

Coriolis and Centrifugal Forces

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Derivation of Coriolis in a Rotating Framework

- Imagine a position vector \mathbf{r} (**C in the figure**), rotating at a constant angular velocity $\boldsymbol{\Omega}$
- From a stationary (Eulerian) view, we are viewing \mathbf{r} rotate around the vertical axis at a constant speed $\boldsymbol{\Omega}$ from some fixed point away from the system
- From a Lagrangian view, we are viewing the system from the moving position vector \mathbf{r}
- Depending on frame (Eulerian or Lagrangian) of reference you view \mathbf{r} from, ***we observe different motions: (WHY?)***
- The motion of the vector should be equivalent between the frames, so there should be an apparent force in one frame that equates them
- The change of a vector in the inertial (sub i) frame can be seen in the rotating frame (sub r) as
 - $(\delta \mathbf{B})_i = (\delta \mathbf{B})_r + \boldsymbol{\Omega} \times \mathbf{B}$
 - \mathbf{B} is an arbitrary vector and $\boldsymbol{\Omega} \times \mathbf{B}$ is the velocity imparted by the rotating frame

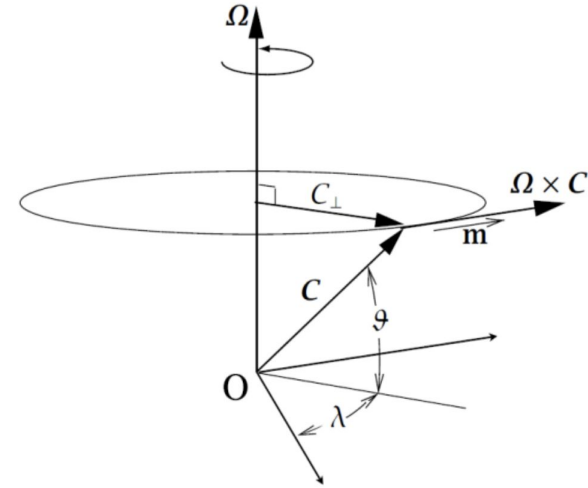


Fig. 2.1 A vector \mathbf{C} rotating at an angular velocity $\boldsymbol{\Omega}$. It appears to be a constant vector in the rotating frame, whereas in the inertial frame it evolves according to $(d\mathbf{C}/dt)_I = \boldsymbol{\Omega} \times \mathbf{C}$.

- Now we can look at how a velocity vector \mathbf{v}_i in the inertial frame changes in the rotating frame

$$\begin{aligned}
 \left(\frac{d\mathbf{v}_i}{dt}\right)_i &= \left(\frac{d\mathbf{v}_i}{dt}\right)_r + \boldsymbol{\Omega} \times \mathbf{v}_i && \text{Recall: } (\partial\mathbf{B})_i = (\partial\mathbf{B})_r + \boldsymbol{\Omega} \times \mathbf{B} \\
 &= \left(\frac{d}{dt}(\mathbf{v}_r + \boldsymbol{\Omega} \times \mathbf{r})\right)_r + \boldsymbol{\Omega} \times (\mathbf{v}_r + \boldsymbol{\Omega} \times \mathbf{r}) \\
 &= \left(\frac{d\mathbf{v}_r}{dt}\right)_r + \boldsymbol{\Omega} \times \left(\frac{d\mathbf{r}}{dt}\right)_r + \boldsymbol{\Omega} \times \mathbf{v}_r + \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}) \\
 &= \left(\frac{d\mathbf{v}_r}{dt}\right)_r + \boldsymbol{\Omega} \times \mathbf{v}_r + \boldsymbol{\Omega} \times \mathbf{v}_r + \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}) \\
 &= \left(\frac{d\mathbf{v}_r}{dt}\right)_r + 2\boldsymbol{\Omega} \times \mathbf{v}_r + \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})
 \end{aligned}$$

- Rearranging our in terms of \mathbf{v}_i in the rotating reference frame, we see two new terms on the RHS which are the **Coriolis and Centrifugal Forces**
- There are **apparent** forces imparted by the rotating frame!

$$\left(\frac{d\mathbf{v}_r}{dt}\right)_r = \left(\frac{d\mathbf{v}_i}{dt}\right)_i - 2\boldsymbol{\Omega} \times \mathbf{v}_r - \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})$$

Magnitude of Coriolis and Centrifugal Forces

- Coriolis $\sim 2\Omega U = 2 * 7 \times 10^{-5} \text{ 1/s} * 10 \text{ m/s} = \mathbf{1.4e-3 \text{ m/s}}$

$$\left(\frac{d\mathbf{v}_r}{dt}\right)_r = \left(\frac{d\mathbf{v}_i}{dt}\right)_i - 2\boldsymbol{\Omega} \times \mathbf{v}_r - \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})$$

What Does Coriolis do?

- In Cartesian coordinates, this is the vector components of Coriolis:

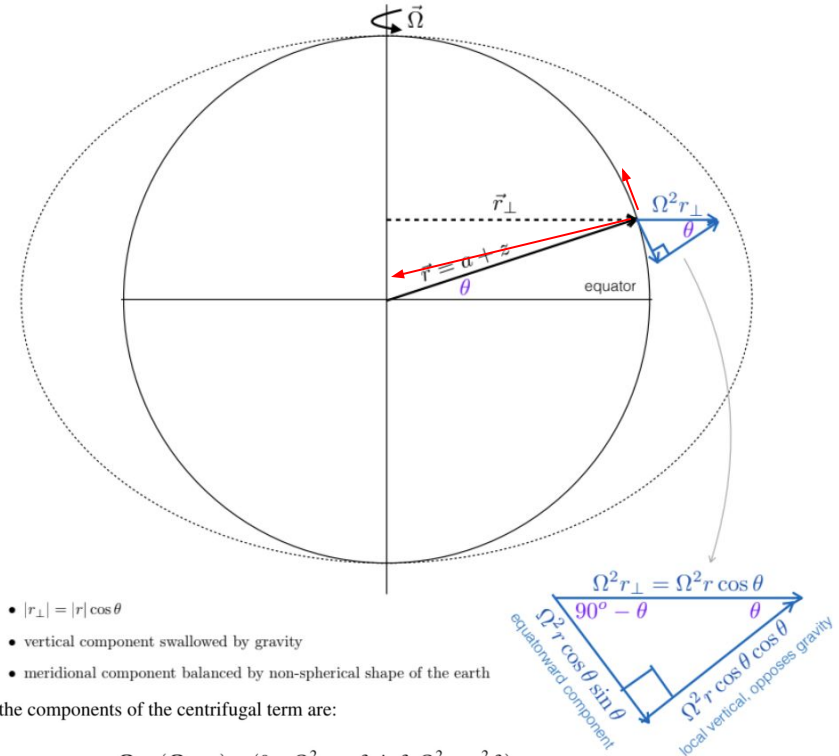
$$-2\boldsymbol{\Omega} \times \mathbf{v} = (-2\Omega v \cos \theta + 2\Omega v \sin \theta)\hat{\mathbf{i}} + -2\Omega u \sin \theta \hat{\mathbf{j}} + 2\Omega u \cos \theta \hat{\mathbf{k}}$$

- We can largely ignore the first term in i-hat-direction and the k-hat term due to their small magnitudes compared to large-scale horizontal motions and gravity, respectively
- Therefore with:
 - +u (eastward motion) → effect of coriolis is southward motion
 - +v (northward motion) → effect of coriolis is eastward motion
 - This is why we see large systems being steering to the right in the Northern Hemisphere
 - And so on...
- Consider the acceleration the Coriolis force has on moving objects by
 - $F = d^2x/dt^2 = dv/dt = 2\Omega v_y \sin(\Theta) \rightarrow$ Integrate equation twice to get distance coriolis affects objects $x = 2\Omega v_y \sin(\Theta)$; where v_y is velocity travelled in y-direction and a choose $\Theta = 45^\circ$
 - A naval war gun is shot northward is swayed 40 whole meters to right east and even a rifle shot is affected by almost 4 meters

Sporting event	y	v _y	v _y	t	x
	(m)	(km h ⁻¹)	(m s ⁻¹)	(s)	(cm)
Projectile type cases					
Cricket: slow bowler's delivery ^b	18	80	22.2	0.81	0.0752
Archery	65	300	83.3	0.78	0.261
Cricket: throw from boundary ^b	80	160	44.4	1.80	0.742
Football: 60 m kick	60	54	15	4.00	1.237
Golf: drive from the tee ^b	250	300	83.3	3.00	3.866
Rifle shot	500	1200	333	1.50	3.866
World War 1 Naval gun	24 000	2696	749	32.05	39.65 m

Derivation of Centrifugal in a Rotating Framework

- The centrifugal force is given by $\mathbf{F}_{ce} = -\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}) = -\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}_{\perp}) = \Omega^2 \mathbf{r}_{\perp}$, where \mathbf{r}_{\perp} is given by $r \cos(\theta)$ where it is perpendicular to the axis of rotation
- This total vector is broken down into two components, one in the \hat{j} -unit vector and one in the \hat{k} -unit vector
- We notice a \hat{j} vector pointing equatorward
 - So why are we sliding towards the Equator?
- This is because the Earth is an oblate spheroid and the mass has been distributed according to the long-term effect of the centrifugal force with there being more mass at the Equator
- The shape of the Earth now results in gravity that generates a term in the $+\hat{j}$ direction which opposes the $-\hat{j}$ centrifugal term
- For simplicity, the meridional terms (\hat{j} hat) are considered to cancel out and are omitted
- By this we assume that the local vertical is now in the opposite direction of the effective gravity



Centrifugal Force: Where is it?

- Centrifugal $\sim \Omega^2 a = (7 \times 10^{-5})^2 \times 6 \times 10^6 \text{m} = 3 \times 10^{-2} \text{ m/s}^2$
- Coriolis $\sim 2\Omega U = 2 \times 7 \times 10^{-5} \text{ 1/s} \times 10 \text{ m/s} = 1.4 \times 10^{-3} \text{ m/s}^2$
- **As we can see Centrifugal > Coriolis by one order of magnitude**

So where is the Centrifugal term in the momentum equations?

- Centrifugal force is actually hidden in the potential on the RHS
- g is no longer the gravity constant 9.8 m/s^2 rather it is what is called effective gravity $\mathbf{g} \equiv \mathbf{g}_{\text{eff}} \equiv \mathbf{g}_{\text{grav}} + \Omega^2 \mathbf{r}_{\perp} = -\nabla \Phi$
- Since gravity (9.8 m/s^2) is so much larger than the centrifugal force magnitude, it is conveniently placed in it

$$\frac{D\mathbf{v}}{Dt} + 2\boldsymbol{\Omega} \times \mathbf{v} = -\frac{1}{\rho} \nabla p - \nabla \Phi$$
