11

12

13

14

15

16

17

18

19

20

21

22

Chemically Self-Consistant Modeling of the Globular Cluster NGC 2808 and its Effects on the Inferred Helium abundance of Multiple Stellar Populations.

EMILY M. BOUDREAUX, BRIAN C. CHABOYER, AMANDA ASH, REHNATA HOH, AND GREGORY FEIDEN

¹Department of Physics and Astronomy, Dartmouth College, Hanover, NH 03755, USA
 ²Department of Physics and Astronomy, University of North Georgia, Dahlonega, GA 30533, USA

ABSTRACT

The Helium abundances in the multiple populations which are now known to comprise all closely studied Milky Way globular clusters are often inferred by fitting isochohrones generated from stellar evolutionary models to globular cluster photometry. It is therefore important to build stellar models that are chemically self-consistent in terms of their structure, atmosphere, and opacity. In this work we present the first chemically self-consistent stellar models of the Milky Way Globular Cluster NGC 2808 using MARCS model atmospheres, OPLIB high-temperature radiative opacities, and AESOPUS low-temperature radiative opacities. These stellar models were fit to the NGC 2808 photometry using Fidanka , a new software tool that was developed optimally fit cluster photometry to isochrones and for population synthesis. Fidanka can determine, in a relatively unbiased way, the ideal number of distinct populations which exist within a dataset and then fits isochrones to each population. We achieve this through a combination of Bayesian Gaussian Mixture Modeling and a novel number density estimation algorithm. Using Fidanka and F275W-F814W photometry from the Hubble UV Globular Cluster Survey we find that the helium abundance of the second generation of stars in NGC 2808 is higher than the first generation by $15\pm3\%$. This is in agreement with previous studies of NGC 2808.

Keywords: Globular Clusters (656), Stellar evolutionary models (2046)

1. INTRODUCTION Globular clusters (GCs) are among the oldest observ-

²³ able objects in the universe (Peng et al. 2011). They are characterized by high densities with typical half²⁵ light radii of ≤10 pc (van den Bergh 2010), and typi²⁶ cal masses ranging from 10^4 – 10^5 M_☉ (Brodie & Strader 2006) — though some GCs are significantly larger than these typical values (e.g. ω Cen, Richer et al. 1991). GCs provide a unique way to probe stellar evolution (Baumgardt & Makino 2003), galaxy formation models (Boylan-Kolchin 2018; Kravtsov & Gnedin 2005), and dark matter halo structure (Hudson & Robison 2018). The traditional view of Globular Clusters was that they consisted of a single stellar population (SSP, in some publications this is referred to as a Simple Stel²⁶ lar Population). This view was supported by spectro-

Corresponding author: Emily M. Boudreaux emily.m.boudreaux.gr@dartmouth.edu, emily@boudreauxmail.com

37 scopically uniform heavy element abundances (Carretta 38 et al. 2010; Bastian & Lardo 2018) across most clus- $_{39}$ ters (M54 and ω Cen are notable exceptions, see Marino 40 et al. (2015) for further details), and the lack of ev-41 idence for multiple stellar populations (MPs) in past 42 color-magnitude diagrams of GCs (i.e. Sandage 1953; 43 Alcaino 1975). However, over the last 40 years non-44 trivial star-to-star light-element abundance variations 45 have been observed (i.e. Smith 1987) and, in the last 46 two decades, it has been definitively shown that most if 47 not all Milky Way GCs have MPs (Gratton et al. 2004, 48 2012; Piotto et al. 2015). The lack of photometric evi-49 dence for MPs prior to the 2000, can be attributed to the 50 more narrow color bands available, until very recently, to ₅₁ ground based photometric surveys (Milone et al. 2017). The prevalence of multiple populations in GCs is so 53 distinct that the proposed definitions for what consti-54 tutes a globular cluster now often center the existence 55 of MPs (e.g. Carretta et al. 2010). Whereas, people have 56 have often tried to categorized objects as GCs through 57 relations between half-light radius, density, and surface 58 brightness profile, in fact many objects which are generBoudreaux et al.

59 ally thought of as GCs don't cleanly fit into these cuts 60 (Peebles & Dicke 1968; Brown et al. 1991, 1995; Bekki 61 & Chiba 2002). Consequently, Carretta et al. (2010) 62 proposed a definition of GC based on observed chem-63 ical inhomogeneities in their stellar populations. The 64 modern understanding of GCs then is not simply one of 65 a dense cluster of stars that may have chemical inho-66 mogeneities and multiple populations; rather, it is one 67 where those chemical inhomogeneities and multiple pop-68 ulations themselves are the defining element of a GC.

All Milky Way globular clusters older than 2 Gyr studied in detail show populations enriched in He, N, and Na while also being deplete in O and C (Piotto et al. 2015; Bastian & Lardo 2018). These light element abundance patterns also are not strongly correlated with variations in heavy element abundance, resulting in spectroscopically uniform Fe abundances between populations. Further, high-resolution spectral studies reveal anti-correlations between N-C abundances, Na-O abundances, and potentially Al-Mg (Sneden et al. 1992; Gratton et al. 2012). Typical stellar fusion reactions can deplete core oxygen; however, the observed abundances of Na, Al, and Mg cannot be explained by the CNO cycle (Prantzos et al. 2007). Consequently, globular cluster populations must be formed by some novel means.

Formation channels for these multiple populations remain a point of debate among astronomers. Most proposed formation channels consist of some older, more
massive, population of stars polluting the pristine cluster media before a second population forms, now enriched in heavier elements which they themselves could
not have generated (for a detailed review see Gratton
et al. 2012). The four primary candidates for these polluters are asymptotic giant branch stars (AGBs, Ventura
et al. 2001; D'Ercole et al. 2010), fast rotating massive stars (FRMSs, Decressin et al. 2007), super massive stars (SMSs, Denissenkov & Hartwick 2014), and
massive interacting binaries (MIBs, de Mink et al. 2009;
Bastian & Lardo 2018).

Hot hydrogen burning (i.e. proton capture), material transport to the surface, and material ejection into the intra-cluster media are features of each of these models and consequently they can all be made to qualitatively agree with the observed elemental abundances. Howver, none of the standard models can currently account for all specific abundances (Gratton et al. 2012). AGB and FRMS models are the most promising; however, both models have difficulty reproducing severe O depletion (Ventura & D'Antona 2009; Decressin et al. 2007). Moreover, AGB and FRMS models require significant mass loss (~ 90%) between cluster formation and the current epoch — implying that a significant fraction of

111 halo stars formed in GCs (Renzini 2008; D'Ercole et al. 112 2008; Bastian & Lardo 2015).

In addition to the light-element anti-correlations ob-

114 served, it is also known that younger populations are significantly enhanced in Helium (Piotto et al. 2007, 2015; 116 Latour et al. 2019). Depending on the cluster, helium mass fractions as high as Y = 0.4 have been inferred (e.g. 118 Milone et al. 2015a). However, due to both the relatively 119 high and tight temperature range of partial ionization 120 for He and the efficiency of gravitational settling in core 121 helium burning stars, the initial He abundance of glob-122 ular cluster stars cannot be observed; consequently, the 123 evidence for enhanced He in GCs originates from com-124 parison of theoretical stellar isochrones to the observed 125 color-magnitude-diagrams of globular clusters. There-126 fore, a careful handling of chemistry is essential when 127 modeling with the aim of discriminating between MPs; 128 yet, only a very limited number of GCs have been stud-129 ied with chemically self-consistent (structure and atmo-130 sphere) isochrones (e.g. Dotter et al. 2015, NGC 6752). NGC 2808 is the prototype globular cluster to host 132 Multiple Populations. Various studies since 2007 have 133 identified that it may host anywhere from 2-5 stellar 134 populations. These populations have been identified both spectroscopically (i.e. Carretta et al. 2004; Carretta 136 2006; Carretta et al. 2010; Gratton et al. 2011; Carretta 137 2015; Hong et al. 2021) and photometrically (i.e. Piotto 138 et al. 2007, 2015; Milone et al. 2015a, 2017; Pasquato & 139 Milone 2019). Note that recent work (Valle et al. 2022) 140 calls into question the statistical significance of the de-141 tections of more than 2 populations in the spectroscopic 142 data. Here we present new, chemically self-consistent 143 modeling of the photometry of the two extreme populations of NGC 2808 identified by Milone et al. (2015a), 145 populations A and E. We use archival photometry from 146 the Hubble UV Globular Cluster Survey (HUGS) (Pi-147 otto et al. 2015; Milone et al. 2017) in the F275W and 148 F814W passbands to characterize multiple populations 149 in NGC 2808 (Milone et al. 2015a,b). Additionally, we 150 present a likelihood analysis of the photometric data 151 of NGC 2808 to determine the number of populations 152 present in the cluster.

2. CHEMICAL CONSISTENCY

153

There are three primary areas in which must the stel-155 lar models must be made chemically consistent: the at-156 mospheric boundary conditions, the opacities, and inte-157 rior abundances. The interior abundances are relatively 158 easily handled by adjusting parameters within our stel-159 lar evolutionary code. However, the other two areas 160 are more complicated to bring into consistency. Atmo-161 spheric boundary conditions and opacities must both be

260

calculated with a consistent set of chemical abundances outside of the stellar evolution code. For evolution we use the Dartmouth Stellar Evolution Program (DSEP) (Dotter et al. 2008), a well tested 1D stellar evolution code which has a particular focus on modelling low mass stars ($\leq 2~{\rm M}_{\odot}$)

2.1. Atmospheric Boundary Conditions

168

204

Certain assumptions, primarily that the radiation field 169 170 is at equilibrium and radiative transport is diffusive (Salaris & Cassisi 2005), made in stellar structure codes, 172 such as DSEP, are valid when the optical depth of a star 173 is large. However, in the atmospheres of stars, the num-174 ber density of particles drops low enough and the opti-175 cal depth consequently becomes small enough that these 176 assumptions break down, and separate, more physically motivated, plasma modeling code is required. Generally 178 structure code will use tabulated atmospheric boundary 179 conditions generated by these specialized codes, such 180 as ATLAS9 (Kurucz 1993), PHOENIX (Husser et al. 181 2013), MARCS (Gustafsson et al. 2008), and MPS-182 ATLAS (Kostogryz et al. 2023). Often, as the boundary 183 conditions are expensive to compute, they are not up-184 dated as interior abundances vary.

One key element when chemically consistently mod-186 eling NGC 2808 modeling is the incorporation of new 187 atmospheric models with the same elemental abun-188 dances as the structure code. We use atmospheres 189 generated from the MARCS grid of model atmospheres 190 (Plez 2008). MARCS provides one-dimensional, hydro-191 static, plane-parallel and spherical LTE atmospheric 192 models (Gustafsson et al. 2008). Model atmospheres 193 are made to match the spectroscopically measured elemental abundances of populations A and E. Moreover, 195 for each population, atmospheres with various helium mass fractions are generated. These range from Y=0.24 197 to Y=0.36 in steps of 0.03. All atmospheric models are 198 computed to an optical depth of $\tau = 100$ where their 199 temperature and pressures serves as boundary condi-200 tions for the structure code. In general, enhancing he-201 lium in the atmosphere has only a small impact on the atmospheric temperature profile, while leading to a drop 203 in the pressure by $\sim 10-20\%$.

2.2. Opacities

In addition to the atmospheric boundary conditions, both the high and low temperature opacities used by DSEP must be made chemically consistent. Here we use OPLIB high temperature opacity tables (Colgan et al. 2016) retrieved using the TOPS web-interface. Retrival of High temperature opacities is done using pyTOPSScrape, first introduced in Boudreaux &

212 Chaboyer (2023). Low temperature opacity tables are retrieved from the Aesopus 2.0 web-interface (Marigo & Aringer 2009; Marigo et al. 2022). Ideally, these opacities would be the same used in the atmospheric models. However, the opacities used in the MARCS models are not publicly available. As such, we use the opacities provided by the TOPS and Aesopus 2.0 web-interfaces.

3. STELLAR MODELS

We use the Dartmouth Stellar Evolution Program (DSEP, Dotter et al. 2008) to generate stellar mod222 els. DSEP is a well-tested, one-dimensional stellar evo223 lution code which includes a mixing length model of
224 convection, gravitational settling, and diffusion. Using
225 the solar composition presented in (Grevesse et al. 2007)
226 (GAS07), MARCS model atmosphers, OPLIB high tem227 perature opacities, and AESOPUS 2.0 low temperature
228 opacities we find a solar calibrated mixing length pa229 rameter, α_{MLT} , of $\alpha_{MLT} = 1.901$.

We use DSEP to evolve stellar models ranging in mass 231 from 0.3 to 2.0 solar masses from the fully convective 232 pre-main sequence to the tip of the red giant branch. 233 Below 0.7 M_{\odot} we evolve a model every 0.03 M_{\odot} and ₂₃₄ above 0.7 M_{\odot} we evolve a model every 0.05 M_{\odot} . We 235 evolve models over a grid of mixing length parameters 236 from $\alpha_{MLT}=1.0$ to $\alpha_{MLT}=2.0$ in steps of 0.1. For 237 each mixing length, a grid of models and isochrones were 238 calculated, with chemical compositions consistent with 239 Populations A and E (see Table 1) and a range of helium 240 abundances (Y=0.24, 0.27, 0.30, 0.33, 0.36, and 0.39). 241 In total, 144 sets of isochrones, each with a unique com-242 position and mixing length were calculated. Each model 243 is evolved in DSEP with typical numeric tolerences of 244 one part in 10⁷. Each model is allowed a maximum 245 time step of 50 Myr.

For each combination of population, Y, and α_{MLT} we use the isochrone generation code first presented in Dotter (2016) to generate a grid of isochrones. The isochrone generation code identified equivalent evolutionary points (EEPs) over a series of masses and interpolates between them. The grid of isochrones generated for this work is avalible as a digital supplement to this paper. Given the complexity of the parameter space when fitting multiple populations along with the recent warnings in the liteerature regarding overfitting datasets (e.g. Valle et al. 2022) we want to develop a more objective way of fitting isochrones to photometry than if we were to mark median ridge line positions by hand.

4. FIDANKA

265

266

267

268

269

270

271

283

Table 1. Population Composition

Element Pop A Pop E Element Pop A Pop E Li -0.08 — In -1.46 — Be 0.25 — Sn -0.22 — B 1.57 — Sb -1.25 — C 6.87 5.91 Te -0.08 — N 6.42 6.69 I -0.71 — O 7.87 6.91 Xe -0.02 — F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Al 5.21 6.67 Nd 0.29 — S 6.31 5.89 Sm -1.3 — C						
Be 0.25 — Sn -0.22 — B 1.57 — Sb -1.25 — C 6.87 5.91 Te -0.08 — N 6.42 6.69 I -0.71 — O 7.87 6.91 Xe -0.02 — F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar <t< td=""><td>Element</td><td>Pop A</td><td>Pop E</td><td>Element</td><td>Pop A</td><td>Pop E</td></t<>	Element	Pop A	Pop E	Element	Pop A	Pop E
B 1.57 — Sb -1.25 — C 6.87 5.91 Te -0.08 — N 6.42 6.69 I -0.71 — O 7.87 6.91 Xe -0.02 — F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K	Li	-0.08	_	In	-1.46	_
C 6.87 5.91 Te -0.08 — N 6.42 6.69 I -0.71 — O 7.87 6.91 Xe -0.02 — F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca	Be	0.25	_	Sn	-0.22	_
N 6.42 6.69 I -0.71 — O 7.87 6.91 Xe -0.02 — F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Al 5.21 6.61 Pr -1.54 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K	В	1.57	_	Sb	-1.25	
O 7.87 6.91 Xe -0.02 — F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 —	\mathbf{C}	6.87	5.91	Te	-0.08	
F 3.43 — Cs -1.18 — Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Cl -1.13 4.37 Eu -0.61 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 —	N	6.42	6.69	I	-0.71	
Ne 7.12 6.7 Ba 1.05 — Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.	O	7.87	6.91	Xe	-0.02	
Na 5.11 5.7 La -0.03 — Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3	\mathbf{F}	3.43	_	Cs	-1.18	
Mg 6.86 6.42 Ce 0.45 — Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — T	Ne	7.12	6.7	Ba	1.05	
Al 5.21 6.61 Pr -1.54 — Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W </td <td>Na</td> <td>5.11</td> <td>5.7</td> <td>La</td> <td>-0.03</td> <td>_</td>	Na	5.11	5.7	La	-0.03	_
Si 6.65 6.77 Nd 0.29 — P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re	$_{ m Mg}$	6.86	6.42	Се	0.45	_
P 4.28 — Pm -99.0 — S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re	Al	5.21	6.61	Pr	-1.54	
S 6.31 5.89 Sm -1.3 — Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os	Si	6.65	6.77	Nd	0.29	
Cl -1.13 4.37 Eu -0.61 — Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt	P	4.28	_	Pm	-99.0	
Ar 5.59 5.17 Gd -1.19 — K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Hf -1.41 — Co 3.86 — Th -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os	\mathbf{S}	6.31	5.89	Sm	-1.3	_
K 3.9 — Tb -1.96 — Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -	Cl	-1.13	4.37	Eu	-0.61	
Ca 5.21 — Dy -1.16 — Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg <td< td=""><td>Ar</td><td>5.59</td><td>5.17</td><td>Gd</td><td>-1.19</td><td></td></td<>	Ar	5.59	5.17	Gd	-1.19	
Sc 2.02 — Ho -1.78 — Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Zn 2.3 — Os -0.86 — Ba -2.0 — — Zn -0.88 — Ge 1.39 — Pt -0.64 — An -1.34 — Se 1.08 — Hg -1.09 — Bn -1.36 — Nt -0.51 — Nt -0.51 <	K	3.9	_	Tb	-1.96	_
Ti 3.82 — Er -1.34 — V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — TI -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi <t< td=""><td>Ca</td><td>5.21</td><td>_</td><td>Dy</td><td>-1.16</td><td>_</td></t<>	Ca	5.21	_	Dy	-1.16	_
V 2.8 — Tm -2.16 — Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — T1 -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi <td< td=""><td>Sc</td><td>2.02</td><td>_</td><td>Но</td><td>-1.78</td><td></td></td<>	Sc	2.02	_	Но	-1.78	
Cr 4.51 — Yb -1.42 — Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At <	Ti	3.82	_	Er	-1.34	_
Mn 4.3 — Lu -2.16 — Fe 6.37 — Hf -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At <	V	2.8	_	$_{ m Tm}$	-2.16	
Fe 6.37 — Hff -1.41 — Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Y 1.08 — At -99.0 — Nb -0.8 — Fr	Cr	4.51	_	Yb	-1.42	_
Co 3.86 — Ta -2.38 — Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Nb -0.8 — Fr -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra	Mn	4.3	_	Lu	-2.16	_
Ni 5.09 — W -1.41 — Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Y 1.08 — At -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Ru -0.51 — Th	Fe	6.37	_	Hf	-1.41	_
Cu 3.06 — Re -2.0 — Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Y 1.08 — At -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa	Co	3.86	_	Ta	-2.38	_
Zn 2.3 — Os -0.86 — Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Y 1.08 — At -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Ni	5.09	_	W	-1.41	_
Ga 0.78 — Ir -0.88 — Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Y 1.08 — At -99.0 — Nb -0.8 — Fr -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Cu	3.06	_	Re	-2.0	_
Ge 1.39 — Pt -0.64 — As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Zn	2.3	_	Os	-0.86	_
As 0.04 — Au -1.34 — Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Ga	0.78	_	Ir	-0.88	_
Se 1.08 — Hg -1.09 — Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Ge	1.39	_	Pt	-0.64	
Br 0.28 — Tl -1.36 — Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	As	0.04	_	Au	-1.34	_
Kr 0.99 — Pb -0.51 — Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Se	1.08	_	Hg	-1.09	
Rb 0.26 — Bi -1.61 — Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Br	0.28	_	Tl	-1.36	
Sr 0.61 — Po -99.0 — Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Kr	0.99	_	Pb	-0.51	_
Y 1.08 — At -99.0 — Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Rb	0.26	_	Bi	-1.61	
Zr 1.45 — Rn -99.0 — Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Sr	0.61	_	Po	-99.0	
Nb -0.8 — Fr -99.0 — Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Y	1.08	_	At	-99.0	_
Mo -0.38 — Ra -99.0 — Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Zr	1.45	_	Rn	-99.0	_
Tc -99.0 — Ac -99.0 — Ru -0.51 — Th -2.2 — Rh -1.35 — Pa -99.0 —	Nb	-0.8	_	Fr	-99.0	
Ru -0.51 — Th -2.2 — Pa -99.0 —	Mo	-0.38	_	Ra	-99.0	_
Rh -1.35 — Pa -99.0 —	Tc	-99.0	_	Ac	-99.0	_
	Ru	-0.51	_	Th	-2.2	_
Pd -0.69 — U -2.8 —	Rh	-1.35	_	Pa	-99.0	_
	Pd	-0.69	_	U	-2.8	

Note—Relative Metal composition used where a(H) = 12. Where the relative composition is the the same for both populations A and E it is only listed in the population A colum for the sake of visual clarity.

When fitting isochrones to the clusters with multiple populations we have four main criteria for any method

- The method must be robust enough to work along the entire main sequence, turn off, and much of the subgiant and red giant branch.
- Any method should consider photometric uncertainty in the fitting process.
- The method should be model independent, weighting any n number of populations equally.
- The method should be automated and require minimal intervention from the user.

We do not believe that any currently available software is a match for our use case. Therefore, we elect
to develop our own software suite, Fidanka. Fidanka
fi is a python package designed to automate much of the
process of measuring fiducial lines in CMDs, adhering to
the four criteria we lay out above. Primary features of
Fidanka may be separated into three categories: fiducial line measurement, stellar population synthesise, and
isochrone optimization/fitting. Additionally, there are
utility functions that are detailed in the Fidanka documentation.

4.1. Fiducial Line Measurement

Fidanka takes a iterative approach to measuring fidu-285 cial lines, the first step of which is to make a "guess" as to the fiducial line. This initial guess is calculated 287 by splitting the CMD into magnitude bins, with uni-288 form numbers of stars per bin (so that bins are cover a 289 small magnitude range over densely populated regions 290 of the CMD while covering a much larger magnitude ²⁹¹ range in sparsely populated regions of the CMD, such 292 as the RGB). A unimodal Gaussian distribution is then 293 fit to the color distribution of each bin, and the resulting ²⁹⁴ mean color is used as the initial fiducial line guess. This 295 rough fiducial line will approximately trace the area of 296 highest density. The initial guess will be used to verti-297 calze the CMD so that further algorithms can work in ²⁹⁸ 1-D magnitude bins without worrying about weighting 299 issues caused by varying projections of the evolutionary 300 sequence onto the magnitude axis. Verticalization is pre-301 formed taking the difference between the guess fiducial 302 line and the color of each star in the CMD.

If Fidanka were to simply apply the same algorithm to the verticalized CMD then the resulting fiducial line would likely be a re-extraction of the initial fiducial line guess. To avoid this, we take a more robust, number density based approach, which considers the distribution of stars in both color and magnitude space simultaneously. For each star in the CMD we first using

an introselect partitioning algorithm to select the 50 nearest stars in F814W vs. F275W-F814W space. To 312 account for the case where the star is at an extreme 313 edge of the CMD, those 50 stars include the star it- $_{314}$ self (such that we really select 49 stars + 1). We use 315 ghull¹(Barber et al. 1996) to calculate the convex hull 316 of those 50 points. The number density at each star 317 then is defined as $50/A_{hull}$, where A_{hull} is the area of the convex hull. Because we use a fixed number of points per star, and a partitioning algorithm as opposed to a sorting algorithm, this method scales like $\mathcal{O}(n)$, where is the number of stars in the CMD. This method also intrinsically weights the density of of each star equally 323 as the counting statistics per bin are uniform. We are 324 left with a CMD where each star has a defined number 325 density (Figure 1).

Fidanka can now exploit this density map to fit a

327 better fiducial line to the data, as the density map is far more robust to outliers. There are multiple algorithms we implement to fit the fiducial line to the color-density 330 profile in each magnitude bin (Figure 2); they are explained in more detail in the Fidanka documentation. 332 However, of most relevance here is the Bayesian Gaussian Mixture Modeling (BGMM) method. BGMM is a 334 clustering algorithm which, for some fixed number of n- $_{335}$ dimensional Gaussian distributions, K, determines the 336 mean, covariance, and mixing probability (somewhat 337 analogous to amplitude) of each k^{th} distribution, such 338 that the local lower bound of the likelyhood of each star belonging strongly to a single distribution is maximized. 339 Maximization is preformed using the Dirichlet pro-340 341 cess, which is a non-parametric Bayesian method of $_{342}$ determining the number of Gaussian distributions, K, which best fit the data (Ferguson 1973; Pedregosa et al. 2011). Use of the Dirichlet process allows for dynamic variation in the number of inferred populations from 346 magnitude bin to magnitude bin. Specifically, popula-347 tions are clearly visually separated from the lower main 348 sequence through the turn off; however, at the turn off 349 and throughout much of the subgiant branch, the two visible populations overlap due to their extremely simi-351 lar ages (i.e. Jordán et al. 2002). The Dirichlet process 352 allows for the BGMM method to infer a single popula-353 tion in these regions, while inferring two populations in 354 regions where they are clearly separated. More gener-355 ally, the use of the Dirichlet process removes the need 356 for a prior on the exact number of populations to fit.

Rather, the user specifies a upper bound on the num-

 $_{358}$ ber of populations within the cluster. An example bin $_{359}$ (F814W = 20.6) is shown in Figure 3.

Fidanka 's BGMM method first breaks down the ver-361 ticalized CMD into magnitude bins with uniform num-362 bers of stars per bin (here we adopt 250). Any stars 363 left over are placed into the final bin. For each bin a 364 BGMM model with a maximum of 5 populations is fit 365 to the color density profile. The number of populations 366 is then inferred from the weighting parameter (the mix-367 ing probability) of each population. If the weighting parameter of any k^{th} components less than 0.05, then that 369 component is considered to be spurious and removed. 370 Additionally, if the number of populations in the bin 371 above and the bin below are the same, then the num-372 ber of populations in the current bin is forced to be the 373 same as the number of populations in the bin above. Fi-374 nally, the initial guess fiducial line is added back to the 375 BGMM inferred line. Figure 4 shows the resulting fidu-376 cial line(s) in each magnitude bin for both a verticalized 377 CMD and a non verticalized CMD.

This method of fiducial line extraction effectively discriminated between multiple populations long the main sequence and RGB of a cluster, while simultaneously allowing for the presence of a single population along the MSTO and subgiant branch.

We can adapt this density map based BGMM method to consider photometric uncertainties by adopting a simple Monte Carlo approach. Instead of measuring the fiducial line(s) a single time, Fidanka can measure the fiducial line(s) many times, resampling the data with replacement each time. For each resampling Fidanka adds a random offset to each filter based on the photometric uncertainties of each star. From these n measurements the mean fiducial line for each sequence can be identified along with upper and lower bound confidence intervals in each magnitude bin.

4.2. Stellar Population Synthesis

While not extensively used in this paper Fidanka can, in addition to measuring fiducial lines, preform stellar population synthesise. Fidanka 's population synthesis module can generate synthetic stellar population from a set of MIST formatted isochrones. This is of primary importance for binary population modeling. The module is also used to generate synthetic CMDs for the purpose of testing the fiducial line extraction algorithms against priors.

Fidanka uses MIST formatted isochrones (Dotter 2016) as input along with distance modulus, B-V color excess, binary mass fraction, and bolometric corrections. An arbitrarily large number of isochrones may be used to define an arbitrary number of populations. Synthetic

¹ https://www.qhull.com

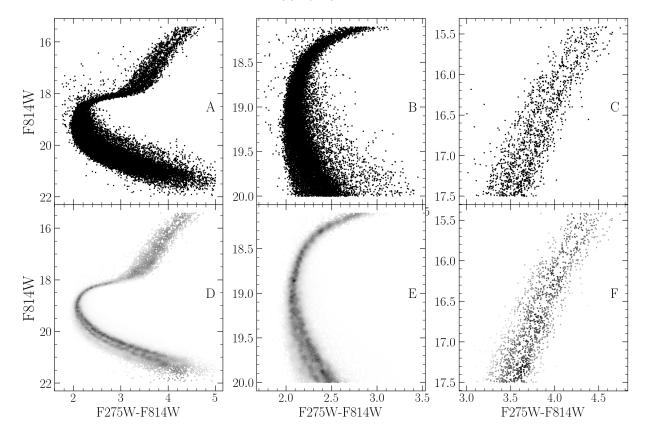


Figure 1. Density map demo showing density estimate over different parts of the evolutionary sequence. The left panel shows the density map over the entire evolutionary sequence, while the middle panel shows the density map over the main sequence and the right most panel shows the density map over the RGB. Figures in the top row are the raw CMD, while figures in the bottom row are colored by the density map.

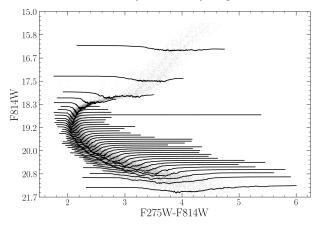


Figure 2. CMD where points are colored by density. Lines show the density-color profile in each magnitude bin. In this figure adaptive binning targeted 1000 stars per bin

 $_{409}$ stars are samples from each isochrone based on a definable probability (for example it is believed that $\sim 90\%$ $_{411}$ of stars in globular clusters are younger population (e.g. Suntzeff & Kraft 1996; Carretta 2013)). Based on the metallicity, $\mu,$ and E(B-V) of each isochrone, bolometric corrections are taken from bolometric correction tables.

Where bolometric correction tables do not include exact metallicities or extinctions a linear interpolation is preformed between the two bounding values.

4.3. Isochrone Optimization

The optimization routines in Fidanka will find the best fit distance modulus, B-V color excess, and binary number fraction for a given set of isochrones. If a sin- gle isochrone is provided then the optimization is done by minimizing the χ^2 of the perpendicular distances between an isochrone and a fiducial line. If multiple isochrones are provided then those isochrones are first used to run stellar population synthesis and generate a synthetic CMD. The optimization is then done by min- imizing the χ^2 of both the perpendicular distances between and widths of the observed fiducial line and the fiducial line of the synthetic CMD.

4.4. Fidanka Testing

In order to validate fidanka we have run an series of injection recovery tests using Fidanka's population synthesis routines to build various synthetic populations and Fidanka's fiducial measurement routines to recover

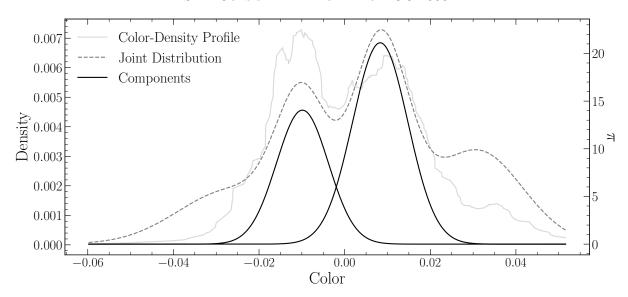


Figure 3. Example of BGMM fit to a magnitude bin. The grey line shows the underlying color-density profile, while the black dashed-line shows the joint distribution of each BGMM component. The solid black lines show the two selected components.

these populations. Each population was generated using the initial mass function given in (Milone et al. 2012) for the redmost population ($\alpha=-1.2$). Further, every population was given a binary population fraction of 10%, distance uniformly sampled between 5000pc and 15000pc, and a B-V color excess uniformly sampled between 0 and 0.1. Finally, each synthetic population was generated using a fixed age uniformly sampled between 7 Gyr and 14 Gyr. An example synthetic population along with its associated best fit isochrone are shown in Figure 5.

For each trial we use Fidanka to measure the fidu-448 cial line and then optimize that fiducial line against the 449 originating isochrone to esimate distance modulus, age, 450 and color B-V excess. Figure 6 is built from 1000 runs of 451 these trials and show the mean and width of the percent 452 error distributions for μ , E(B-V), and age. In general 453 Fidanka is able to recover distance modulii effectively 454 with age and E(B-V) reovery falling in line with other 455 literature that does not cosider the CMD outside of the 456 main sequence, main sequence turn off, sub giant, and 457 red giant branches; specifically, it should be noted that 458 Fidanka is not setup to model the horizontal branch.

5. ISOCHRONE FITTING

We fit pairs of isochrones to the HUGS data for NGC 2808 using Fidanka , as described in §4. Two isochrones, one for Population A and one for Population E are fit simultaneously. These isochrones are constrained to have distance modulus, μ , and color excess, E(B-V) which agree to within 0.5% and an ages which agree to within 1%. Moreover, we constrain the mixing length, α_{ML} , for any two isochrones in a set to be within 0.5 of one

⁴⁶⁸ and other. For every isochrone in the set of combina-⁴⁶⁹ tion of which fulfilling these constraints μ , E(B-V), ⁴⁷⁰ Age_A, and Age_B are optimized to reduce the χ^2 distance ⁴⁷¹ ($\chi^2 = \sum \sqrt{\Delta \text{color}^2 + \Delta \text{mag}^2}$) between the fiducial lines ⁴⁷² and the isochrones. Because we fit fiducial lines directly, ⁴⁷³ we do not need to consider the binary population frac-⁴⁷⁴ tion, f_{bin} , as a free parameter.

The best fit isochrones are shown in Figure 7 and optimized parameters for these are presented in Table 1. We find helium mass fractions that are consistent with those identified in past literature (e.g. Milone et al. 2015a). Note that our helium mass fraction grid has a spacing of 0.03 between grid points and we are therefore unable to resolve between certain proposed helium mass fractions for the younger sequence (for example between 0.37 and 0.39).

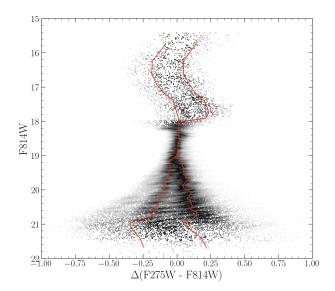
Past literature (e.g. Milone et al. 2015a, 2018) have found helium mass fraction variation from the low red-most to bluemost populations of ~ 0.12 . Here we find a helium mass fraction variation of 0.15 which, given the spacing of the helium grid we use is consistent with these past results.

5.1. The Number of Populartions in NGC 2808

In order to estimate the number of populations which ideally fit the NGC 2808 F275W-F814W photometry without overfitting the data we make use of silhouette analysis (Rousseeuw 1987, and in a similar manner to how Valle et al. (2022) preform their analysis of spectroscopic data). We find the average silhouette score for all tagged clusters identified using BGMM in all magnitude bins over the CMD using the standar python module sklearn. Figure 8 shows the silhouette analysis results

Population	Age	Distance Modulus	Extinction	Y	α_{ML}	χ^2_{ν}
	[Gyr]		[mag]			
A	$12.996^{+0.87}_{-0.64}$	15.021	0.54	0.24	2.050	0.021
\mathbf{E}	$13.061^{+0.86}_{-0.69}$	15.007	0.537	0.39	1.600	0.033

Table 1. Best fit parameters derived from fitting isochrones to the fiducual lines derived from the NCG 2808 photometry. The one sigma uncertainty reported on population age were determined from the 16th and 84th percentiles of the distribution of best fit isochrones ages.



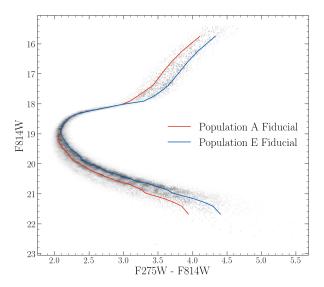


Figure 4. CMD where points are colored by density. Lines trace the infered fiducial line(s) in each magnitude bin.

500 and that two populations fit the photometry most ide-501 ally. This is in line with what our BGMM model predicts 502 for the majority of the the CMD.

5.2. ACS-HUGS Photometric Zero Point Offset

The Hubble legacy archive photometry used in this work is calibrated to the Vega magnitude system. How-

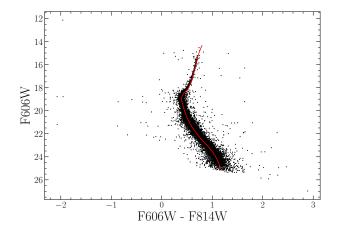


Figure 5. Synthetic population generated by fidanka at 10000pc with E(B-V) = 0, and an age of 12 Gyr along with the best fitting isochrone. The best fit paremeters are derived to be mu = 15.13, E(B-V) = 0.001, and an age of 12.33 Gyr.

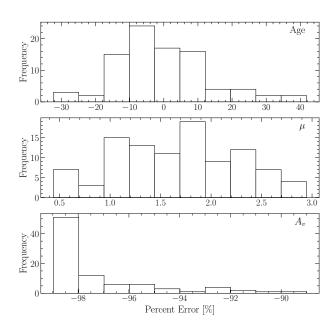


Figure 6. Percent Error distribution for each of the three deriver parameters. Note that these values will be sensitive to the magnitude uncertainties of the photometry. Here we made use of the ACS artificial star tests to estimate the uncertanties. Note that currently this is built with 100 runs, these take a long time so currently re running with 1000 runs.

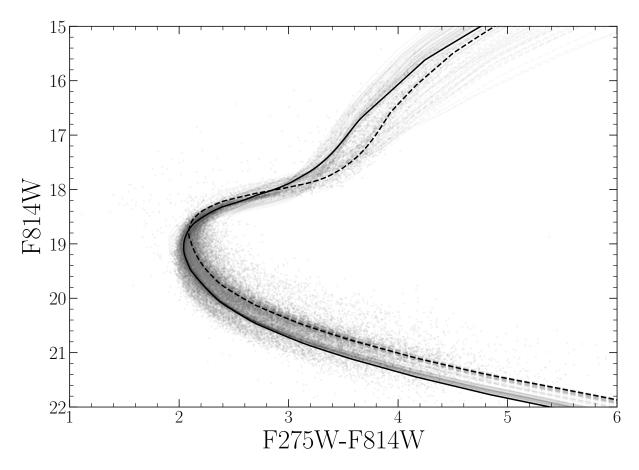


Figure 7. Best fit isochrone results for NGC 2808. The best fit population A and E models are shown as black lines. The following 50 best fit models are presented as grey lines. The solid black line is fit to population A, while the dashed black line is fit to population E.

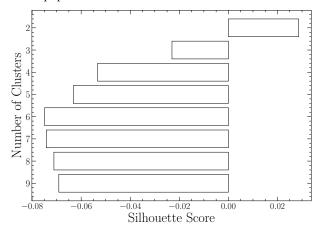


Figure 8. Silhouette analysis for NGC 2808 F275W-F814W photometry. The Silhouette scores are an average of score for each magnitude bin. Positive scores incidate that the clustering algorithm produced well distinguised clusters while negative scores indicate clusters which are not well distinguised.

506 ever, we have found that the photometry has a system-

sor atic offset of ~ 0.026 magnitudes in the F814W band when compared to the same stars in the ACS survey (Figure 9). The exact cause of this offset is unknown, but it is likely due to a difference in the photometric zero point between the two surveys. A full correction of this offset would require a careful re-reduction of the HUGS photometry, which is beyond the scope of this work. We instead recognize a 0.02 inherent uncertainty in the inferred magnitude of any fit when comparing to the ACS survey. This uncertainty is small when compared to the uncertainty in the distance modulus and should not affect the conclusion of this paper.

The oberved photometric offset between ACS and HUGS reductions introduces a systematic uncertainity when comparing parameters derived from isochrone fits to ACS data vs those fit to HUGS data. Specifically, this offset introduces a ~AGE Gyr uncertainity. Moreover, for two isochrone of the same age, only seperated by helium mass fraction, a shift of the main sequence turn off of is also expected. Figure 10 shows this shift.
Note a change in the helium mass fraction of a model by 0.03 results in an approximate 0.08 magnitude shift

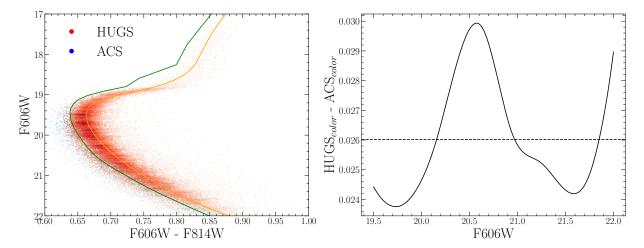


Figure 9. (left) CMD showing the photometric offset between the ACS and HUGS data for NGC 2808. CMDs have been randomly subsampled and colored by point density for clarity. (right) Mean difference between the color of the HUGS and ACS fiducual lines at the same magnitude. Note that the ACS data is systematically bluer than the HUGS data.

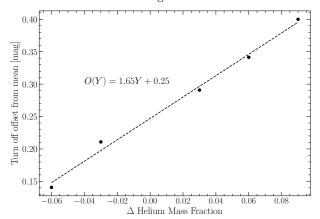


Figure 10. Main sequence turn off magnitude offset from a guage helium mass fraction (Y=0.30 chosen). All main sequence turn off locations are measured at 12.3 Gyr Should I make these contour surfaces for various ages?

529 to the main sequence turn off location. This means that 530 the mean 0.026 magnitude offset we find in between ACS 531 and HUGS data corresponds to an additional approaxi-532 mate 0.01 uncertainity in the derived helium mass frac-533 tion when comparing between these two datasets.

6. CONCLUSION

Here we have preformed the first chemically selfconsistnent modeling of the Milky Way Globular Cluster
NGC 2808. We find that, updated atmospheric boundary conditions and opacity tables do not have a significant effect on the inferred helium abundances of multiple
populations. Specifically, we find that the older stellar
population, populatio A, has a helium mass fraction of
20.24, while the older population, E, has a helium mass
fraction of 0.39.

Further, we introduce a new software suite for globular cluster science, Fidanka, which has been released under a permissive open source liscense. Fidanka aims to provide a statistically robust set of tools for estimating the parameters of multiple populations within globular clusters.

This work has made use of the NASA astrophysical data system (ADS). We would like to thank Elisabeth New-ton and Aaron Dotter for their support and for useful disscusion related to the topic of this paper. Addition-stally, we would like to thank Kara Fagerstrom, Aylin Garcia Soto, and Keighley Rockcliffe for their useful disscusion related to in this work. We acknowledge the strong support of a NASA grant (No. 80NSSC18K0634).

REFERENCES

Alcaino, G. 1975, A&AS, 21, 15
Barber, C. B., Dobkin, D. P., & Huhdanpaa, H. 1996, ACM
Transactions on Mathematical Software (TOMS), 22, 469
Bastian, N., & Lardo, C. 2015, MNRAS, 453, 357,
doi: 10.1093/mnras/stv1661

Bastian, N., & Lardo, C. 2018, Annual Review of
 Astronomy and Astrophysics, 56, 83
 Baumgardt, H., & Makino, J. 2003, MNRAS, 340, 227,
 doi: 10.1046/j.1365-8711.2003.06286.x

```
567 Bekki, K., & Chiba, M. 2002, The Astrophysical Journal,
     566, 245, doi: 10.1086/337984
568
569 Boudreaux, E. M., & Chaboyer, B. C. 2023, ApJ, 944, 129,
     doi: 10.3847/1538-4357/acb685
570
571 Boylan-Kolchin, M. 2018, MNRAS, 1423,
     doi: 10.1093/mnras/sty1490
572
573 Brodie, J. P., & Strader, J. 2006, Annu. Rev. Astron.
     Astrophys., 44, 193
574
575 Brown, J. H., Burkert, A., & Truran, J. W. 1991, ApJ, 376,
     115, doi: 10.1086/170260
576
   —. 1995, ApJ, 440, 666, doi: 10.1086/175304
577
578 Carretta, E. 2006, AJ, 131, 1766, doi: 10.1086/499565
    -. 2013, A&A, 557, A128,
     doi: 10.1051/0004-6361/201322103
    -. 2015, ApJ, 810, 148, doi: 10.1088/0004-637X/810/2/148
582 Carretta, E., Bragaglia, A., & Cacciari, C. 2004, ApJL,
     610, L25, doi: 10.1086/423034
583
584 Carretta, E., Bragaglia, A., Gratton, R. G., et al. 2010,
     Astronomy & Astrophysics, 516, A55
585
586 Colgan, J., Kilcrease, D. P., Magee, N. H., et al. 2016, in
     APS Meeting Abstracts, Vol. 2016, APS Division of
587
     Atomic, Molecular and Optical Physics Meeting
588
     Abstracts, D1.008
589
590 de Mink, S. E., Pols, O. R., Langer, N., & Izzard, R. G.
     2009, A&A, 507, L1, doi: 10.1051/0004-6361/200913205
591
592 Decressin, T., Meynet, G., Charbonnel, C., Prantzos, N., &
     Ekström, S. 2007, A&A, 464, 1029,
593
     doi: 10.1051/0004-6361:20066013
594
595 Denissenkov, P. A., & Hartwick, F. D. A. 2014, MNRAS,
     437, L21, doi: 10.1093/mnrasl/slt133
596
  D'Ercole, A., D'Antona, F., Ventura, P., Vesperini, E., &
597
     McMillan, S. L. W. 2010, MNRAS, 407, 854,
598
     doi: 10.1111/j.1365-2966.2010.16996.x
600 D'Ercole, A., Vesperini, E., D'Antona, F., McMillan, S.
     L. W., & Recchi, S. 2008, MNRAS, 391, 825,
601
     doi: 10.1111/j.1365-2966.2008.13915.x
602
603 Dotter, A. 2016, ApJS, 222, 8,
     doi: 10.3847/0067-0049/222/1/8
604
605 Dotter, A., Chaboyer, B., Jevremović, D., et al. 2008, The
     Astrophysical Journal Supplement Series, 178, 89
606
607 Dotter, A., Ferguson, J. W., Conroy, C., et al. 2015,
     MNRAS, 446, 1641, doi: 10.1093/mnras/stu2170
608
609 Ferguson, T. S. 1973, The annals of statistics, 209
610 Gratton, R., Sneden, C., & Carretta, E. 2004, ARA&A, 42,
     385, doi: 10.1146/annurev.astro.42.053102.133945
611
612 Gratton, R. G., Carretta, E., & Bragaglia, A. 2012,
```

Astronomy and Astrophysics Reviews, 20, 50,

Gratton, R. G., Lucatello, S., Carretta, E., et al. 2011,

A&A, 534, A123, doi: 10.1051/0004-6361/201117690

doi: 10.1007/s00159-012-0050-3

613

614

615

```
617 Grevesse, N., Asplund, M., & Sauval, A. J. 2007, SSRv,
     130, 105, doi: 10.1007/s11214-007-9173-7
618
619 Gustafsson, B., Edvardsson, B., Eriksson, K., et al. 2008,
     A&A, 486, 951, doi: 10.1051/0004-6361:200809724
620
621 Hong, S., Lim, D., Chung, C., et al. 2021, AJ, 162, 130,
     doi: 10.3847/1538-3881/ac0ce6
622
623 Hudson, M. J., & Robison, B. 2018, Monthly Notices of the
     Royal Astronomical Society, 477, 3869,
624
     doi: 10.1093/mnras/sty844
625
626 Husser, T. O., Wende-von Berg, S., Dreizler, S., et al. 2013,
     A&A, 553, A6, doi: 10.1051/0004-6361/201219058
628 Jordán, A., Côté, P., West, M. J., & Marzke, R. O. 2002,
     ApJL, 576, L113, doi: 10.1086/343759
630 Kostogryz, N., Shapiro, A. I., Witzke, V., et al. 2023,
     Research Notes of the AAS, 7, 39,
631
     doi: 10.3847/2515-5172/acc180
632
633 Kravtsov, A. V., & Gnedin, O. Y. 2005, The Astrophysical
     Journal, 623, 650
634
635 Kurucz, R.-L. 1993, Kurucz CD-Rom, 13
636 Latour, M., Husser, T. O., Giesers, B., et al. 2019, A&A,
     631, A14, doi: 10.1051/0004-6361/201936242
638 Marigo, P., & Aringer, B. 2009, A&A, 508, 1539,
     doi: 10.1051/0004-6361/200912598
639
640 Marigo, P., Aringer, B., Girardi, L., & Bressan, A. 2022,
     ApJ, 940, 129, doi: 10.3847/1538-4357/ac9b40
641
642 Marino, A. F., Milone, A. P., Karakas, A. I., et al. 2015,
     Monthly Notices of the Royal Astronomical Society, 450,
643
     815, doi: 10.1093/mnras/stv420
644
645 Milone, A. P., Piotto, G., Bedin, L. R., et al. 2012, ApJ,
     744, 58, doi: 10.1088/0004-637X/744/1/58
647 Milone, A. P., Marino, A. F., Piotto, G., et al. 2015a, ApJ,
     808, 51, doi: 10.1088/0004-637X/808/1/51
649 —. 2015b, MNRAS, 447, 927, doi: 10.1093/mnras/stu2446
650 Milone, A. P., Piotto, G., Renzini, A., et al. 2017, MNRAS,
     464, 3636, doi: 10.1093/mnras/stw2531
652 Milone, A. P., Marino, A. F., Renzini, A., et al. 2018,
     MNRAS, 481, 5098, doi: 10.1093/mnras/sty2573
653
654 Pasquato, M., & Milone, A. 2019, arXiv e-prints,
     arXiv:1906.04983, doi: 10.48550/arXiv.1906.04983
655
656 Pedregosa, F., Varoquaux, G., Gramfort, A., et al. 2011,
     Journal of Machine Learning Research, 12, 2825
657
658 Peebles, P. J. E., & Dicke, R. H. 1968, ApJ, 154, 891,
     doi: 10.1086/149811
659
660 Peng, E. W., Ferguson, H. C., Goudfrooij, P., et al. 2011,
     The Astrophysical Journal, 730, 23
661
662 Piotto, G., Bedin, L. R., Anderson, J., et al. 2007, The
     Astrophysical Journal Letters, 661, L53,
663
     doi: 10.1086/518503
664
  Piotto, G., Milone, A. P., Bedin, L. R., et al. 2015, AJ, 149,
     91, doi: 10.1088/0004-6256/149/3/91
```

12 BOUDREAUX ET AL.

```
667 Plez, B. 2008, Physica Scripta Volume T, 133, 014003,
                                                                   681 Sandage, A. R. 1953, AJ, 58, 61, doi: 10.1086/106822
     doi: 10.1088/0031-8949/2008/T133/014003
669 Prantzos, N., Charbonnel, C., & Iliadis, C. 2007, A&A, 470,
     179, doi: 10.1051/0004-6361:20077205
671 Renzini, A. 2008, Monthly Notices of the Royal
                                                                   686
     Astronomical Society, 391, 354,
672
                                                                        doi: 10.1086/117930
     doi: 10.1111/j.1365-2966.2008.13892.x
673
674 Richer, H. B., Fahlman, G. G., Buonanno, R., et al. 1991,
     ApJ, 381, 147, doi: 10.1086/170637
                                                                   690
676 Rousseeuw, P. J. 1987, Journal of Computational and
                                                                   691
     Applied Mathematics, 20, 53,
                                                                   692
     doi: https://doi.org/10.1016/0377-0427(87)90125-7
                                                                   693
678
                                                                      Ventura, P., D'Antona, F., Mazzitelli, I., & Gratton, R.
679 Salaris, M., & Cassisi, S. 2005, Evolution of stars and
                                                                   694
                                                                        2001, ApJL, 550, L65, doi: 10.1086/319496
     stellar populations (John Wiley & Sons)
```

```
682 Smith, G. H. 1987, Publications of the Astronomical
    Society of the Pacific, 99, 67, doi: 10.1086/131958
684 Sneden, C., Kraft, R. P., Prosser, C. F., & Langer, G. 1992,
    The Astronomical Journal, 104, 2121
  Suntzeff, N. B., & Kraft, R. P. 1996, AJ, 111, 1913,
  Valle, G., Dell'Omodarme, M., & Tognelli, E. 2022, A&A,
    658, A141, doi: 10.1051/0004-6361/202142454
  van den Bergh, S. 2010, The Astronomical Journal, 140,
     1043, doi: 10.1088/0004-6256/140/4/1043
  Ventura, P., & D'Antona, F. 2009, A&A, 499, 835,
    doi: 10.1051/0004-6361/200811139
```