to efficiently perform, in a straightforward and user-friendly fashion, sophisticated physics analyses of event files such as those generated by a large class of Monte Carlo event generators. Madanalysis 5 comes with two modes of running. The first one, easier to handle, uses the strengths of a powerful Python interface in order to implement physics analyses by means of a set of intuitive commands. The second one requires to implement the analyses in the C++ programming language, directly within the core of the analysis framework. This opens unlimited possibilities concerning the level of complexity which can be reached, being only limited by the programming skills and the originality of the user.

Keywords: Particle physics phenomenology, Monte Carlo event generators, hadron colliders.

PROGRAM SUMMARY

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Program Title: Madanalysis 5

Licensing provisions: Permission to use, copy, modify and distribute this program is granted under the terms of the GNU General Public License.

Programming language: Python, C++.

Computer: All platforms on which Python version 2.7, Root version 5.27 and the G++ compiler are available. Compatibility with newer versions of these programs is also ensured. However, the Python version must be below version 3.0. Operating system: Unix, Linux and Mac OS operating systems on which the above-mentioned versions of Python and Root, as well as G++, are available. Keywords: Particle physics phenomenology, Monte Carlo event generators, hadron colliders.

Classification: 11.1 General, High Energy Physics and Computing.

Nature of problem: Implementing sophisticated phenomenological analyses in highenergy physics through a flexible, efficient and straightforward fashion, starting from event files as those produced by Monte Carlo event generators. The events files can have been matched or not to parton-showering and can have been processed or not by a (fast) simulation of a detector. According to the sophistication level of the event files (parton-level, hadron-level, reconstructed-level), one must note that several input formats are possible.

Solution method: We implement an interface allowing to produce predefined as well as user-defined histograms for a large class of kinematical distributions after applying a set of event selection cuts specified by the user. This therefore allows to devise robust and novel search strategies for collider experiments, such as those currently running at the Large Hadron Collider at CERN, in a very efficient way. Restrictions: Unsupported event file format.

Unusual features: The code is fully based on object representations for events, particles, reconstructed objects and cuts, which facilitates the implementation of an analysis.

Running time: It depends on the purposes of the user and on the number of events to process. It varies from a few seconds to the order of the minute for several millions of events.

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1. Introduction

Among the key topics of the present experimental program of high-energy physics lies the quest for new physics and the identification of the fundamental building blocks of matter, together with their interactions. In these prospects, the Large Hadron Collider (LHC) is currently exploring the TeV scale and multi-purpose experiments, such as ATLAS or CMS, are currently pushing the limits on beyond the Standard Model physics to a further and further frontier. Discoveries from these experiments, together with their interpretation, are in general challenging and strongly rely on our ability to accurately simulate both the possible candidate signals and the backgrounds. This task is however rendered quite complicated due to the complexity of the typical final states to be produced at the LHC, which contain large numbers of light and heavy flavor jets, charged leptons and missing transverse energy. Consequently, the overwhelming sources of Standard Model background require the development of robust, and possibly novel, search strategies. In this context, tools allowing to compute predictions for large classes of models are central.

This has triggered, during the last twenty years, a lot of efforts dedicated to the development of multi-purpose matrix-element based event generators such as Alpgen [1], Comix [2], Comphep/Calchep [3, 4, 5], Helac [6], Madgraph/Madevent [7, 8, 9, 10, 11], Sherpa [12, 13] and Whizard [14, 15]. As a result, the problem of the generation of parton-level events, at the leading-order accuracy, for many renormalizable or non-renormalizable new physics theories has been solved. More recently, progress has also been achieved in the automation of next-to-leading order computations. On the one hand, the generation of the real emission contributions with the appropriate subtraction terms has been achieved in an automatic way [16, 17, 18, 19, 20, 21]. On the other hand, several algorithms addressing the numerical calculation of loop amplitudes have been proposed [22, 23, 24, 25] and successfully applied to the computation of Standard Model processes of physical interest [26, 27, 28, 29, 30].

Even if each of the above-mentioned tools is based on a different philosophy, uses a specific programming language and requires a well-defined input format for the physics models under consideration, programs such as FEYNRULES [31, 32, 33, 34, 35] and LANHEP [36, 37] have alleviated the time-consuming and error-prone task of implementing new physics theories in the Monte Carlo tools. Furthermore, the introduction of the Universal FEYN-

RULES Output (UFO) format [38] and the development of an automated tool computing helicity amplitudes [39] have also streamlined the communication between the construction of a new physics theory and its implementation in the matrix-element generators through a standardized fashion.

Parton-level physics analyses based on event samples produced by matrixelement generators are far from describing the reality of what is observed in any existing detector. This kind of phenomenological work can however be useful in the prospects of investigating new types of signature and devise original search strategies, preliminary to more complete phenomenological investigations. In order to facilitate the information transfer from the matrixelement generators, a generic format for storing (parton-level) events and their properties have been proposed ten years ago, the so-called Les Houches Event (LHE) file format [40, 41]. Following the LHE standards, events are stored in a single file, following an XML-like structure, together with additional information related to the way in which the events have been generated.

Monte Carlo generator-based physics analyses at the parton-level then consist in first parsing LHE event files related to both the signal and the different sources of background, then implementing various selection cuts on the objects contained in the events, *i.e.*, quarks, leptons, neutrinos, new stable particles, *etc.*, and finally in creating histograms representing several kinematical quantities. By means of signal over background ratios, optimization of the selection cuts can be achieved in order to maximize the chances to unveil the signal of interest. Of course, the only sensible conclusions which could be stated at this stage would be to motivate (or not) a more realistic analysis, including at least parton showering and hadronization.

An accurate simulation of the collision to be observed at hadron colliders indeed requires a proper modeling of the strong interaction, including parton showering, fragmentation and hadronization. This is efficiently provided by packages such as Pythia [42, 43] and Herwig [44, 45] and several algorithms matching parton showering to hard scattering matrix elements have been recently developed [46, 47, 48, 49].

Even if not directly necessary for the analysis of the event samples, a large part of the information present in a parton-level LHE event file allows for a further matching procedure. Consequently, the generation of hadron-level events for a multitude of Standard Model and beyond the Standard Model processes can be (and has been) done in a systematic fashion. This starts from parton-level events stored into LHE format-compliant files, as

generated by matrix-element based event generators. These files are then further processed by a parton showering and hadronization code which allows to match the strengths of both the description of the physics embedded into the matrix elements and the one modeled by the parton showering. Consequently, physics analyses at the hadron-level are far more sophisticated than their counterparts at the partonic level.

When analyzing hadron-level events, one must note that the key difference with respect to the parton-level case lies in the objects which are contained in the event files. In contrast to partonic events, hadronic events consist in general in a huge collection of hadrons given together with their four-momentum. The particle content of the final state is thus much richer and the storing of the information at the time of parsing the event file is hence rendered rather cumbersome.

In order to streamline this procedure, dedicated event class libraries have been developed and several common formats for outputting hadron-level event samples exist. Event files compliant with such formats can in general be straightforwardly read by fast detector simulation programs, necessary for a more advanced level of sophistication of the phenomenological analysis. As examples, one finds the Stdhep [50] or HepMC [51] structure for event files, both widely used in the high-energy physics community. According to these conventions, an event contains, in addition to the final state particles and their properties, the whole event history as a set of mother-particle to daughter-particle relations. Let us also note that the task of parsing hadron-level event files can be highly simplified after clustering the hadrons into jets, using jet algorithms as, e.g., those included in the FAStJet program [52]. This simplified picture allows for the usage of a simpler event format, such as the LHE format introduced above, which subsequently renders an analysis easier to implement.

State-of-the-art phenomenological analyses require in general fast (and realistic) detector simulation in order to correctly estimate both the signals under consideration and the backgrounds. Starting from hadronized event samples, several frameworks, such as PGS 4 [53] and Delphes [54] are dedicated to this task. They produce event samples containing reconstructed objects such as photons, jets, electrons, muons or missing energy, together with their properties. Whereas the description of a specific signature as provided by a fast detector simulation tool is still far from what could have been obtained using a full detector simulation, including among others the transport of the particles through the detector material, fast simulation of

collider experiments is often sufficient to study the feasibility of a specific analysis on realistic grounds. The results then motivate (or not) the associated analysis in the context of a full detector simulation, the latter being embedded in generally complex experimental software such as those used by large collaborations.

As stated above, after their processing by a fast detector simulation, the original sets of hadrons are replaced by high-level reconstructed objects stored together with properties such as their (smeared) four-momentum or if they have been tagged as a b-jet or not. The structure of the events and their properties can be very efficiently saved into a file following the LHCO conventions [55]. Let us note that this format has in particular been designed for that purpose.

As outlined above, the study of the hadron collider phenomenology of a given signature can be performed at several levels of sophistication. Sometimes, very preliminary works at the parton-level are necessary in order to motivate further investigations. Hadron-level analyses give a clearer hint about the expected sensitivity of colliders to the signal under consideration. This assumes a perfect and ideal detector. In contrast, state-of-the art phenomenology includes a fast and semi-realistic simulation of the detectors. This allows to derive conclusions closer to the expectations estimated with the help of a full simulation of the detector as performed by the large collaboration software. The major drawback of the latter is that such tools consist in very heavy, complex, time-consuming and non-public algorithms. This therefore motivates pioneering works under the parton-level, hadronic-level or reconstructed-level assumptions¹. It is however important to keep in mind that the final word is always included in the data.

Performing a phenomenological analysis on the basis of the results provided by Monte Carlo generators always starts with the reading of several event samples. Since partonic or hadronic events (the latter including or not a fast detector simulation) are in general stored under different formats, this step demands to use appropriate routines capable to understand the file structure. Selection cuts on the objects included in the events, *i.e.*, partons, hadrons or high-level objects, according to the level of sophistication, are

¹In this paper, we denote by *reconstructed-level* events which have already been processed by a fast detector simulation, in contrast to *hadron-level* events which have only been showered, fragmented and hadronized and *parton-level* events where the matching to a parton showering algorithm is fully absent.

then implemented and applied to both the signal and background samples. This allows for the creation of histograms of various quantities in order to be able to extract (or not, in the worst case scenario) information about a signal in general swamped by backgrounds.

This procedure is most of the time based on home-made and non-public programs, especially due to the lack of a dedicated framework. As consequences, this can lead to various problems in the validation and the traceability of the analyses as well as possibly in the interpretation of the results. In this work, we are alleviating this issue by proposing a single efficient framework for phenomenological analyses at any level of sophistication. We introduce the package MADANALYSIS 5, an open source program based on a multi-purpose C++ kernel, denoted SAMPLEANALYZER, which uses the ROOT platform [56].

Using the strengths of a PYTHON interface, the user can define his own physics analysis in an efficient, flexible and straightforward way. Similarly to the older FORTRAN and PERL version of MADANALYSIS which was linked to the version 4 of the MADGRAPH program, MADANALYSIS 5 can either be run within MADGRAPH 5 or as a standalone package. It includes a complete reorganization of the code and the implementation of many novel functionalities such as an efficient method to implement cuts as well as to generate histograms and cut-flow charts in an automated fashion. Moreover, care has been taken for developing a fast and optimized code.

Consequently, the procedure for performing a phenomenological analysis has been drastically simplified since the only task left to the user is to define the corresponding selection cuts and the distributions to be computed. Sometimes, these embedded features might however not be sufficient according to the needs of the user. In order to overcome this limitation, MADANALYSIS 5 offers an expert mode of running with unlimited possibilities. It is unlimited in the sense that the user directly implements, within the C++ kernel, his own analysis.

Finally, in order to release a single framework for the analysis of parton-level, hadron-level or reconstructed-level based event simulations, several event file formats are supported as input, from the LHE files which could describe partonic or hadronic events to the more complex Stdhep, HepMC and LHCO file formats. Let us note that, according to the needs of the users, interfacing additional event formats, as, e.g., the Exrootanalysis format [57], can be easily achieved.

This paper documents the fifth version of Madanalysis and consists in

its user-guide. An up-to-date version of this document, together with the program can be found on the web page

http://madanalysis.irmp.ucl.ac.be

Moreover, an extended, more pedestrian, version of this manual, is also available at the same Internet address. The outline of this paper is as follows: In Section 2, we give a general overview of the program, its philosophy and its structure. Section 3 illustrates the capabilities of MADANALYSIS 5 without entering into the details, the latter being given in the next Sections. Hence, Section 4 provides the guidelines for implementing basic (but still professional) physics analyses, *i.e.*, implementing simple selection cuts and creating histograms for various kinematical distributions, whilst Section 5 is dedicated to a more expert usage of the program which allows to design more sophisticated analyses. Our conclusions are presented in Section 6. Finally, a technical Appendix on the installation of the program and its dependencies follows.

2. Overview of Madanalysis 5

2.1. Madanalysis 5 in a nutshell

The MADANALYSIS 5 package is an open source program allowing to perform physics analyses of Monte Carlo event samples in an efficient, flexible and straightforward way. It relies on a C++ kernel, named SAMPLEANALYZER, which uses the ROOT platform and interacts with the user by means of a Python command line interface.

The distribution of the program includes a home-made reader of event files created by Monte Carlo event generators. The reader is compliant with Monte Carlo samples containing events either at the parton level, at the hadron level or at the reconstructed level. Hence, a unique framework can be used for analyzing events embedded equivalently into simplified LHE files or into more complex StdHep, HepMC and LHCO files.

From the information included in the event files, the user can ask MAD-ANALYSIS 5 to generate histograms illustrating various properties of the generated physics processes. These properties range from observables related to the whole content of the events, such as the multiplicity of a given particle species or the missing transverse-energy distribution, to observables associated to a specific particle such as, for instance, the transverse-momentum distribution of the final-state muon with the highest energy or the angular

separation between the final-state jets. The user has also the possibility to compare event samples related to different physical processes in a straightforward fashion. On the one hand, the investigated distributions can be stacked on the same histogram, after automatic normalization of the samples to an integrated luminosity which the user can specify. On the other hand, the distributions can be normalized to unity in order to facilitate the design of event selection cuts allowing for an efficient background rejection based on the shape of the distribution. In this case, the histograms could be superimposed rather than stacked in the aim of improving the readability.

Moreover, if available, the whole event history, *i.e.*, all the event information from the hard-scattering process to the final-state hadrons, including the mother-particle to daughter-particle relations among the different stages yielding the final state particle content, is stored. One can emphasize that this feature is important for sanity and consistency checks of the generated event samples.

Event selection can be easily performed by means of a series of intuitive PYTHON commands yielding the application of selection cuts to the samples under consideration. For instance, it is possible to only consider, in the current analysis, events which contain at least a certain amount of missing transverse energy or a given number of final-state leptons. In the same way, one can access a specific particle in the event, such as the jet with the hardest transverse momentum, and require that it fulfills a given geometrical or kinematical criterion. As a last example, the user could decide to only select events containing at least two muons whose the invariant mass is compatible with the Z-boson mass. After applying a given cut, MADANALYSIS 5 proceeds with an automatic computation of the associated cut efficiency.

The ultimate goal of a physics analysis in particle physics is to extract some signal from an usually swamping background so that one could study its properties. The MADANALYSIS 5 framework offers the possibility to tag a specific event sample as a background or signal sample. This tagging allows for an automatic treatment of the signal over background ratio, or of any other similar observable which can be specified by the user, together with the associated uncertainty. This quantity is recomputed after each of the different selection cuts implemented by the user. With these pieces of information available, optimizing selection cuts consequently becomes easier, which allows to investigate in a fast and efficient way whether a given signature could be observable at colliders.

To display the results in a human-readable form, MADANALYSIS 5 can

collect them either into a LATEX document to be further compiled or under the form of an HTML webpage.

2.2. Basic concepts

The Madanalysis 5 program provides to the user a platform including a wide class of functionalities allowing to perform sophisticated physics analyses. Even if the existing possibilities are rather large, see Section 4 and Section 5, implementing an analysis within the Madanalysis 5 framework always follows the same steps, each of these steps being linked to one or several of the key features of the program. These key features are briefly described in this Section, whilst additional and more detailed information can be found in the rest of this manual.

Sample declaration - datasets.

When implementing an analysis within the framework of MADANALYSIS 5, the first task which is asked to the user is to indicate the Monte Carlo event files to be processed on run-time. In general, a full Monte Carlo event generation requires the simulation of several physical processes, such as the signal under consideration and the associated sources of background. Furthermore, for the processes with the highest cross sections, more than one event samples may have to be generated in order to get enough statistical significance. The user then requires these samples, describing the same physics, to be treated on the same footings. In MADANALYSIS 5, this can be performed in a very natural way. The event files to be merged are gathered under the form of an object dubbed dataset. When the analysis is effectively executed by the C++ core of the program, datasets, i.e., collections of event files, are processed rather than the event files individually.

The particle content of the events - particles and multiparticles.

According to the conventions of the different event formats introduced above, the particles described in the events are defined, in an unambiguous way, through their Particle Data Group identifier (PDG-id) [58]. Similarly, the algorithms generated by MADANALYSIS 5 are widely based on this concept to distinguish the various particle species. The drawback is obviously that this identifier is most of the time non-intuitive for the user and can even become unreadable when the set of particles included in the events becomes large. For instance, the particle content of hadron-level events consists in

hundreds of different particles, each of them being identified by a different PDG-id.

In the spirit of the MADGRAPH program, MADANALYSIS 5 alleviates this issue by allowing to associate a particle label to a given PDG-id. Hence, instead of representing an electron and a positron by the integer numbers 11 and -11, one can create the (more intuitive) labels e- and e+ and associate them to the PDG-ids 11 and -11, respectively. It is also possible to collect labels together through the concept of multiparticles. Hence, one could define a label e referring to both the electron and the positron.

In order to implement an analysis in an efficient way, it is recommended to the user to define, in a first step, a series of particle and multiparticle labels facilitating the readability (and then the validation) of his analysis.

Definition of the analysis - selections.

Implementing the analysis is the core task of the user. It consists in defining the histograms that have to be generated and the selection cuts that need to be applied. These two types of objects, *i.e.*, histograms and cuts, are uniquely dubbed *selections*. It can be noted that even if asking for the generation of several histograms is a commutative operation, the ordering of the selection cuts is in contrast important. Applying the cuts in a different order indeed leads to the production of different (intermediate) histograms and efficiency tables.

Running the analysis - jobs.

After having created a series of dataset objects and defined the analysis to be performed, *i.e.*, having implemented the histograms and selection cuts of interest, the user needs to execute the analysis on the datasets. Contrary to the previous steps which rely on the Python module of Madanalysis 5, this task is built, for efficiency reasons, upon a C++ code which is generated by Madanalysis 5 and denoted as a *job*. After that Madanalysis 5 creates the job, it has to be executed by the Sampleanalyzer kernel for each of the predefined datasets. Once the analysis has been performed, the results are subsequently imported in the Python module and re-interpreted by Madanalysis 5.

Display of the results - reports.

The generated results can be collected and displayed in a synthetic *report*. This report is given either under the HTML format, as a webpage, or as a

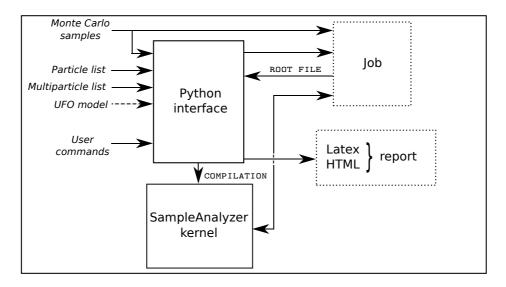


Figure 1: The Madanalysis 5 flowchart.

LATEX document that has to be further compiled. The report contains the histograms which have been generated and the efficiency tables related to the implemented selection cuts. Furthermore, a signal over background table is generated provided the datasets defined by the user have been tagged as signal and background samples.

2.3. Logical architecture of the program

The Madanalysis 5 program has two components, a Python command line interface and a C++ kernel dubbed Sampleanalyzer. In the normal mode of running of the program, these two modules interact at different stages, as illustrated in Figure 1. The way to perform physics analyses in this mode is presented in details in Section 3 and Section 4. For more sophisticated analyses, users with advanced skills in programming can bypass the Python interface and directly implement their analysis within the Sampleanalyzer framework. This expert mode of running the code is described in details in Section 5.

The Python command line interface of Madanalysis 5 consists in a command prompt where the user accesses all the functionalities of the program through a set of commands. This allows to implement an analysis in a very user-friendly way. Each command entered by the user is first checked from the point of view of the syntax and if necessary, an error message is printed to the screen. The issued command is then stored in the memory of the computer.

The interface can import different types of external information. Hence, a predefined list of particle and multiparticle labels could be imported if given as a text file or as a UFO model [38]. In a similar way, the list of the Monte Carlo event files to be analyzed can be easily loaded in the current session of MADANALYSIS 5.

After that the user has typed in all the commands defining his analysis (definition of the datasets, histograms and selection cuts), the PYTHON interface creates a C++ code using the SAMPLEANALYZER framework, the previously introduced MADANALYSIS 5 job. This C++ code comes with a Makefile and is kept available to the user for further modifications or improvements of the analysis without having to be regenerated. After the compilation and execution of the job, the PYTHON interface loads the results and uses the ROOT library of functions [56] to normalize and draw the histograms according to the requirements of the user. The HTML and LATEX reports are eventually generated.

The SampleAnalyzer C++ kernel of Madanalysis 5 is, strictly speaking, the part of the program dedicated to the analysis of the Monte Carlo event samples itself. It consists in a framework built upon an adaptive data format common for all types of Monte Carlo event samples and which contains a series of well-suited functionalities. In addition, it includes a reader compliant with the LHE, Stdhep, Hepmc and LHCO event formats and a library of specific functions facilitating particle physics analyses. Among the latter, one finds, e.g., methods allowing for boosting four-momenta or testing if a particle is a final-state particle or not. Let us finally note that the storage of information within the context of the SampleAnalyzer platform is widely based on the functions implemented within the Root library.

3. First steps with Madanalysis 5

In this Section, we present the philosophy, main features and the user-friendliness of MADANALYSIS 5 by means of a simple example. In contrast, the full list of capabilities of MADANALYSIS 5, which are much broader than what is shown in this Section, are described in Section 4.

For the sake of the illustration, we decide to perform a toy-analysis at the parton-level. We consider several samples of 1000 events each describing various Standard Model hard-scattering processes at the LHC collider running

at a center-of-mass energy of 7 TeV. Events are generated with the Monte Carlo generator MADGRAPH 5 [11]. Neglecting all quark masses (but the top mass), we employ the leading order set of the CTEQ6 parton density fit [59] and identify both the renormalization and factorization scales as the transverse mass of the produced particles.

We consider four different event samples describing three different physical processes. They are stored into the directory samples of Madanalysis 5. If this directory is not present on the system of the user or if the event files are not there for any reason, one can type, once the command line interface of Madanalysis 5 has been started (see Section 3.1),

install samples

This leads to the creation of the directory samples, if relevant, and to the download from the Internet of the four samples

```
ttbar_sl_1.lhe.gz
ttbar_sl_2.lhe.gz
ttbar_fh.lhe.gz
zz.lhe.gz
```

The first two samples are related to the production of a semi-leptonically decaying top-antitop pair where the lepton is an electron or a muon, and the third one is related to the production of a fully hadronically decaying top-antitop pair. The last sample describes the production of a Z-boson pair, including also diagrams with virtual photons, where each of the bosons is decaying either to an electron pair or to a muon pair. At the time of event generation, we demand that the produced parton-level jets have a transverse momentum $p_T > 20$ GeV, a pseudorapidity $|\eta| < 2.5$ and a relative distance $\Delta R > 0.4$. Leptons are required to have a pseudorapidity $|\eta| < 2.5$ and we ask the invariant mass of a pair of two leptons of the same flavor to be higher than 20 GeV.

3.1. Starting the command interface

Once downloaded from the web and unpacked, the MADANALYSIS 5 package does not require any compilation² or configuration and its command in-

²The C++ core of Madanalysis 5 has in fact to be compiled but this task is performed automatically, behind the scene, without requiring any interaction from the user.

terface, consisting in a command prompt ma5>, can immediately be launched by issuing

bin/ma5

from the directory where Madanalysis 5 has been installed. For the installation procedure of the program and all its dependencies, we refer to Appendix A. The user is now able to access all the functionalities of the tool and can start implementing an analysis.

When launched, the program firstly checks that all the required dependencies, such as the Root header files and libraries and a C++ compiler, are present on the system and correctly installed. If not, an error message is printed, so that the user has enough information to solve the issue, and the program exits. On its first run, MADANALYSIS 5 also compiles a static library which is stored into the directory lib of the distribution. This library is further used by the SAMPLEANALYZER kernel when analyses are executed.

Secondly, two lists of labels corresponding to standard definitions of particles and multiparticles are loaded into the memory of the current session of the program. By default, when MADANALYSIS 5 is used as a standalone package, they are imported from the corresponding files stored in the input directory of the distribution of the program. In contrast, if MADANALYSIS 5 has been installed in the directory where MADGRAPH 5 has been unpacked, the lists of particle and multiparticle labels are directly imported from MADGRAPH 5. Let us note that several multiparticles, such as hadronic or invisible, are essential for a correct running of MADANALYSIS 5. Therefore, if not included in the imported files, they are automatically created.

3.2. Particles and multiparticles

As soon as all the particle and multiparticle labels have been imported, MADANALYSIS 5 creates links pointing to lists containing all the declared labels. This allows to access them in an easy way at any time of the analysis, by issuing in the command interface

```
display_particles
display_multiparticles
```

New labels can be created on run-time with the command define, as, e.g.,

```
define mu = mu+ mu-
```

This command creates a multiparticle label mu which is associated to both the muon and the antimuon. Moreover, new labels are automatically added to the relevant list of (multi)particles.

The definition of a specific label can be retrieved with the help of the display command which outputs the PDG-id of the associated particle(s). For example, after invoking the two commands

```
display b
display l+
```

the command interface outputs information about the (multi)particle labels b and 1+,

```
The particle 'b' is defined by the PDG-id 5. The multiparticle 'l+' is defined by the PDG-ids -11 -13.
```

As can be seen from the screen output above, the two symbols **b** and **1+** define a *b*-quark and a positively charged lepton different from a tau, respectively.

3.3. Importing event samples

Preliminary to perform any analysis, event files under consideration must be parsed and loaded into the memory. The command import has been designed in that purpose. As a mandatory (single) argument, it requires the name of the Monte Carlo sample(s) to be parsed. In order to import several files at one time, the wildcard characters * and ? are allowed when typing-in the argument of the function. Hence, the four samples introduced above can be loaded simultaneously by issuing the command

```
import samples/*.lhe.gz
```

which is equivalent to the set of four commands

```
import samples/ttbar_sl_1.lhe.gz
import samples/ttbar_sl_2.lhe.gz
import samples/ttbar_fh.lhe.gz
import samples/zz.lhe.gz
```

The result is the creation of a unique event sample denoted by defaultset containing all the imported files. We are now ready to define selection cuts which the SampleAnalyzer kernel will apply to the events included in the dataset defaultset. The name as well as the way how to merge the samples can be tuned according to the needs of the user, but this goes beyond the scope of this Section and will be addressed in Section 4.

3.4. Selection cuts and creation of histograms

Creating an histogram representing a specific kinematical distribution has been made very efficient in the framework of Madanalysis 5 through the command plot. This command requires one mandatory argument, the observable to be computed, and a set of optional arguments containing, among others, the number of bins of the histogram to be created and the lower and upper bounds of its x-axis. In the setup of the bounds of an histogram, one must note that the standard unit of energy used in Madanalysis 5 is the GeV.

In the toy-analysis which is implemented in the sequel, we focus on the missing transverse-energy distribution as well as on the transverse-momentum distribution of the final state muons. We recall that we are considering a unique event sample resulting from the merging of three $t\bar{t}$ and one diboson event files, as explained in Section 3.3. In order to create the associated histograms, it is sufficient to issue the two commands

```
plot MET
plot PT(mu) 20 0 100
```

where we remind that the multiparticle mu has been defined in Section 3.2 and represents both the muon and the antimuon. The symbol MET is associated to the missing transverse energy whilst the function PT stands for the transverse momentum of a given (multi)particle provided as its argument. The next pieces of information to be passed to the command plot are optional and related to the binning of the histograms. By default, *i.e.*, in the case the binning information is not specified as in the first example above, MADANALYSIS 5 uses hard-coded values which depend on the observable under consideration. In contrast, as in the second example above, the user can provide, at the time of typing-in the command, the number of bins (20 here) together with the values of the lowest and highest bins (which are chosen equal to 0 GeV and 100 GeV, respectively, in our example).

We now turn to illustrating the implementation of the two types of selection cuts which is possible to employ in Madanalysis 5. These cuts will be further applied, by the Sampleanalyzer kernel, to the events contained in the defaultset dataset. We recall that the program contains a set of possibilities much broader than what is presented in this Section and we refer to Section 4 for more information. In a first step, we decide to select events where the missing transverse energy (E_T) is not too large, *i.e.*, $E_T < 100$

GeV. Rejecting events not fulfilling this criterion can be very efficiently and easily performed by issuing the command

The function reject tells MADANALYSIS 5 to remove from the event selection any event where the missing energy (MET) is larger than 100 GeV.

As a second illustration about the implementation of selection cuts in Madanalysis 5, we focus on the kinematical properties of the muons contained in the selected events. In particular, a part of the events contain a muon-antimuon pair and we would like to represent the corresponding invariant-mass distribution through an histogram. However, many events do not exactly contain one muon and one antimuon. This is taken into account by Madanalysis 5 at the time of the generation of the histogram, where one entry is included for each different muon-antimuon pair which can be formed from the event particle content. For a given event, the number of entries can hence be zero, one or bigger than one according to the (anti)muon multiplicity of the final state.

Realistic detectors are in general not capable to correctly reconstruct too soft particles. Even at the parton level, this effect can (and should) be implemented. In our case, we could consider as (anti)muon candidate only final state (anti)muons with, e.g., a transverse momentum $p_T > 20$ GeV. This cut can be implemented in MADANALYSIS 5 through the command reject, but following a syntax different from above,

The effect of this command is to consider, for each of the selected events, as muon and antimuon candidates only muons and antimuons with a transverse momentum (PT) harder than 20 GeV. For all the histograms which are generated after having typed the line above in the command interface, only the selected (anti)muon candidate's are considered. As stated above, we choose, as an example, to represent the muon-pair invariant mass distribution computed from the analyzed event samples included in the dataset defaultset. The command to create this histogram is similar as those presented in the beginning of this Section,

plot M(mu+ mu-) 20 0 100

where the arguments of the function M, associated to the invariant mass, are multiple. This denotes that we are summing the four-momenta of the muon and the antimuon to derive the invariant-mass to be represented. As it can be seen from the optional arguments of the command plot, we once again ask for an histogram of 20 bins with its lower and upper bounds being fixed to 0 GeV and 100 GeV, respectively. We remind that for a specific event, the number of entries which are included in the histogram corresponds to the number of different muon-antimuon pairs that can be formed from the final state particle content. This can then be any integer number.

Once the selection is defined, it has to be executed through a job to be run by SampleAnalyzer. This is done through the command submit which takes as an argument the name of a directory which will be created,

submit <dirname>

The created directory <dirname> contains a series of C++ source and header files that are necessary for SampleAnalyzer to properly run. The compilation, linking to the external static library of Madanalysis 5 (see Section 3.1) and the execution of the resulting code is handled by Madanalysis 5. The screen output indicates the status of these different tasks and various information such as the detected format of the event samples or the number of processed events. In the case anything is not going as smoothly as it should, Madanalysis 5 also prints warning and/or error messages to the user and the program exists in the worst case scenario.

3.5. Displaying the results

After having performed the analysis as indicated in the previous Section, MADANALYSIS 5 offers several ways to present the results under a human-readable report. This report can be either generated under the HTML format or under a TeX format that could be compiled with a LATEX or a PDFLATEX compiler. The histograms are saved under a Portable Network Graphics file (.png) for HTML and PDFLATEX reports, and under an Encapsulated PostScript file (.eps) for LATEX reports. The MADANALYSIS 5 commands generating the reports read

generate_latex <dirname>
generate_pdflatex <dirname>
generate_html <dirname>

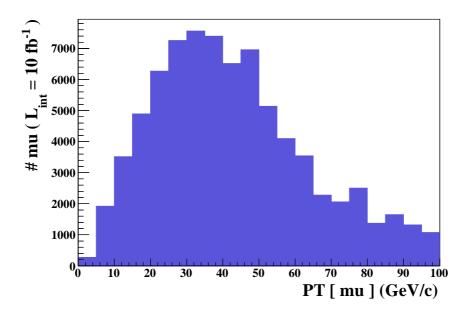


Figure 2: Transverse-momentum distribution of the muons for a dataset consisting in the four event samples introduced in Section 3.

where **dirname** stands for the directory in which the report is generated.

The structure of the report is similar whatever is the adopted format. It starts with a table of contents and firstly displays all the information necessary to ensure the reproducibility of the analysis. One hence finds the list of commands which have been issued, in the current session of the program, before the generation of the report, followed by the employed version of the code as well as the values of all the setup parameters. In this last category, one has, e.g., the list of the event samples which have been analyzed, given together with the corresponding cross sections and the number of events contained in the event files.

The core of the report contains the results of all the commands related to histogram creation and to the application of a selection cut. These commands have been treated one by one by SAMPLEANALYZER and the report follows this pattern. By default, histograms are normalized to an integrated luminosity of 10 fb⁻¹ and comes together with a summary table containing information on the mean value of the computed distribution and the associated root mean square (RMS), as well as on the presence of possible

Dataset	Integral	Entries / event	Mean	RMS	Underflow	Overflow
defaultset	82747	0.752	42.8177	21.36	0.0	6.181

Table 1: Statistics associated to the histogram of Figure 2.

Cuts	Signal (S)	Background (B)	S vs B
Initial (no cut)	109999 +/- 789		
cut 1	82994 +/- 375443		
cut 2	82994 +/- 612		

Table 2: Efficiencies of the cuts applied in Section 3.4.

underflow and overflow bins in the histogram. The latter are given as a ratio of the value of the integral of the represented distribution between its lower and upper bounds. A green-orange-red color code indicates their relative importance, an orange or red color suggesting more or less strongly to the user to modify the bounds of the histograms, if relevant. We refer to Section 4 for the description of all the properties of the histograms that can be modified by the use, such as the way to modify the luminosity or the binning of a given histogram.

In Figure 2, we take the example of the transverse-momentum distribution of the muons implemented in Section 3.4 and present the histogram generated by MADANALYSIS 5. As stated above, if several muons are contained in one specific event, they each correspond to a different entry in the histogram with the same weight. The summary table generated together with the histograms is given in Table 1. In the first column of the table, the name of the dataset is printed. In the second column, one finds the number of events normalized to an integrated luminosity of 10 fb⁻¹, since this is the default value which has not been modified. In the third column, the average number of histogram entries for each event is provided, with the mean and the root mean square of the distribution under consideration shown in the fourth and fifth column. From the orange cells in the table, one immediately notices that the number of underflow and overflow entries given in the last two columns are only under a roughly reasonable control.

Let us finally note that for each selection cut, a summary table of the cut efficiency is also given, together with the corresponding effect on the signal over background ratio. In our toy-example, we have not tagged any sample as signal or background. Therefore, this feature is irrelevant here and we refer to Section 4 for more information. The generated summary of all cuts is provided in Table 2, where all the events are (by default) considered as signal events.

4. Implementing analyses in an efficient and user-friendly way

4.1. Starting a Madanalysis 5 session

The Madanalysis 5 package is able to handle several types of Monte Carlo samples. Supported event files can describe partonic, hadronic or reconstructed events. The key difference between these three types of events lies in the basic objects in which the event content is expressed in terms of. In the first case, an event consists in a set of partons (quark, gluons, charged leptons, neutrinos, new states, etc.), whilst in the second case, it consists in hadrons (pions, kaons, baryons, etc.). Finally, a reconstructed event contains a set of high-level reconstructed objects such as electrons, muons, jets, etc.. For each of these classes of events, Madanalysis has a specific running mode which the user is required to specify when launching the code,

bin/ma5 [level]

Typing in a shell bin/ma5 --partonlevel or bin/ma5 -P allows to run MADANALYSIS 5 in the mode required to analyze parton-level events, whilst issuing bin/ma5 --hadronlevel or bin/ma5 -H makes the program running in the hadron-level mode. In a similar fashion, the two shell commands bin/ma5 --recolevel and bin/ma5 -R allow to start MADANALYSIS 5 in a readily way to analyze reconstructed-level events. In the case no argument is provided, as in the example of Section 3, the parton-level mode is automatically selected. Of course, the different modes cannot be combined, since the levels of sophistication are self-excluding. Once one of the previously introduced command is issued, the command line interface of MADANALYSIS 5 is started. It consists in a command prompt ma5> where the user can access all the functionalities of the program and directly implement a physics analysis by issuing a set of commands. We refer to Section 4.2 for more information about the whole set of existing commands.

The list of commands to be typed can also be provided under the form of a script, *i.e.*, a simple text file. In order to execute the script, its path has to be provided as the last argument when typing the bin/ma5 command in a shell. Hence,

```
bin/ma5 --recolevel <filename>
```

starts the command interface of Madanalysis 5 in the reconstructed-level mode. All the commands included in the file <filename> are then sequentially applied, one after the other. When having such a script executed by Madanalysis 5, it can be useful to bypass all the confirmation questions usually asked by the program. This can be done by issuing in a shell one of the two commands

```
bin/ma5 --script --recolevel <filename>
bin/ma5 -s --recolevel <filename>
```

This *script mode* also enforces the automatic exit of MADANALYSIS 5 after the completion of the script <filename>, in contrast to the *forced mode* described below. It is also possible to execute several scripts by providing various filenames,

```
bin/ma5 -s -R <filename1> <filename2> <filename3>
```

The different files are handled as concatenated in one single file, *i.e.*, all the commands included in filename1 are processed sequentially, followed by the commands included in filename2, *etc.*. Let us note that the character '#' can be included when writing down script files. Everything standing to the right of this character is considered as a comment by MADANALYSIS 5 and ignored by the command line interpreter.

For highly sophisticated analyses, the features which are presented in the sequel of this Section might not be sufficient. Therefore, Madanalysis 5 can be run in an expert mode by issuing one of the three commands

```
bin/ma5 -- expert bin/ma5 -e bin/ma5 -E
```

The possibilities to implement any analysis are here unlimited. More information about the expert mode is provided in Section 5. Let us however note that in this case, any other option which could be passed when starting the interface is ignored.

Among the other options which could be provided when starting the interpreter, one can note that the version of the distribution installed on the system of the user can be accessed through one of the following commands,

Table 3: Syntax to launch the MADANALYSIS 5 program:

bin/ma5 -h or bin/ma5 --help

This allows to display to in a shell a summary of the content of this Table.

bin/ma5 -v, bin/ma5 --version or bin/ma5 --release This allows to display in a shell the number of the version of MAD-ANALYSIS 5 present on the system.

bin/ma5 -f , bin/ma5 --forced

This allows to run Madanalysis 5 in a mode where confirmation questions are ignored.

bin/ma5 -e , bin/ma5 -E or bin/ma5 --expert This indicates that MADANALYSIS 5 has to run in expert mode. In this case, any other option is ignored.

bin/ma5 [level] [files]

[level]

When provided, this optional argument allows to select the type of event files to analyze. If absent, parton-level events are assumed. The user can choose one of the following options:

-P or --partonlevel Parton-level events.
 -H or --hadronlevel Hadron-level events.
 -R or --recolevel Reconstructed events.

[files]

When provided, this optional argument consists in a sequence of filenames, separated by blank characters, containing each a set of MAD-ANALYSIS 5 commands. The files are handled as concatenated and the commands are applied sequentially. If the option pattern -s or --script is included, MADANALYSIS 5 exits after having executed the script and does not ask any confirmation question.

and that typing-in

ensures a running mode of MADANALYSIS 5 where confirmation questions, such as, e.g., those printed to the screen when a directory is about to be removed, are not asked to the user.

The different running modes of MADANALYSIS 5 and the way to cast them are summarized in Table 3. They can also be displayed by the program when the user types in a shell one of the following commands,

4.2. The command line user interface of Madanalysis 5

The command line interface of Madanalysis 5 is built upon the Python module cmd which allows for a flexible processing of commands issued by the user. It features, among others, text help, tab completion and shell access. In addition, the history of commands issued by the user can be obtained by means of the up and down keys of the keyboard. As presented in Section 4.1, the command interface can be started by issuing in a shell the command

bin/ma5 [options]

from the directory where MADANALYSIS 5 has been installed. However, MADANALYSIS 5 could also be started from any location on the computer of the user, the command above needing only to be modified accordingly.

The user interface consists in a command prompt ma5> where the user can implement a physics analysis very efficiently and easily by means of a series of (case sensitive) commands. The related syntax is inspired by the Python programming language and a command therefore always starts with an action. Several actions can be typed together in a single command line after separating them with a semi-column ';'. The number of different implemented actions has been made as small as possible in order to guarantee a quick handling of the code by any physicist. Their list has been collected in Table 4, with basic instructions about how to use each of these actions. The information included in this Table can also be displayed to the screen at run-time by issuing

help [<action>]

where **<action>** is the action under consideration. Moreover, if **help** is typed-in without any argument, the list of all the available actions is printed out. Let us emphasize that tab completion could also be used in the aim of obtaining that list. We now dedicate the rest of this Section 4 to a detailed description of each of those actions.

As stated above, the history of all the actions undertaken by the user can be accessed through the up and down keys of the keyboard, as for standard shells. They are read from the file .ma5history stored in the directory where Madanalysis 5 has been unpacked. This file is updated every time that the user types a command, and the list of the last 100 commands executed by the interpreter is saved. In addition, if one restricts ourselves to the current session of Madanalysis 5, the list of typed commands can be accessed by issuing in the command interface

history

Besides actions, the language handled by the interpreter contains *objects* that actions is acting on. There exist various types of objects representing particles, multiparticles, datasets, selection cuts or even the configuration panel of the current session of the program, denoted by main. In addition, any object is described by several attributes related to its properties, which are generically denoted as *options*. These attributes can be accessed through the syntax *object.option*.

Finally, the language of MADANALYSIS 5 contains a set of reserved keywords, such as all, or or as, whose the usage is described in the sequel. Moreover, it is important to keep in mind that both predefined and user-defined objects have a name which has to be unique. Therefore, in an analysis, any object which is created, such as, e.g., a new instance of the multiparticle class, must have a name different from those of the existing actions (see Table 4) or those of the objects already defined and/or used in the current analysis.

The syntax introduced above has been developed with a focus on uniformization and intuition. For the sake of the example, we now focus on the three commands display, set and remove. To display to the screen an object denoted by <object> together with its properties, it is enough to type the command

display <object>

Similarly, the display of a specific property, denoted by <option>, of the object under consideration can be performed by issuing

```
display <object>.<option>
```

On the same footings, the value of the attribute <object>.<option> can be easily modified via the action set

```
set <object>.<option> = x
```

where in the example above, the option under consideration is set to the value x. Finally, an object can be deleted from the computer memory with the help of the command remove,

```
remove <object>
```

However, three objects, the objects main, hadronic and invisible, as well as all the particles and multiparticles currently used in the analysis, i.e., being related to an histogram or to a selection cut, cannot be deleted.

As already stated in the beginning of this Section, the cmd module of PYTHON allows for an efficient access to shell commands directly from inside the command line interface of MADANALYSIS 5. This requires to start the command line either with an exclamation mark '!' or with the command shell. Therefore, the syntax

```
!<command>
```

or equivalently

```
shell <command>
```

has to be followed in order to execute the command <command> from an external shell.

Finally, the configuration of the command line interpreter can be restored as for a new session of Madanalysis 5 by issuing the command

reset

Consequently, all the user-defined particles and multiparticles as well as all the implemented histograms and selection cuts are removed from the computer memory.

Table 4:	Actions available from the command
line	user interface of Madanalysis 5

define <name> = <def> This creates a new (multi)particle la-

bel <name> defined through the (list of)

PDG-id(s) <def>.

display <var> This displays to the screen either the

properties of an object or the value of one of its options, both generically de-

noted by **<var>**.

display_datasets This displays to the screen the list of

the imported datasets.

display_multiparticles This displays to the screen the list of

defined multiparticle labels.

display_particles This displays to the screen the list of

defined particle labels.

EOF This closes the current session of MAD-

Analysis 5.

exit This closes the current session of MAD-

Analysis 5.

generate_html <dir> The report of the current analysis is

generated, under the HTML format, and

stored in the directory <dir>.

generate_latex <dir> The report of the current analysis is

generated, under the TEX format, and stored in the directory <dir>. Compilation requires a LATEX compiler.

generate_pdflatex <dir> The report of the current analysis is

generated, under the TEX format, and stored in the directory <dir>. Compilation requires a PDFLATEX compiler.

help This displays to the screen the list of

all the actions implemented in MAD-

Analysis 5.

history This displays to the screen the history

of the commands typed in the current

session of the program.

Table 4	(continued):	Actions	a vailable	from t	he command
	line user i	nterface	of MADA	NALVSIS	5

import <obj> This allows to import in the current ses-

sion of Madanalysis 5 external information, generically denoted by <obj>, such as Monte Carlo samples, a configuration used by SampleAnalyzer in a previous analysis or a UFO model.

install <obj> This allows to install the external ob-

ject <obj> from the Internet. Currently, the only allowed choice for the value of the variable <obj> consists in the keyword samples which yields the download of the four example samples

presented in Section 3.

plot <obs> [options] This creates, in an analysis, an his-

togram associated to the distribution of the observable **<obs>**. We refer to Section 4.5 for the list of available options.

open <rep> This opens the report <rep> of a given

analysis, with the appropriate program.

preview <hist> This displays an histogram <hist> in a

graphical window.

quit This closes the current session of MAD-

Analysis 5.

reject (<prt>) <cond> This adds a selection cut. In an anal-

ysis, a candidate to a given particle species cprt> is ignored if the condi-

tion <cond> is fulfilled.

reject <cond> This adds a selection cut. An event is

rejected if the condition <cond> is ful-

filled.

remove <obj> This removes the object <obj> from the

memory.

Table 4 (continued): Actions available from the command line user interface of Madanalysis 5

reset This reinitializes Madanalysis 5 as

when a new session starts.

resubmit This allows, for an analysis which

has already been run by SAMPLEAN-ALYZER and further modified, to be re-

executed.

select (<prt>) <cond> This adds a selection cut. In order to

consider, in an analysis, a candidate to a given particle species <prt>, the con-

dition <cond> has to be fulfilled.

select <cond> This adds a selection cut. Events are

selected only if the condition <cond> is

fulfilled.

set <obj>.<opt> = <val> This allows to set the attribute <opt>

of the object <obj> to the value <val>.

shell <com> This allows to run the shell command <com> from the command interface of

MADANALYSIS 5. The action shell can be replaced by an exclamation

mark '!'.

submit <dir> This allows to execute an analysis as

a job run by SAMPLEANALYZER. The C++ source and header files related to the job are created in the directory

<dir>.

swap <sel1> <sel2> This allows to permute the sequence of

two instances of the class selection, <sel1> and <sel2>. We refer to Sec-

tion 4.5 for more information.

Before moving on, some remarks on tab completion are in order. This feature allows for an easy typing of commands, attributes of the objects, etc.. Typing on the 'tab' key has the effect of printing to the screen all the allowed possibilities for completing the command about to be written. It works equivalently for actions, objects and attributes. Let us emphasize that this makes tab completion in MADANALYSIS 5 much more advanced than its counterpart in standard command shells.

4.3. Datasets

In Section 3, we have shown that four example Monte Carlo samples can be downloaded from the Internet by employing the command install,

install samples

and stored into a directory samples of the current distribution of MAD-ANALYSIS 5. Actions, such as their gathering into datasets, the creation of histograms and the definition of selection cuts, have been subsequently executed on those samples. This Section is dedicated to the dataset class of objects, whilst Section 4.5 and Section 4.6 concerns histograms and cuts, respectively.

In order to load Monte Carlo samples in the computer memory, they have to be imported from outside the current session of MADANALYSIS 5. This requires the use of the command import, as briefly presented in both Section 3 and Table 4,

import <path-to-sample>

The syntax above allows to import a Monte Carlo sample stored at a location <path-to-sample> on the computer of the user. Several samples can be imported at one time by employing the wildcard characters * and ?. As a generic example, the command

import <directory>/*

imports all the samples stored in the directory <directory>. Moreover, the tilde character ~ can be used when typing the path to a sample. As in standard Linux shells, it points to the location of the home directory of the user.

According to their event format, MADANALYSIS 5 handles differently the imported event samples. The key of the procedure lies in the extension of the

event files, up to a possible packing with GZIP. The formats currently supported are the LHE event file format, whose the corresponding file extensions are .lhe or .lhe.gz, the STDHEP event file format, whose the corresponding file extensions are .hep or .hep.gz, the HEPMC event file format, whose the corresponding defining extensions are .hepmc or .hepmc.gz and the LHCO event file format, whose the corresponding event file extensions are .lhco or .lhco.gz.

Of course, all the different formats are not appropriate for any level of sophistication of the analysis. For instance, using the parton-level mode of MADANALYSIS 5 to analyze an event file compliant with the LHCO format is clearly inadequate. More into details, events stored in files compliant with the LHE format can be imported whatever is the level of sophistication of the analysis, i.e., equally for parton-level, hadron-level or reconstructed-level analyses. A remark is in order here. In principle, the LHE format is not supposed to describe reconstructed events. However, there exist routines external to Madanalysis 5, such as the Hep2LHE converter included in MadGraph, which allow to efficiently convert events as produced by hadronization algorithms or a fast detector simulation tool into reconstructed events. Those routines use a jet algorithm in order to reconstruct light and b-tagged jets, and the total missing transverse energy of the event is eventually computed. In contrast, STDHEP and HEPMC event files can only store parton-level and hadron-level events, whilst complementary, LHCO files can only be used after that high-level objects have been reconstructed, i.e. at the reconstructed-level after detector simulation.

The effect of the action import is to unify all the imported samples into one single object dubbed dataset. If nothing is specified by the user, events are gathered together into a dataset called defaultset. The possibility to import events and collect them into different datasets allows, for instance, to differentiate background event samples from signal event samples, as with the following commands

```
import <path-to-signal-events>/* as signalset
import <path-to-background-events>/* as backgroundset
```

The first command imports all the event samples present in the directory <path-to-signal-events> and collects them into one single dataset labeled as signalset. Similarly, the second command loads into the current session of MADANALYSIS 5 all the event files located in the directory

<path-to-background-events> and gathers them together into a dataset labeled as backgroundset. On the same footings, this feature can be used to collect efficiently event files corresponding to different sources of background,

```
import <path-to-events>/ttbar* as ttbar
import <path-to-events>/zz* as zz
import <path-to-events>/drellyan* as drellyan
```

These three commands import in Madanalysis 5 three series of samples, related respectively to top-antitop pair, Z-boson pair and Drell-Yan events. Three different datasets, ttbar, zz and drellyan, are created so that these events can be treated separately when performing the analysis.

A dataset has several properties which are implemented as options of the dataset class and whose the value can be modified with the command set. These attributes can be grouped into three categories of properties which are summarized in Table 5.

The first category of options consists in fact in one single property which is related to the attribute type. It allows to tag the corresponding event sample(s) as a part of the signal or the background. Taking the example of a generic dataset <dataset>, the two commands

```
set <dataset>.type = background
set <dataset>.type = signal
```

tag it as a dataset belonging to the series of background or signal event samples, respectively. By default, a dataset is always considered as of the type signal. In an analysis, this distinction between signal and background is mandatory for a correct (automated) computation of the cut efficiencies by Madanalysis 5, as well as for the corresponding derivation of the signal over background ratio.

When creating histograms, their overall normalization is related to both the luminosity, which can be specified by the user (see Section 4.7) and the cross sections associated to the different datasets included in the histograms. These cross sections are included in LHE event files and are directly read and imported in the current session of the program, in the case LHE samples are analyzed. In contrast, the numerical value of the cross sections are absent from Stdhep, HepMC and LHCO files. Therefore, in most of the cases, the user has to indicate the cross section manually. Moreover, in the case of LHE files, one could also want to modify the values which have been read.

For instance, the cross section associated to an event sample which has been generated with a leading order Monte Carlo tool could be modified by the user in order to account for next-to-leading order normalization effects.

The second category of options related to the dataset class offers two ways to perform the above-mentioned task, either through the weight attribute or through the xsection attribute. A weight different from one can be assigned to each event of a dataset by modifying the value of the attribute weight of the dataset class,

set <dataset>.weight = <weight>

The command line above leads to the assignment of a weight <weight> to each event included in a generic dataset which has been defined as <dataset>. Consequently, the default value of one has been superseded by the value <weight>. In contrast, the user can also decide to leave the weight of each event unchanged and equal to unity, and modify instead the value of the cross section. The new value of the cross section has to be stored through the attribute xsection of the dataset class.

set <dataset>.xsection = <value>

If the value <value> is different from the default choice of zero, MADANAL-YSIS 5 uses it at the time of the creation of the histogram when calculating its normalization, ignoring the value of the cross section possibly included in the event file.

The last class of attributes associated to dataset objects concerns the layout of the histograms generated by MADANALYSIS 5, and more in particular the style of the curves associated to the datasets which can be customized. On the one hand, styles and colors can be modified via the value of the attributes linestyle, linewidth, linecolor, backstyle and backcolor. On the other hand, the text used in histogram legends can be set by the user through the attribute title. The first three attributes above are associated to the type of lines (solid, dashed or dotted) used in the histograms, their width and their color (blue, green, none, purple, white, black, cyan, gray, orange, red or yellow). The two attributes backstyle and backcolor refer to the surface under an histogram and the style (solid, dotted, hatched lines) and color (blue, green, none, purple, white, black, cyan, gray, orange, red or yellow) employed when it is drawn.

Table 5: List of the attributes of the dataset class
<pre>set <dataset>.<option> = <value></value></option></dataset></pre>

backcolor In an histogram, this changes the color filling the surface under the curve associated to the dataset under consideration. The allowed choices are auto (default), blue, green, none, purple, white, black, cyan, gray, orange, red and yellow. Integer numbers between one and four can be added or subtracted to these values (see backstyle In an histogram, this changes the style employed when filling the surface under the curve associated to the dataset under consideration. The allowed choices are auto, solid, dotted, hline, dline and vline (see Figure 4). linecolor In an histogram, this changes the color of the line of the curve associated to the dataset under consideration. For the allowed choices, we refer to the attribute backcolor. In an histogram, this changes the style of the line of linestyle the curve associated to the dataset under consideration. The allowed choices are solid, dashed, dotted and dash-dotted. linewidth In an histogram, this changes the width of the line of the curve associated to the dataset under consideration. It takes an integer value between one and ten. title This change the string used in histogram legends for the dataset under consideration. The possible choices are either auto (the name of the dataset) or any string under a valid T_FX form. type This modifies the type of a dataset, associating it to the set of either signal (signal) or background (background) samples. weight This allows to change the weight of each event included in the dataset. The default value is one. This allows to modify the total cross section associated xsection to the events included in the dataset under considera-

tion. The value can be any real number.

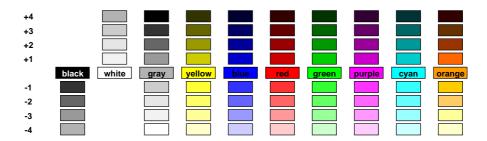


Figure 3: Complete list of allowed choices for the values taken by the attributes backcolor and linecolor of the dataset class.

The default value for the two attributes related to colors in histograms, linecolor and backcolor, is auto. This means that MADANALYSIS 5 handles the color features automatically, assigning different colors to the datasets created by the user. For a specific dataset object denoted by <dataset>, the values of the two color attributes can be, as usual, superseded by employing the command set,

```
set <dataset>.linecolor = <color>
set <dataset>.backcolor = <color>
```

The supported values for the variable <color> are intuitive and read auto, blue, green, none, purple, white, black, cyan, gray, orange, red or yellow. For histograms employing a large number of datasets, this panel of colors might however be insufficient. Shades of these basic colors can be used by adding or subtracting an integer number between one and four to those colors. For instance, the set of commands

```
set <dataset1>.backcolor = red+1
set <dataset2>.backcolor = red+2
set <dataset3>.backcolor = red+3
set <dataset4>.backcolor = red+4
```

allows to assign four different shades of red to the four datasets <dataset1>, <dataset2>, <dataset3> and <dataset4>. Consequently, when a stacked histogram containing these four datasets is drawn, the associated surfaces will be all filled with different color derived from the basic red. The complete list of colors integrated in MADANALYSIS 5 is illustrated on Figure 3.

The style of the lines of the curves drawn in the histograms can be modified through the attribute linestyle of the dataset class,



Figure 4: Complete list of styles allowed for the values possibly taken by the attribute backstyle of the dataset class.

set <dataset>.linestyle = <value>

This command changes the default employed solid style to the value <value>, which can be either solid, dashed, dotted or dash-dotted. Similarly, the attribute linewidth allows to change the width of the drawn lines,

set <dataset>.linewidth = <value>

where **<value>** is an integer number between one and ten, the default value being unity. The option **backstyle** of the **dataset** class is linked to the style employed to fill the surface under an histogram and can be set to a new value by issuing in the command line interface

set <dataset>.backstyle = <value>

The allowed choices for the parameter <value> are either the default value auto or solid, dotted, hline, dline or vline, as presented in Figure 4. The last three choices, less intuitive, stand for horizontal, diagonal and vertical lines, respectively.

Finally, when creating an histogram where the curves related to several datasets are stacked or superimposed, MADANALYSIS 5 includes a legend with the explanation of the color/style code employed to distinguish the datasets. By default, the name of the dataset is used in this legend, but this can be modified by the user once setting the attribute title to a string consisting in a valid TeX expression,

set <dataset>.title = <string>

where $\langle string \rangle$ could stand, e.g., in the case of a dataset describing events related to the production of a pair of W-bosons, for "W^{+} W^{-}" (including the quotation marks).

The effects of the customization of the histograms are only taken into account when the reports containing the results are generated. Therefore, the user does not have to issue the command submit each time he modifies the style of a curve or how the area under a curve is filled inside an histogram.

Before moving on, let us remind that the command display_datasets introduced in Table 4 allows to display to the screen the list of all the instances of the dataset class which have been created in the current session of Madanalysis 5. Furthermore, the action display can also be used on datasets. This leads to the printing to the screen of the current values of all the attributes of a given dataset. Hence, for a dataset labeled by ttbar where not any attribute has been modified by the user, the effect of the command

display ttbar

is to print to the screen the following pieces of information,

```
Name of the dataset = ttbar (signal)
Title = 'ttbar'
User-imposed cross section = 0.0
User-imposed weight of the set = 1.0
Line color in histograms = auto
Line style in histograms = solid
Line width in histograms = 1
Background color in histograms = auto
Background style in histograms = solid
List of event files included in this dataset:
- ttbar.lhe.gz
```

This indeed summarizes the status of the instance of the dataset class labeled ttbar and displays the value of all the options.

Finally, a dataset can be deleted from the memory of the computer by employing the command remove. If several event files are included in a single dataset, it is however not possible to remove a particular file. In this case, the whole dataset has to be removed and then recreated without including this file.

4.4. Particles and multiparticles

When starting a session of MADANALYSIS 5, lists of predefined particle and multiparticle labels are loaded into the memory, as it is mentioned in the

example of Section 3. These lists are taken from the content of the directory madanalysis/input which contains several files with label definitions. Not all the files present in this directory are loaded when a new session of MAD-ANALYSIS 5 starts. Indeed, according to a running in the aim of analyzing parton-level, hadron-level or reconstructed-level events, only some of the files are appropriate, and only these are in imported when MADANALYSIS 5 is launched.

For parton-level analyses, the particle labels are imported from the file particles_name_default.txt whilst the multiparticle labels are read from the file multiparticles_default.txt. These two files contains standard names for the Standard Model and the Minimal Supersymmetric Standard Model particles (e.q., an antimuon is denoted by mu+ and the lightest neutralino by n1) as well as appropriate definitions for the multiparticle hadronic and invisible (see below). For hadron-level analyses, only a single file with a list of 415 hadrons (see Ref. [58]), hadron_default.txt, is imported. Finally, for analyses at the reconstructed-level, the two lists reco_default.txt and multiparticles_reco_default.txt are imported. They contain labels for high-level reconstructed objects necessary for such a level of sophistication of the Monte Carlo simulations. The set of predefined labels is rather short and consists in e+, e-, mu+, mu-, ta+, ta- for charged leptons, j and b and nb for jets, b-tagged jets and non-b-tagged jets. In addition, several (intuitive) multiparticle labels are included, e, mu, ta, 1+, 1- and 1 for various combinations of charged leptons, as well as two special labels dedicated to muons and antimuons, mu-isol and mu+isol, in the case they are isolated.

Let us note that if MADANALYSIS 5 has been installed in the directory where MADGRAPH 5 is unpacked, the predefined particle and multiparticle labels are directly imported from MADGRAPH 5 rather than from the directory madanalysis/input of MADANALYSIS 5.

Two multiparticles, denoted by the labels hadronic and invisible, have a special role and are essential in any analysis. They are respectively related to the computation of observables related to the missing energy and to the hadronic activity, respectively. Therefore, if they are not included in the imported files, they are automatically created by MADANALYSIS 5 when a new session starts. Moreover, their special role prevent them from being deleted with the command remove.

Full support is also offered to load entire UFO models [38]. The UFO format is automatically detected by MADANALYSIS 5 so that, in order to load the model, it is enough to issue in the interpreter

import <path-to-UFO-files>

where path-to-UFO-files> is the location of the UFO model under consideration. The list of particle labels is derived from the particles.py file containing the UFO implementation of the particles of the model. Moreover, any stable, electrically and color neutral particle, with the exception of the photon, is automatically added to the invisible multiparticle object.

As already presented in Section 3 and in Table 4, the command interpreter of MADANALYSIS 5 contains two actions for printing to the screen the list of all the particle and multiparticle labels which have been defined in the current session,

display_particles display_multiparticles

Moreover, as for any instance of any class, the properties of a given particle or multiparticle object can be obtained with the action display,

display <label>

where <label> is the label of the (multi)particle under consideration. The effect of the command above is to print to the screen, in the case of a particle, the PDG-id linked to the label <label>. Similarly, for multiparticle labels, the whole list of associated PDG-ids is displayed.

As already shown in Section 3 where we have taken the example of muons and antimuons, new particle and multiparticle labels can be created, according to the needs of the user, by means of the command define,

define <label> = <identifiers>

The command above creates a label denoted by <label> and associates to it the content of the parameter <identifiers>. If the value of <identifiers> consists in one single integer number, a new particle label related to the corresponding PDG-id is created. In the case we have either several integer numbers separated by spaces or several other particle and multiparticle labels separated by spaces, a new multiparticle label is created and linked to the list of corresponding PDG-ids. The define command also allows for redefining existing labels or to add extra (multi)particles to a definition. In this last case, the user is allowed to issue commands such as, e.g.,

define <label> = <label> <identifiers>

The effect of the command above is to add the particle(s) contained in the variable <identifiers> to the definition of the label <label>.

Even if, in principle, the user is allowed to choose any name for the labels to be created, MADANALYSIS 5 contains a set of reserved keywords that cannot be used as labels, such as the symbols all, or and and (see below). In addition, the names of the different actions (see Table 4) cannot be used for (multi)particle labels.

At any time, a label not already used for the definition of an histogram or of a cut can be deleted from the memory of the computer by issuing

remove <label>

where the label to be deleted is denoted by <label>. Due to their particular nature, the labels hadronic and invisible can however never be removed.

4.5. Creating histograms

There are two types of observables that can be represented as histograms by Madanalysis 5, global observables (see Table 7), related to the entire event content, and observables related to a given type of particle (see Table 8), thus specific to only a part of the event. In addition, a few additional observables are dedicated to analyses at the reconstructed-level and are collected in Table 9.

We first address the description of the global observables. Two of them are related to the missing energy. The corresponding symbols implemented in MADANALYSIS 5 are denoted by MET and MHT and correspond to the missing transverse energy E_T and the missing hadronic transverse energy E_T , respectively. The definitions of these two quantities are not unique in the literature, and we decide to adopt the choice

$$E_T = \left\| \sum_{\text{visible particles}} \vec{p}_T \right\| \quad \text{and} \quad E_T = \left\| \sum_{\text{hadronic particles}} \vec{p}_T \right\|, \quad (1)$$

where \vec{p}_T stands for the particle transverse momentum. At the parton-level, the generic name hadronic particles stands for gluons and light and b-quarks, but not for top quarks (decaying most of the time to a W-boson and a b-quark). At the hadronic level and reconstructed level, hadronic particles are trivially hadrons and jets, respectively. These definitions are related to the default values of the multiparticle labels hadronic and invisible. We

remind that the latter can be modified by the user by means of the command define (see Section 4.4).

Even if particle objects are not explicitly tagged as visible, any non-invisible object, *i.e.*, an object whose the PDG-id is not included in the definition of the label invisible, is considered by MADANALYSIS 5 as visible. Similarly, any object which is not tagged as hadronic is considered as non-hadronic

Concerning the missing transverse energy $\not E_T$, a remark is in order. Detector simulation tools are in general internally computing the missing energy, following a built-in definition which might be different from the one of Eq. (1). This value is stored under a special tag in the event files, compliant with the LHCO format relevant for analyses at the reconstructed-level. When this is read by MADANALYSIS 5, the value of the missing transverse energy is subsequently imported and always supersedes the one which would have been computed by employing Eq. (1).

Parallel to those observables related to the invisible sector of the events, two important global observables are related to the visible objects, the total transverse energy E_T and the total transverse hadronic energy H_T . In Madanalysis 5, these kinematical quantities are defined by

$$E_T = \sum_{\text{visible particles}} ||\vec{p}_T|| \quad \text{and} \quad H_T = \sum_{\text{hadronic particles}} ||\vec{p}_T||, \quad (2)$$

and are associated to the symbols TET and THT, respectively.

Creating histograms associated to the four variables introduced above follows the standard syntax of the command plot,

where the value of the variable <observable> corresponds here either to MET, MHT, TET or THT. The number of bins, the value of the lowest bin and the one of the highest bin are optional information which can be passed through the symbols <nbins>, <min> and <max>, respectively. If left unspecified, MADANALYSIS 5 uses built-in values.

The effect of the command plot is to create a new instance of the class selection which has been designed to handle histograms (and cuts, as it is shown in Section 4.6). The labeling of the selection objects is internally handled by MADANALYSIS 5. The first time that the command plot is issued in the current session of MADANALYSIS 5, the label selection[1] is assigned

to the corresponding histogram. The second occurrence of the command plot leads to the creation of the object selection[2], and so on. The full list of instances of the selection class which have been created can be displayed to the screen with the display command,

display selection

As for particle and multiparticle labels, typing in the interpreter

display selection[<i>]

allows to display the properties of the <i>th created selection object. This <i>th selection can always be deleted from the memory of the computer with the action remove,

remove selection[<i>]

After the removal of a specific selection, the numbering of the different created selections is automatically adapted by MADANALYSIS 5 so that the sequence of the integer numbers is respected, without any hole. Even if handled automatically, the ordering of the selections can be modified by the user by means of the command swap. Issuing

swap selection[<i>] selection[<j>]

results then in the exchange of the $<i>^{th}$ and $<j>^{th}$ instances of the selection class

Objects of the selection class have several attributes, as shown in Table 6. In this section, we however only focus on selection objects related to histograms. In the case of cuts, we refer to Section 4.6. The number of bins, the value of the lowest bin of an histogram and the one of its highest bin are stored in the attributes nbins, xmin and xmax, respectively. Their values are fixed at the time the command plot is issued in the interpreter. They can however be further modified by the use of the command set, as for any other attribute of an object, by typing in the command interface, e.g.,

This commands allows to set the value of the highest bin of the histogram associated to the object selection[<i>] to 100.

Two of the other attributes of the selection class are related to the scales used when drawing the axes of the histograms. By default, a linear

Table 6: List of the attributes of the selection class

logX If set to true, this enforces a logarithmic scale for the x-axis. If set to false (default), a linear scale is used.

logY If set to true, this enforces a logarithmic scale for the y-axis. If set to false (default), a linear scale is used.

nbins Number of bins of the histogram. The value taken by

this attribute must be an integer.

This refers to the observable to

This refers to the observable to be used for ordering particles of the same type. The value of this option has to be anything among ETAordering (pseudorapidity ordering), ETordering (transverse-energy ordering), Eordering (energy ordering), Pordering (momentum ordering), PTordering (transverse-momentum ordering, the default choice), PXordering (ordering according to the x-component of the momentum), PYordering (ordering according to the y-component of the momentum) or PZordering (ordering according to the z-component of the momentum).

stacking_method

xmax

When several datasets are represented on an histogram, the different curves can be stacked (stack), superimposed (superimpose) or normalized to unity (normalize2one), including superimposing. The default value is auto and then refers to the properties of the attribute stacking_method of the object main (see Section 4.7).

statuscode By default, only final state particles are considered in histograms. This attribute allows to consider instead initial or intermediate particles, or all the particle content. The (intuitive) allowed values are allstate, initialstate, finalstate (default) and interstate.

titleX The title of the x-axis of the histogram, given as a string.

titleY The title of the y-axis of the histogram, given as a string.

xmin Central value of the lowest bin of the histogram. This must be a real number.

Central value of the highest bin of the histogram. This

must be a real number.

scale is employed. Logarithmic scales can be enforced through the boolean attributes logX and logY which can then take the values true or false (default choice). For instance, issuing

```
set selection[<i>].logY = true
set selection[<j>].logX = false
```

ensures the usage of a logarithmic scale for the y-axis of the histogram related to the object $selection[\langle i \rangle]$ and the one of a linear scale for the x-axis of the histogram associated to the object $selection[\langle j \rangle]$.

Another important attribute related to the layout of the histograms concerns the stacking method used for the curves related to the different datasets created by the user. By default, the curves are drawn as stacked one above each other, but this can be modified through the attribute stacking_method of the class selection. For instance, focusing on the object selection[<i>], the command

```
set selection[<i>].stacking_method = <value>
```

allows to change the employed stacking method to the value <value>. By default, this attribute is set to the value auto. This means that the value of the attribute stacking_method of the object main (see Section 4.7) is employed. The other allowed choices are stack, superimpose and normalize2one. In the first case, the curves are all stacked, whilst in the second case, they are superimposed. The last possibility for the attribute stacking_method, i.e., normalize2one, has been designed for comparing the shapes of the curves related to the different datasets. Here, the normalization of each curve is set to one and they are drawn superimposed. Let us note that this consists in an helpful feature when optimizing selection cuts.

To achieve the description of the functions allowing to tune the layout of an histogram, the user is allowed to change the titles of the x-axis and y-axis by means of the attributes titleX and titleY of the selection class,

```
set selection[<i>].titleX = <string>
set selection[<i>].titleY = <string'>
```

The effect of the commands above leads to the use of the strings <string> and <string'> as titles for the x-axis and y-axis of the histogram represented by the object selection[<i>], respectively.

The attributes of an instance of the **selection** class related to an histogram can also be directly set when issuing the command **plot**,

Table 7: List of the global observables that can be represented by an histogram		
NAPID	The particle multiplicity of the events after mapping particles	
	and antiparticles.	
NPID	The particle multiplicity of the events.	
MET	The missing transverse energy as defined in Eq. (1). At the	
	reconstructed-level, the missing energy is however directly	
	read from the LHCO file.	
MHT	The missing transverse hadronic energy as defined in Eq. (1).	
SQRTS	Partonic center of mass energy. Only available for parton-level	
	and hadron-level event samples.	
TET	The visible transverse energy as defined in Eq. (2).	
THT	The visible transverse hadronic energy as defined in Eq. (2).	

plot <observable> <nbins> <min> <max> [options]

The optional pattern [options] stands for logX, logY, stack, superimpose, normalize2one or the value of any of the other attributes presented below, statuscode and rank. Multiple keywords are allowed (but however only one for each attribute). In this way, the values of several options can be passed at one time, separated each by a space character. For instance, creating an histogram representing the visible transverse hadronic energy with a y-axis represented with a logarithmic scale and where all the curves are drawn superimposed can be done by issuing the command

plot THT [logY superimpose]

The binning is then automatically handled by MADANALYSIS 5, since not provided by the user.

In addition to the kinematical global observables E_T , $\not\!E_T$, H_T and $\not\!H_T$ introduced in the beginning of this Section, the partonic center-of-mass energy can be represented by an histogram by typing the command

plot SQRTS

Finally, two other global observables are available. The latter are related to the particle multiplicity of the events. In order to draw the associated histograms, it is enough to enter in the command interface the two commands plot NPID
plot NAPID

The first command, plot NPID, generates an histogram where one bin is associated to each possible type of final state particle, the height of the bin being related to the multiplicity of the corresponding particle within the whole sample. Hence, if several particles of the same type are present in one specific event, they correspond to several entries in the histogram. The second command above, plot NAPID, produces a similar histogram after mapping antiparticles and particles. For both types of histograms, the labels of the x-axis correspond to the particle label imported in MADANALYSIS 5. If non-existing, the PDG-ids are used instead. We emphasize that these histograms are special histograms dedicated to the task of getting an idea about the particles present in the input sample. To compute the multiplicity of a given particle species, we refer to the observable N described below.

The second large class of observables that can be represented by histograms in Madanalysis 5 refers to the kinematical properties of the particles contained in the events. Hence, distributions such as the invariant mass or the transverse-momentum of given particle species can be computed. The complete list of implemented observables can be found in Table 8.

Creating histograms associated to a given property **<observable>** of a specific particle represented by the label **<label>** is also based on the command **plot**. In this case, the syntax is however slightly different as for the global observables. The symbol **<observable>** has to be seen as a function which takes as argument the label associated to the particle under consideration,

plot <observable>(<label>)

For instance, if mu+ stands for the particle label related to antimuons, the command

plot PT(mu+)

results in representing by an histogram the transverse-momentum distribution of all the antimuons in the sample. As above, if several antimuons are included in one single event, they contribute to several entries in the histogram. In the case of the relative distance between two particles (denoted by DELTAR), two particle labels <label1> and <label2> are required,

Table 8: List of the kinematical observables that can		
be represented by histograms		
BETA	Velocity $\beta = v/c$.	
DELTAR	Relative distance between two objects in the $\eta - \phi$ plane.	
E	Energy.	
ET	Transverse energy.	
ETA	Pseudorapidity.	
GAMMA	Lorentz factor.	
M	Invariant mass.	
N	Particle multiplicity.	
MT	Transverse mass.	
P	Norm of the momentum.	
PHI	Azimuthal angle of the momentum.	
PT	Norm of the transverse momentum.	
PX	Projection of the momentum on the x-axis.	
PY	Projection of the momentum on the y -axis.	
PZ	Projection of the momentum on the z -axis.	
R	Position of the object in the $\eta - \phi$ plane.	
THETA	Angle between the momentum and the beam axis.	
Y	Rapidity.	

plot DELTAR(<label1>, <label2>)

As illustrated above, when studying the kinematical properties of a given particle species, the normalization of the histograms reflects the total number of particles of this type included in the full sample. This can be different from the total number of events, since one single event could describe a final state with zero, one or several particle of the considered type, which then corresponds to zero, one or several entries in the produced histograms³.

The syntax detailed above is also valid for multiparticle objects. In this case, the label <label> is related to an instance of the multiparticle class. Histograms are created by treating on the same footings all the particles

³In this Section, we do not consider the luminosity which also affects the normalization of the histograms. This feature is described more in details in Section 4.7.

linked to the label <label>. For example, the set of commands

```
define <multi> = <particle1> <particle2>
plot <observable>(<multi>)
```

defines, in a first step, a multiparticle label denoted by <multi> and linked to the particle species <particle1> and <particle2> (see Section 4.4). In a second step, a property represented by the observable <observable> is investigated. For a specific event, each particle of the type <particle1> or <particle2> is associated to one entry in the histogram.

As shown above, if several particles of the considered type appear in a specific event, they always correspond to several entries in the histograms. In phenomenological analyses, it is however often more relevant to only consider the leading particle, *i.e.*, the one which has the highest value of a kinematical variable such as the transverse momentum or the energy. The squared brackets '[]' allow to access leading, next-to-leading, *etc.*, particles. For instance, the histogram resulting from the command

plot <observable>(<label>[<i>])

represents the distribution of the observable <observable> for the particle of the type <label> with the <i>th largest transverse momentum, <i> being a positive integer. Similarly, a negative value of the parameter <i> corresponds to a pointer to the particle with the <i >th smallest transverse momentum. Events where the number of particles of the type under consideration is smaller than the absolute value of <i > are ignored at the time of the creation of the histogram.

By default, the ordering variable employed in Madanalysis 5 is the transverse momentum. Other possible choices exist and the information is passed, for a given histogram, through the value of the attribute rank of the selection class. Hence, considering the instance of the class selection [<i>], one can modify the ordering variable by issuing

The parameter <value> can take any value among ETAordering (pseudorapidity ordering), ETordering (transverse-energy ordering), Eordering (energy ordering), Pordering (momentum ordering), PTordering (transverse-momentum ordering), PXordering (ordering according to the x-component of the momentum), PYordering (ordering according to the y-component of

the momentum) and PZordering (ordering according to the z-component of the momentum).

For all the histograms produced so far, MADANALYSIS 5 only uses information related to the final state particles. However, for phenomenological purposes, it is often interesting to investigate the properties of the intermediate particles or those of final state particles issued from the decays of a specific type of intermediate particle. This requires, of course, that the relevant information is available. This is not the case for event samples at the reconstructed-level since these features are totally absent from event files under the LHCO format. Contrary, the user can access, for hadron-level or parton-level event files, to the whole particle history by a set of mother-to-daughter relations. To benefit from those relations, there exist two special functions in MADANALYSIS 5.

Firstly, the symbol '<' links one particle or set of particles to their direct mother. The command line

plot <observable>(<type1> < <type2>)

allows hence to study a given property, which is represented by the symbol <observable>, of the particles of type <type1> included in the final state of the events under consideration. However, in order to correspond to an entry in the histogram, the particles of type type1 must be issued from the direct decay of a particle of type <type2>. Doubling the symbol '<', i.e., replacing it by '<<', allows to remove the restriction of a direct decay. Finally, the usage of these symbols recursively

allows to focus on entire decay chains. Here, we investigate the properties of the particle species <type1>, but only for particles of type type1 issued from a direct decay of a particle of type <type2>, which is itself issued from a decay of a particle represented by the label <type3>, this last decay possibly occurring in several steps.

Secondly, kinematical properties of the intermediate particles can be directly investigated through the option statuscode of the selection class. As suggested above, by default, only final state particles are considered. This corresponds to the value finalstate of the attribute statuscode. Issuing in the interpreter

set selection[<i>].statuscode = interstate

indicates that, for the selection selection[<i>], we are not considering final state particles anymore, but only intermediate states. On the same footings, setting statuscode to the value initial state allows to focus on the initial state particles only, whilst setting it to allstate allows to consider equivalently initial state, final state and intermediate state particles.

Key observables to design efficient selection cuts in an analysis are in general related more than one single particle. For instance, highlighting a new Z' gauge boson in Drell-Yan events and estimating its mass with a good accuracy rely on the invariant-mass distribution of the produced lepton pair. All the functions of Table 8, but the DELTAR observable which requires exactly two arguments, can take an arbitrary number of arguments. This allows to combine particles before computing the kinematical distribution to be represented. Hence, the two equivalent command lines

```
plot <observable>(<prtcl1> <prtcl2>)
plot v<observable>(<prtcl1> <prtcl2>)
```

lead to the creation of an histogram showing the distribution of the observable <observable>. The observable is computed on the basis of the combined four-vector built from the sum of the four-momentum of the particles cprtcl1> and combined vectorially (other options are shown below). If, for a given event, several pairs of particles cprtcl1> and cprtcl2> can be formed, each possible pair leads to one different entry in the histogram. Hence, the histogram describing the invariant mass of a muon pair can be created by issuing

```
plot M(mu+ mu-)
```

where, as before, the binning is automatically handled by MADANALYSIS 5. The observable that is computed corresponds to the norm of the sum of the four-momentum of the muon and the one of the antimuon. This syntax can straightforwardly be generalized to multiparticles or to combinations of more than two particles. A remark is however in order here. If <multi> denotes the label associated to an instance of the multiparticle class, issuing

```
plot <observable>(<multi> <multi>)
```

generates an histogram where each entry corresponds to one different combination of the particles represented by the multiparticle <multi> MAD-ANALYSIS 5 indeed forbids double-counting any combination.

By default, the particles are combined by adding vectorially their four-momentum, i.e.,

$$p_{\mu} = \sum_{i} p_{\mu}^{i} ,$$

where p_{μ} is the resulting four-momentum and p_{μ}^{i} are the four-momenta of the particles to be combined. The four-vector p^{μ} is the one which is used when computing the value of the observable to be represented on the corresponding histogram. Madanalysis 5 offers additional ways to perform this combination. The sum could be, in contrast, done scalarly, *i.e.*, by firstly computing the considered observable for each of the particles to be combined and secondly adding the results. To select this option, the user must add a prefix 's' in front of the name of the observable when typing the command in the interpreter,

```
plot s<observable>(<prtcl1> ...)
```

where the dots stand for the list of particles to be combined. In the case the user wants to combine all particles of a given type cprtcl> in an event, the reserved keyword all can be used in order to simplify the syntax,

plot <observable(all <prtcl>)

If two (and only two) particles are combined, differences can also be computed, rather than sums. To allow for this option, it is enough to include one of the the prefixes 'dv', 'vd', 'd', 'ds', 'sd' or 'r' in front of the symbol of the observable to be computed. For the options 'dv', 'vd' and 'd', vectorial differences are considered. Hence, the result of one of the equivalent commands

```
plot dv<observable>(<prtcl1> <prtcl2>)
plot vd<observable>(<prtcl1> <prtcl2>)
plot d<observable>(<prtcl1> <prtcl2>)
```

is to subtract the two four-momenta from each other,

$$p_{\mu} = p_{\mu}^1 - p_{\mu}^2 \; ,$$

 p_{μ}^{1} being the four-momentum of the particle associated to the label <prtcl1> and p_{μ}^{2} the one of the particle associated to the label <prtcl2> and then compute the observable <observable> from the resulting four-vector p^{μ} . In the case of the prefixes 'ds' and 'sd',

Table 9: List of the additional observables that can be represented by histograms for reconstructed events

EE_HE Ratio of the electromagnetic and hadronic energy for a

given object.

HE_EE Ratio of the hadronic and electromagnetic energy for a

given object.

NTRACKS Number of tracks in a jet. This returns zero for non-jet

objects.

```
plot ds<observable>(<prtcl1> <prtcl2>)
plot sd<observable>(<prtcl1> <prtcl2>)
```

scalar differences are computed, *i.e.*, the observable **<observable>** is computed for each of the particles **<prtcl1>** and **<prtcl2>** and the results are subtracted from each other. Relative differences can also be computed by means of the prefix 'r',

```
plot r<observable>(<prtcl1> <prtcl2>)
```

In this case, the results consists in the scalar difference of the values of the observable computed for the two particles cprtcl1> and cprtcl2> taken individually, which is however given relatively to the value of the observable for the first particle represented by the label cprtcl1>. This corresponds then to the quantity

$$\frac{\texttt{()} - \texttt{()}}{\texttt{()}}$$

If the considered observable vanishes for the particle labeled by <prtcl1>, the value zero is returned.

A last way for combining observables is related to the reserved word and. When the considered observable has to be evaluated for several possible combination of particles, the user can use the keyword and to efficiently implement the corresponding command in MADANALYSIS 5,

The command line above computes the observable represented by the symbol <observable>, firstly, for the vectorial combination of the particles <p1> and <p2> and secondly, for the vectorial combination of the particles <p3> and <p4>. The two distributions are eventually summed.

Three additional observables can be represented by histograms in the case of fully-reconstructed objects. To allow for generating histograms for these observables, MADANALYSIS 5 must be run in the reconstructed-level mode. These observables consist in the ratio between the hadronic and the electromagnetic energy for a given object (the ratio between the energy deposited in the electromagnetic and hadronic calorimeters of a detector), the inverse ratio and the number of tracks within a (reconstructed) jet. For objects different from a jet, this last observable always returns the number zero. The associated symbols are HE_EE, EE_HE and NTRACKS and their definitions are collected in Table 9. The corresponding histograms can be created by following the usual syntax,

plot <observable>(<label>)

Let us note that for these three observables, combining objects is not supported, *i.e.*, only one single (multi)particle label can be passed as an argument.

4.6. Selection cuts

In Madanalysis 5, the process of event selection is based on two equivalent classes of kinematical cuts which can be applied to the imported datasets. The program offers to the user the two choices of either selecting or rejecting events in the case a certain condition is fulfilled. This task can be performed at the level of the command line interpreter of the program by means of the two actions select and reject. The associated syntax is very intuitive and reads

```
select <condition> [<options>]
reject <condition> [<options>]
```

For the first (second) command, events are selected (rejected) if the condition <condition> is satisfied. Hence, the command

```
reject PT(mu) > 50
```

leads to the rejection of all the events where at least one muon with a transverse momentum $p_T > 50$ GeV is found, whilst issuing

select M(e+e-) > 100

allows for the selection of all the events where we have a electron-positron pair with an invariant-mass larger than 100 GeV. Internally, the effects of the commands above are to create instances of the selection class with special properties. Contrary to the command plot which is related to the creation of histograms, the commands select and reject lead to the production of tables of cut efficiencies. Consequently, the only attributes of the selection class which are relevant are rank and statuscode (see Table 6). They can be either passed directly at the time of typing-in the commands, by including the desired values as the optional parameter <options> above, or at a later stage by means of the command set (see Section 4.5).

Modifying the rank attribute plays only a role if the ordering of the particles is a necessary information for a good application of the condition, as e.g., if we are constraining some observable related to the leading or next-to-leading particles. Furthermore, setting the option statuscode to a non-default value allows to apply, if needed by the user, cuts on initial or intermediate states rather than on final states only.

The condition < condition > must be given according to the pattern

<observable> <logical-operator> <value>

where the observable <observable> can be any observable from Tables 7, 8 and 9, computed for a given particle or for any combination of particles, with the exception of the global variables related to the symbols NPID and NAPID⁴. The supported logical operators are '>' (greater than), '>=' (greater than or equal to), '<' (smaller than), '<=' (smaller than or equal to), '==' (equal to) and '!=' (different from).

Conditions can also be combined by using one or several of the connecting keywords and (logical and) and or (logical inclusive or), such as in

<cond1> <connector1> <cond2> <connector2> <cond3>

The condition above consists in the combination of the three conditions <cond1>, <cond2> and <cond3> by employing the two connecting keywords <connector1> and <connector2> being and or or. Moreover, brackets are

⁴To implement selection cuts on the multiplicity of a given particle species, cuts on the observable N have to be performed, rather than on the global observable NPID and NAPID.

also authorized by the syntax for handling more complex conditions. In the special case both upper and lower limits are imposed on a given observable, the condition can be easily implemented by means of the keyword and. However, there exists a more compact syntax

```
<value1> <logical-operator> <obs> <logical-operator> <value2>
```

which could be employed. The logical operator <logical-operator> is here used twice. The parameter <obs> denotes the observable and <value1> and <value2> the imposed bounds.

So far, we have supposed that all the objects present in each event can be used for applying selection cuts. However, this is barely the case in most realistic phenomenological analyses, since, for example, too soft particles are in general omitted. This feature can also be implemented in analyses performed with MADANALYSIS 5 by means of the commands select and reject, but following a slightly different syntax as the one shown before,

```
select (<particle>) <condition> [<options>]
reject (<particle>) <condition> [<options>]
```

In the case the particle candidate is rejected, the event is considered as without containing this particle. Let us note that whilst selecting and rejecting events have a direct influence on the signal over background ratio, selection or rejecting candidates to a particle type only affects the number of entries for one event in the histograms.

4.7. Executing an analysis and displaying the results

Once an analysis has been implemented, it must be passed to Sample-Analyzer, the C++ kernel of Madanalysis 5, for execution by means of the command

submit <dirname>

A directory named dirname is created and all the files necessary for SAMPLE-ANALYZER to properly run are generated and included in this directory. The code is further compiled and linked to the external static library of MAD-ANALYSIS (see Section 3.1), and the execution of the resulting program is eventually managed by MADANALYSIS 5.

This execution starts with the reading of the event samples under consideration and their storing in the memory of the computer according to a format internal to MADANALYSIS 5. All the histograms required by the user are then sequentially created, including the application to all the events of the defined selection cuts. As an output, a ROOT file [56] is generated so that the analysis can be accessed later, as, e.g., directly in the ROOT framework or in a new session of MADANALYSIS 5. In this last case, the SAMPLEANALYZER (executed) job can be imported by means of the command import

import <dirname>

where the directory <dirname> contains the analysis previously performed.

After modifications such as asking for the creation of a new histogram or the application of a new selection cut, the user does not have to submit the SampleAnalyzer job entirely again. A much faster option, saving a sensible amount of computing time, consists in the command

resubmit

This updates the already existing SampleAnalyzer directory and only the differences with respect to the original analysis are executed.

Once the ROOT file has been created by SAMPLEANALYZER, the command preview allows for the display of a single histogram,

preview selection[<i>]

where in the generic example above, the <i>th histogram is asked to be displayed to the screen, in a ROOT popup window. The command preview only works for displaying histograms. Consequently, previewing the efficiency table associated to a selection cut is not possible.

A more complete report, with all the selection cuts, efficiency tables and histograms can be generated by issuing in the command line interface one of the three commands

```
generate_html <html-dirname>
generate_latex <tex-dirname>
generate_pdflatex <pdftex-dirname>
```

The first command, <code>generate_html</code>, <code>generates</code> the report under the HTML format and stores the files in the directory <code><html-dirname></code>. The second and third commands, <code>generate_latex</code> and <code>generate_pdflatex</code>, are related to the creation of <code>TeX</code> files to be compiled with the help of the shell commands <code>latex</code> and <code>pdflatex</code>, respectively. In this two cases, the <code>TeX</code> files are already compiled by <code>Madanalysis 5</code>, if the <code>latex</code> and <code>pdflatex</code> commands are available on the system of the user.

Histograms are exported to (non-ROOT) figures under either the Encapsulated PostScript (.eps) format, required for a proper compilation with latex, or to the Portable Network Graphics (.png) format for HTML files or TEX files to be compiled with pdflatex. In addition, each figure is also saved as a ROOT macro which can be modified by the user for further processing. Once generated, the report can be immediately displayed by typing-in the command

open <report-dirname>

which opens the report in a web browser.

The analysis, and thus the generation of the histograms and the efficiency tables, is so far based on the default configuration of Madanalysis 5. This configuration can be modified by superseding the default values of the attributes of the object main, which are listed in Table 10 and Table 11.

Two options allow to control the normalization of the histogram, lumi and normalize. This last attributes defines the way the histograms are normalized. The allowed values are none, lumi and lumi_weight and can be set as for any other attribute of any class,

set main.normalize = <value>

For the first choice (<value> = none), the total number of events included in each dataset is kept. The two other options imply that the histograms are normalized with respect to the integrated luminosity. The difference lies in the weight which can be possibly associated to each event (see Section 4.3), which can be ignored (<value> = lumi) or included (<value> = lumi_weight). This last possibility consists in the default choice.

By default, all the histograms are normalized to an integrated luminosity of 10 fb⁻¹. This value can be updated by modifying the attribute lumi of the object main,

set main.lumi = <new-value>

Table 10: List of the attributes of the object main set main.coption> = <value>

currentdir Current directory in which any directory created by MADANALYSIS 5 is stored.

lumi This allows to modify the value of the integrated luminosity, in fb^{-1} , used for the normalization of the histograms. The default value is 10 fb^{-1} .

normalize This defines the way the histograms are normalized. The allowed choices are the number of events included in the different datasets (none), the integrated luminosity without (lumi) or accounting for the event weights associated to each dataset (the default choice, lumi_weight).

This fixes the way the uncertainty on the signal over background ratios is computed by MADANALYSIS 5. The attribute <value> is the corresponding analytical formula passed as a valid Python expression given as a string. It has to depend on S (number of signal events), B (number of background events), ES (uncertainty on the number of signal events) and EB (uncertainty on the number of background events). It is automatically handled for the most simple cases (see Eq. (3)).

SBratio This fixes the way signal over background ratios are computed by Madanalysis 5. The attribute <value> is the corresponding analytical formula passed as a valid Python expression given as a string. It has to depend on S (number of signal events) and B (number of background events).

stacking_method

SBerror

When several datasets are represented on histograms, the different curves can be stacked (stack, default), superimposed (superimpose) or normalized to unity (normalize2one), including superimposing. When the stacking_method attribute of an instance of the class selection is set to auto, the value of main.stacking_method is employed.

where the new value of the integrated luminosity, $\langle new-value \rangle$, is given in fb⁻¹.

Once all the datasets have been defined as part of the signal or background samples, MADANALYSIS 5 can compute automatically the signal (S) over background (B) ratio, and thus the efficiency of each selection cut. This feature is related to an attribute of the object main denoted by SBratio. It refers to an analytical formula, expressed as a valid PYTHON expression, which indicates how the signal over background ratio must be calculated⁵. When implementing this formula, the symbols related to the signal and background number of events are S and B, respectively. By default, the signal over background ratio is computed according to S/B. This can be modified through the command set, by issuing in the interpreter,

where <formula> is, as sketched above, a valid Python expression depending on the two variables S and B. For instance, the command

enforces the signal over background ratio r to be computed according to

$$r = \frac{S}{\sqrt{S+B}} \ .$$

In the case the signal over background ratio is undefined due, e.g., to the evaluation of the squared root of a negative number or to a division by zero, the value zero is returned. Let us emphasize that it is safer for the user to use real numbers when typing-in the analytical expression formula, rather than integer numbers. We indeed recall that 1/2 is evaluated as 0 whilst 1./2. returns 0.5.

It is fundamental to associate an uncertainty to the signal over background ratio. The way to compute this quantity is related to the SBerror attribute of the object main. As for SBratio, it refers to an analytical formula, given as a valid Python expression, which indicates how the uncertainty on the signal over background ratio must be computed. For the three

⁵The formula is internally stored as an instance of the TFormula class. This structure is defined in the ROOT library linked to MADANALYSIS 5, and all the associated attributes can therefore be used. We refer to the ROOT manual for more information [56].

choices

$$r_1 = \frac{S}{B}$$
, $r_2 = \frac{S}{S+B}$ and $r_3 = \frac{S}{\sqrt{S+B}}$, (3)

as well as for the three additional cases obtained when S and B are exchanged, MADANALYSIS 5 automatically detects the formula stored in the SBratio attribute. It then updates accordingly the uncertainty, setting the SBerror attribute to the values Δr_i given by

$$\Delta r_1 = \frac{\sqrt{B^2(\Delta S)^2 + S^2(\Delta B)^2}}{B^2} ,$$

$$\Delta r_2 = \frac{\sqrt{B^2(\Delta S)^2 + S^2(\Delta B)^2}}{(S+B)^2} ,$$

$$\Delta r_3 = \frac{\sqrt{(S+2B)^2(\Delta S)^2 + S^2(\Delta B)^2}}{2(S+B)^{3/2}} ,$$

where ΔS and ΔB are the uncertainties on the signal and on the background number of events. If the user wants to use another formula or to set himself the SBerror attribute, the command set has to be employed,

```
set main.SBerror = '<formula>'
```

where $\langle formula \rangle$ depends this time on the number of signal and background events S and B as well as on the associated uncertainties ES ($\equiv \Delta S$) and EB ($\equiv \Delta B$). For instance, implementing by hand the error Δr_1 above would give

```
set main.SBerror = 'sqrt(B**2*ES**2+S**2*EB**2)/B**2'
```

Four specific attributes of the object main concern muon isolation when MADANALYSIS 5 is run in order to analyze reconstructed-level events. The user has the possibility to choose the algorithm which is employed by MADANALYSIS 5 when defining an isolated muon. Two choices are implemented and can be adopted by issuing one of the two commands

```
set main.isolation.algo = DELTAR
set main.isolation.algo = SUMPT
```

In the first case, a muon is tagged as isolated when no track lies inside a cone around the muon. The size of this cone can be specified by the user by typing in the command interface

Table 11: List of the attributes of the object main associated to the isolation of the muons.

set main.isolation.<option> = <value>

This specifies the algorithm to be employed for muon isolation. The allowed choices are DELTAR (default) and SUMPT. In the first case, we require that no tracks lies inside a cone around the muon. In the second case, the sum of the transverse-momentum of all the tracks inside the cone and the ratio of the summed transverse energy of these tracks over their summed transverse momentum must be lower than values specified by the user. For the second algorithm, the size of the cone is fixed by the detector simulation tool.

deltaR This specifies the radius of the isolation cone to be used

by the DELTAR isolation algorithm.

This specifies the value of the ratio of the sum of the transverse energy of the tracks lying in the isolation cone over the sum of their transverse-momentum to be used

by the SUMPT algorithm.

algo

ET_PT

sumPT This specifies the transverse-momentum threshold to be

used by the SUMPT algorithm.

```
set main.isolation.deltaR = <value>
```

where <value> is a floating-point number. For the second algorithm, a muon is considered as isolated when the sum of the transverse momentum of all the tracks lying in a cone around the muon is lower than a value <value>. The latter can be specified by typing

```
set main.isolation.sumPT = <value>
```

In addition, the ratio of the sum of the transverse energy of these tracks over the sum of their transverse momentum has also to be lower than a value <value'> to be specified. This value is provided by means of the command

```
set main.isolation.ET_PT = <value'>
```

On a fairly different line, the last attribute of the object main which can be modified by the user is denoted by currentdir

```
set main.currentdir = <dirname>
```

It fixes the path to the directory in which any file and/or directory created by MADANALYSIS 5 is stored.

5. Madanalysis 5 for expert users

Besides its user-friendliness, the way of using Madanalysis 5 described in Section 4 is (obviously) restricted by the set of functionalities that have been implemented. The latter allow, in general, to perform rather traditional and standard analyses but might not be sufficient for more sophisticated and/or exotic investigations. For example, the normal running mode of the program does not allow to create an histogram related to either the distribution of a new observable or to the one of an existing observable that needs to be computed in a reference frame different from the laboratory reference frame. Along the same lines, selection cuts must match the pattern presented in Section 4.6, which forbids any other selection than cutting on the implemented kinematical variables by means of assigning a lower and/or upper bound on the result of the computation of the associated observable. Finally, as a last example, two- or three-dimensional histograms, necessary for correlation studies, are not included.

In order to overcome the above-mentioned restrictions as well as any type of, even unforeseen, limitations, MADANALYSIS 5 comes with an expert mode

of running. In this case, the possibilities are only limited by the programming skills of the user and his originality in designing the analysis. The user is asked to implement himself the entire analysis, knowing that he is able to benefit from all the strengths of the SampleAnalyzer framework. The latter comes indeed with its own set of reading routines for the event samples, its own data format and a large class of functions and methods ready to be employed. In addition, an automated compilation of the analysis C++ files as well as a programming error management system are included.

As already presented in Section 4.1, the expert mode of MADANALYSIS 5 can be cast by issuing in a shell one of the three commands

5.1. The SampleAnalyzer framework

In the expert mode, there are two possibilities in order to create an analysis, *i.e.*, to implement the C++ source and header files which are automatically generated by MADANALYSIS 5 in its normal mode of running. Either one can start from and extend an existing analysis already generated by MADANALYSIS 5, or one can design a new analysis from scratch.

In the first case, a working directory has already been created through the issue, in the command line interface of MADANALYSIS 5, of the command submit (see Section 4.7). Consequently, at least one dataset has been imported (see Section 4.3) and at least one histogram has been created (see Section 4.5). The created working directory contains three subdirectories, SampleAnalyzer, lists and root. The first of these directories, SampleAnalyzer, includes all the files necessary for having the C++ kernel of MADANALYSIS 5 properly running and executing the analysis implemented by the user. When issuing the command submit in the command interface, all the PYTHON commands entered by the user is translated into a set of three C++ files, user.cpp, user.h and analysisList.cpp, stored in the sub-directory Analysis of SampleAnalyzer.

The two other sub-directories included in the working directory, lists and root, are dedicated to the input and output files, respectively. As mentioned in Section 4.3, the imported events are gathered into different datasets. For each dataset, a single list, containing the paths to the various associated event files, is stored in the directory lists. After being executed, Sample-Analyzer creates a Root file for each of the defined datasets. These files,

stored into the directory root, contain the necessary information to create the histograms requested by the user.

When the user starts an analysis from scratch, the situation is similar to the one described in the paragraphs above, with the exception that the working directory is initialized directly from the shell, when MADANALYSIS 5 is launched. When typing in a shell the command

bin/ma5 --expert

or equivalently any of the three commands recalled in the introduction of this Section 5, MADANALYSIS 5 indeed asks a series of questions to the user such as the name of the working directory to be created or the chosen label for the analysis to be created (see below).

As a result, a working directory is created, together with the sub-directory SampleAnalyzer which contains a blank analysis. The implementation of this analysis, as the one of any analysis, is divided into the implementation of three core functions. The latter have to be provided by the user in the file user.cpp which comes together with the corresponding header file user.h. Both files are stored in the sub-directory SampleAnalyzer/Analysis, together with an additional file, analysisList.cpp, that contains the list of all the analyses which are/will be implemented in the working directory under consideration. After the creation of a fresh working directory, this list only contains a single analysis, the blank analysis pre-implemented in the file user.cpp. However, there is not any limitation on the number of analyses which can be included and nothing prevents this list from becoming very large.

Once the three C++ files introduced above have been created, MAD-ANALYSIS 5 exits and the user can then start implementing his analysis by modifying these files. In this Section 5.1, we do not address the way in which an analysis has to be implemented, since this is the scope of Section 5.3, Section 5.4 and Section 5.5, but we only describe the steps leading to the execution of the analysis and the generation of the output files.

In order to properly compile and execute the implemented analysis, the linking to the external dependencies, such as the ROOT header files and libraries, must be performed appropriately. This is allowed by a correct setting of the environment variables LD_LIBRARY_PATH (or rather DYLD_LIBRARY_PATH for MACOS systems), LIBRARY_PATH and CPLUS_INCLUDE_PATH prior to the compilation. This task has been rendered automatic by means of the scripts setup.sh and setup.csh included in the SampleAnalyzer sub-directory. All

the above-mentioned environment variables can be appropriately set at once by issuing

source setup.sh

in a bash shell or

source setup.csh

in a tcsh shell before compiling and/or executing SampleAnalyzer.

After this step, the analysis can be compiled with the help of the Makefile present in the SampleAnalyzer directory, assuming that the GNU Make package has been installed on the computer of the user. As for any program to be compiled with GNU Make, it is enough to issue in a shell the command

make

which also takes care of linking the external libraries to the SampleAna-Lyzer program. In addition, the command

make clean

has also been implemented and leads to the cleaning of the various subdirectories included in the current working directory.

Once compiled, the SAMPLEANALYZER core, containing the user's analysis, is ready to be executed. The only input left to be provided consists in the location paths of the event samples. They have, for a given dataset, to be collected into a single text file, denoted in the following by <datasetlist>. The SAMPLEANALYZER package is then simply run by issuing in a shell

SampleAnalyzer [options] <datasetlist>

If the event samples are collected into several datasets, SampleAnalyzer has to be run once for each of the datasets to be included in the analysis, with a different text file with the paths to the relevant event samples.

The only option supported by Sampleanalyzer consists in specifying the label of the analysis to be performed. In particular, this allows to include several analyses, each of them specified by a unique label, in one single working directory of Sampleanalyzer. The user can hence indicate which analysis to perform on run-time. If the option pattern is not provided, Sampleanalyzer lists to the screen all the analyses included in the file

analysisList.cpp, together with the corresponding labels, and asks to the user to make his choice.

The label of an analysis is defined in the corresponding header file. In the case of a fresh working directory just created by MADANALYSIS 5, the chosen name for the label is asked by the program and exported to the file user.h. We refer to Section 5.2 for more information about the way to declare labels independently from MADANALYSIS 5. Assuming that the label under consideration is denoted by <label>, SAMPLEANALYZER is launched by issuing in a shell

SampleAnalyzer --analysis=<label> <datasetlist>

As a result, the analysis defined by the label <label> is executed by SAM-PLEANALYZER and the output files are generated according to what the user has implemented in the C++ files related to the corresponding analysis.

5.2. Implementing new analyses using the analysis template

As mentioned in Section 5.1, the implementation of an analysis within the SampleAnalyzer framework consists in the writing of three files. Two of them are related to the analysis itself, *i.e.*, one C++ source file together with the associated header file. Since a given working directory of Sample-Analyzer can include many analyses, their list must be provided. This is the aim of the third file, analysisList.cpp, which is common to all the present analyses.

As presented in the previous Section, a pair of such analysis source/header files is automatically created when MADANALYSIS 5 is run in expert mode. These files are denoted by user.cpp and user.h. However, we adopt, in the following, the generic names name.cpp and name.h since the user has in fact the freedom to choose the name of the files. The only requirement is that all the filenames are different. Moreover, all these files have to be stored, together with the list of the implemented analyses included in the file analysisList.cpp, in the sub-directory SampleAnalyzer/Analysis.

The pair of the (generic) header and source analysis files name.cpp and name.h contains the declaration of a class denoted by name, i.e., having the same name as the files. The class name is a daughter class inheriting from the base class AnalysisBase that contains (empty) analysis methods. These methods are then specified at the level of the definition of the daughter classes.

The structure of the header file name.h follows, for any of the analyses included in the working directory,

The only pieces among these predefined lines to be modified by the user are the name tag name, related to the filename, and the label of the analysis label which is a string. This label is the one that can be provided as an option when running the SAMPLEANALYZER code (see Section 5.1).

In the C++ code above, the INIT_ANALYSIS macro automatically creates the constructor and destructor methods associated to the class name and consistently links the filename to the name of the analysis, which must be the same. As for any class derived from the mother class AnalysisBase, the class name must contain the three methods Initialize, Execute and Finalize. A part from this, the user is free to include his own set of additional functions and variables as required to perform his analysis.

The role of the three functions above is intuitive and related to their names. When SampleAnalyzer is executed in a shell, the program starts by calling (once) the function Initialize

```
void Initialize()
```

Its aim is to initialize all the internal variables, as well as the entire set of user-defined variables such as the histograms that have been requested (and that must be properly declared in the header file name.h).

After initialization, SAMPLEANALYZER loops over all the events which are passed to the code. As shown in Section 5.1, the list of event files is collected into a single file which is passed as an argument when executing the program from a shell. For each event, the function Execute is called in order to perform the analysis. Among others, the histograms are filled event by event and the cuts are applied. The method,

takes as arguments, on the one hand, the event event under consideration, which is passed as a set of particles with specific properties, and on the other hand, general information on the entire event sample sample currently analyzed, such as the corresponding integrated cross section necessary for a proper normalization of the histograms to be produced.

Once all the events have been processed, the function Finalize is eventually called (once) by SampleAnalyzer

Its aim is twofold. Firstly, it allows for the creation of the histograms and cutflow charts requested by the user. To implement the latter in an easy way, the
user can employ all the functionalities included in the ROOT library and we
therefore refer to the ROOT manual for more information [56]. Secondly, the
method Finalize stores the results of the analysis as an output ROOT files
which can be further processed or modified. The method takes as arguments
general information related to the event samples, summary, such has the total
cross section and the associated uncertainty. However, this time, instead of
having this information linked to a specific event sample, an average over
the whole list of samples included in the dataset under consideration has
been performed. In the case the user wants to compute additional quantities
when implementing the method Finalize, the same information, related to
each of the samples individually, is passed as the second argument which is
denoted, in the example above, as files.

These three predefined functions have to be implemented by the user in the corresponding source file name.cpp. To this aim, we refer to the description of the internal data format used by SAMPLEANALYZER for event processing (see Section 5.3 and Section 5.4) and to the one of the methods included upon the mother class AnalysisBase (see Section 5.5).

As mentioned above, several analyses can be included in the same working directory. The only requirement consists in implementing them in different files and classes since each analysis is unambiguously defined by its (file)name, which must therefore be unique. In order for SampleAnalyzer to correctly handle them, the user must refer to the various classes and files in the C++ source file analysisList.cpp, located in the sub-directory SampleAnalyzer/Analysis. This file contains a link to each of the header files associated to the analyses to be included. In addition, one instance of each analysis class is created. The architecture of this file follows the structure

```
#include "Analysis/name1.h"
#include "Analysis/name2.h"
...
#include "Core/AnalysisManager.h"
#include "Core/logger.h"

void AnalysisManager::BuildTable()
{
   Add(new name1);
   Add(new name2);
   ....
}
```

At least two analyses, denoted by name1 and name2 are implemented and the corresponding header files have been included in SampleAnalyzer. In the example above, the dots stand for possible additional analyses that the user might want to embed too. In the case the user has only implemented one single analysis, the file analysisList.cpp must still be present. It however then only includes the header file of this analysis and creates an instance of the related class.

In order to facilitate the implementation of new analyses, SAMPLEANA-LYZER comes with a PYTHON script newAnalysis.py located in the directory SampleAnalyzer. As shown above, creating a new analysis with the name newname requires the implementation of the two C++ files newname.h and newname.cpp and the update the file analysisList.cpp in order to include

the new analysis. The analysis-independent part of this task has been automated through this PYTHON script. The user can use it from a shell, by typing

newAnalysis.py newname

The script starts by asking to the user to type-in the label of the new analysis. As a result, the two files containing the declaration of the new analysis class are created (with a blank analysis included) and the list of the existing analyses in analysisList.cpp is updated. The user must now start implementing the analysis itself.

To this aim, he has to declare the variables necessary for the analysis and implement the three methods Initialize, Execute and Finalize, together with possible additional user-defined functions. This step requires the knowledge of the methods already included in the base class AnalysisBase and the one of the data format used internally by SAMPLEANALYZER. This is the scope of Section 5.3, Section 5.4 and Section 5.5.

5.3. The data format used by SampleAnalyzer

The function Execute introduced in the previous Section takes as a second argument an object of the type EventFormat, which points to the current event being analyzed denoted by event in the sequel. We dedicate this section to the description of the EventFormat class, which can be also found as a DOXYGEN documentation on the MADANALYSIS 5 website,

http://madanalysis.irmp.ucl.ac.be

In order to implement his analysis and to apply it to the event under consideration, the user generally needs to access various pieces of information related to this event. For example, the momentum of a specific type of object could be needed to implement some sophisticated cuts. To this aim, the Sampleanalyzer framework contains many built-in methods. They are split into two categories, the first one being related to hadron-level and parton-level events, which are generically named, in the following, as Monte-Carlo-level events, and the second one to reconstructed events.

When processing the event, SAMPLEANALYZER automatically creates, when reading the event, an object with a structure appropriate to store the information included in the event under consideration,

event.mc()

These two objects are C++ pointers to all the methods implemented to facilitate the design of an analysis at the Monte-Carlo-level and at the reconstructed level, respectively. These methods are extensively described in the rest of this Section.

5.3.1. The data format for parton-level or hadron-level events

The pointer event.mc() introduced above allows to access general information related to a (Monte Carlo) event denoted by event such as the weight of the event or the employed value of the strong coupling constant. In addition, the whole set of initial, intermediate and final state particles, together with their kinematical properties, is available. These properties form the so-called data format of SampleAnalyzer for Monte-Carlo-level events and are summarized in Table 12.

For Monte Carlo samples which include events related to several physical processes, matrix-element generators usually assign to each event a tag, *i.e.*, an unsigned integer number, that allows to identify the physical process which the event is originating from. If the user needs to access this tag in the analysis, it is available through the function

```
event.mc()->processId()
```

which returns an unsigned integer. With a similar syntax, the weight associated to the event event can be used in the C++ source file of the analysis through,

```
event.mc()->weight()
```

which returns a floating-point number. If the implementation of the analysis requires the evaluation of the factorization scale, it can be obtained from the method

```
event.mc()->scale()
```

which gives the results as a floating-point number. In contrast, the renormalization scale is not available but the values of the strong and electromagnetic coupling constants (including a possible running) can also be employed in the analysis,

```
event.mc()->alphaQCD()
event.mc()->alphaQED()
```

Table 12: Methods related to the parton-level and hadron-level event format

Let ${\tt ev}$ be an ${\tt EventFormat}$ object, i.e., an instance of the class related to events issued from a partonic or hadronic Monte Carlo sample.

ev.mc()->alphaQCD()	This returns, as a floating-point number, the value of the strong coupling constant used in the event.
ev.mc()->alphaQED()	This returns, as a floating-point number, the value of the electromagnetic coupling constant used in the event.
ev.mc()->particles()	This returns a vector of MCParticleFormat objects whose each entry consists in one of the particles present in the event, together with its properties. All initial, intermediate and final state particles are included.
ev.mc()->processId()	This returns an unsigned integer number related to the tag of the physical process the event is originating from. It is especially useful when several physical processes are merged into one event sample.
ev.mc()->scale()	This returns, as a floating-point number, the factorization scale.
ev.mc()->weight()	This returns, as a floating-point number, the event weight.

These last two methods also return floating-point numbers.

More importantly, implementing histograms or selection cuts requires to investigate the particle content of the event, as well as the properties of one or several of these particles. All the initial, intermediate and final state particles included in the event event are stored in a vector of MCParticleFormat objects, which can be called in the analysis through

```
event.mc()->particles()
```

The syntax above returns a vector that each entry consists in a particle present in the event, together with its properties, given as an instance of the MCParticleFormat class.

In order to implement a loop over all the particle content of the event event, in, e.g., the function Execute of the analysis source file, it is sufficient to program

```
unsigned int n = event.mc()->particles().size();
for (unsigned int i=0; i<n; i++) { ... }</pre>
```

where we remind that at this stage, all the initial, intermediate and final state particles are considered equivalently. We will show in Section 5.5 how to implement loops over, e.g., the final state particles only. Similarly, denoting by the object prt an instance of the MCParticleFormat class, the ith particle is given by

```
MCParticleFormat* prt = &event.mc()->particles()[i];
```

In the rest of this Section, we focus on the attributes of the object prt.

The class MCParticleFormat contains seven methods which allow to extract and use the properties of a given particle within the analysis. These methods are summarized in Table 13.

The PDG-id of the particle prt can be accessed through

```
prt->pdgid()
```

which returns the corresponding signed integer number. The **statuscode** of the particle, *i.e.*, its tag as an initial state, intermediate state or final state particle, can be obtained through the function

```
prt->statuscode()
```

Table 13: Methods related to the MCParticleFormat class

Let ${\tt prt}$ be a ${\tt MCParticleFormat}$ object.

<pre>prt.ctau()</pre>	This returns, as a floating-point number, the
_	decay length of the particle, assuming that it
	moves at the speed of light.
<pre>prt.momentum()</pre>	This returns a TLorentzVector containing
P= 0 · 0 ()	the four-momentum of the particle prt.
prt.mother1()	This returns a MCParticleFormat object
pro.modiferr()	pointing to the mother particle of the par-
	ticle prt, <i>i.e.</i> , either the particle which de-
	cays into prt (plus other particles) or a parti-
	cle which interacts with another particle (see
	mother2()) to produce the particle prt.
<pre>prt.mother2()</pre>	This returns a MCParticleFormat object
	pointing to the mother particle of the par-
	ticle prt, but only in the case prt is pro-
	duced from the interaction of two particles.
	These particles can be accessed through the
	two methods mother1() and mother2().
<pre>prt.pdgid()</pre>	This returns an integer number standing for
	the PDG-id of the particle.
<pre>prt.spin()</pre>	This returns a floating-point number, the co-
	sine of the angle between the momentum of
	the particle prt and its spin vector, com-
	puted in the laboratory reference frame.
<pre>prt.statuscode()</pre>	This returns an integer depending on the ini-
	tial state, intermediate state or final state na-
	ture of the particle. The conventions on this
	integer number are taken from Ref. [40].
	i i

which returns an integer associated to the initial state, intermediate state or final state nature of the particle. For the conventions on this integer number we refer to Ref. [40].

For non-initial state particles, information on mother particle(s) can be extracted (and used in the analysis) through the two methods

```
prt->mother1()
prt->mother2()
```

which return the MCParticleFormat objects related to the particle(s) from which the current particle prt is issued. In the case of a decay chain, only the mother1 method has to be used. In contrast, when the particle prt is produced from the interaction of two particles, both the mother1 and mother2 give results, each of them pointing to one of the initial particles.

The most important property of the MCParticleFormat class to be used in physics analyses consists in the particle four-momentum, accessible from

```
prt->momentum()
```

The result of this function is given as a TLorentzVector, a ROOT class appropriate to store four-momentum. In addition, this class contains various methods to compute a large set of observables which can be derived from the knowledge of the four-momentum, such as the energy or the transverse momentum. The complete list of these observables is given, together with the associated syntax, in Table 14.

An additional observable related to the four-momentum, and more in particular to the three-momentum component of the four-momentum, which can be employed in an analysis reads

```
prt->spin()
```

This returns a floating-point number which stands for the cosine of the angle between the momentum of the particle prt and its spin vector, evaluated in the laboratory reference frame.

Finally, the decay length of the particle (assuming that the particle is moving at the speed of light) can also be easily obtained through the function

```
prt->ctau()
```

which returns a floating-point number too.

Table 14: Common methods related to the MCParticleFormat and RecParticleFormat classes

Let P be a MCParticleFormat or RecParticleFormat object.

- P.angle(P2) Angle between the momenta of the objects P and P2, where P2 is an instance of the MCParticleFormat or RecParticleFormat class.
- P.beta() Velocity $\beta = v/c$.
- P.dr(P2) Relative distance between the objects P and P2 in the $\eta-\phi$ plane, where P2 is an instance of the MCParticleFormat or RecParticleFormat class.
- P.e() Energy.
- P.et() Transverse energy.
- P.eta() Pseudorapidity.
- P.gamma() Lorentz factor.
- P.m() Invariant mass.
- P.mt() Transverse mass.
- P.p() Norm of the momentum.
- P.phi() Azimuthal angle of the momentum.
- P.pt() Norm of the transverse momentum.
- P.px() Projection of the momentum on the x-axis.
- P.py() Projection of the momentum on the y-axis.
- P.pz() Projection of the momentum on the z-axis.
- P.r() Position of the object in the $\eta \phi$ plane.
- P.theta() Angle between the momentum and the beam axis.
- P.y() Rapidity.

Table 15: Methods related to the reconstructed event format

Let ev be an EventFormat object, *i.e.*, an instance of the class related to events issued from a reconstructed Monte Carlo sample.

ev.rec()->electrons()	This returns a vector with all the
	electrons of the event, encoded as
	RecLeptonFormat objects.
ev.rec()->jets()	This returns a vector with all the jets of
	the event, given as RecJetFormat objects.
ev.rec()->met()	This points to the missing energy of the
	event, encoded as a RecMETFormat object.
ev.rec()->muons()	This returns a vector with all the muons of
	the event, encoded as RecLeptonFormat
	objects.
ev.rec()->taus()	This returns a vector with all the taus of
	the event, given as RecTauFormat objects.

5.3.2. The data format for reconstructed events

At the beginning of this Section, we have introduced two types of data format which are used for event processing by SampleAnalyzer. They consist in the two sides of the more general EventFormat class. In the previous Section, we have focused on the description of an event object event read from a parton-level or hadron-level sample. We have shown that its properties are embedded, in the SampleAnalyzer framework, into a structure which can be browsed from the C++ pointer event.mc().

Since the properties of reconstructed events are in general very different from those of Monte Carlo events, SampleAnalyzer embeds the latter into another structure which is, this time, linked to the C++ pointer event.rec(). It consists in five methods, collected in Table 15, associated to the five types of objects that can be reconstructed, *i.e.*, electrons, muons, taus, jets and missing energy⁶. In the SampleAnalyzer framework, each

⁶By electrons, muons and taus, we are considering both the corresponding particles

reconstructed objects is associated to a specific class derived from the generic RecParticleFormat mother class and depending on its species. In the following, we adopt the choice of describing the daughter classes, more relevant for the user, rather than the generic class.

Two methods are associated to first and second generation charged leptons, *i.e.*, to the reconstructed electrons and muons, present in the event event,

```
event.rec()->electrons()
event.rec()->muons()
```

They return vectors that the entries are the different reconstructed electrons and muons of the event, respectively. Each reconstructed electron or muon is encoded as a RecLeptonFormat object. This last class has six associated methods, gathered in Table 16, related to the attributes of the reconstructed leptonic electron and muon objects.

As for Monte Carlo event samples, the four-momentum of the reconstructed objects is one of the most incontrovertible variables to be employed in analyses. For a reconstructed lepton lep, its four-momentum can be included and used in the analysis source file by calling the function

```
lep.momentum()
```

the syntax being similar to the one employed for particles included in parton-level and hadron-level events. This method returns a TLorentzVector object containing the corresponding four-momentum. Therefore, the methods associated to all the observables which can be derived from the knowledge of the four-momentum, collected in Table 14, are also methods of the RecLeptonFormat class.

Reconstructed leptons are considered as electrons or muons regardless of their charge. In the case the user wants to select them according to the electric charge, the method

```
lep.charge()
```

of the RecLeptonFormat class can be used. It returns, as a floating-point number, the electric charge of the reconstructed lepton which can be equal to either -1 or +1 according to its antiparticle or particle nature.

and antiparticles.

Table 16: Methods related to the ${\tt RecLeptonFormat}$ class

Let lep be an RecLeptonFormat object, i.e., an instance of the class describing the reconstructed electrons and muons.

lep.charge()	This returns the electric charge of the reconstructed object lep as a floating-point number. The returned value consists in -1 or $+1$.
lep.EEoverHE()	This returns the ratio of the electromagnetic and hadronic energy for the reconstructed object lep, given as a floating-point number.
<pre>lep.ET_PT_isol()</pre>	This returns the ratio of the values of the functions sumET_isol() and sumPT_isol().
lep.HEoverEE()	This returns the ratio of the hadronic and electromagnetic energy for the reconstructed object lep, given as a floating-point number.
<pre>lep.momentum()</pre>	This returns a TLorentzVector containing the four-momentum of the reconstructed object lep.
<pre>lep.sumET_isol()</pre>	This returns, if the object lep is a muon, the sum of the transverse energy of all the tracks lying in a cone around the muon. The size of the cone is fixed by the detector simulation tool. If lep is an electron, this function returns zero.
<pre>lep.sumPT_isol()</pre>	This returns, if the object lep is a muon, the sum of the transverse momentum of all the tracks lying in a cone around the muon. The size of the cone is fixed by the detector simulation tool. If lep is an electron, this function returns zero.

All the methods presented in Table 14 can also be used.

Two methods allow to get information on the splitting of the reconstructed electron or muon energy between the electromagnetic and the hadronic parts of the detector,

```
lep.EEoverHE()
lep.HEoverEE()
```

These methods compute the ratio between the energy deposited in the electromagnetic calorimeter and the one deposited in the hadronic calorimeter, and *vice-versa*, and return them as floating-point numbers.

Finally, the RecLeptonFormat class contains three specific methods related to muon isolation (see also Section 5.5). The algorithms to be employed for deciding if a muon is considered as isolated or not require in general the evaluation of two quantities, the sum of the transverse momentum of all tracks lying in a cone around the muon candidate and the sum of their transverse energy. These two observables can be accessed by typing

```
lep.sumPT_isol()
lep.sumET_isol()
```

which return a zero value in the case the lepton lep is an electron. In contrast, for muons, the value read from the event file is employed. The size of the cone is fixed by the fast detector simulation tool and not available in the LHCO event format. The ratio of the values of these two functions can be obtained via the function

```
lep.ET_PT_isol()
```

Tau leptons being unstable, they always decay either into a narrow jet, into a muon or into an electron, each time in association with missing energy. Therefore, a specific class, different from the RecLeptonFormat class, exists in order to embed reconstructed taus. This class is denoted by RecTauFormat. In the internal data format used by SAMPLEANALYZER, the pointer to the reconstructed event, event.rec(), contains a specific method to access all the taus present in the event

```
event.rec()->taus()
```

This returns, as a vector of RecTauFormat objects, all the reconstructed tau leptons of the event (regardless of their electric charge).

Table 17: Methods related to the RecTauFormat class

Let tau be an instance of the RecTauFormat class, *i.e.*, the class describing the reconstructed taus.

tau.charge()	This returns the electric charge of the object tau as a floating-point number. The returned value consists in -1 or $+1$.
tau.EEoverHE()	This returns the ratio of the electromagnetic and hadronic energy for the object tau, given as a floating-point number.
tau.HEoverEE()	This returns the ratio of the hadronic and electromagnetic energy for the object tau, given as a floating-point number.
tau.momentum()	This returns a TLorentzVector containing the four-momentum of the object tau.
tau.ntracks()	This returns, as a short integer, the number of charged tracks contained in the object tau.

All the methods presented in Table 14 can also be used.

In addition to the four methods momentum(), charge(), EEoverHE(), HEoverEE() and the methods given in Table 14 already present for charged leptons of the first and second generations, the RecTauFormat class includes an additional method extracting the number of charged tracks issued from the decaying tau and included in the reconstructed object. Denoting by tau the reconstructed object, this number, given as a short integer, can be obtained and further used in the implementation of the analysis by

tau.ntracks()

The entire list of methods associated to the RecTauFormat class can be found in Table 17.

As for electrons, muons and taus, the entire set of reconstructed jets

Table 18: Methods related to the RecJetFormat class

Let j be an RecJetFormat object, *i.e.*, an instance of the class describing the reconstructed jets.

j.btag()	This returns true or false according to the <i>b</i> -tagging (or not) of the object j.
j.EEoverHE()	This returns the ratio of the electromagnetic and hadronic energy for the object
	j, given as a floating-point number.
j.HEoverEE()	This returns the ratio of the hadronic and
	electromagnetic energy for the object j,
	given as a floating-point number.
j.momentum()	This returns a TLorentzVector contain-
	ing the four-momentum of the object j.
j.ntracks()	This returns, as a short integer, the num-
	ber of charged tracks contained in the ob-
	ject j.

All the methods presented in Table 14 can also be used.

included in the event event can be obtained through through the intuitive method of the event.rec() structure,

This returns a vector of instances of the RecJetFormat class whose the properties are collected in Table 18.

Reconstructed jets are characterized by their momentum, the number of charged tracks induced by the parton showering, fragmentation and hadronization of the initial partons and the ratio of the hadronic and electromagnetic parts of the jet energy. These properties are related to the methods momentum(), ntracks(), EEoverHE(), HEoverEE() as well as all those included in Table 14, of the RecJetFormat class. In addition, an extra method returns true or false according to the fact that the jet has been tagged as a b-jet or not,

Table 19: Methods related to the RecMetFormat class

Let miss be a RecMetFormat object, *i.e.*, the variable containing the reconstructed missing energy associated to an event.

miss.mag()	This returns the magnitude of the missing energy represented by the object miss as
	a floating-point number.
miss.phi()	This returns the azimuthal angle of the
	missing energy represented by the object
	miss as a floating-point number.
miss.x()	This returns the x -component of the miss-
	ing energy represented by the object miss
	as a floating-point number.
miss.y()	This returns the y -component of the miss-
	ing energy represented by the object miss
	as a floating-point number.

j.btag()

where j denotes an instance of the RecJetFormat class.

The last type of objects which is included in reconstructed events consists in the associated missing transverse energy. The structure event.rec() comes with the related method

which returns a two-dimensional vector implemented as a TVector2 object. It contains then the direction and magnitude of the missing energy in the transverse plane as shown in Table 19.

The x-component and y-component of the missing transverse energy can be used when implementing an analysis through the two methods

miss.x()
miss.y()

which return floating-point numbers associated to the two components of the two-dimensional vector describing the missing energy, the object miss being an instance of the RecMETFormat class. In the case the user prefers to use polar coordinates instead of Cartesian coordinates when implementing his analysis, he can use the two methods

```
miss.mag()
miss.phi()
```

which return the magnitude and the azimuthal angle of the two-dimensional transverse-momentum vector as floating-point numbers.

5.4. The sample format used by SampleAnalyzer

In Section 5.2, we have introduced the function Execute as the main function of the analysis class. We have shown that it requires two arguments, the event being currently analyzed, stored as an EventFormat object, and general information about the sample which this event belongs to, passed as a SampleFormat object. This format is also the one to be used for the arguments of the function Finalize, dedicated to the creation of the output files containing the results of the analysis. In this case, the user must pass two arguments, an instance of the SampleFormat class associated to all the event samples included in the dataset under consideration, *i.e.*, information averaged over all the samples, and one vector of SampleFormat objects with one entry for each of the samples included in the dataset, *i.e.*, the same information, but given for each individual sample.

The SampleFormat class contains several methods allowing to access general information about the sample under consideration, provided the information is available. If not, the associated entries in the SampleFormat structure are left non-initialized and the corresponding quantity cannot consequently be used in an analysis. A full DOXYGEN documentation is available on the MADANALYSIS 5 website,

```
http://madanalysis.irmp.ucl.ac.be
```

Since information supplementing the events is only available within LHE and Stdhep files, the SampleFormat structure is then only relevant for the analysis of such event files. In particular, LHCO files with reconstructed events do not include anything but the events. In this case, it is however still possible to pass additional information such as cross sections to SampleAnalyzer through the attributes of the dataset class (see Section 4.3).

In Section 5.3, we have shown that SAMPLEANALYZER is creating an EventFormat structure each time an event is processed, resulting into a C++ pointer pointing to the whole methods available for the EventFormat structure. Similarly, when processing a new sample generically denoted by sample, SAMPLEANALYZER creates a C++ pointer

```
sample.mc()
```

which points to a structure containing all the methods allowing to access the global information of the event file. These methods are collected in Table 20 and Table 21.

As shown by its name, sample.mc() is related to parton-level or hadron-level events. The counterpart of this object in the case of a sample containing reconstructed events, sample.rec(), has been implemented within the SAM-PLEANALYZER framework. However, the only event file format appropriate for reconstructed events, the LHCO format, does not leave a possibility for including additional information to the events. Therefore, the pointer sample.rec() points to a set of null information. If in the future, a new format for reconstructed events is designed, with the room for global information about the event sample, the related structure in the SAMPLEANALYZER framework will be ready for the new format.

The first series of methods available within the SampleFormat structure, collected in Table 20, are related to the description of the initial colliding beams. The two members of the SampleFormat class

```
sample.mc()-> beamPDGID().first
sample.mc()-> beamPDGID().second
```

return, as integer numbers, the PDG-id of the first and second beams, respectively, whilst the four class members

```
sample.mc()-> beamPDFauthor().first
sample.mc()-> beamPDFID().first
sample.mc()-> beamPDFauthor().second
sample.mc()-> beamPDFID().second
```

return, as four unsigned integer numbers, information with respect to the set of parton density functions used for the beams. These identifying numbers are exported from the event samples (if available) and correspond to the numbering scheme employed by matrix element generators to identify

Table 20: Methods related to the SampleFormat class giving information on the colliding beams

Let sample be a SampleFormat object.

sample.mc()->beamE().first

This returns, as a floating-point number, the energy of the first of the colliding beams.

sample.mc()->beamPDFauthor().first

This returns the author group of the parton density set used for the first of the colliding beams, as an unsigned integer number.

sample.mc()->beamPDFID().first

This returns the identifier, as an unsigned integer number, of the parton density set used for the first of the colliding beams, within a given author group of parton densities.

sample.mc()->beamPDGID().first

This returns the PDG-id of the first of the colliding beams, as an integer number.

sample.mc()->beamE().second

This returns, as a floating-point number, the energy of the second of the colliding beams.

sample.mc()->beamPDFauthor().second

This returns the author group of the parton density set used for the second of the colliding beams, as an unsigned integer number.

sample.mc()->beamPDFID().second

This returns the identifier, as an unsigned integer number, of the parton density set used for the second of the colliding beams, within a given author group of parton densities.

sample.mc()->beamPDGID().second

This returns the PDG-id of the second of the colliding beams, as an integer number.

Table 21: Other methods related to the SampleFormat class

Let sample be a SampleFormat object.

sample.mc()->weightingmode()

This returns, as an integer number, information on the type of events present in the sample (e.g., weighted versus unweighted events).

sample.mc()->processes()

This returns a vector where the entries are ProcessFormat objects containing basic information about each of the physical processes included in the sample sample.

sample.mc()->xsection()

This returns the cross section associated to the event sample sample, as a floating-point number.

sample.mc()->xsection_error()

This returns the Monte Carlo uncertainty related to the cross section associated to the event sample sample, as a floating-point number.

the set of parton densities employed. According to the Les Houches conventions, this scheme is based on the PDFLIB [60] and LHAPDF [61] packages. Hence, the method beamPDFauthor() is related to the author group which has released the parton density relevant for the sample under consideration and the method beamPDFID() indicates which specific set of parton densities of the corresponding group has been used.

The last pieces of information available with respect to the beams which can be stored in the event files, and thus exported to the SAMPLEANALYZER framework, consist in their energy. This can be used in the implementation of an analysis by means of the two methods

```
sample.mc()-> beamE().first
sample.mc()-> beamE().second
```

for the first and second beams, respectively.

The second series of methods available from the SampleFormat class are related to the sample itself. When available, information about the weighting of the events is passed to SAMPLEANALYZER and can be accessed through

```
sample.mc()-> weightingmode()
```

According to the LHE format [40, 41], this type of information is stored as an integer number which indicates how event weights and cross sections have to be interpreted, and the weightingmode() method returns this integer number. Even if presently, SampleAnalyzer does not fully support weighted events, the architecture has already been implemented with respect to future developments of the program. Finally, the most important quantities included in this second series of methods concern the cross section associated to the sample sample, together with the corresponding uncertainty. Both can be called at the level of the analysis by using the methods

```
sample.mc()->xsection()
sample.mc()->xsection_error()
```

which return floating-point numbers.

As mentioned in Section 5.3.1, several physical processes can be included within the same event sample. In this case, each of the processes is associated with an identifying tag (see the processId method introduced in Section 5.3.1). The SampleFormat structure allows to easily access detailed information about the processes individually. This information is stored into a new structure, the ProcessFormat class, whose all the associated methods are summarized in Table 22. The method

```
sample.mc()->processes()
```

returns a vector that each entry is a ProcessFormat object. For an instance of this class denoted by process, information on the identifier of the process can be obtained through

```
process.processId()
```

the corresponding cross section together with its associated uncertainty through the methods

```
process.xsection()
process.xsection_error()
```

Table 22: Methods related to the ProcessFormat class

Let proc be an instance of the ProcessFormat class.

proc.processId()

This returns an unsigned integer number related to the tag of the physical process the event is originating from.

proc.xsection()

This returns the cross section associated to the process proc, as a floating-point number.

proc.xsection_error()

This returns the Monte Carlo uncertainty related to the cross section associated to the process proc, as a floating-point number.

proc.maxweight()

This returns the maximum weight carried by an event associated to the process proc, as a floating-point number.

and finally, the maximum weight carried by a single event originating from the subprocess process through

process.maxweight()

5.5. Framework services

In this Section, we describe the various services included in the SampleAnalyzer framework. These services are components of the program that are initialized (internally or by the user) at the beginning of the execution of SampleAnalyzer (within the Initialize method) and can then be further called within the analysis as many times as necessary. Two series of services are currently available, message services and physics services and are described in the rest of this Section. We remind that a full Doxygen documentation is available on the Madanalysis 5 website,

http://madanalysis.irmp.ucl.ac.be

5.5.1. Message services

This class of functions has been implemented and can be used by the users in order to print text to the screen during an on-going analysis in a rather sophisticated fashion.

As for the implementation of any C++ program, the user can, when designing his analysis, use the standard C++ streamers std::cout and std:cerr in order to implement messages to be printed to the screen. However, these methods only include two levels of messages, *i.e.*, normal and error messages. It is hence rather cumbersome to handle multi-level messages and give information both on the reason leading to the printing of the message and on its location in the analysis code in a clear and useful way.

Therefore, SAMPLEANALYZER includes its own message services with four different levels of printing, DEBUG, INFO, WARNING and ERROR. The way to use these four modes mimics the one of the C++ commands std::cout and std::cerr,

```
DEBUG << "Debug message." << std::endl;
INFO << "Information message." << std::endl;
WARNING << "Warning message." << std::endl;
ERROR << "Error message." << std::endl;</pre>
```

Each message level is associated to a different color. Debugging messages are printed in yellow, information messages in white, warning messages in purple and error messages in red. However, if the user is not interested by the color of the messages, this message can be switched off by including, in the source code of the analysis, the line

```
LEVEL->DisableColor()
```

The color can be enabled again through the command

```
LEVEL->EnableColor()
```

In the command lines above, the keyword LEVEL stands for any of the four levels of message, DEBUG, INFO, WARNING or ERROR.

Warning and error messages have a special role concerning the debugging of the analysis code. In addition to the message, the line number of the code having generated the message is also printed in order to facilitate the debugging. Finally, let us note that messages associated to a given level can be fully switched off, if this is needed by the user, by including in the analysis code the command

```
LEVEL->Mute()
```

Messages can be restored by implementing

```
LEVEL->UnMute()
```

where the keyword LEVEL again stands for any of the four levels of message.

5.5.2. Physics services

By the name *physics services*, we collect a series of methods and functions aiming to facilitate the writing of an analysis by the user. It includes, among others, functions to compute global observables related to the entire event, such as the transverse missing energy, and this for any type of event (parton level, hadron level and reconstructed level). All these functions can be called through the C++ pointer PHYSICS.

Before moving on with the description of these functions, one must note that in the case the user wants to use, within his analysis, one or several of the methods related to the invisible or hadronizing particles, the definition of these invisible and hadronic particles must be first provided. This task is performed within the Initialize function introduced in Section 5.2 with the help of two different functions, one being associated to the invisible particles and another one being associated to the particles taking part to the hadronic activity. However, before declaring new particle as invisible or hadronizing, the physics services class must be initialized by means of one of the commands

```
PHYSICS->mcConfig()->Reset()
PHYSICS->recConfig()->Reset()
```

In the case the event samples which are analyzed consist in parton-level and hadron-level event files, the method mcConfig is used whilst for reconstructed-level event files, recConfig is instead used. Then, a particle whose the PDG-id is PDGID can be declared as an invisible particle by means of

PHYSICS->mcConfig()->AddInvisibleId(PDGID)

In the non-expert mode of Madanalysis 5, the definition of the multiparticle invisible (see Section 4.4) corresponds to several calls of the function above behind the scene. Taking the example of an analysis at the parton-level in the context of the Minimal Supersymmetric Standard Model, the invisible particles consist in the neutrinos and the lightest superpartner assumed to be the lightest neutralino. The corresponding declaration within the Sampleanalyzer framework reads

```
PHYSICS->mcConfig().AddInvisibleId(-16);
PHYSICS->mcConfig().AddInvisibleId(-14);
PHYSICS->mcConfig().AddInvisibleId(-12);
PHYSICS->mcConfig().AddInvisibleId(12);
PHYSICS->mcConfig().AddInvisibleId(14);
PHYSICS->mcConfig().AddInvisibleId(16);
PHYSICS->mcConfig().AddInvisibleId(1000022);
Similarly, the method
PHYSICS->mcConfig()->AddHadronicId(PDGID)
```

declares the particle whose the PDG-id is given by PDGID as a particle related to the hadronic activity in an event. Again, in the non-expert mode of MADANALYSIS 5, the definition of the multiparticle hadronic (see Section 4.4) corresponds to a series of calls to this last function. For the example of a parton-level analysis within the Standard Model, the related declaration would be given by

```
PHYSICS->mcConfig().AddHadronicId(-5);
PHYSICS->mcConfig().AddHadronicId(-4);
PHYSICS->mcConfig().AddHadronicId(-3);
PHYSICS->mcConfig().AddHadronicId(-2);
PHYSICS->mcConfig().AddHadronicId(-1);
PHYSICS->mcConfig().AddHadronicId(1);
PHYSICS->mcConfig().AddHadronicId(2);
PHYSICS->mcConfig().AddHadronicId(3);
PHYSICS->mcConfig().AddHadronicId(4);
PHYSICS->mcConfig().AddHadronicId(4);
PHYSICS->mcConfig().AddHadronicId(5);
PHYSICS->mcConfig().AddHadronicId(5);
```

Table 23: Methods related to the initialization of the physics services

PHYSICS->mcConfig()->AddHadronic(PDGID)

This adds the particle whose the PDG-id is PDGID to the list of the invisible particles.

PHYSICS->mcConfig()->AddInvisible(PDGID)

This adds the particle whose the PDG-id is PDGID to the list of the particles taking part of the hadronic activity of an event.

PHYSICS->mcConfig()->Reset()

This initializes physics services when analyzing parton-level or hadron-level events.

PHYSICS->recConfig()->Reset()

This initializes physics services when analyzing reconstructed-level events.

In the case of events at the reconstructed-level, the user may need to specify the algorithm to be employed when testing the isolation of a muon. This can be done through the two (self-excluding) methods

```
PHYSICS->recConfig().UseDeltaRIsolation(deltaR=0.5);
PHYSICS->recConfig().UseSumPTIsolation(sumPT, ET_PT);
```

In the first case, a muon is tagged as isolated when no track lies inside a cone of size deltaR around the muon. The default size of this cone is set to 0.5. In the second case, the muon is tagged as isolated when the sum of the transverse momentum of all the tracks lying in a cone around the muon is lower than sumPT and in addition, when the sum of the transverse energy of these tracks over the sum of their transverse momentum is lower than ET_PT. By default, the first algorithm is adopted with deltaR being set to 0.5.

The usage of all the methods associated to the initialization of the physics services, *i.e.*, to the PHYSICS->mcConfig() and PHYSICS->recConfig() objects, is summarized in Table 23. All the other methods included in the physics services are in general called within the function Execute, at the core of the analysis and are summarized in Table 24 and Table 25.

Three boolean functions exist in order to test if a particle prt is invisible or visible as well as if this particle takes part to the hadronic activity of the event or not.

```
PHYSICS->IsHadronic(prt)
PHYSICS->IsInvisible(prt)
PHYSICS->IsVisible(prt)
```

When analyzing partonic or hadronic event samples, the particle prt is passed as a MCParticleFormat object. In contrast, for analyses at the reconstructed level, prt is an instance of the RecParticleFormat class⁷. In order for these three methods to correctly work, it is necessary to declare which particle is invisible and which particle takes part to the hadronic activity as shown above. In the reconstructed-level case, an additional method exists to test whether a muon is isolated within an event event,

```
PHYSICS->IsIsolatedMuon(prt,event)
```

which returns always false for particles which are not muon candidates. In the case of muons, this method applies the isolation algorithm chosen by the user when setting the PHYSICS->recConfig() properties.

Another set of three methods checks whether a given particle prt is a final state, initial state or intermediate state particle,

```
PHYSICS->IsFinalState(prt)
PHYSICS->IsInitialState(prt)
PHYSICS->IsInterState(prt)
```

These functions all return a boolean value according to the final state, initial state or intermediate state nature of the particle under consideration. As above, those methods work equivalently for analyses at the hadronic, partonic and reconstructed levels, the particle prt being hence either an instance of the MCParticleFormat or of the RecParticleFormat classes. Since status codes are defined in a different fashion according to the event file format, we have adopted the choice to include these features within the physics services rather than the data format itself, which allows to have a unified way to probe the initial, intermediate or final state nature of the particles included

⁷The RecParticleFormat class is the mother class of all the classes defining reconstructed objects, *i.e.*, the RecLeptonFormat, RecJetFormat and RecTauFormat classes.

Table 24: Boolean methods included in the physics services

Let prt be an instance of either the MCParticleFormat or the RecParticleFormat classes.

```
PHYSICS->IsFinalState(prt)
```

This checks if the particle prt is a final state particle.

PHYSICS->IsHadronic(prt)

This checks if the particle prt takes part to the hadronic activity in an event.

PHYSICS->IsInitinalState(prt)

This checks if the particle prt is an initial state particle.

PHYSICS->IsInterState(prt)

This checks if the particle prt is an intermediate state particle.

PHYSICS->IsInvisible(prt)

This checks if the particle prt is an invisible particle.

IsIsolatedMuon(prt,evt)

This checks if the particle prt is a muon isolated from the other particles of the event evt.

PHYSICS->IsVisible(prt)

This checks if the particle prt is a visible particle.

in an event. These last series of functions allows to implement efficiently a loop over, e.g., the final state particles of the event, and not on all the particles, in contrast to the example given in Section 5.3.1,

```
unsigned int n = event.mc()->particles().size();
for (unsigned int i=0; i<n; i++)
{
   MCParticleFormat* prt = &event.mc()->particles()[i];
   if(PHYSICS->IsFinalState(prt))
      { ... }
}
```

Five methods included in the physics services are dedicated to the computation of global observables related to the entire event content: the missing

transverse energy $\not \!\! E_T$, the missing hadronic energy $\not \!\! H_T$, the total transverse energy E_T and the total hadronic energy H_T as defined in Eq. (1) and Eq. (2), as well as the partonic center-of-mass energy. The associated functions read, respectively,

```
PHYSICS->EventMET(evt->mc())
PHYSICS->EventMET(evt->rec())
PHYSICS->EventMHT(evt->mc())
PHYSICS->EventTET(evt->mc())
PHYSICS->EventTET(evt->rec())
PHYSICS->EventTHT(evt->mc())
PHYSICS->EventTHT(evt->mc())
PHYSICS->SqrtS(event->mc())
PHYSICS->SqrtS(event->rec())
```

where evt is an instance of the EventFormat class. These five methods work for any level of analysis (partonic, hadronic and reconstructed) and then takes accordingly as argument either an event->mc() or an event->rec() object. The value, in GeV, of the corresponding observable is returned as a floating-point number.

For most analyses, it is important to be able to order the particles according to a specific variable. Therefore, in order to avoid requiring the user to implement his own sorting algorithm, physics services include one dedicated function,

```
PHYSICS->sort(prtVect, OrderingObs)
```

The command above allows to sort a vector of particles prtVect, that each entry is either a MCParticleFormat or RecParticleFormat object, according to the ordering observable OrderingObs. The implemented choices for the latter are ETAordering (pseudorapidity ordering), ETordering (transverse-energy ordering), Eordering (energy ordering), Pordering (momentum ordering), PTordering (transverse-momentum ordering), PXordering (ordering according to the x-component of the momentum), PYordering (ordering according to the y-component of the momentum) and PZordering (ordering according to the z-component of the momentum).

The last method included in the physics services that can be useful when implementing an analysis is related to the reference frame in which the four-momenta of the particles included in the event are computed. When reading the event files, all the four-momenta are, by conventions, given in the laboratory reference frame. However, for specific observables, it is necessary to recalculate one or several of the four-momenta in the rest frame of a given particle, which is equivalent to apply a Lorentz boost to these four-vectors. This task can be done automatically by means of the method

Table 25: Other methods included in the physics services

Let evt be an instance of the EventFormat class assuming to point to an evt->mc() or a evt->rec() structure.

PHYSICS->ToRestFrame(prt,prt1)

This recalculates (and modifies) the four-momentum of the particle prt after a Lorentz boost to the rest frame of the particle prt1. The objects prt and prt1 are two instances of the MCParticleFormat or RecParticleFormat classes.

PHYSICS->EventMET(evt)

This computes the missing transverse energy E_T associated to the event evt.

PHYSICS->EventMHT(evt)

This computes the missing transverse hadronic energy H_T associated to the event evt.

PHYSICS->EventTET(evt)

This computes the total transverse energy E_T associated to the event ext.

PHYSICS->EventTHT(evt)

This computes the total transverse hadronic energy H_T associated to the event evt.

PHYSICS->sort(prtvector,orderobs)

This sorts a vector of MCParticleFormat or MCRecFormat objects according to the ordering observable orderobs. The allowed choices for the ordering variable are ETAordering (pseudorapidity ordering), ETordering (transverse-energy ordering), Eordering (energy ordering), Pordering (momentum ordering), PTordering (transverse-momentum ordering), PXordering (ordering according to the x-component of the momentum), PYordering (ordering according to the y-component of the momentum) and PZordering (ordering according to the z-component of the momentum).

PHYSICS->SqrtS(evt)

This returns the partonic center-of-mass energy of the event in GeV.

PHYSICS->ToRestFrame(prt,prt1)

where the objects prt and prt1 are two instances of the MCParticleFormat or RecParticleFormat classes. The four-momentum of the particle prt is boosted to the rest frame of the particle prt1 and overwritten. Let us note that the Lorentz transformation is chosen in such a way that the tridimensional (x, y, z) axes are kept unchanged.

5.6. A detailed example

In this Section, we give a detailed example about how to implement an analysis within the expert mode of Madanalysis 5. We choose to investigate the polarization of the W-boson issued from a top quark decaying leptonically,

$$t \to W^+ b \to \ell^+ \ \nu_l \ b \ . \tag{4}$$

This property of the W-boson is usually studied through the shape a particular angular distribution, $d\sigma/d\cos\theta^*$. The θ^* angle is defined as the angle between the momentum of the W-boson evaluated in the rest frame of the originating top quark and the one of the lepton ℓ^+ evaluated in the rest frame of the W-boson.

To investigate this observable, we focus on the production of a top-antitop pair where one of the final state top quarks undergoes a leptonic decay and the other one decays hadronically,

$$pp \to t\bar{t} \to (b\ j\ j)(\bar{b}\ \ell^-\ \bar{\nu}_l)$$
 or $pp \to t\bar{t} \to (b\ \ell^+\ \nu_l)(\bar{b}\ j\ j)$,

where ℓ stands for an electron or a muon and j for a light jet. From the parton-level samples stored into the directory samples of MADANALYSIS 5⁸, ttbar_sl_1.lhe.gz and ttbar_sl_2.lhe.gz, we illustrate the implementation, within the SAMPLEANALYZER framework, of an analysis code leading to the creation of an histogram representing the $d\sigma/d\cos\theta^*$ distribution introduced above. We refer to Section 3 concerning details on the generation of these two Monte Carlo samples.

When launching MADANALYSIS 5 in expert mode by issuing in a shell

bin/ma5 -e

⁸If the samples directory is absent, we remind that it can be created by issuing, in the command line interface of MADANALYSIS 5, the command install samples.

the user is asked to enter the name of the directory to be created and the label of the analysis. For the sake of the example, the name of the directory is chosen to be Wpol whilst the string W polarization from a top decay is entered as the label, or title, of the analysis. The three files analysisList.cpp, user.cpp and user.h are automatically created and stored in the directory Wpol/SampleAnalyzer/Analysis, as mentioned in Section 5.1. Whilst the first of these files does not have to be modified by the user, the two others must be updated to include the analysis which has to be implemented.

The file user.h strictly follows the structure presented in Section 5.2,

The label of the analysis has been automatically included into the arguments of the INIT_ANALYSIS function at the execution of MADANALYSIS 5, together with the declaration of the three main functions Initialize, Execute and Finalize. Since we are interested in the creation of a single histogram, we therefore need to declare, as a private member of the user class, an instance

of the TH1F class⁹ which we denote by myHisto.

The C++ source file user.cpp contains the implementation of the three core functions Initialize, Execute and Finalize. Its structure reads thus

where the dots are specified in the sequel of this Section. The function Initialize contains on the one hand the initialization of the physics services as well as, on the other hand, the one of the histogram to be drawn, following the ROOT syntax,

```
void user::Initialize()
{
    // Initializing Physics services
    PHYSICS->mcConfig().Reset();

    // Initializing the histogram
    myHisto = new TH1F("myHisto","cos#theta^{*}",15,-1.,1.);
}
```

We here ask for the creation of an histogram containing 15 bins ranging from -1 to +1, the minimal and maximal values for the cosine of the θ^* -angle.

The implementation of the function Execute is a bit more complicated. First of all, the observable of interest can only be computed once the three relevant particles have been identified, *i.e.*, the final state lepton ℓ , the intermediate W-boson decaying into this lepton ℓ and the originating top quark.

⁹We refer to the ROOT manual [56] for the definition of the TH1F class and its attributes.

Then, the four-momentum of the W-boson has to be boosted into the rest frame of the top quark and the one of the lepton into the rest frame of the W-boson. Once this is done, the cosine of the angle θ^* can be computed and the value filled into the histogram. The skeleton of the function Execute reads thus

```
void user::Execute(const SampleFormat& sample,
    const EventFormat& event)
{
    // Initialization of three pointers to the lepton,
    // W and top quark
    { ... }

    // Identification of the three particles of interest
    { ... }

    // Computing the observable; filling the histogram
    { ... }
}
```

The first series of dots concerns the declaration of the three objects which contain the three particles of interest, once identified, and which are necessary to compute the observable $\cos \theta^*$,

```
const MCParticleFormat* top = 0;
const MCParticleFormat* w = 0;
const MCParticleFormat* lepton = 0;
```

In the lines above, three pointers to a MCParticleFormat structure are created and initialized to the zero value.

The second series of dots are related to the identification of the top quark, the W-boson and the lepton. In particular, the analysis code needs to check that the mother-to-daughter relations given by the decay chain of Eq. (4) are fulfilled. Otherwise, the event must be rejected. In practice, this task is performed through a loop over the particle content of the event. The lepton is firstly selected as a final state particle whose the PDG-id reads ± 11 (electron) or ± 13 (muon). The W-boson is then further identified by requiring that the mother particle of this lepton has a PDG-id equal to ± 24 . The top quark is finally identified by requiring that the 'grand-mother' of the lepton has a PDG-id equal to ± 6 . The implemented code is given by

```
const MCParticleFormat* prt=&(event.mc()->particles()[i]);
    // The lepton is a final state particle
    if (!PHYSICS->IsFinalState(prt)) continue;
    // Lepton selection based on the PDG-id
    if ( std::abs(prt->pdgid())!=11 &&
         std::abs(prt->pdgid())!=13 ) continue;
    // Getting the mother of the lepton and
    // checking if it is a W-boson
    const MCParticleFormat* mother = prt->mother1();
    if (mother==0) continue;
    if (std::abs(mother->pdgid())!=24) continue;
    // Getting the grand-mother of the lepton and
    // checking if it is a top quark
    const MCParticleFormat* grandmother = mother->mother1();
    if (grandmother==0) continue;
    if (std::abs(grandmother->pdgid())!=6) continue;
    // Saving the selected particles
    lepton = prt;
           = mother;
           = grandmother;
    top
    // Particles are found: breaking the loop
    break;
In the case the three particles are not identified, the event must be rejected,
which is done by implementing
 // Rejection of the event if the decay chain is not found
  if (lepton==0)
  ₹
    WARNING << "(t > b W > b l nu) decay chain not found!"
```

for (unsigned int i=0;i<event.mc()->particles().size();i++)

```
 << std::endl;
 return;
}</pre>
```

where we have used the message services included in the SampleAnalyzer framework to print a warning message to the screen.

In the third and last series of dots included in the skeleton of the user.cpp file, the observable $\cos \theta^*$ is computed and the histogram filled, this last task being performed by means of standard ROOT commands,

```
// Boosting the lepton four-momentum to the W rest frame
MCParticleFormat lepton_new = *lepton;
PHYSICS->ToRestFrame(lepton_new,w);

// Boosting the W four-momentum to the top rest frame
MCParticleFormat w_new = *w;
PHYSICS->ToRestFrame(w_new,top);

// Computing the observable; filling the histogram
myHisto -> Fill( cos(lepton_new.angle(w_new)) );
```

The histogram must now be created and exported to a human-readable format. This is done at the level of the function Finalize by means of a series of ROOT commands. For the sake of the example, we have decided to normalize the histogram to the number of events \mathcal{N} expected for an integrated luminosity of 10 fb⁻¹,

$$\mathcal{N} = \sigma \mathcal{L}_{\rm int}/N$$
,

where σ is the cross section, in pb, corresponding to the process under consideration, $\mathcal{L}_{\mathrm{int}}$ the integrated luminosity (thus set to 10000 pb⁻¹) and N the number of events included in the samples. As stated above, the value of the cross section read from the LHE files is stored in the SampleFormat object summary (see Section 5.4) where the average over both input samples has been performed. The corresponding C++ implementation of the function Finalize reads

```
void user::Finalize(const SampleFormat& summary,
  const std::vector<SampleFormat>& files)
{
```

```
// Color of the canvas background
gStyle->SetCanvasColor(0);
// Turning off the border lines of the canvas
gStyle->SetCanvasBorderMode(0);
// Configuring statistics printing
gStyle->SetOptStat(111110);
// Creating the output root file
TCanvas* myCanvas = new TCanvas("myCanvas","");
// Setting background color
myHisto->SetFillColor(kBlue);
// Normalization of the histogram: L = 10 fb-1
double nrm = summary.mc()->xsection() * 10000. /
   static_cast<float>(summary.nevents());
myHisto->Scale(nrm);
// Setting axis title
myHisto->GetXaxis()->SetTitle("cos#theta^{*}");
// Drawing histogram
myHisto->Draw();
// Saving plot
myCanvas->SaveAs((outputName_+".eps").c_str());
```

After compiling the analysis and linking it to the external libraries with the help of the provided Makefile located in the SampleAnalyzer directory, the analysis can be executed,

}

```
SampleAnalyzer --analysis="W polarization from a top decay" \
   list.txt
```

where list.txt is a text file containing the absolute paths to the two event samples under consideration. A figure, named list.eps, is generated in the

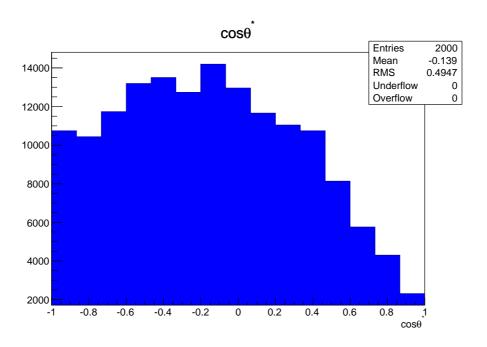


Figure 5: $\cos \theta^*$ distribution for the production of a top-antitop pair semi-leptonically decaying. The θ^* angle is defined as in the text.

directory where the executable is stored. The results, shown in Figure 5, agrees, for instance, with Ref. [62].

6. Conclusions

The task of performing a phenomenological analysis based on event files as those generated by Monte Carlo generators can be divided into three stages. Firstly, the event samples must be read and loaded into the memory of the computer. The format of these files depends in general on the level of sophistication of the analysis (at the parton level, hadron level or reconstructed level). Secondly, the analysis itself must be performed, *i.e.*, selection cuts are applied on the signal and background event samples in the aim of being able to extract information on the signal from the often overwhelming background. Finally, the results are outputted as histograms and/or cut-flow charts to improve their readability and interpretation.

In this work, we have presented MADANALYSIS 5, a new user-friendly and efficient framework aiming to facilitate the implementation of phenomenological analyses such as the one described above. We have explained how to

implement and run an analysis within this framework in a straightforward way and have given several detailed examples addressing the different facets of the program.

MADANALYSIS 5 is based on a multi-purpose C++ kernel, SAMPLEAN-ALYZER, which uses the ROOT platform. Compatible with most of the event file formats commonly used for analyzing parton-level, hadron-level and reconstructed events, MADANALYSIS 5 offers two modes of running according to the needs and the expertise of the user.

A highly-user friendly mode, the normal running mode of the program, uses the strengths of a Python interface to reduce the implementation of an analysis to a set of intuitive commands whose the syntax has been inspired by the Python programming language. Therefore, rather complex analyses can be achieved without too much efforts.

For users with more advanced C++ and Root programming skills, Madanalysis 5 can also be run in its expert mode. This overcomes the limitations of the normal mode of running which the scope is restricted by the set of functionalities that have been implemented. The user is in this case required to directly implement his analysis in C++ within the Sampleanalyzer framework, rendering the possibilities, at the analysis level, only limited by the imagination and the skills of the user. However, even if this mode of running is in principle more complicated to handle, the existence of many built-in functions and methods renders the task of implementing the analysis easier and more straightforward.

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Appendix A. Installation of the program

Appendix A.1. Requirements

For a proper running, Madanalysis 5 requires a couple of mandatory and optional external dependencies to be installed on the computer of the user. Concerning the optional packages, in the case they are not installed, several functionalities of Madanalysis 5 are deactivated but analyses of event files can still be performed. In contrast, the absence of one of the mandatory external library does simply not allow to use the program.

In order to run Madanalysis 5 locally, Python 2.6 or a more recent version (however not from the 3.x series) must be installed on the computer of the user [63]. We remind that in order to check the version of Python present on a system, it is enough to type in a shell

```
python --version
```

The installation of the Python package is one of the three mandatory requirements without which the program cannot run. The two other external packages that have to be installed are related to C++ and Root.

The Sampleanalyzer kernel requires the installation of a C++ compiler together with the associated Standard Template Libraries (STL). Since the validation procedure of Madanalysis 5 has been only achieved within the context of the GNU GCC compiler and since this compiler is available on most operating systems [64], we have adopted the choice of requiring the installation of GCC, version 4.3.x or 4.4.x, in order to have Madanalysis 5 properly running. We remind that the version of the GCC installed on a system can be obtained by issuing in a shell

```
g++ --version
```

Let us however stress that compatibility with any other C++ compiler is in principle ensured, but requires the modification of several core files of Madanalysis 5. Therefore, it is currently not supported.

Concerning the Root package, a version more recent than version 5.27 has to be installed [65], and the user has to check that the Python functionalities of the Root library are available. We recall that the version of Root present on a computer, together with the compatibility with the Python language, can be checked by issuing in a shell the commands

```
root-config --version
root-config --has-python
```

The three above-mentioned programs, *i.e.*, PYTHON, the GCC compiler and ROOT, must be available from any location of the computer. If this is not the case, the user has to modify the system variable \$PATH. Moreover, it is also left to the user to check that the paths to the associated header and library files are included in the environment variables of the GCC compiler, \$CPLUS_INCLUDE_PATH and \$LIBRARY_PATH. Finally, for a proper linking of the external libraries, the variable \$LD_LIBRARY_PATH (or, on MACOS operating systems, \$DYLD_LIBRARY_PATH) must also contain the paths to the libraries to be linked.

We now turn to the optional libraries that can be linked to MADANALYSIS 5. In order to handle zipped Monte Carlo event samples, the ZLIB library has to be installed on the system [66]. However, if ZLIB is not locally installed, the possibility of analyzing zipped event sample is simply deactivated, *i.e.*, the user will have to unzip the event files manually before running the code.

Appendix A.2. Downloading the program

It is recommended to always use the latest stable version of the MAD-ANALYSIS 5 package, which contains, when downloaded from the web, both the Python command line interface and the SampleAnalyzer framework. It can always be found together with an up-to-date manual on the webpage

```
http://madanalysis.irmp.ucl.ac.be
```

The package does not require any compilation or configuration. After having downloaded the tar-ball from the website, it can be unpacked either in the directory where MADGRAPH 5 is installed,

```
cd <path-to-madgraph>
tar -xvf ma5_xxxx.tgz
```

or in any location on the computer of the user,

```
mkdir <ma5-dirname>
cd <ma5-dirname>
tar -xvf ma5_xxxx.tgz
```

where xxxx stands for the version number of the MADANALYSIS 5 release which has been downloaded. When MADANALYSIS 5 is installed as a dependency of MADGRAPH 5, the list of predefined particle and multiparticle labels is directly updated from the MADGRAPH 5 standards. Moreover, this

allows to perform analyses at the time of the event generation by MAD-GRAPH 5 which, in this case, pilots MADANALYSIS 5. We emphasize that several predefined analyses exist, but the user has the freedom to tune the desired analysis to be performed according to his own aims.

Once installed and unpacked, MADANALYSIS 5 can be immediately launched by issuing in a shell

bin/ma5

as presented in Section 4.1.

Appendix A.3. Running Madanalysis 5

When started, Madanalysis 5 first checks that all the dependencies (GCC, Python, Root, ZLIB) are present on the system and that compatibility is ensured with the installed versions. In the case of any problem, a message is printed to the screen and the code exists if it is found that it cannot properly run. On the first session of Madanalysis 5, the Sample-Analyzer core is compiled behind the scene as a static library stored in the directory lib. For the next sessions, the kernel is only recompiled if the configuration of the system has changed (new version of the dependencies or of the main program).

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