

## Correlations between Ca II H&K Emission and the Gaia M dwarf Gap

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### ABSTRACT

Previous work has demonstrated a paucity of Hydrogen Alpha emission at the same Gaia G Magnitude as the Jao Gap in the solar neighborhood. The exact mechanism which results in this paucity is as yet unknown; however, the authors of the originating paper suggestion that it may be the result of complex variations to a stars magnetic topology driven by the Jao Gaps characteristic formation and breakdown of stars radiative transition zones. Here I present a brief summary of a potential extension to this work looking at Ca II H&K emission lines. Preliminary work with archival data shows a much stronger correlation between the calcium emission lines and the Jao Gap than was observed between hydrogen emission and the Jao Gap. If this observation withstands further testing then it may provide a new way to locate the gap in populations which lack the large counting statistics currently required — and which are only practically available from Gaia data at moment.

*Keywords:* Stellar Evolution (1599) — Stellar Evolutionary Models (2046)

### 1. INTRODUCTION

Due to the initial mass requirements of the molecular clouds which collapse to form stars, star formation is strongly biased towards lower mass, later spectral class stars when compared to higher mass stars. Partly as a result of this bias and partly as a result of their extremely long main-sequence lifetimes, M Dwarfs make up approximately 70 percent of all stars in the galaxy. Moreover, some planet search campaigns have focused on M Dwarfs due to the relative ease of detecting small planets in their habitable zones (e.g. [Nutzman & Charbonneau 2008](#)). M Dwarfs then represent both a key component of the galactic stellar population as well as the possible set of stars which may host habitable exoplanets. Given this key location M Dwarfs occupy in modern astronomy it is important to have a thorough understanding of their structure and evolution.

[Jao et al. \(2018\)](#) discovered a novel feature in the Gaia Data Release 2 (DR2)  $G_{BP} - G_{RP}$  color-magnitude-diagram. Around  $M_G = 10$  there is an approximately

17 percent decrease in stellar density of the sample of stars [Jao et al. \(2018\)](#) considered. Subsequently, this has become known as either the Jao Gap, or Gaia M Dwarf Gap. Following the initial detection of the Gap in DR2 the Gap has also potentially been observed in 2MASS ([Skrutskie et al. 2006](#); [Jao et al. 2018](#)); however, the significance of this detection is quite weak and it relies on the prior of the Gap’s location from Gaia data. Further, the Gap is also present in Gaia Early Data Release 3 (EDR3) ([Jao & Feiden 2021](#)). These EDR3 and 2MASS data sets then indicate that this feature is not a bias inherent to DR2.

The Gap is generally attributed to convective instabilities in the cores of stars straddling the fully convective transition mass (0.3 - 0.35  $M_\odot$ ) ([Baraffe & Chabrier 2018](#)). These instabilities interrupt the normal, slow, main sequence luminosity evolution of a star and result in luminosities lower than expected from the main sequence mass-luminosity relation ([Jao & Feiden 2020](#)).

The Jao Gap, inherently a feature of M Dwarf populations, provides an enticing and unique view into the interior physics of these stars ([Feiden et al. 2021](#)). This is especially important as, unlike more massive stars, M Dwarf seismology is infeasible due to the short periods and extremely small magnitudes which both radial

and low-order low-degree non-radial seismic waves are predicted to have in such low mass stars (Rodríguez-López 2019). The Jao Gap therefore provides one of the only current methods to probe the interior physics of M Dwarfs.

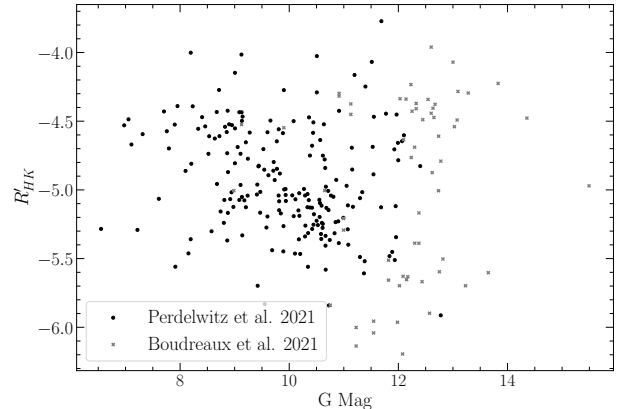
(Jao et al. 2023) identify the Jao gap as a strong discontinuity point for activity in M dwarfs. Two primary observations from their work are that the Gap serves as a boundary where very few active stars in their sample of 640 M dwarfs exist below the gap and that the overall downward trend of activity moving to fainter magnitudes is anomalously high in within the 0.2 mag range of the gap. Their figures 3 and 13 are of particular relevance here and have been included below for convince. Based on previous work from Spada & Lanzafam 2020, Curtis et al. 2020, and Dungee et al. 2022 the authors propose that the mechanism resulting in the reduced fraction of active stars within the gap is that as the radiative zone disipates due to core expansion, angular momentum from the outter convective zone is dumped into the core resulting in a faster spin down than would otherwise be possible. Effectively the core of the star acts as a sink, reducing the amount of angular momentum which needs to be lost by magnetic breaking for the outer convective region to reach the same angular velocity. Given that  $H\alpha$  emission is strongly coupled magnetic activity in the lower photosphere and that a stars angular velocity is a primary factor in its magnetic activity, a faster spin down will serve to more quickly dampen  $H\alpha$  activity.

## 2. CORRELATION

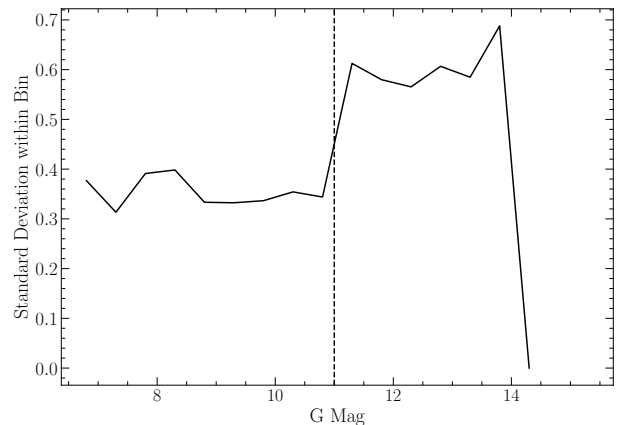
Using Ca II H&K emission data from Boudreaux et al. (2022) and Perdelwitz et al. (2021) (quantified using the  $R_{HK}$  metric) we investigate the correlation between the Jao Gap magnitude and stellar magnetic activity. We are more statistically limited here than past authors have been due to therequirment for high resolution spectroscopic data when measuring Calcium emission; however, this is balanced by the aparent stronger correlation between Calcium emission and the Jao gap when compared to  $H\alpha$  emission.

The merged dataset is presented in Figure 1. There is a visual discontinuity around the Jao Gap mangitude; however, this manifests as an increase in the spread of the emission measurments rather than a change in the mean value. In order to quantify the signifigance of this discontinuity we measure the false alarm probability of the change in standard deviation.

First bin the merged dataset into bins with a width of 0.5 mag. In each bin we measure the standard deviation about the mean of the data. The results of this are



**Figure 1.** Merged Dataset from Boudreaux et al. (2022); Perdelwitz et al. (2021). Note the increase in the spread of  $R_{HK}$  around the Jao Gap Magnitude.

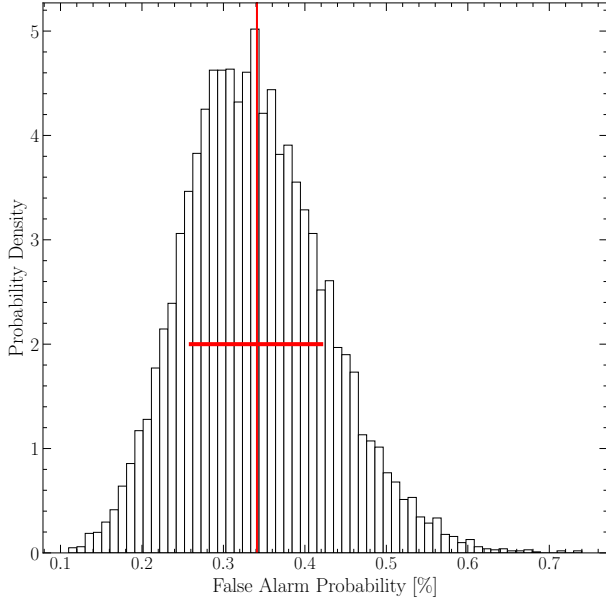


**Figure 2.** Standard deviation of Calcium emission data within each bin. Note the discontinuity near the Jao Gap Magnitude.

shown in Figure 2. In order to measure the false alarm probability of this discontinuity we first resample the merged calcium emission data based on the associated uncertainties for each datapoint as presented in their respective publications. Then, for each of these “resample trials” we measure the probability that a change in the standard deviation of the size seen would happen purley due to noise. Results of this test are show in in Figure 3.

This rapid increase star-to-star variability would only arise due purley to noise  $0.3 \pm 0.08$  percent of the time and is therefore likeley either a true effect or an alias of some sample bias. **COME BACK TO HERE TO FLUSH OUT SAMPLE BIAS SECTION.**

If the observed increase in variability is not due to a sample bias and rather is a physically driven effect then there is an obvious similarity between these findings and those of (Jao et al. 2023). Specifically we find a increase

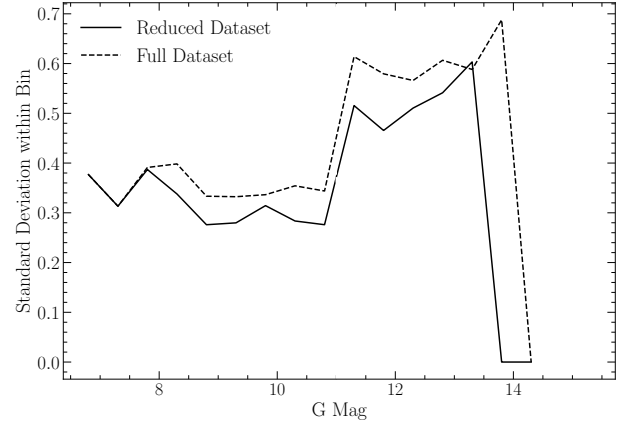


**Figure 3.** Probability distribution of the false alarm probability for the discontinuity seen in Figure 2. The mean of this distribution is  $0.341\% \pm 0.08$ .

in variability just below the magnitude of the gap. Moreover, this variability increase is primarily driven by an increase in the number of low activity stars (as opposed to an increase in the number of high activity stars). We can further investigate the observed change in variability for only low activity stars by filtering out those stars at or above the saturated threshold for magnetic activity. Boudreaux et al. (2022) identify  $\log(R'_{HK}) = -4.436$  as the saturation threshold. We adopt this value and filter out all stars where  $\log(R'_{HK}) \geq -4.436$ . Applying the same analysis to this reduced dataset as was done to the full dataset we still find a discontinuity at the same location (Figure 4). This discontinuity is of a smaller magnitude and consequently is more likely to be due purely to noise, with a  $7 \pm 0.2$  percent false alarm probability. This false alarm probability is however only concerned with the first point after the jump in variability. If we consider the false alarm probability of the entire high variability region then the probability that the high variability region is due purely to noise drops to  $1.4 \pm 0.04$  percent.

We observe a strong, likely statistically significant, discontinuity in the star-to-star variability of Ca II K & K emission coincident with the magnitude of the Jao Gap. However, modeling is required to determine if this discontinuity may be due to the same underlying physics.

### 3. MODELING



**Figure 4.** Spread in the magnetic activity metric for the merged sample with any stars  $\log(R'_{HK}) > -4.436$  filtered out.

One of the most pressing questions related to this work is whether or not the increased star-to-star variability in the activity metric and the Jao Gap, which are coincident in magnitude, are driven by the same underlying mechanism. The challenge when addressing this question arises from current computational limitations. Specifically, the kinds of three dimensional magneto-hydrodynamical simulations — which would be needed to derive the effects of convective kissing instabilities on the magnetic field of the star — are infeasible to run over gigayear timescales while maintaining thermal timescale resolutions needed to resolve periodic mixing events.

In order to address this and answer the specific question of *could kissing instabilities result in increased star-to-star variability of the magnetic field*, we adopt a very simple toy model. Kissing instabilities result in transient radiative zone separating the core of a star (convective) from its envelope (convective). When this radiative zone breaks down two important things happen: one, the entire star becomes mechanically coupled, and two, convective currents can now move over the entire radius of the star. Jao et al. (2023) propose that this mechanical coupling may allow the stars core to act as an angular momentum sink thus accelerating a stars spin down and resulting in anomalously low H $\alpha$  emission.

Regardless of the exact mechanism by which the magnetic field may be effected, it is reasonable to expect that both the mechanical coupling and the change to the scale of convective currents will have some effect on the stars magnetic field. On a microscopic scale both of these will change how packets of charge within a star move and may serve to disrupt a stable dynamo. Therefore, in the model we present here we make only one primary assumption: *every mixing event may modify the stars magnetic field by some amount*. Within our model

this assumption manifests as a random linear perturbation applied to some base magnetic field at every mixing event. The strength of this perturbation is sampled from a normal distribution with some standard deviation,  $\sigma_B$ .

Synthetic stars are sampled from a grid of stellar models evolved using the Dartmouth Stellar Evolution Program (DSEP). Each stellar model was evolved using a high temporal resolution (timesteps no larger than 10,000 years [Check this](#)) and typical numerical tolerances of one part in  $10^5$ . Each model was based on a GS98 ([CITATION](#)) solar composition with a mass range from  $0.3 M_\odot$  to  $0.4 M_\odot$ . Finally, models adopt OPLIB high temperature radiative opacities, Ferguson 2004 low temperature radiative opacities, and include both atomic diffusion and gravitational settling. A Kippenhan-Iben diagram showing the structural evolution of a model within the gap is shown in Figure 5.

Each synthetic star is assigned some base magnetic activity ( $B_0 \sim \mathcal{N}(1, \sigma_B)$ ) and then the number of mixing events before some age  $t$  are counted based on local maxima in the core temperature. The toy magnetic activity at age  $t$  for the model is given in Equation 1. An example of the magnetic evolution resulting from this model is given in Figure 6. Fundamentally, this model presents magnetic activity variation due to mixing events as a random walk and therefore results will increasingly diverge over time.

$$B(t) = B_0 + \sum_i B_i \sim \mathcal{N}(1, \sigma_B) \quad (1)$$

Applying the same analysis to these models as was done to the observations as described in Section [X.X](#) we find that this simple model results in a qualitatively similar trend in the standard deviation vs. magnitude graph. The interpretation here is important, what this qualitative similarity demonstrates is that it may be reasonable to expect kissing instabilities to result in the observed increased star-to-star variation. Importantly, we are not able to claim that kissing instabilities *do* lead to these increased variations, only that they reasonably could. Further modeling, observational, and theoretical efforts will be needed to more definitively answer this question.

#### 4. CONCLUSION

The Jao Gap provides an intriguing probe into the interior physics of M Dwarfs stars where traditional methods of studying interiors break down. However, before detailed physics may be inferred it is essential to have

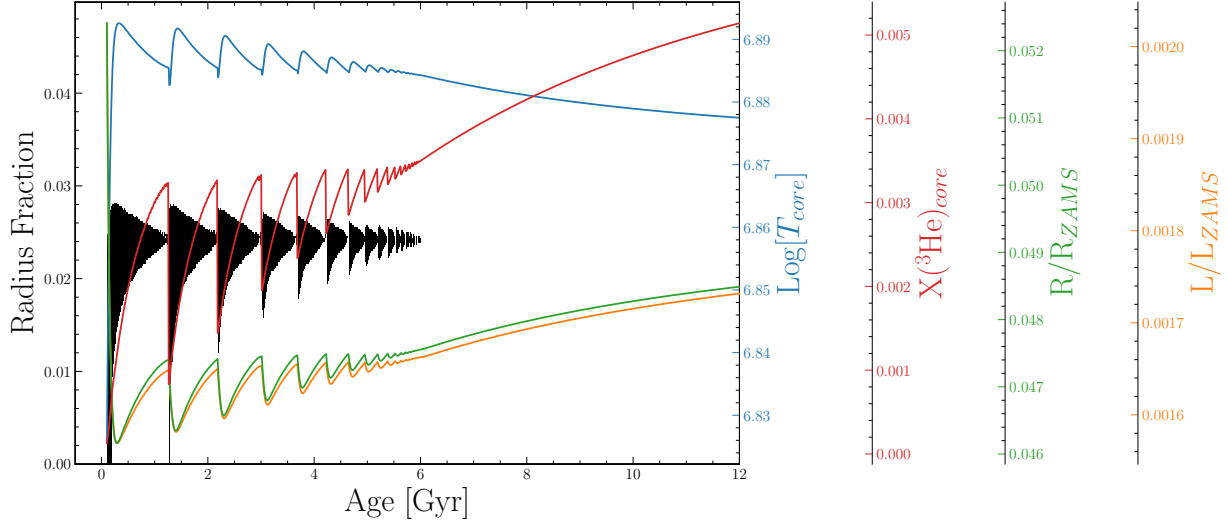
models which are well matched to observations. Here we investigate whether the OPLIB opacity tables reproduce the Jao Gap location and structure more accurately than the widely used OPAL opacity tables. We find that while the OPLIB tables do shift the Jao Gap location more in line with observations, by approximately 0.05 magnitudes, the shift is small enough that it is likely not distinguishable from noise due to population age and chemical variation. However, future measurement of  $[\text{Fe}/\text{H}]$  for stars within the gap will be helpful in constraining the degree to which the gap should be smeared by these theoretical models.

We also find that both the color and magnitude of the Jao Gap are correlated to the convective mixing length parameter. Specifically, a lower mixing length parameter will bring the gap in the populations presented in this paper more in line with the current best estimate for the actual gap magnitude. Using this relation it may be possible for mixing length to be calibrated for low mass stars such that models match the Jao Gap location. Further, the Jao gap location may provide a test of alternative convection models such as entropy calibrated convection ([Spada et al. 2021](#)). Both of these potential uses require careful handling of other uncertainties such as the uncertainties in bolometric correction, population composition, and population age. As we currently do not have reason to suspect that the mixing length for the low mass stars in the DR2 and ERD3 CMD is substantially lower than that of the sun we leave the investigation of these potential additional uses for future work.

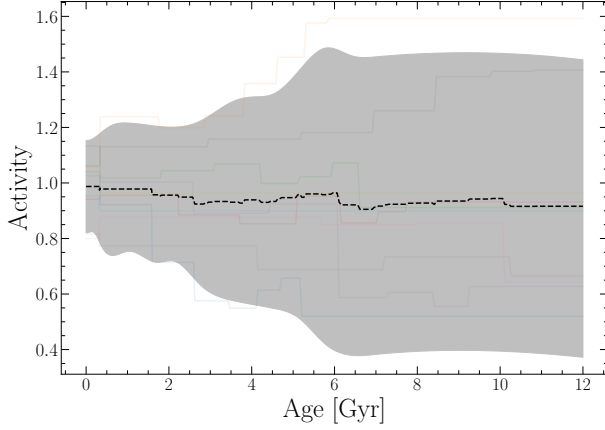
Finally, we do not find that the OPLIB opacity tables help in reproducing the as yet unexplained wedge shape of the observed Gap.

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*Software:* The Dartmouth Stellar Evolution Program (DSEP) ([Dotter et al. 2008](#)), `BeautifulSoup` ([Richardson 2007](#)), `mechanize` ([Chandra & Varanasi 2015](#)), `FreeEOS` ([Irwin 2012](#)), `pyTOPSScrape` ([Boudreaux 2022](#))



**Figure 5.** Kippenhan-Iben diagram for a 0.345 solar mass star. Note the periodic mixing events (where the plotted curves peak).



**Figure 6.** Example of the toy model presented here resulting in increased divergence between stars magnetic fields. The shaded region represents the maximum spread in the two point correlation function at each age.

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