11

12

13

14

15

16

17

18

19

20

21

22

23

24

25

# Chemically Self-Consistant Modeling of the Globular Cluster NGC 2808 and its Effects on the Inferred Helium abundance of Multiple Stellar Populations.

EMILY M. BOUDREAUX, BRIAN C. CHABOYER, AMANDA ASH, RENATA EDAES HOH, AND GREGORY FEIDEN

<sup>1</sup>Department of Physics and Astronomy, Dartmouth College, Hanover, NH 03755, USA
<sup>2</sup>Department of Physics and Astronomy, University of North Georgia, Dahlonega, GA 30533, USA

#### ABSTRACT

The Helium abundances in the multiple populations which are now known to comprise all closely studied Milky Way globular clusters are often inferred by fitting isochohrones generated from stellar evolutionary models to globular cluster photometry. It is therefore important to build stellar models that are chemically self-consistent in terms of their structure, atmosphere, and opacity. In this work we present the first chemically self-consistent stellar models of the Milky Way Globular Cluster NGC 2808 using MARCS model atmospheres, OPLIB high-temperature radiative opacities, and AESOPUS low-temperature radiative opacities. These stellar models were fit to the NGC 2808 photometry using Fidanka, a new software tool that was developed optimally fit cluster photometry to isochrones and for population synthesis. Fidanka can determine, in a relatively unbiased way, the ideal number of distinct populations which exist within a dataset and then fits isochrones to each population. We achieve this through a combination of Bayesian Gaussian Mixture Modeling and a novel number density estimation algorithm. Using Fidanka and F275W-F814W photometry from the Hubble UV Globular Cluster Survey we find that the helium abundance of the second generation of stars in NGC 2808 is higher than the first generation by  $15 \pm 3\%$ . This is in agreement with previous studies of NGC 2808. This work, along with previous work by Dotter (2016) focused on NGC 6752 demonstrates that chemically self-consistent models of globular clusters do not signifigantly alter infered helium abundances and are therefor unlikeley to be worth the signifigant additional time investment.

Keywords: Globular Clusters (656), Stellar evolutionary models (2046)

#### 1. INTRODUCTION

Globular clusters (GCs) are among the oldest observable objects in the universe (Peng et al. 2011). They are characterized by high densities with typical half-light radii of  $\leq$ 10 pc (van den Bergh 2010), and typical masses ranging from  $10^4$ – $10^5~{\rm M}_{\odot}$  (Brodie & Strader 2006) — though some GCs are significantly larger than these typical values (e.g.  $\omega$  Cen, Richer et al. 1991). GCs provide a unique way to probe stellar evolution (Baumgardt & Makino 2003), galaxy formation models (Boylan-Kolchin 2018; Kravtsov & Gnedin 2005), and dark matter halo structure (Hudson & Robison 2018).

Corresponding author: Emily M. Boudreaux emily.m.boudreaux.gr@dartmouth.edu, emily@boudreauxmail.com

The traditional view of Globular Clusters was that 38 they consisted of a single stellar population (SSP, in 39 some publications this is referred to as a Simple Stel-40 lar Population). This view was supported by spectro-41 scopically uniform heavy element abundances (Carretta 42 et al. 2010; Bastian & Lardo 2018) across most clus-43 ters (M54 and  $\omega$ Cen are notable exceptions, see Marino 44 et al. (2015) for further details), and the lack of ev-45 idence for multiple stellar populations (MPs) in past 46 color-magnitude diagrams of GCs (i.e. Sandage 1953; 47 Alcaino 1975). However, over the last 40 years non-48 trivial star-to-star light-element abundance variations 49 have been observed (i.e. Smith 1987) and, in the last 50 two decades, it has been definitively shown that most if 51 not all Milky Way GCs have MPs (Gratton et al. 2004, 52 2012; Piotto et al. 2015). The lack of photometric evi-53 dence for MPs prior to the 2000, can be attributed to the Boudreaux et al.

54 more narrow color bands available, until very recently, to <sub>55</sub> ground based photometric surveys (Milone et al. 2017). The prevalence of multiple populations in GCs is so 57 distinct that the proposed definitions for what consti-58 tutes a globular cluster now often center the existence <sup>59</sup> of MPs (e.g. Carretta et al. 2010). Whereas, people have 60 have often tried to categorized objects as GCs through 61 relations between half-light radius, density, and surface 62 brightness profile, in fact many objects which are gener-63 ally thought of as GCs don't cleanly fit into these cuts 64 (Peebles & Dicke 1968; Brown et al. 1991, 1995; Bekki 65 & Chiba 2002). Consequently, Carretta et al. (2010) 66 proposed a definition of GC based on observed chem-67 ical inhomogeneities in their stellar populations. The  $_{\rm 68}$  modern understanding of GCs then is not simply one of dense cluster of stars that may have chemical inho-70 mogeneities and multiple populations; rather, it is one 71 where those chemical inhomogeneities and multiple populations themselves are the defining element of a GC.

All Milky Way globular clusters older than 2 Gyr studied in detail show populations enriched in He, N, and Na while also being deplete in O and C (Piotto et al. 2015; Bastian & Lardo 2018). These light element abundance patterns also are not strongly correlated with variations in heavy element abundance, resulting in spectroscopically uniform Fe abundances between populations. Further, high-resolution spectral studies reveal anti-correlations between N-C abundances, Na-O abundances, and potentially Al-Mg (Sneden et al. 1992; Gratton et al. 2012). Typical stellar fusion reactions can deplete core oxygen; however, the observed abundances of Na, Al, and Mg cannot be explained by the CNO cycle (Prantzos et al. 2007). Consequently, globular cluster populations must be formed by some novel means.

Formation channels for these multiple populations remain a point of debate among astronomers. Most proposed formation channels consist of some older, more massive, population of stars polluting the pristine cluster media before a second population forms, now enriched in heavier elements which they themselves could not have generated (for a detailed review see Gratton et al. 2012). The four primary candidates for these polluters are asymptotic giant branch stars (AGBs, Ventura et al. 2001; D'Ercole et al. 2010), fast rotating massive stars (FRMSs, Decressin et al. 2007), super massive stars (SMSs, Denissenkov & Hartwick 2014), and massive interacting binaries (MIBs, de Mink et al. 2009; Bastian & Lardo 2018).

Hot hydrogen burning (i.e. proton capture), material transport to the surface, and material ejection into the intra-cluster media are features of each of these models and consequently they can all be made to *qualitatively* 

agree with the observed elemental abundances. Howver, none of the standard models can currently account
for all specific abundances (Gratton et al. 2012). AGB
and FRMS models are the most promising; however,
both models have difficulty reproducing severe O depletion (Ventura & D'Antona 2009; Decressin et al. 2007).
Moreover, AGB and FRMS models require significant
mass loss (~ 90%) between cluster formation and the
current epoch — implying that a significant fraction of
halo stars formed in GCs (Renzini 2008; D'Ercole et al.
2008; Bastian & Lardo 2015).

In addition to the light-element anti-correlations ob-118 served, it is also known that younger populations are significantly enhanced in Helium (Piotto et al. 2007, 2015; 120 Latour et al. 2019). Depending on the cluster, helium mass fractions as high as Y = 0.4 have been inferred (e.g. 122 Milone et al. 2015a). However, due to both the relatively 123 high and tight temperature range of partial ionization 124 for He and the efficiency of gravitational settling in core 125 helium burning stars, the initial He abundance of glob-126 ular cluster stars cannot be observed; consequently, the 127 evidence for enhanced He in GCs originates from com-128 parison of theoretical stellar isochrones to the observed 129 color-magnitude-diagrams of globular clusters. There-130 fore, a careful handling of chemistry is essential when modeling with the aim of discriminating between MPs; 132 yet, only a very limited number of GCs have been stud-133 ied with chemically self-consistent (structure and atmo-134 sphere) isochrones (e.g. Dotter et al. 2015, NGC 6752). NGC 2808 is the prototype globular cluster to host 136 Multiple Populations. Various studies since 2007 have 137 identified that it may host anywhere from 2-5 stellar 138 populations. These populations have been identified 139 both spectroscopically (i.e. Carretta et al. 2004; Car-140 retta 2006; Carretta et al. 2010; Gratton et al. 2011; 141 Carretta 2015; Hong et al. 2021) and photometrically 142 (i.e. Piotto et al. 2007, 2015; Milone et al. 2015a, 2017; Pasquato & Milone 2019). Note that recent work (Valle 144 et al. 2022) calls into question the statistical signifi-145 cance of the detections of more than 2 populations in 146 the spectroscopic data. Here we present new, chemi-147 cally self-consistent modeling of the photometry of the 148 two extreme populations of NGC 2808 identified by 149 Milone et al. (2015a), populations A and E. We do 150 not consider populations B, C, or D identified 151 in Milone et al. (2015a) as the purpose of this 152 work is to identify if chemically self-consistent 153 modelling results in a statistically signifigant de-154 viation in the infered helium abundance when 155 compared to non chemically self-consistent mod-Use of the two populations in the NGC 157 2808 with the highest identified difference be-

234

tween their helium populations is sufficent for to answer this question. We use archival photometry from the Hubble UV Globular Cluster Survey (HUGS) (Piotto et al. 2015; Milone et al. 2017) in the F275W and F814W passbands to characterize multiple populations in NGC 2808 (Milone et al. 2015a,b) (This data is avalible at MAST: 10.17909/T9810F). Additionally, we present a likelihood analysis of the photometric data of NGC 2808 to determine the number of populations present in the cluster.

#### 2. CHEMICAL CONSISTENCY

168

183

There are three primary areas in which must the stel170 lar models must be made chemically consistent: the at171 mospheric boundary conditions, the opacities, and inte172 rior abundances. The interior abundances are relatively
173 easily handled by adjusting parameters within our stel174 lar evolutionary code. However, the other two areas
175 are more complicated to bring into consistency. Atmo176 spheric boundary conditions and opacities must both be
177 calculated with a consistent set of chemical abundances
178 outside of the stellar evolution code. For evolution we
179 use the Dartmouth Stellar Evolution Program (DSEP)
180 (Dotter et al. 2008), a well tested 1D stellar evolution
181 code which has a particular focus on modelling low mass
182 stars ( $\leq 2 M_{\odot}$ )

# 2.1. Atmospheric Boundary Conditions

Certain assumptions, primarily that the radiation field 184 185 is at equilibrium and radiative transport is diffusive (Salaris & Cassisi 2005), made in stellar structure codes, 187 such as DSEP, are valid when the optical depth of a star is large. However, in the atmospheres of stars, the number density of particles drops low enough and the optical depth consequently becomes small enough that these 191 assumptions break down, and separate, more physically motivated, plasma modeling code is required. Generally 193 structure code will use tabulated atmospheric boundary 194 conditions generated by these specialized codes, such as ATLAS9 (Kurucz 1993), PHOENIX (Husser et al. 196 2013), MARCS (Gustafsson et al. 2008), and MPS-<sup>197</sup> ATLAS (Kostogryz et al. 2023). Often, as the boundary 198 conditions are expensive to compute, they are not updated as interior abundances vary.

One key element when chemically consistently modeling NGC 2808 modeling is the incorporation of new atmospheric models with the same elemental abundances as the structure code. We use atmospheres generated from the MARCS grid of model atmospheres (Plez 2008). MARCS provides one-dimensional, hydrostatic, plane-parallel and spherical LTE atmospheric models (Gustafsson et al. 2008). Model atmospheres

<sup>208</sup> are made to match the spectroscopically measured ele-<sup>209</sup> mental abundances of populations A and E. Moreover, <sup>210</sup> for each population, atmospheres with various helium <sup>211</sup> mass fractions are generated. These range from Y=0.24 <sup>212</sup> to Y=0.36 in steps of 0.03. All atmospheric models are <sup>213</sup> computed to an optical depth of  $\tau=100$  where their <sup>214</sup> temperature and pressures serves as boundary condi-<sup>215</sup> tions for the structure code. In general, enhancing he-<sup>216</sup> lium in the atmosphere has only a small impact on the <sup>217</sup> atmospheric temperature profile, while leading to a drop <sup>218</sup> in the pressure by  $\sim 10-20\%$ .

## 2.2. Opacities

In addition to the atmospheric boundary conditions, both the high and low temperature opacities used by DSEP must be made chemically consistent. Here we use OPLIB high temperature opacity tables (Colgan et al. 2016) retrieved using the TOPS web-interface. Retrival of High temperature opacities is done using pyTOPSScrape, first introduced in Boudreaux & Chaboyer (2023). Low temperature opacity tables are retrieved from the Aesopus 2.0 web-interface (Marigo & Aringer 2009; Marigo et al. 2022). Ideally, these opacities would be the same used in the atmospheric models. However, the opacities used in the MARCS models are not publicly available. As such, we use the opacities provided by the TOPS and Aesopus 2.0 web-interfaces.

# 3. STELLAR MODELS

We use the Dartmouth Stellar Evolution Program (DSEP, Dotter et al. 2008) to generate stellar models. DSEP is a one-dimensional stellar evolution code which includes a mixing length model of convection, gravitational settling, and diffusion. Using the solar composition presented in (Grevesse et al. 2007) (GAS07), MARCS model atmosphers, OPLIB high temperature opacities we find a solar calibrated mixing length parameter,  $\alpha_{MLT}$ , of  $\alpha_{MLT}=1.901$ .

We use DSEP to evolve stellar models ranging in mass from 0.3 to 2.0 solar masses from the fully convective pre-main sequence to the tip of the red giant branch. Below 0.7  $M_{\odot}$  we evolve a model every 0.03  $M_{\odot}$  and above 0.7  $M_{\odot}$  we evolve a model every 0.05  $M_{\odot}$ . We evolve models over a grid of mixing length parameters from  $\alpha_{MLT}=1.0$  to  $\alpha_{MLT}=2.0$  in steps of 0.1. For each mixing length, a grid of models and isochrones were calculated, with chemical compositions consistent with Populations A and E (see Tables 1 and 1) and a range of helium abundances (Y=0.24, 0.27, 0.30, 0.33, 0.36, and 0.39). In total,144 sets of isochrones, each with a unique composition and mixing length were calculated.

278

279

280

281

282

283

284

285

286

Each model is evolved in DSEP with typical numeric tolerences of one part in  $10^7$ . Each model is allowed a maximum time step of 50 Myr.

For each combination of population, Y, and  $\alpha_{MLT}$  we use the isochrone generation code first presented in Dotter (2016) to generate a grid of isochrones. The isochrone generation code identified equivalent evolutionary points (EEPs) over a series of masses and interpolates between them. The grid of isochrones generated for this work is avalible as a digital supplement to this paper 10.5281/zenodo.10631439. Given the complexity of the parameter space when fitting multiple populations along with the recent warnings in the liteerature regarding overfitting datasets (e.g. Valle et al. 2022) we want to develop a more objective way of fitting isochrones to photometry than if we were to mark median ridge line positions by hand.

#### 4. FIDANKA

When fitting isochrones to the clusters with multiple populations we have four main criteria for any method

- The method must be robust enough to work along the entire main sequence, turn off, and much of the subgiant and red giant branch.
- Any method should consider photometric uncertainty in the fitting process.
- The method should be model independent, weighting any n number of populations equally.
- The method should be automated and require minimal intervention from the user.

We do not believe that any currently available software is a match for our use case. Therefore, we elect
to develop our own software suite, Fidanka. Fidanka
is a python package designed to automate much of the
process of measuring fiducial lines in CMDs, adhering to
the four criteria we lay out above. Primary features of
Fidanka may be separated into three categories: fiducial line measurement, stellar population synthesise, and
isochrone optimization/fitting. Additionally, there are
utility functions that are detailed in the Fidanka documentation.

#### 4.1. Fiducial Line Measurement

Fidanka takes a iterative approach to measuring fidu-300 cial lines, the first step of which is to make a "guess" 301 as to the fiducial line. This initial guess is calculated

Table 1. Population Composition

Element	Pop A	Pop E	Element	Pop A	Pop E
Li	-0.08		In	-1.46	
Be	0.25	_	Sn	-0.22	_
В	1.57	_	Sb	-1.25	_
$^{\mathrm{C}}$	6.87	5.91	Те	-0.08	
N	6.42	6.69	I	-0.71	
O	7.87	6.91	Xe	-0.02	_
$\mathbf{F}$	3.43	_	Cs	-1.18	
Ne	7.12	6.7	Ba	1.05	_
Na	5.11	5.7	La	-0.03	_
Mg	6.86	6.42	Се	0.45	
Al	5.21	6.61	Pr	-1.54	
Si	6.65	6.77	Nd	0.29	
P	4.28	_	Pm	-99.0	_
$\mathbf{S}$	6.31	5.89	Sm	-1.3	_
Cl	-1.13	4.37	Eu	-0.61	_
$\operatorname{Ar}$	5.59	5.17	Gd	-1.19	_
K	3.9	_	Tb	-1.96	_
Ca	5.21	_	Dy	-1.16	_
$\operatorname{Sc}$	2.02	_	Но	-1.78	_
Ti	3.82	_	Er	-1.34	_
V	2.8	_	Tm	-2.16	
$\operatorname{Cr}$	4.51	_	Yb	-1.42	
Mn	4.3	_	Lu	-2.16	_
Fe	6.37	_	Hf	-1.41	
Co	3.86	_	Ta	-2.38	
Ni	5.09	_	W	-1.41	_
Cu	3.06	_	Re	-2.0	_
Zn	2.3	_	Os	-0.86	_
Ga	0.78	_	Ir	-0.88	_
Ge	1.39	_	Pt	-0.64	
As	0.04	_	Au	-1.34	_
Se	1.08	_	Hg	-1.09	_
$\operatorname{Br}$	0.28	_	Tl	-1.36	_
Kr	0.99	_	Pb	-0.51	
Rb	0.26	_	Bi	-1.61	
$\operatorname{Sr}$	0.61	_	Po	-99.0	
Y	1.08	_	At	-99.0	
$\operatorname{Zr}$	1.45	_	Rn	-99.0	
Nb	-0.8	_	Fr	-99.0	_
Мо	-0.38	_	Ra	-99.0	_
$\operatorname{Tc}$	-99.0	_	Ac	-99.0	_
Ru	-0.51	_	Th	-2.2	_
Rh	-1.35	_	Pa	-99.0	
Pd	-0.69	_	U	-2.8	

Note—Relative Metal composition used where a(H) = 12. Where the relative composition is the the same for both populations A and E it is only listed in the population A colum for the sake of visual clarity.

References—Milone et al. (2015a)

Table 2. Population Abundance Ratios

Population	[Fe/H]	$[\alpha/\mathrm{Fe}]$	[C/Fe]	[N/Fe]	[O/Fe]	[r/Fe]	[s/Fe]	C/O	X	Y	Z
A	-1.13	0.32	-0.43	-0.28	0.31	-1.13	-1.13	0.10	0.7285	0.2700	0.00154
$\mathbf{E}$	-1.13	-0.11	-1.39	-0.02	-0.66	-1.13	-1.13	0.10	0.7594	0.240	0.00063

Note—Abundance Ratios for populations A and E in NGC 2808.

References—Milone et al. (2015a)

302 by splitting the CMD into magnitude bins, with uni-303 form numbers of stars per bin (so that bins are cover a 304 small magnitude range over densely populated regions 305 of the CMD while covering a much larger magnitude 306 range in sparsely populated regions of the CMD, such 307 as the RGB). A unimodal Gaussian distribution is then 308 fit to the color distribution of each bin, and the resulting mean color is used as the initial fiducial line guess. This 310 rough fiducial line will approximately trace the area of 311 highest density. The initial guess will be used to verti-312 calze the CMD so that further algorithms can work in 313 1-D magnitude bins without worrying about weighting 314 issues caused by varying projections of the evolutionary 315 sequence onto the magnitude axis. Verticalization is pre-316 formed taking the difference between the guess fiducial 317 line and the color of each star in the CMD.

If Fidanka were to simply apply the same algorithm 319 to the verticalized CMD then the resulting fiducial line 320 would likely be a re-extraction of the initial fiducial line 321 guess. To avoid this, we take a more robust, number 322 density based approach, which considers the distribu-323 tion of stars in both color and magnitude space simul-324 taneously. For each star in the CMD we first using an introselect partitioning algorithm to select the 50 326 nearest stars in F814W vs. F275W-F814W space. To 327 account for the case where the star is at an extreme 328 edge of the CMD, those 50 stars include the star itself (such that we really select 49 stars + 1). We use 330 ghull¹(Barber et al. 1996) to calculate the convex hull 331 of those 50 points. The number density at each star 332 then is defined as  $50/A_{hull}$ , where  $A_{hull}$  is the area of the convex hull. Because we use a fixed number of points 334 per star, and a partitioning algorithm as opposed to a sorting algorithm, this method scales like  $\mathcal{O}(n)$ , where 336 n is the number of stars in the CMD. This method also 337 intrinsically weights the density of of each star equally 338 as the counting statistics per bin are uniform. We are

339 left with a CMD where each star has a defined number 340 density (Figure 1).

Fidanka can now exploit this density map to fit a 342 better fiducial line to the data, as the density map is far 343 more robust to outliers. There are multiple algorithms 344 we implement to fit the fiducial line to the color-density 345 profile in each magnitude bin (Figure 2); they are ex-346 plained in more detail in the Fidanka documentation. 347 However, of most relevance here is the Bayesian Gaus-348 sian Mixture Modeling (BGMM) method. BGMM is a 349 clustering algorithm which, for some fixed number of n- $_{350}$  dimensional Gaussian distributions, K, determines the 351 mean, covariance, and mixing probability (somewhat analogous to amplitude) of each  $k^{th}$  distribution, such 353 that the local lower bound of the likelyhood of each star 354 belonging strongly to a single distribution is maximized. Maximization is preformed using the Dirichlet pro-356 cess, which is a non-parametric Bayesian method of  $_{357}$  determining the number of Gaussian distributions, K, 358 which best fit the data (Ferguson 1973; Pedregosa et al. 359 2011). Use of the Dirichlet process allows for dynamic 360 variation in the number of inferred populations from 361 magnitude bin to magnitude bin. Specifically, popula-362 tions are clearly visually separated from the lower main 363 sequence through the turn off; however, at the turn off 364 and throughout much of the subgiant branch, the two 365 visible populations overlap due to their extremely simi-366 lar ages (i.e. Jordán et al. 2002). The Dirichlet process 367 allows for the BGMM method to infer a single popula-368 tion in these regions, while inferring two populations in 369 regions where they are clearly separated. More gener-370 ally, the use of the Dirichlet process removes the need 371 for a prior on the exact number of populations to fit. 372 Rather, the user specifies a upper bound on the num-373 ber of populations within the cluster. An example bin  $_{374}$  (F814W = 20.6) is shown in Figure 3. Fidanka 's BGMM method first breaks down the ver-

376 ticalized CMD into magnitude bins with uniform num-377 bers of stars per bin (here we adopt 250). Any stars

378 left over are placed into the final bin. For each bin a 379 BGMM model with a maximum of 5 populations is fit

<sup>&</sup>lt;sup>1</sup> https://www.qhull.com

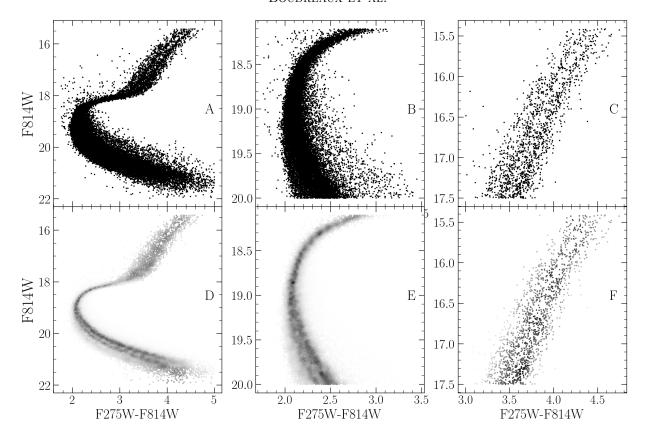
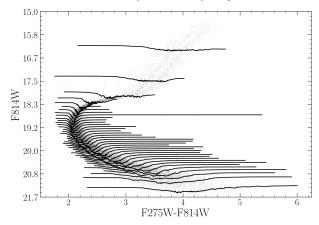


Figure 1. Density map demo showing density estimate over different parts of the evolutionary sequence. The left panel shows the density map over the entire evolutionary sequence, while the middle panel shows the density map over the main sequence and the right most panel shows the density map over the RGB. Figures in the top row are the raw CMD, while figures in the bottom row are colored by the density map.



**Figure 2.** CMD where point brightness is determined by local density. Lines show the density-color profile in each magnitude bin. In this figure adaptive binning targeted 1000 stars per bin

 $_{380}$  to the color density profile. The number of populations  $_{381}$  is then inferred from the weighting parameter (the mix-  $_{382}$  ing probability) of each population. If the weighting pa-  $_{383}$  rameter of any  $k^{th}$  components less than 0.05, then that  $_{384}$  component is considered to be spurious and removed.

Additionally, if the number of populations in the bin above and the bin below are the same, then the number of populations in the current bin is forced to be the same as the number of populations in the bin above. Finally, the initial guess fiducial line is added back to the BGMM inferred line. Figure 4 shows the resulting fiducial line(s) in each magnitude bin for both a verticalized CMD and a non verticalized CMD. In contrast to other work in the literature where evidence for up to 5 distinct populations has been found; we only find evidence for two stellar populations.

This method of fiducial line extraction effectively discriminated between multiple populations along the main sequence and RGB of a cluster, while simultaneously allowing for the presence of a single population along the MSTO and subgiant branch.

We can adapt this density map based BGMM method to consider photometric uncertainties by adopting a simmap ple Monte Carlo approach. Instead of measuring the fiducial line(s) a single time, Fidanka can measure the fiducial line(s) many times, resampling the data with remap placement each time. For each resampling Fidanka adds a random offset to each filter based on the photometric

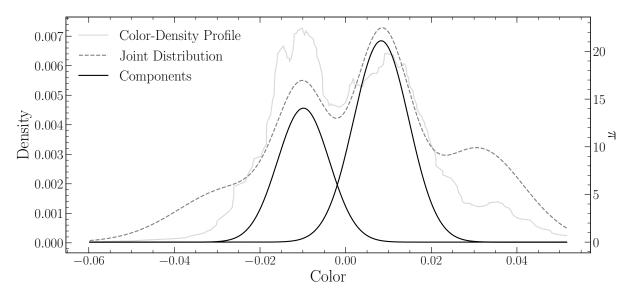


Figure 3. Example of BGMM fit to a magnitude bin. The grey line shows the underlying color-density profile, while the black dashed-line shows the joint distribution of each BGMM component. The solid black lines show the two selected components.

uncertainties of each star. From these n measurements the mean fiducial line for each sequence can be identified along with upper and lower bound confidence intervals in each magnitude bin.

#### 4.2. Stellar Population Synthesis

412

While not extensively used in this paper Fidanka can, in addition to measuring fiducial lines, preform stellar population synthesise. Fidanka's population synthesis module can generate synthetic stellar population from a set of MIST formatted isochrones. This is of primary importance for binary population modeling. The module is also used to generate synthetic CMDs for the purpose of testing the fiducial line extraction algorithms against priors.

Fidanka uses MIST formatted isochrones (Dotter 2016) as input along with distance modulus, B-V color excess, binary mass fraction, and bolometric corrections. An arbitrarily large number of isochrones may be used to define an arbitrary number of populations. Synthetic stars are samples from each isochrone based on a definable probability (for example it is believed that  $\sim 90\%$  of stars in globular clusters are younger population (e.g. Suntzeff & Kraft 1996; Carretta 2013)). Based on the metallicity,  $\mu$ , and E(B-V) of each isochrone, bolometric corrections are taken from bolometric correction tables. Where bolometric correction tables do not include exact metallicities or extinctions a linear interpolation is preformed between the two bounding values.

## 4.3. Isochrone Optimization

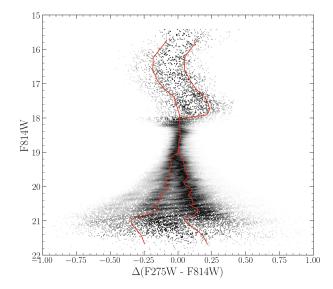
The optimization routines in Fidanka will find the best fit distance modulus, B-V color excess, and binary

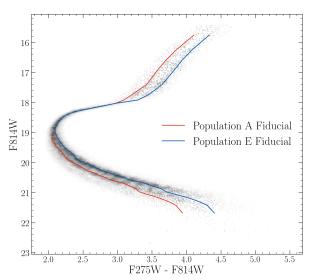
439 number fraction for a given set of isochrones. If a sin-440 gle isochrone is provided then the optimization is done 441 by minimizing the  $\chi^2$  of the perpendicular distances 442 between an isochrone and a fiducial line. If multiple 443 isochrones are provided then those isochrones are first 444 used to run stellar population synthesis and generate a 445 synthetic CMD. The optimization is then done by min-446 imizing the  $\chi^2$  of both the perpendicular distances be-447 tween and widths of the observed fiducial line and the 448 fiducial line of the synthetic CMD.

#### 4.4. Fidanka Testing

In order to validate fidanka we have run an series of injection recovery tests using Fidanka's population synthesis routines to build various synthetic populations and Fidanka's fiducial measurement routines to recover these populations. Each population was generated using the initial mass function given in (Milone et al. 2012) for the redmost population ( $\alpha = -1.2$ ). Further, every population was given a binary population fraction of 10%, distance uniformly sampled between 5000pc and 15000pc, and a B-V color excess uniformly sampled between 0 and 0.1. Finally, each synthetic population was generated using a fixed age uniformly sampled between 7 Gyr and 14 Gyr. An example synthetic population along with its associated best fit isochrone are shown in Figure 5.

For each trial we use Fidanka to measure the fidu-466 cial line and then optimize that fiducial line against the 467 originating isochrone to esimate distance modulus, age, 468 and color B-V excess. Figure 6 is built from 1000 runs of 469 these trials and show the mean and width of the percent 470 error distributions for  $\mu$ , E(B-V), and age. In general





**Figure 4.** Verticalized CMD where point brightness is determined by density (top). CMD where point brightness is determined by density, calculated fiducial lines are shown (bottom). The data used is from the Hubble Space Telescope UV Legacy Survey of Galactic Globular Clusters.

<sup>471</sup> Fidanka is able to recover distance modulii effectively with age and E(B-V) reovery falling in line with other literature that does not cosider the CMD outside of the main sequence, main sequence turn off, sub giant, and red giant branches; specifically, it should be noted that Fidanka is not setup to model the horizontal branch.

#### 5. ISOCHRONE FITTING

477

We fit pairs of isochrones to the HUGS data for NGC 2808 using Fidanka, as described in  $\S 4$ . Two isochrones, one for Population A and one for Population E are fit simultaneously. These isochrones are constrained to have distance modulus,  $\mu$ , and color excess, E(B-V) which

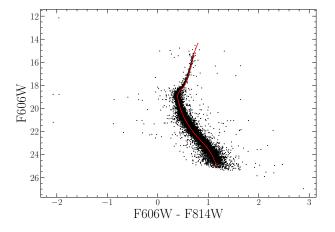
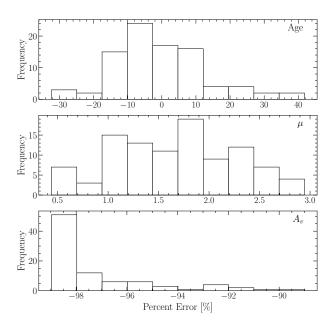


Figure 5. Synthetic population generated by fidanka at 10000pc with E(B-V) = 0, and an age of 12 Gyr along with the best fitting isochrone. The best fit paremeters are derived to be mu = 15.13, E(B-V)=0.001, and an age of 12.33 Gyr.



**Figure 6.** Percent Error distribution for each of the three deriver parameters. Note that these values will be sensitive to the magnitude uncertainties of the photometry. Here we made use of the ACS artificial star tests to estimate the uncertanties.

 $^{483}$  agree to within 0.5% and an ages which agree to within  $^{484}$  1%. Moreover, we constrain the mixing length,  $\alpha_{ML},$   $^{485}$  for any two isochrones in a set to be within 0.5 of one  $^{486}$  and other. For every isochrone in the set of combina- $^{487}$  tion of which fulfilling these constraints  $\mu,~E(B-V),$   $^{488}$  Age\_A, and  $Age_B$  are optimized to reduce the  $\chi^2$  distance  $^{489}$  ( $\chi^2=\sum\sqrt{\Delta {\rm color}^2+\Delta {\rm mag}^2}$ ) between the fiducial lines and the isochrones. Because we fit fiducial lines directly,  $^{491}$  we do not need to consider the binary population fraction,  $f_{bin}$ , as a free parameter.

The best fit isochrones are shown in Figure 7 and optimized parameters for these are presented in Table 1. We find helium mass fractions that are consistent with those identified in past literature (e.g. Milone et al. 2015a). Note that our helium mass fraction grid has a spacing of 0.03 between grid points and we are therefore unable to resolve between certain proposed helium mass fractions for the younger sequence (for example between 0.37 and 0.39).

Past literature (e.g. Milone et al. 2015a, 2018) have found helium mass fraction variation from the low red-most to bluemost populations of  $\sim 0.12$ . Here we find a helium mass fraction variation of 0.15 which, given the spacing of the helium grid we use is consistent with these past results.

# 5.1. The Number of Populartions in NGC 2808

In order to estimate the number of populations which ideally fit the NGC 2808 F275W-F814W photometry without overfitting the data we make use of silhouette analysis (Rousseeuw 1987, and in a similar manner to how Valle et al. (2022) preform their analysis of spectroscopic data). We find the average silhouette score for all tagged clusters identified using BGMM in all magnitude bins over the CMD using the standar python module sklearn. Figure 8 shows the silhouette analysis results and that two populations fit the photometry most ideally. This is in line with what our BGMM model predicts for the majority of the the CMD.

While we make use a purley CMD based ap-521 522 proach in this work, other literature has made 523 use of Chromosome Maps. These consist of im-524 plicitly verticalized pseudo colors. In the chro-525 mosome map for NGC 2808 there may be evi-526 dence for more than two populations; however, 527 the process of transforming magnitude measure-528 ments into chromosome space results in dramat-529 ically increased uncertanties for each star. We 530 find a mean fractional uncertantie for chromosome parameters of  $\approx 1$  when starting with mag-532 nutude measurements having a mean best-case (i.e. uncertainty assumed to only be due to Poisson statistics) fractional uncertainty of  $\approx 0.0005$ . Because of how Fidanka operates, i.e. resampling 536 a probabilty distribution for each star in order to 537 idenfify clusters, we are unable to make statisit-538 cally meaningful statements from the chromo-539 some map

## 5.2. ACS-HUGS Photometric Zero Point Offset

The Hubble legacy archive photometry used in this work is calibrated to the Vega magnitude system. How-

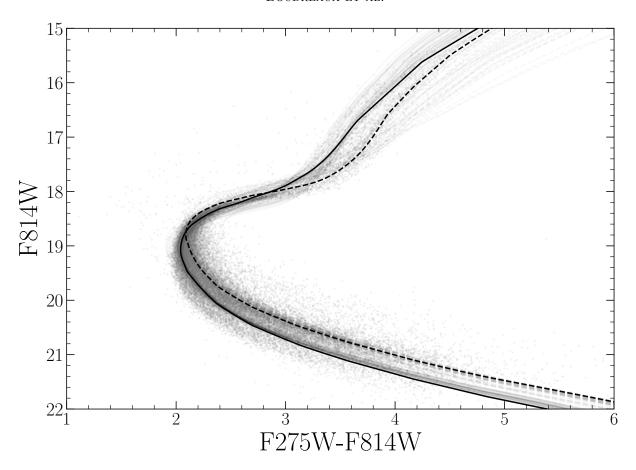
 $_{543}$  ever, we have found that the photometry has a system- $_{544}$  atic offset of  $\sim 0.026$  magnitudes in the F814W band  $_{545}$  when compared to the same stars in the ACS survey  $_{546}$  (Figure 9). The exact cause of this offset is unknown,  $_{547}$  but it is likely due to a difference in the photometric  $_{548}$  zero point between the two surveys. A full correction  $_{549}$  of this offset would require a careful re-reduction of the  $_{550}$  HUGS photometry, which is beyond the scope of this  $_{551}$  work. We instead recognize a 0.02 inherent uncertainty  $_{552}$  in the inferred magnitude of any fit when comparing to  $_{553}$  the ACS survey. This uncertainty is small when com- $_{554}$  pared to the uncertainty in the distance modulus and  $_{555}$  should not affect the conclusion of this paper.

The oberved photometric offset between ACS and 557 HUGS reductions introduces a systematic uncertainity 558 when comparing parameters derived from isochrone fits 559 to ACS data vs those fit to HUGS data. Specifically, this offset introduces a  $\sim 2Gyr$  uncertainty when compar-561 ing ages between ACS and HUGS. Moreover, for two 562 isochrone of the same age, only seperated by helium 563 mass fraction, a shift of the main sequence turn off of 564 is also expected. Figure 10 shows this shift. Note a 565 change in the helium mass fraction of a model by 0.03 566 results in an approximate 0.08 magnitude shift to the 567 main sequence turn off location. This means that the 568 mean 0.026 magnitude offset we find in between ACS 569 and HUGS data corresponds to an additional approaxi-570 mate 0.01 uncertainty in the derived helium mass frac-571 tion when comparing between these two datasets.

# 6. CONCLUSION

Here we have preformed the first chemically selfconsistent modeling of the Milky Way Globular Cluster
NGC 2808. We find that, updated atmospheric boundary conditions and opacity tables do not have a significant effect on the inferred helium abundances of multiple
populations. Specifically, we find that population has a
helium mass fraction of 0.24, while population E has a
helium mass fraction of 0.39. Additionally, we find that
the ages of these two populations agree within uncertainties. Further, we only find evidence for two distinct
stellar populations, which is in agreement with recent
work studying the number of populations in NGC 2808
spectroscopic data.

Further, we introduce a new software suite for glob-1587 ular cluster science, Fidanka, which has been released 1588 under a permissive open source license. Fidanka aims to 1589 provide a statistically robust set of tools for estimating 1590 the parameters of multiple populations within globular 1591 clusters.



**Figure 7.** Best fit isochrone results for NGC 2808. The best fit population A and E models are shown as black lines. The following 50 best fit models are presented as grey lines. The solid black line is fit to population A, while the dashed black line is fit to population E.

Population	Age	Distance Modulus	Extinction	Y	$\alpha_{ML}$	$\chi^2_{\nu}$
	[Gyr]		[mag]			
A	$12.996^{+0.87}_{-0.64}$	15.021	0.54	0.24	2.050	0.021
E	$13.061^{+0.86}_{-0.69}$	15.007	0.537	0.39	1.600	0.033

**Table 1.** Best fit parameters derived from fitting isochrones to the fiducual lines derived from the NCG 2808 photometry. The one sigma uncertainty reported on population age were determined from the 16th and 84th percentiles of the distribution of best fit isochrones ages.

This work has made use of the NASA astrophysical data system (ADS). We would like to thank Elisabeth Newton and Aaron Dotter for their support and for useful disscusion related to the topic of this paper. Additionally, we would like to thank Kara Fagerstrom, Aylin Garcia Soto, and Keighley Rockcliffe for their useful disscusion related to in this work. We acknowledge the support of a NASA grant (No. 80NSSC18K0634).

# REFERENCES

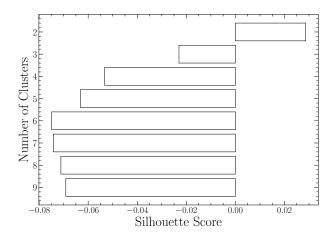


Figure 8. Silhouette analysis for NGC 2808 F275W-F814W photometry. The Silhouette scores are an average of score for each magnitude bin. Positive scores incidate that the clustering algorithm produced well distinguised clusters while negative scores indicate clusters which are not well distinguised.

```
603 Bastian, N., & Lardo, C. 2015, MNRAS, 453, 357,
     doi: 10.1093/mnras/stv1661
605 Bastian, N., & Lardo, C. 2018, Annual Review of
     Astronomy and Astrophysics, 56, 83
606
607 Baumgardt, H., & Makino, J. 2003, MNRAS, 340, 227,
     doi: 10.1046/j.1365-8711.2003.06286.x
608
  Bekki, K., & Chiba, M. 2002, The Astrophysical Journal,
609
     566, 245, doi: 10.1086/337984
610
611 Boudreaux, E. M., & Chaboyer, B. C. 2023, ApJ, 944, 129,
     doi: 10.3847/1538-4357/acb685
613 Boylan-Kolchin, M. 2018, MNRAS, 1423,
     doi: 10.1093/mnras/sty1490
615 Brodie, J. P., & Strader, J. 2006, Annu. Rev. Astron.
     Astrophys., 44, 193
616
617 Brown, J. H., Burkert, A., & Truran, J. W. 1991, ApJ, 376,
     115, doi: 10.1086/170260
618
    -. 1995, ApJ, 440, 666, doi: 10.1086/175304
619
620 Carretta, E. 2006, AJ, 131, 1766, doi: 10.1086/499565
     -. 2013, A&A, 557, A128,
621
     doi: 10.1051/0004-6361/201322103
622
    -. 2015, ApJ, 810, 148, doi: 10.1088/0004-637X/810/2/148
  Carretta, E., Bragaglia, A., & Cacciari, C. 2004, ApJL,
     610, L25, doi: 10.1086/423034
625
  Carretta, E., Bragaglia, A., Gratton, R. G., et al. 2010,
626
     Astronomy & Astrophysics, 516, A55
627
628 Colgan, J., Kilcrease, D. P., Magee, N. H., et al. 2016, in
```

APS Meeting Abstracts, Vol. 2016, APS Division of

Atomic, Molecular and Optical Physics Meeting

632 de Mink, S. E., Pols, O. R., Langer, N., & Izzard, R. G.

2009, A&A, 507, L1, doi: 10.1051/0004-6361/200913205

629

630

631

Abstracts, D1.008

```
634 Decressin, T., Meynet, G., Charbonnel, C., Prantzos, N., &
     Ekström, S. 2007, A&A, 464, 1029,
635
     doi: 10.1051/0004-6361:20066013
636
  Denissenkov, P. A., & Hartwick, F. D. A. 2014, MNRAS,
637
     437, L21, doi: 10.1093/mnrasl/slt133
638
639 D'Ercole, A., D'Antona, F., Ventura, P., Vesperini, E., &
     McMillan, S. L. W. 2010, MNRAS, 407, 854,
640
     doi: 10.1111/j.1365-2966.2010.16996.x
641
  D'Ercole, A., Vesperini, E., D'Antona, F., McMillan, S.
     L. W., & Recchi, S. 2008, MNRAS, 391, 825,
     doi: 10.1111/j.1365-2966.2008.13915.x
645 Dotter, A. 2016, ApJS, 222, 8,
     doi: 10.3847/0067-0049/222/1/8
647 Dotter, A., Chaboyer, B., Jevremović, D., et al. 2008, The
     Astrophysical Journal Supplement Series, 178, 89
648
649 Dotter, A., Ferguson, J. W., Conroy, C., et al. 2015,
     MNRAS, 446, 1641, doi: 10.1093/mnras/stu2170
650
651 Ferguson, T. S. 1973, The annals of statistics, 209
  Gratton, R., Sneden, C., & Carretta, E. 2004, ARA&A, 42,
     385, doi: 10.1146/annurev.astro.42.053102.133945
653
  Gratton, R. G., Carretta, E., & Bragaglia, A. 2012,
654
     Astronomy and Astrophysics Reviews, 20, 50,
655
     doi: 10.1007/s00159-012-0050-3
656
  Gratton, R. G., Lucatello, S., Carretta, E., et al. 2011,
657
     A&A, 534, A123, doi: 10.1051/0004-6361/201117690
658
  Grevesse, N., Asplund, M., & Sauval, A. J. 2007, SSRv,
659
     130, 105, doi: 10.1007/s11214-007-9173-7
660
  Gustafsson, B., Edvardsson, B., Eriksson, K., et al. 2008,
661
     A&A, 486, 951, doi: 10.1051/0004-6361:200809724
663 Hong, S., Lim, D., Chung, C., et al. 2021, AJ, 162, 130,
     doi: 10.3847/1538-3881/ac0ce6
  Hudson, M. J., & Robison, B. 2018, Monthly Notices of the
     Royal Astronomical Society, 477, 3869,
666
     doi: 10.1093/mnras/sty844
667
  Husser, T. O., Wende-von Berg, S., Dreizler, S., et al. 2013,
668
     A&A, 553, A6, doi: 10.1051/0004-6361/201219058
669
670 Jordán, A., Côté, P., West, M. J., & Marzke, R. O. 2002,
     ApJL, 576, L113, doi: 10.1086/343759
671
672 Kostogryz, N., Shapiro, A. I., Witzke, V., et al. 2023,
     Research Notes of the AAS, 7, 39,
673
     doi: 10.3847/2515-5172/acc180
674
675 Kravtsov, A. V., & Gnedin, O. Y. 2005, The Astrophysical
     Journal, 623, 650
676
677 Kurucz, R.-L. 1993, Kurucz CD-Rom, 13
678 Latour, M., Husser, T. O., Giesers, B., et al. 2019, A&A,
     631, A14, doi: 10.1051/0004-6361/201936242
679
680 Marigo, P., & Aringer, B. 2009, A&A, 508, 1539,
     doi: 10.1051/0004-6361/200912598
681
682 Marigo, P., Aringer, B., Girardi, L., & Bressan, A. 2022,
     ApJ, 940, 129, doi: 10.3847/1538-4357/ac9b40
683
```

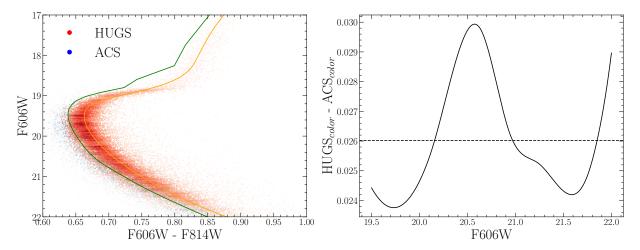


Figure 9. (left) CMD showing the photometric offset between the ACS and HUGS data for NGC 2808. CMDs have been randomly subsampled and colored by point density for clarity. (right) Mean difference between the color of the HUGS and ACS fiducual lines at the same magnitude. Note that the ACS data is systematically bluer than the HUGS data.

705

706

707

708

710

712

714

700 Peebles, P. J. E., & Dicke, R. H. 1968, ApJ, 154, 891,

The Astrophysical Journal, 730, 23

Astrophysical Journal Letters, 661, L53,

91, doi: 10.1088/0004-6256/149/3/91

179, doi: 10.1051/0004-6361:20077205

Astronomical Society, 391, 354,

713 Renzini, A. 2008, Monthly Notices of the Royal

702 Peng, E. W., Ferguson, H. C., Goudfrooij, P., et al. 2011,

Piotto, G., Bedin, L. R., Anderson, J., et al. 2007, The

709 Plez, B. 2008, Physica Scripta Volume T, 133, 014003, doi: 10.1088/0031-8949/2008/T133/014003

711 Prantzos, N., Charbonnel, C., & Iliadis, C. 2007, A&A, 470,

Piotto, G., Milone, A. P., Bedin, L. R., et al. 2015, AJ, 149,

doi: 10.1086/149811

doi: 10.1086/518503

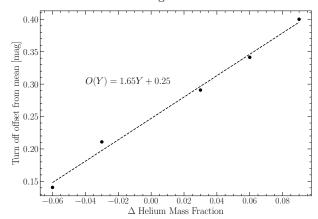


Figure 10. Main sequence turn off magnitude offset from a guage helium mass fraction (Y=0.30 chosen). All main sequence turn off locations are measured at 12.3 Gyr Should I make these contour surfaces for various ages?

```
doi: 10.1111/j.1365-2966.2008.13892.x
                                                                    715
                                                                    716 Richer, H. B., Fahlman, G. G., Buonanno, R., et al. 1991,
684 Marino, A. F., Milone, A. P., Karakas, A. I., et al. 2015,
                                                                         ApJ, 381, 147, doi: 10.1086/170637
                                                                    717
     Monthly Notices of the Royal Astronomical Society, 450,
685
                                                                    718 Rousseeuw, P. J. 1987, Journal of Computational and
     815, doi: 10.1093/mnras/stv420
686
                                                                         Applied Mathematics, 20, 53,
                                                                    719
  Milone, A. P., Piotto, G., Bedin, L. R., et al. 2012, ApJ,
687
                                                                         doi: https://doi.org/10.1016/0377-0427(87)90125-7
                                                                    720
     744, 58, doi: 10.1088/0004-637X/744/1/58
688
                                                                    721 Salaris, M., & Cassisi, S. 2005, Evolution of stars and
  Milone, A. P., Marino, A. F., Piotto, G., et al. 2015a, ApJ,
689
                                                                         stellar populations (John Wiley & Sons)
                                                                    722
     808, 51, doi: 10.1088/0004-637X/808/1/51
690
                                                                    723 Sandage, A. R. 1953, AJ, 58, 61, doi: 10.1086/106822
     -. 2015b, MNRAS, 447, 927, doi: 10.1093/mnras/stu2446
691
                                                                       Smith, G. H. 1987, Publications of the Astronomical
  Milone, A. P., Piotto, G., Renzini, A., et al. 2017, MNRAS,
692
                                                                         Society of the Pacific, 99, 67, doi: 10.1086/131958
     464, 3636, doi: 10.1093/mnras/stw2531
                                                                    725
693
                                                                       Sneden, C., Kraft, R. P., Prosser, C. F., & Langer, G. 1992,
                                                                    726
694 Milone, A. P., Marino, A. F., Renzini, A., et al. 2018,
                                                                         The Astronomical Journal, 104, 2121
     MNRAS, 481, 5098, doi: 10.1093/mnras/sty2573
                                                                    727
695
                                                                    728 Suntzeff, N. B., & Kraft, R. P. 1996, AJ, 111, 1913,
696 Pasquato, M., & Milone, A. 2019, arXiv e-prints,
                                                                         doi: 10.1086/117930
     arXiv:1906.04983, doi: 10.48550/arXiv.1906.04983
                                                                    729
697
                                                                       Valle, G., Dell'Omodarme, M., & Tognelli, E. 2022, A&A,
    Pedregosa, F., Varoquaux, G., Gramfort, A., et al. 2011,
698 F
                                                                    730
                                                                         658, A141, doi: 10.1051/0004-6361/202142454
     Journal of Machine Learning Research, 12, 2825
690
```

 $_{732}$  van den Bergh, S. 2010, The Astronomical Journal, 140,

733 1043, doi: 10.1088/0004-6256/140/4/1043

734 Ventura, P., & D'Antona, F. 2009, A&A, 499, 835,

735 doi: 10.1051/0004-6361/200811139

736 Ventura, P., D'Antona, F., Mazzitelli, I., & Gratton, R.

<sup>737</sup> 2001, ApJL, 550, L65, doi: 10.1086/319496