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Chemically Self-Consistant Modeling of the Globular Cluster NGC 2808 and its Effects on the Inferred Helium abundance of Multiple Stellar Populations.

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ABSTRACT

Globular Clusters (GCs) provide a unique astrophysical laboratory for studying the formation and evolution of stars. GCs are old, dense, and it has historically been believed that they have a single stellar population (SSP, in some publications also simple stellar population). However, in the last two decades, it has been definitively shown that most if not all Milky Way GCs have multiple stellar populations (MPs). These MPs are chemically distinct from one another, primarily separated by light element abundance variations without the standard accompanying heavy element abundance variations. In addition to the light element abundance variations, MPs also have different inferred helium abundances. The origin of these MPs remains an open question, and one which is sensitive to the population - population compositional differences; consequently, the extent of the composition variations is a key parameter to constrain. Metal abundances may be directly measured spectroscopically; however, helium abundances are not directly observable in GCs. Instead, helium abundances are inferred from stellar models. It is therefore important to understand build stellar models that are self-consistent in the compositions of the structure, atmosphere, and opacity. In this work we present the first chemically self-consistent stellar models of the Milky Way Globular Cluster NGC 2808. We find that the helium abundance of the second generation of stars is higher than the first generation by SOME AMOUNT. This is in agreement with previous studies of NGC 2808.

Keywords: Globular Clusters (656), Stellar evolutionary models (2046)

1. INTRODUCTION

Globular clusters (GCs) are among the oldest observable able objects in the universe (Peng et al. 2011). They are characterized by high densities with typical half-light radii of ≤ 10 pc (van den Bergh 2010), and typical masses ranging from 10^4 – 10^5 M $_{\odot}$ (Brodie & Strader 2006) — though some GCs are significantly larger than these typical values EXAMPLE. GCs provide a unique way to probe stellar evolution (Baumgardt & Makino 2003), galaxy formation models (Boylan-Kolchin 2018; Kravtsov & Gnedin 2005), and dark matter halo structure (Hudson & Robison 2018). BRING IN SOME MORE RECENT CITATIONS.

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The traditional view of Globular Clusters was, for a 37 long time, that they consisted of a single stellar popu-38 lation (SSP, in some publications this is referred to as 39 a Simple Stellar Population). This view was supported 40 by spectroscopically uniform heavy element abundances 41 (?) [CHECK], and the lack of evidence for multiple stel-42 lar populations (MPs) in the color-magnitude diagrams 43 of GCs (?). However, in the last two decades, it has 44 been definitively shown that most if not all Milky Way 45 GCs have MPs (Piotto et al. 2015). The lack of photo-46 metric evidence for MPs can be attributed to the short 47 color throw available to ground based photometric sur-48 veys (?). While MPs are chemically distinct from one 49 another, that distinction is most prominent when ob-50 serving with U and B filters (Sbordone et al. 2011). The prevalence of multiple populations in GCs is so 52 distinct that the proposed definitions for what consti-53 tutes a globular cluster now often center the existence 54 of MPs. Whereas, people have have often tried to cate-55 gorized objects as GCs through relations between halfBoudreaux et al.

56 light radius, density, and surface brightness profile, in 57 fact many objects which are generally thought of as 58 GCs don't cleanly fit into these cuts EXAMPLE + 59 CITATION. Consequently, Carretta et al. (2010) pro- 60 posed a definition of GC based on observed chemical inhomogeneities in their stellar populations. The mod- 62 ern understanding of GCs then is not simply one of a 63 dense cluster of stars which may have chemical inhomogeneities and multiple populations; rather, it is one 64 where those chemical inhomogeneities and multiple populations themselves are the defining element of a GC.

All Milky Way globular clusters older than 2 Gyr studied in detail show populations enriched in He, N, and
Na while also being deplete in O and C (Piotto et al.
2015; Bastian & Lardo 2018). These light element abundance patterns also are not strongly correlated with variations in heavy element abundance, resulting in spectroscopically uniform Fe abundances between populations. Further, high-resolution spectral studies reveal
anti-correlations between N-C abundances, Na-O abundances, and potentially Al-Mg (Sneden et al. 1992; Gratton et al. 2012). Typical stellar fusion reactions can deplete core oxygen; however, the observed abundances of
Na, Al, and Mg cannot be explained by the likes of the
CNO cycle (Prantzos et al. 2007).

Formation channels for these multiple populations remain a point of debate among astronomers. Most proposed formation channels consist of some older, more massive, population of stars polluting the pristine cluster media before a second population forms, now enriched in heavier elements which they themselves could not have generated (for a detailed review see Gratton et al. 2012). The four primary candidates for these polluters are asymptotic giant branch stars (AGBs, Ventura et al. 2001; D'Ercole et al. 2010), fast rotating massive stars (FRMSs, Decressin et al. 2007), super massive stars (SMSs, Denissenkov & Hartwick 2014), and massive interacting binaries (MIBs, de Mink et al. 2009; Bastian & Lardo 2018).

Hot hydrogen burning (proton capture), material fransport to the surface, and material ejection into the intra-cluster media are features of each of these models and consequently they can all be made to qualitatively agree with the observed elemental abundances. Howtop ever, none of the standard models can currently account for all specific abundances (Gratton et al. 2012). AGB and FRMS models are the most promising; however, both models have difficulty reproducing severe O depletion (Ventura & D'Antona 2009; Decressin et al. 2007). Moreover, AGB and FRMS models require significant mass loss (~ 90%) between cluster formation and the current epoch — implying that a significant fraction of

halo stars formed in GCs (Renzini 2008; D'Ercole et al. 2008; Bastian & Lardo 2015).

In addition to the light-element anti-correlations ob-111 served it is also known that younger populations are sig-112 nificantly enhanced in Helium (Piotto et al. 2007, 2015; 113 Latour et al. 2019). Depending on the cluster, Helium mass fractions as high as Y = 0.4 have been inferred 115 (e.g Milone et al. 2015). However, due to the relatively 116 high and tight temperature range of partial ionization 117 for He it cannot be observed in globular clusters; conse-118 quently, the evidence for enhanced He in GCs originates 119 from comparison of theoretical stellar isochrones to the 120 observed color-magnitude-diagrams of globular clusters. 121 Therefore, a careful handling of chemistry is essential 122 when modeling with the aim of discriminating between 123 MPs; yet, only a very limited number of GCs have yet 124 been studied with chemically self-consistent (structure and atmosphere) isochrones (e.g. Dotter et al. 2015, 126 NGC 6752).

NGC 2808 is the prototype globular cluster to host Multiple Populations. Various studies since 2007 have identified that it may host anywhere from 2-5 stellar populations. These populations have been identified both spectroscopically (i.e.) and photometrically (i.e. 132). Note that recent work (Valle et al. 2022) calls into question the statistical significance of the detections of more than 2 populations in the spectroscopic data. Here we present new, chemically self-consistent modeling of the two extreme populations of NGC 2808 identified by Milone et al. (2015), populations A and E. Additionally, we present a likelihood analysis of the photometric data of NGC 2808 to determine the number of populations present in the cluster.

2. CHEMICAL CONSISTENCY

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There are three primary areas in which must the stellar models must be consistent: the atmospheric boundlar ary conditions, the opacities, and interior abundances.
The interior abundances are relatively easily handled
lar by adjusting parameters within our stellar evolutionary
lar code. However, the other two areas are more complilar cated to bring into consistency. Atmospheric boundary
conditions and opacities must both be calculated with a
consistent set of chemical abundances outside of DSEP.
The

2.1. Atmospheric Boundary Conditions

153 Certain assumptions [WHICH ONES, CITATION]
154 made in stellar structure codes, such as DSEP, are valid
155 when the optical depth of a star is small. However, in
156 the atmospheres of stars, the number density of par157 ticles drops low enough and the optical depth conse158 quently becomes large enough that these assumptions

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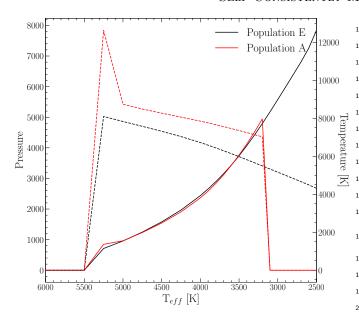


Figure 1. Comparison of the MARCS model atmospheres generated for the two extreme populations of NGC 2808. These lines shows population A and E with the same Helium abundance; though, we fit a grid of models over various helumn abundances. Dashed lines show the temperature of the boundary condition while sold lines show the pressure.

break down, and more separate plasma modeling code is required. Generally structure code will use tabulated atmospheric boundary conditions generated by these specialized codes PHEONIX, ATLAS, MARCS + [CITA-TIONS]. Often, as the boundary conditions are both expensive to compute and not the speciality of stellar structure researchers, the boundary conditions are not updated as as light-element interior abundance varies.

One key element when chemically consistently mod-167 168 eling NGC 2808 modeling is the incorporation of new 169 atmospheric models with the same elemental abun-170 dances as the structure code. We use atmospheres generated from the MARCS grid of model atmospheres 172 (Plez 2008). MARCS provides one-dimensional, hydro-173 static, plane-parallel and spherical LTE atmospheric models (Gustafsson et al. 2008). Model atmospheres 175 are made to match the spectroscopically measured ele-176 mental abundances of populations A and E. Moreover, 177 for each populations, atmospheres with various helium mass fractions are generated. These range from Y=0.24 179 to Y=0.36 in steps of 0.02. A comparison of the pres- $_{\mbox{\scriptsize 180}}$ sure and temperature throughout the atmospheres of the 181 two populations with helium abundances representative 182 of literature values is shown in Figure 2.1.

In addition to the atmospheric boundary conditions, both the high and low temperature opacities used by DSEP must be made chemically consistent. Here we use OPLIB high temperature opacity tables (Colgan et al. 2016) retrieved using the TOPS web-interface. Low temperature opacity tables are retrieved from the Aesopus 2.0 web-interface (Marigo & Aringer 2009; Marigo et al. 2022). Ideally, these opacities would be the same used in the atmospheric models. However, the opacities used in the MARCS models are not publicly available. As such, we use the opacities provided by the TOPS and Aesopus 2.0 web-interfaces.

3. STELLAR MODELS

We use the Dartmouth Stellar Evolution Program (DSEP, Dotter et al. 2008) to generate stellar models.
DSEP is a well-tested, one-dimensional stellar evolution code which includes a mixing length model of convection, gravitational settling, and diffusion.

We use DSEP to evolve stellar models ranging in mass from 0.3 to 2.0 solar masses from the zero-age main sequence (ZAMS) to the tip of the red giant branch. Below 0.7 M_{\odot} we evolve a model every 0.03 M_{\odot} and above 0.7 M_{\odot} we evolve a model every 0.5 M_{\odot} . Additionally, we evolve models over a grid of mixing length parameters, α_{MLT} , from $\alpha_{MLT}=1.0$ to $\alpha_{MLT}=2.0$ in steps of 0.1. In addition to the mixing length grid the evolved grid of models also has dimensions population (A or E) and helium abundance (Y). Each model is evolved in DSEP's "high resolution" mode and had a maximum allowed time step of 50 Myr.

For each combination of population, Y, and α_{MLT} we use the isochrone generation code first presented in ? to generate a grid of isochrones. The isochrone generation code identified equivalent evolutionary points (EEPs) over a series of masses and interpolates between them. The grid of isochrones generated for this work is avalible as a digital supplement to this paper.

4. FIDANKA

When fitting isochrones to the data we have four main criteria for any method

- The method must be robust enough to work along the entire main sequence, turn off, and much of the subgiant and red giant branchs.
- Any method should consider photometric uncertainty in the fitting process.
- The method should be model independent, weighting any n number of populations equally.
- The method should be automated and require minimal intervention from the user.

We do not believe that any currently available software is a match for our use case. Therefore, we elect
to develop our own software suite, Fidanka. Fidanka
is a python package designed to automate much of the
process of measuring fiducial lines in CMDs, adhering
to the four criteria we lay out above. Primary features
of Fidanka may be separated into three primary categories: fiducial line measurement, stellar population
synthesise, and isochrone optimization/fitting. Additionally, there are utility functions which are detailed
in the Fidanka documentation.

4.1. Fiducial Line Measurement

Fidanka takes a iterative approach to measuring fidu-245 246 cial lines, the first step of which is to make a "guess" 247 as to the fiducial line. This initial guess is calculated 248 by splitting the CMD into magnitude bins, with uni-249 form numbers of stars per bin (so that bins are cover a 250 small magnitude range over densely populated regions of the CMD while covering a much larger magnitude 252 range in sparsely populated regions of the CMD, such 253 as the RGB). A unimodal Gaussian distribution is then 254 fit to the color distribution of each bin, and the resulting ²⁵⁵ mean color is used as the initial fiducial line guess. This 256 rough fiducial line will approximately trace the area of 257 highest density. The initial guess will be used to verti-258 calze the CMD so that further algorithms can work in 259 1-D magnitude bins without worrying about weighting ²⁶⁰ issues caused by varying projections of the evolutionary 261 sequence onto the magnitude axis. Verticalization is preformed taking the difference between the guess fiducial line and the color of each star in the CMD.

If Fidanka were to simply apply the same algorithm 265 to the verticalized CMD then the resulting fiducial line would likely be a re-extraction of the initial fiducial line guess. To avoid this, we take a more robust, number 268 density based approach, which considers the distribu-269 tion of stars in both color and magnitude space simultaneously. For each star in the CMD we first using a 271 introselect partitioning algorithm to select the 50th 272 nearest stars. To account for the case where the star is at an extreme edge of the CMD, those 50 stars include the star itself (such that we really select 49 stars + 1). We use qhull¹(Barber et al. 1996; ?) to calculate the 276 convex hull of those 50 points. The number density at each star then is defined as $50/A_{hull}$, where A_{hull} is the 278 area of the convex hull. Because we use a fixed num-279 ber of points per star, and a partitioning algorithm as 280 opposed to a sorting algorithm, this method scales like

 $\mathcal{O}(n)$, where n is the number of stars in the CMD. This method also intrinsically weights the density of of each star equally as the counting statistics per bin are uniform. We are left with a CMD where each star has a defined number density (Figure 4.1).

Fidanka can now exploit this density map to fit a better fiducial line to the data, as the density map is far more robust to outliers. There are multiple algorithms we implement to fit the fiducial line to the color-density profile in each magnitude bin (Figure 4.1); they are explained in more detail in the Fidanka documentation. However, of most relevance here is the Bayesian Gaussian Mixture Modeling (BGMM) method. BGMM is a clustering algorithm which, for some fixed number of number of number of magnitude distributions, K, determines the mean, covariance, and mixing probability (somewhat analogous to amplitude) of each k^{th} distribution, such that the local lower bound of the evidence of each star belonging strongly to a single distribution is maximized.

Maximization is preformed using the Dirichlet pro-301 cess, which is a non-parametric Bayesian method of $_{302}$ determining the number of Gaussian distributions, K, 303 which best fit the data [CITATION]. Use of the Dirich-304 let process allows for dynamic variation in the number of 305 inferred populations from magnitude bin to magnitude 306 bin. Specifically, populations are clearly visually sepa-307 rated from the lower main sequence through the turn 308 off; however, at the turn off and throughout much of 309 the subgiant branch, the two visible populations over-310 lap due to their extremely similar ages (Citation). The 311 Dirichlet process allows for the BGMM method to infer 312 a single population in these regions, while inferring two 313 populations in regions where they are clearly separated. 314 More generally, the use of the Dirichlet process removes 315 the need for a prior on the exact number of populations 316 to fit. Rather, the user specifies a upper bound on the 317 number of populations within the cluster. An example $_{318}$ bin (F814W = 20.6) is shown in Figure 4.

Fidanka 's BGMM method first breaks down the verticalized CMD into magnitude bins with uniform numbers of stars per bin (here we adopt 250). Any stars
left over are placed into the final bin. For each bin a
BGMM model with a maximum of 5 populations is fit
to the color density profile. The number of populations
is then inferred from the weighting parameter (the mixing probability) of each population. If the weighting parameter of any k^{th} components less than 0.05, then that
component is considered to be spurious and removed.
Additionally, if the number of populations in the bin
above and the bin below are the same, then the number of populations in the current bin is forced to be the

¹ https://www.qhull.com

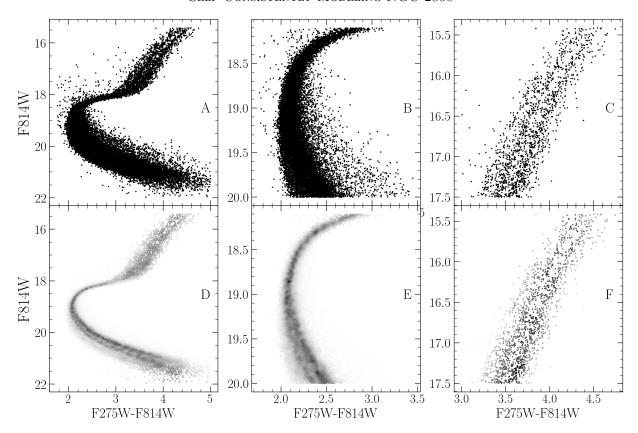


Figure 2. Density map demo showing density estimate over different parts of the evolutionary sequence. The left panel shows the density map over the entire evolutionary sequence, while the middle panel shows the density map over the main sequence and the right most panel shows the density map over the RGB. Figures in the top row are the raw CMD, while figures in the bottom row are colored by the density map.

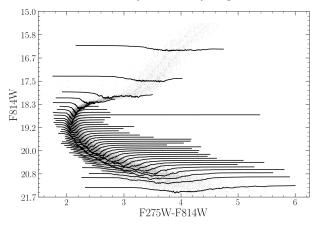


Figure 3. CMD where points are colored by density. Lines show the density-color profile in each magnitude bin. In this figure adaptive binning targeted 1000 stars per bin

333 nally, the initial guess fiducial line is added back to the BGMM inferred line. Figure 5 shows the resulting fidu-335 cial line(s) in each magnitude bin for both a verticalized CMD and a non verticalized CMD.

This method of fiducial line extraction effectively discriminated between multiple populations long the main 339 sequence and RGB of a cluster, while simultaneously al-340 lowing for the presence of a single population along the 341 MSTO and subgiant branch.

We can adapt this density map based BGMM method to consider photometric uncertainties by adopting a simple Monte Carlo approach. Instead of measuring the fiducial line(s) a single time, Fidanka can measure the fiducial line(s) many times, resampling the data with replacement each time. For each resampling Fidanka adds a random offset to each filter based on the photometric uncertainties of each star. From these n measurments the mean fiducial line for each sequence can be identified along with upper and lower bound confidence intervals in each magnitude bin.

4.2. Stellar Population Synthesis

In addition to measuring fiducial lines, Fidanka also includes a stellar population synthesise module. This module is used to generate synthetic CMDs from a given set of isochrones. This is of primary importance for binary population modelling. The module is also used to generate synthetic CMDs for the purpose of testing the fiducial line extraction algorithms against priors.

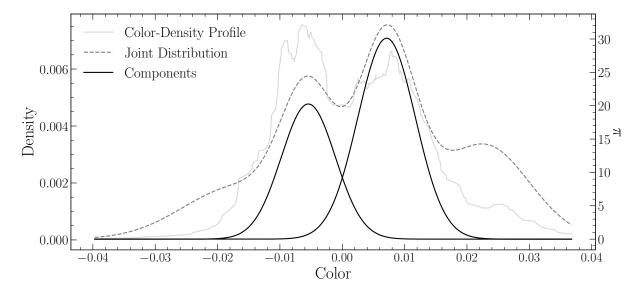


Figure 4. Example of BGMM fit to a magnitude bin. The grey line shows the underlying color-density profile, while the black dashed-line shows the joint distribution of each BGMM component. The solid black lines show the two selected components.

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Fidanka uses MIST formatted isochrones [CITA-TION as input along with distance modulus, B-V color excess, binary mass fraction, and bolometric corrections. An arbitrarily large number of isochrones may be used define an arbitrary number of populations. thetic stars are samples from each isochrone based on definable probability (for example it is believed that $\sim 90\%$ of stars in globular clusters are younger popu-₃₆₉ lation [CITATION]). Based on the metallicity, μ , and 370 E(B-V) of each isochrone, bolometric corrections are taken from bolometric correction tables. Where bolo-372 metric correction tables do not include exact metallic-373 ities or extinctions a linear interpolation is preformed between the two bounding values. [FIGURE] shows an 375 example of a synthetic CMD generated from a set of 2 376 NGC 2808 isochrones as well as a comparison between 377 those isochrones and the measured fiducial line of the 378 synthetic population.

4.3. Isochrone Optimization

The optimization routines in Fidanka will find the best fit distance modulus, B-V color excess, and binary number fraction for a given set of isochrones. If a single isochrone is provided then the optimization is done by minimizing the χ^2 of the perpendicular distances between an isochrone and a fiducial line. If multiple isochrones are provided then those isochrones are first used to run stellar population synthesis and generate a synthetic CMD. The optimization is then done by minimizing the χ^2 of both the perpendicular distances between and widths of the observed fiducial line and the fiducial line of the synthetic CMD.

4.4. Fidanka Testing

In order to validate fidanka we have run a series of tests on synthetic CMDs generated from NGC 2808 isochrones. For each test we use fidanka to both build a synthetic population and then measure the fiducial line of that population. The measured fiducial lines(s) are then compared to the underlying isochrones used.

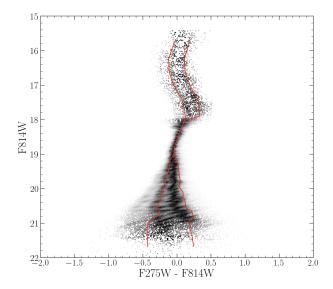
5. ISOCHRONE FITTING

We fit pairs of isochrones to the HUGS data for NGC 2808 using Fidanka, as described in §FIDANKA SEC-TION. Two isochrones, one for Population A and one for 403 Population E are fit simultaneously. These isochrones 404 are constrained to have the same distance modulus, μ , 405 and color excess, E(B-V). Moreover, we constrain the 406 mixing length, α_{ML} , to be constant between any two 407 isochrones in a set. For every isochrone in the set of combination of which fulfilling these constraints μ , E(B-V), Age_A , and Age_B are optimized to reduce the χ^2 distance 410 between the fiducial lines and the isochrones. Because 411 we fit fiducial lines directly, we do not need to consider the binary population fraction, f_{bin} , as a free parameter. Table ?? shows the results of the isochrone fitting for 414 NGC 2808. The best fit isochrones are shown in Figure 415 ??.

Disscus the implications of the specific isochrones we 417 fit.

5.1. ACS-HUGS Photometric Zero Point Offset

The Hubble legacy archive photometry used in this work is calibrated to the Vega magnitude system. However, we have found that the photometry has a system- atic offset of ~ 0.02 magnitudes in the F814W band



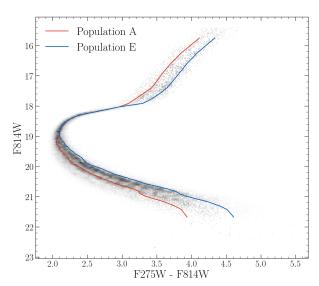


Figure 5. CMD where points are colored by density. Line trace the infered fiducial line(s) in each magnitude bin.

when compared to the same stars in the ACS survey [CITATION]. The exact cause of this offset is unknown, but it is likely due to a difference in the photometric zero point between the two surveys. A full correction of this offset would require a careful re-reduction of the HUGS photometry, which is beyond the scope of this work. We instead recognize a 0.02 inherent uncertainty in the inferred magnitude of any fit when comparing to

431 the ACS survey. This uncertainty is small when com-432 pared to the uncertainty in the distance modulus and 433 should not affect the conclusion of this paper.

Figure Showing the offset

The independent analysis of the Hubble Space Telescope photometry ACS and HUGS, and an additional
ground-based dataset created by Stetson et al. (2009),
of the globular cluster NGC2808 resulted in findings of
a systematic difference in the magnitude of the stars. A
significant offset between the three datasets was found.
HUGS and ACS were found to have consistent data for
Vvega, however an offset of approximately 0.025 was
found in the Ivega. ACS and Stetson were found to
have inconsistent differences in both Vvega and Ivega.
The three photometric studies do not match, and future
studies which compare theoretical models to the data
will need to take into account these systematic uncer-

6. RESULTS

Using Fidanka we fit pairs of Population A + E isochrones to the HUGS data for NGC 2808. Each pair of isochrones is allowed to vary in distance modulus, reddening, relative helium mass fraction (A/E), and age. The χ^2 distribution for the isochrone pairs is shown in Figure [FIGURE]. The best fit isochrones are shown in Figure [FIGURE] and optimized parameters for these are presented in Table [TABLE].

7. CONCLUSION

Here we have preformed the first chemically selfconsistent modeling of the Milky Way Globular Cluster
KGC 2808. We find that, updated atmospheric boundary conditions and opacity tables do not have a significant effect on the inferred helium abundances of multiple
populations.

This work has made use of the NASA astrophysical data system (ADS). We would like to thank Elisabeth Newton and Aaron Dotter for their support and for useful disscusion related to the topic of this paper. Additionally, we would like to thank Kara Fagerstrom, Aylin Garcia Soto, and Keighley Rockcliffe for their useful disscusion related to in this work. We acknowledge the support of a NASA grant (No. 80NSSC18K0634).

REFERENCES

473 Barber, C. B., Dobkin, D. P., & Huhdanpaa, H. 1996, ACM

475 Bastian, N., & Lardo, C. 2015, MNRAS, 453, 357,

Transactions on Mathematical Software (TOMS), 22, 469

doi: 10.1093/mnras/stv1661

8 Boudreaux et al.

```
477 Bastian, N., & Lardo, C. 2018, Annual Review of
```

- Astronomy and Astrophysics, 56, 83
- 479 Baumgardt, H., & Makino, J. 2003, MNRAS, 340, 227,
- doi: 10.1046/j.1365-8711.2003.06286.x
- 481 Boylan-Kolchin, M. 2018, MNRAS, 1423,
- doi: 10.1093/mnras/sty1490
- 483 Brodie, J. P., & Strader, J. 2006, Annu. Rev. Astron.
- 484 Astrophys., 44, 193
- 485 Carretta, E., Bragaglia, A., Gratton, R. G., et al. 2010,
- 486 Astronomy & Astrophysics, 516, A55
- 487 Colgan, J., Kilcrease, D. P., Magee, N. H., et al. 2016, in
- 488 APS Meeting Abstracts, Vol. 2016, APS Division of
- 489 Atomic, Molecular and Optical Physics Meeting
- 490 Abstracts, D1.008
- 491 de Mink, S. E., Pols, O. R., Langer, N., & Izzard, R. G.
- 492 2009, A&A, 507, L1, doi: 10.1051/0004-6361/200913205
- 493 Decressin, T., Meynet, G., Charbonnel, C., Prantzos, N., &
- ⁴⁹⁴ Ekström, S. 2007, A&A, 464, 1029,
- doi: 10.1051/0004-6361:20066013
- 496 Denissenkov, P. A., & Hartwick, F. D. A. 2014, MNRAS,
- 497 437, L21, doi: 10.1093/mnrasl/slt133
- 498 D'Ercole, A., D'Antona, F., Ventura, P., Vesperini, E., &
- 499 McMillan, S. L. W. 2010, MNRAS, 407, 854,
- doi: 10.1111/j.1365-2966.2010.16996.x
- 501 D'Ercole, A., Vesperini, E., D'Antona, F., McMillan, S.
- 502 L. W., & Recchi, S. 2008, MNRAS, 391, 825,
- doi: 10.1111/j.1365-2966.2008.13915.x
- 504 Dotter, A., Chaboyer, B., Jevremović, D., et al. 2008, The
- Astrophysical Journal Supplement Series, 178, 89
- 506 Dotter, A., Ferguson, J. W., Conroy, C., et al. 2015,
- MNRAS, 446, 1641, doi: 10.1093/mnras/stu2170
- 508 Gratton, R. G., Carretta, E., & Bragaglia, A. 2012,
- Astronomy and Astrophysics Reviews, 20, 50,
- doi: 10.1007/s00159-012-0050-3
- 511 Gustafsson, B., Edvardsson, B., Eriksson, K., et al. 2008,
- 512 A&A, 486, 951, doi: 10.1051/0004-6361:200809724

- 513 Hudson, M. J., & Robison, B. 2018, Monthly Notices of the
- Royal Astronomical Society, 477, 3869,
- doi: 10.1093/mnras/sty844
- $_{516}$ Kravtsov, A. V., & Gnedin, O. Y. 2005, The Astrophysical
- 517 Journal, 623, 650
- 518 Latour, M., Husser, T. O., Giesers, B., et al. 2019, A&A,
- 631, A14, doi: 10.1051/0004-6361/201936242
- 520 Marigo, P., & Aringer, B. 2009, A&A, 508, 1539,
- doi: 10.1051/0004-6361/200912598
- 522 Marigo, P., Aringer, B., Girardi, L., & Bressan, A. 2022,
- 523 ApJ, 940, 129, doi: 10.3847/1538-4357/ac9b40
- 524 Milone, A. P., Marino, A. F., Piotto, G., et al. 2015, ApJ,
- 525 808, 51, doi: 10.1088/0004-637X/808/1/51
- 526 Peng, E. W., Ferguson, H. C., Goudfrooij, P., et al. 2011,
- 527 The Astrophysical Journal, 730, 23
- 528 Piotto, G., Bedin, L. R., Anderson, J., et al. 2007, The
- Astrophysical Journal Letters, 661, L53,
- doi: 10.1086/518503
- 531 Piotto, G., Milone, A. P., Bedin, L. R., et al. 2015, AJ, 149,
- 91, doi: 10.1088/0004-6256/149/3/91
- 533 Plez, B. 2008, Physica Scripta Volume T, 133, 014003,
- doi: 10.1088/0031-8949/2008/T133/014003
- 535 Prantzos, N., Charbonnel, C., & Iliadis, C. 2007, A&A, 470,
- 179, doi: 10.1051/0004-6361:20077205
- 537 Renzini, A. 2008, Monthly Notices of the Royal
- Astronomical Society, 391, 354,
- doi: 10.1111/j.1365-2966.2008.13892.x
- 540 Sbordone, L., Salaris, M., Weiss, A., & Cassisi, S. 2011,
- Astronomy & Astrophysics, 534, A9
- 542 Sneden, C., Kraft, R. P., Prosser, C. F., & Langer, G. 1992,
- The Astronomical Journal, 104, 2121
- Valle, G., Dell'Omodarme, M., & Tognelli, E. 2022, A&A,
- 658, A141, doi: 10.1051/0004-6361/202142454
- 546 van den Bergh, S. 2010, The Astronomical Journal, 140,
- 1043, doi: 10.1088/0004-6256/140/4/1043
- ⁵⁴⁸ Ventura, P., & D'Antona, F. 2009, A&A, 499, 835,
- $\qquad \qquad \text{doi: } 10.1051/0004\text{-}6361/200811139$
- 550 Ventura, P., D'Antona, F., Mazzitelli, I., & Gratton, R.
- ⁵⁵¹ 2001, ApJL, 550, L65, doi: 10.1086/319496