

ATS 601: Atmospheric Dynamics I  
**HW2: Momentum Equations**

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Estimate of Time to Completion: 5 hours

Maximum Allotted Time to Completion: 10 hours

Actual Time to Completion: 7 hours

ATS 601: HW#2 DUE September 3<sup>rd</sup>

1.

 (a) Starting From Mass continuity  $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{u}) = 0$ 

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{u}) = 0$$

$$\frac{\partial \rho}{\partial t} + \rho \nabla \cdot \vec{u} + \vec{u} \cdot \nabla \rho = 0 \quad \bullet \text{ given that } \vec{u}(x,t) = u \text{ i.e. } \nabla \cdot \vec{u} = 0$$

$$\frac{\partial \rho}{\partial t} + u \cdot \nabla \rho = 0$$

 $u = u(x,t)$  only x-dimension

$$\boxed{\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} = 0} \quad \checkmark$$

 (b) Show that (a) =  $\frac{D\rho}{Dt} = 0$ 

 From the material derivative:  $\frac{D\rho}{Dt} = \frac{\partial \rho}{\partial t} + \vec{u} \cdot \nabla \rho$ 

$$\therefore \frac{\partial \rho}{\partial t} = \frac{D\rho}{Dt} - u \cdot \nabla \rho$$

Plug into (a)

$$\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} = 0$$

$$\left( \frac{D\rho}{Dt} - u \cdot \nabla \rho \right) + u \frac{\partial \rho}{\partial x} = 0$$

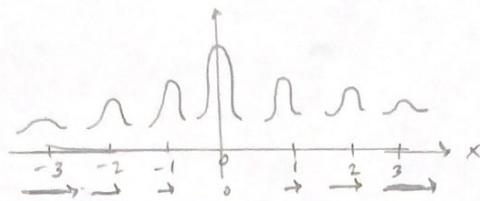
$$\boxed{\frac{D\rho}{Dt} = 0} \quad \checkmark$$

 (c) The solution to the advection equation for  $u(x,t) = u$  is  $\rho(x,t) = \rho_0(x - ut)$  w/ I.C.  $\rho(x,t=0) = \rho_0(x)$ .

$\Rightarrow$  The blob will be advected to the right @ some constant  $u$ , given by  $\rho(x,t) = \rho_0(x - ut)$ . Over some time we see  $\rho(x,1) = \rho_0(x - u)$  and  $\rho(x,2) = \rho_0(x - 2u)$  and so on.

$\Rightarrow$  The eulerian perspective we are watching from a stationary point and see the blob move right wrt us. The lagrangian perspective is when we are moving along w/ the blob, to the right, but we are stationary from this point of view.





(d) consider  $u(x,t) = \Lambda|x|$   $\Lambda \rightarrow \text{constant} > 0$   $u(x,t) = \Lambda|x|$

$\Rightarrow$  the tracer blob will be moving to the right however it will slow down as  $x$  approaches zero, stop @  $x=0$  and then speed up past  $x=0$ . Due to mass conservation, the blob will squish & squeeze

$$\begin{aligned} \Rightarrow \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{u}) &= 0 \\ \frac{\partial \rho}{\partial t} + \rho \frac{\partial u}{\partial x} + u \frac{\partial \rho}{\partial x} &= 0 \\ \frac{\partial \rho}{\partial t} + \rho \frac{\partial (\Lambda|x|)}{\partial x} + (\Lambda|x|) \frac{\partial \rho}{\partial x} &= 0 \end{aligned}$$

$$\left. \begin{aligned} \frac{\partial \rho}{\partial t} + \rho \Lambda + \Lambda|x| \frac{\partial \rho}{\partial x} &= 0 & x > 0 \\ \frac{\partial \rho}{\partial t} - \rho \Lambda + \Lambda|x| \frac{\partial \rho}{\partial x} &= 0 & x < 0 \end{aligned} \right\}$$

$\Rightarrow$  This new continuity equation has an extra  $\rho$  term scaled by  $\Lambda$

and  $\frac{\partial \rho}{\partial x}$  is scaled by a non-constant term  $\Lambda|x|$  which will always

be positive since  $\Lambda > 0$  and  $|x| \geq 0$ . I think its also useful to

equate some terms instead of having it equal to zero:  $\frac{\partial \rho}{\partial t} = -\Lambda|x| \frac{\partial \rho}{\partial x} - \rho \Lambda$

which says the change in density w/ time equals the negative rate of change of density w/ spatial  $x$ , scaled by a  $\Lambda|x|$ , minus some other term  $-\rho \Lambda$ .

(e) Plugging in these terms on the left to our continuity eqn from (d):

$$\begin{aligned} \frac{\partial \rho_j^i}{\partial t} &\approx \frac{\rho_j^{i+1} - \rho_j^i}{\Delta t} \\ \frac{\partial \rho_j^i}{\partial x} &\approx \frac{\rho_j^i - \rho_{j-1}^i}{\Delta x} \end{aligned}$$

$$\Rightarrow \frac{\rho_j^{i+1} - \rho_j^i}{\Delta t} + \rho_j^i \Lambda + \Lambda|x| \frac{(\rho_j^i - \rho_{j-1}^i)}{\Delta x} = 0$$

$$\rho_j^{i+1} - \rho_j^i + \Delta t \rho_j^i \Lambda + \frac{\Delta t \Lambda |x|}{\Delta x} (\rho_j^i - \rho_{j-1}^i) = 0$$

$$\left. \begin{aligned} \rho_j^{i+1} &= \rho_j^i - \Delta t \Lambda \rho_j^i - \frac{\Delta t \Lambda |x|}{\Delta x} (\rho_j^i - \rho_{j-1}^i) & x > 0 \\ \rho_j^{i+1} &= \rho_j^i + \Delta t \Lambda \rho_j^i - \frac{\Delta t \Lambda |x|}{\Delta x} (\rho_j^i - \rho_{j-1}^i) & x < 0 \end{aligned} \right\}$$

# ATS 601: Atmospheric Dynamics I

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```
In [1]: %matplotlib inline
```

```
In [2]: import numpy as np
import math
import xarray as xr
import matplotlib.pyplot as plt
import warnings
```

/anaconda3/lib/python3.7/site-packages/dask/config.py:168: YAMLLoadWarning: calling yaml.load() without Loader=... is deprecated, as the default Loader is unsafe. Please read <https://msg.pyyaml.org/load> for full details.

```
data = yaml.load(f.read()) or {}
/anaconda3/lib/python3.7/site-packages/distributed/config.py:20: YAMLLoadWarning: calling yaml.load() without Loader=... is deprecated, as the default Loader is unsafe. Please read https://msg.pyyaml.org/load for full details.
```

```
defaults = yaml.load(f)
```

```
In [3]: large = 28; med = 24; small = 15
params = {'axes.titlesize': small,
          'legend.fontsize': med,
          'figure.figsize': (8, 4),
          'axes.labelsize': small,
          'axes.titlesize': small,
          'xtick.labelsize': small,
          'ytick.labelsize': small,
          'figure.titlesize': med,
          'axes.titlepad': 10,
          'axes.facecolor': 'black'} # For electric background plots
plt.rcParams.update(params)
```

## Problem 1: Passive tracer advection

1e) Step 2: Convert your equation from (d) into numerical form, and solve for the density  $\psi_j^{i+1}$  at a future time  $i + 1$  at location  $j$ .

*i,j notational form of continuity equation*

$$\psi_j^{i+1} = \psi_j^i - \Delta t \Lambda \psi_j^i - \frac{\Delta t \Lambda |x|}{\Delta x} (\psi_j^i - \psi_{j-1}^i) \text{ for } x > 0 \text{ and}$$

$$\psi_j^{i+1} = \psi_j^i + \Delta t \Lambda \psi_j^i - \frac{\Delta t \Lambda |x|}{\Delta x} (\psi_j^i - \psi_{j-1}^i) \text{ for } x < 0.$$

*Numerical form of continuity equation*

$\text{psi}[xi, ti + 1] = \text{psi}[xi, ti] + (dt * A * \text{psi}[xi, ti]) - (A * \text{np.abs}(x[xi]) * (dt/dx) * (\text{psi}[xi, ti] - \text{psi}[xi - 1$   
for  $x > 0$  and

$\text{psi}[xi, ti + 1] = \text{psi}[xi, ti] - (dt * A * \text{psi}[xi, ti]) - (A * \text{np.abs}(x[xi]) * (dt/dx) * (\text{psi}[xi, ti] - \text{psi}[xi - 1$   
for  $x < 0$ .

```
In [4]: dt = 0.01
dx = 0.10
mu = 5
A = 0.1 #lambda
sigma = 1

# Define time and space and psi arrays
t = np.arange(0,22,dt)
x = np.arange(-100,100,dx)
```

```
In [5]: # assign array for psi
psi = np.zeros([len(x),len(t)])

# intial condition
psi_0 = np.exp(-1*((x-mu)**2)/(2*sigma**2))
psi[:,0] = psi_0

# code up numerical equation
# tmp = psi[xi,ti]+(dt*A*psi[xi,ti])-(A*np.abs(x[xi])*dt/dx*(psi[xi,ti]-p
si[xi-1,ti]))

# enumerate gets the index and value
for ti,tv in enumerate(t[1:]):
    for xi,xv in enumerate(x[1:]):

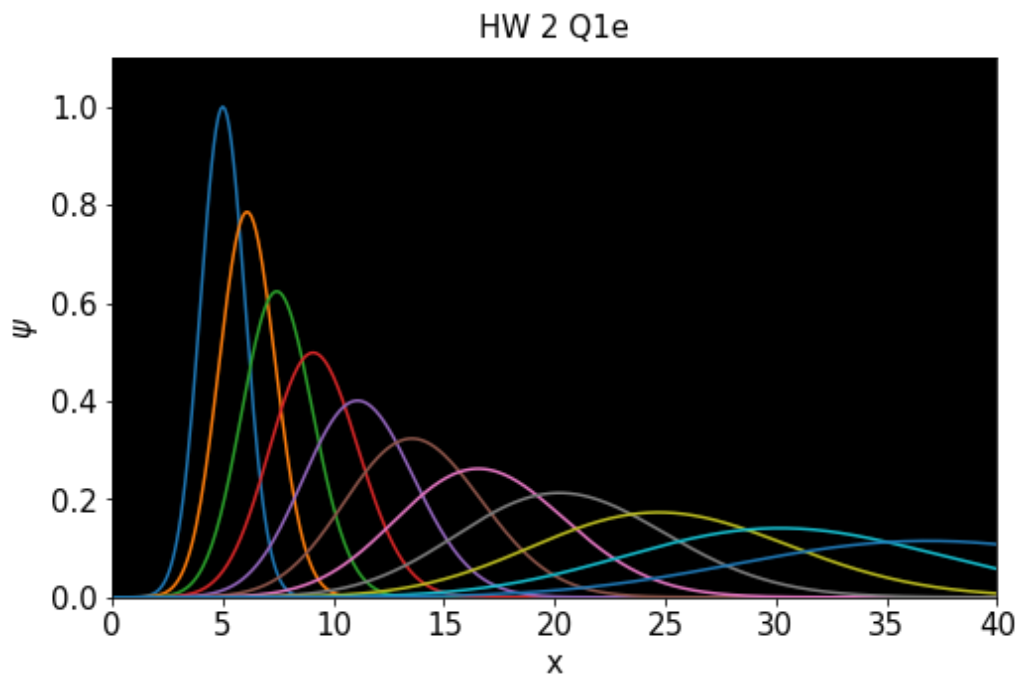
        # calculating psi at each x and at each t index

        # case where value of x < 0
        if xv < 0:
            psi[xi,ti+1] = psi[xi,ti]+(dt*A*psi[xi,ti])-(A*np.abs(x[xi])*
(dt/dx)*(psi[xi,ti]-psi[xi-1,ti]))

        # case where value of x > 0
        elif xv > 0:
            psi[xi,ti+1] = psi[xi,ti]-(dt*A*psi[xi,ti])-(A*np.abs(x[xi])*
(dt/dx)*(psi[xi,ti]-psi[xi-1,ti]))
```

```
In [6]: plt.figure(figsize=(8,5))
        for i in np.arange(0,len(t),200):
            plt.plot(x,psi[:,i]);
            # plt.legend([i])

        plt.xlabel('x');
        plt.ylabel('$\psi$');
        plt.xlim([0,40])
        plt.ylim([0,1.1])
        plt.title('HW 2 Q1e');
```



## Problem 2: Reading-in global atmospheric data

```
In [7]: ds = xr.open_dataset('/Users/tyler/Desktop/601/homework_2_ERA-interim_300
        hPa_2014-Jan1-8.nc')
```

```
In [8]: u = ds.u
        v = ds.v
```

```
In [9]: latitude = ds.latitude
        longitude = ds.longitude
        time = ds.time
```

## Problem 3: Approximations to the momentum equations (gravity, viscosity)

(a) Using the data from Problem 2, estimate typical values\* of the local accelerations of the zonal and meridional wind in the Northern midlatitudes (30-60 N), namely:  $\partial u / \partial t$  and  $\partial v / \partial t$

```
In [10]: u_midlat = u.sel(latitude=slice(60,30))
v_midlat = v.sel(latitude=slice(60,30))
```

```
In [11]: dudt = np.zeros(np.shape(u_midlat))
dvdt = np.zeros(np.shape(v_midlat))

for i in np.arange(1,np.shape(dudt)[0]-1):
    dudt[i-1,:,:] = (u_midlat[i+1,:,:] - u_midlat[i-1,:,:])/(2*6*60*60)
    dvdt[i-1,:,:] = (v_midlat[i+1,:,:] - v_midlat[i-1,:,:])/(2*6*60*60)
```

```
In [12]: # calculate typical values using RMS
def RMS(matrix):
    tmp = matrix**2
    tmp = np.mean(tmp)
    tmp = np.sqrt(tmp)
    return tmp
```

```
In [13]: print('The RMS or typical values of the zonal wind is: '+str(RMS(dudt))+
m/s')
print('The RMS or typical values of the meridional wind is: '+str(RMS(dvd
t))+ ' m/s')
```

The RMS or typical values of the zonal wind is: 0.00024879785447811195  
m/s

The RMS or typical values of the meridional wind is: 0.00034777278712432  
18 m/s

(b) Estimate the magnitude of the wind tendencies due to viscosity in the momentum equations for the atmosphere. You can do this a variety of ways, some easier than others. I would suggest following the scaling method in Vallis (link to textbook on Canvas) that uses the Reynolds number (section 1.11.1, page 44), as this will provide a ballpark solution in record time. Alternatively, you can try and calculate the laplacian of the wind on the sphere and compute the actual term (much more difficult).

- I found a typical value of viscosity for air to be  $\nu = 1.81e - 5$  at 15 degrees C. From this page (<https://www.ncdc.noaa.gov/sotc/global/202004>), NOAA says the average global temperature from the 20th century was 13.7C (I couldnt find global average for the 2000-2021.) It will probably be higher in 2000-2021 so I think this is close enough.
- I chose a typical characteristic length scale as  $L = 10^7 m$  because from lecture I think from lecture it said synoptic scales were between  $10^6$  and  $10^8 m$ . And definitely we're looking at midlatitudes between 30 and 60N which is definitely a synoptic scale.
- For a characteristic velocity scale I took the average of the RMS of the zonal and meridional component of wind speeds. For a typical velocity, I found  $U = 35.7 m/s$ .
- Finally, using the Force due to viscosity scaling in Vallis:  $(\nu \frac{U}{L^2})$

```
In [14]: # Typical value of viscosity (https://www.engineersedge.com/physics/viscosity_of_air_dynamic_and_kinematic_14483.htm)
v = 1.81e-5 # units: kg/(m*s)

# Typical characteristic length scale (L) (between 10^6 and 10^8)
L = 10**7 # units: meters

# Typical characteristic velocity scale (average both du_dt and dv/dt)
U = np.sqrt(RMS(u_midlat)**2 + RMS(v_midlat)**2)

# Contribution due to viscosity is v*U/L^2
F_v = v*U/(L**2)
print('Typical force due to viscosity is: '+str(np.array((F_v))))
```

Typical force due to viscosity is: 6.460452476501464e-18

- The typical time tendencies due to force of viscosity is  $6.5e - 18 \text{ kg/m}^2 \text{ s}^2$ . (F/m = acceleration)
- Compared to time tendencies of wind (du/dt and dv/dt), from a) those values were magnitude  $10^{-4} \text{ m/s}^2$ .
- Therefore, I since the difference in magnitudes is so large, I think it is okay to neglect viscosity in this course.