MA319 — 偏微分方程

Assignment 7

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习题 2.2/3

如果有一长度为 / 的均匀细棒, 其周围及两端 x = 0, x = 1 均为绝热, 初始温度分布为 u(x,0) = f(x), 问以后时刻的温度分布如何? 且证明当 f(x) 等于常数 u_0 时, 恒有 $u(x,t) = u_0$.

题目可归结为求解如下的定解问题:

$$\begin{cases} \frac{\partial u}{\partial t} = a^2 \frac{\partial^2 u}{\partial x^2}, \\ \frac{\partial u}{\partial x}(0, t) = \frac{\partial u}{\partial x}(I, t) = 0, \\ u(x, 0) = f(x). \end{cases}$$

方程的特征值和对应的特征函数为

$$X''(x) + \lambda X(x) = 0,$$

$$X(x) = C_1 \cos \sqrt{\lambda} x + C_2 \sin \sqrt{\lambda} x,$$

代入 X'(0) = 0, X'(I) = 0 得

$$C_2 = 0$$
, $C_1 \sqrt{\lambda} \sin \sqrt{\lambda} I = 0$,

$$\lambda = \lambda_k = \frac{k^2 \pi^2}{I^2}, \quad X_k(x) = C_k \cos \frac{k\pi}{I} x, \quad k = 0, 1, 2, \cdots.$$

方程的通解为

$$u(x,t) = \frac{1}{2}A_0 + \sum_{k=1}^{\infty} A_k e^{-a^2 \lambda t} \cos \sqrt{\lambda} x = \frac{1}{2}A_0 + \sum_{k=1}^{\infty} A_k e^{-\frac{k^2 \pi^2}{J^2} a^2 t} \cos \frac{k\pi}{J} x.$$

代入初值条件得

$$u(x,0) = \frac{1}{2}A_0 + \sum_{k=1}^{\infty} A_k \cos \sqrt{\lambda} x = \frac{1}{2}A_0 + \sum_{k=1}^{\infty} A_k \cos \frac{k\pi}{l} x = f(x).$$

解得

$$A_k = \frac{2}{I} \int_0^I f(x) \cos \sqrt{\lambda} x dx = \frac{2}{I} \int_0^I f(x) \cos \frac{k\pi}{I} x dx.$$

故

$$u(x,t) = \frac{1}{l} \int_0^l f(x) dx + \sum_{k=1}^{\infty} \frac{2}{l} \int_0^l f(x) \cos \frac{k\pi}{l} x dx \cdot e^{-\frac{k^2 \pi^2}{l^2} a^2 t} \cos \frac{k\pi}{l} x.$$

当 f(x) 等于常数 u_0 时

$$A_k = \begin{cases} 2u_0, & k = 0, \\ \frac{u_0 I \sin k\pi}{k\pi} = 0, & k > 0. \end{cases}$$

故

$$u(x, t) = \frac{1}{2}A_0 = u_0.$$

习题 2.2/4

在区域 t > 0, 0 < x < I 中求解如下的定解问题:

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} - \beta(u - u_0), \\ u(0, t) = u(1, t) = u_0, \\ u(x, 0) = f(x), \end{cases}$$

其中 a, β, u_0 均为常数, f(x) 为已知函数

设

$$u(x, t) = u_0 + v(x, t)e^{-\beta t}.$$

则

$$\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} + \beta(u - u_0) = \frac{\partial v}{\partial t} e^{-\beta t} - \beta v e^{-\beta t} - \frac{\partial^2 v}{\partial x^2} e^{-\beta t} + \beta v e^{-\beta t} = e^{-\beta t} \left(\frac{\partial v}{\partial t} - \frac{\partial^2 v}{\partial x^2} \right) = 0,$$

$$v(0, t) e^{-\beta t} = v(l, t) e^{-\beta t} = u(0, t) - u_0 = u(l, t) - u_0 = 0,$$

$$v(x, 0) = u(x, 0) - u_0 = f(x) - u_0.$$

故可以先求解如下定解问题:

$$\begin{cases} \frac{\partial v}{\partial t} = \frac{\partial^2 v}{\partial x^2}, \\ v(0, t) = v(1, t) = 0, \\ v(x, 0) = f(x) - u_0. \end{cases}$$

方程的特征值和对应的特征函数为

$$\lambda = \lambda_k = \frac{k^2 \pi^2}{l^2}, \quad X_k(x) = C_k \sin \sqrt{\lambda} x = C_k \sin \frac{k\pi}{l} x, \quad k = 1, 2, \cdots.$$

方程的通解为

$$v(x,t) = \sum_{k=1}^{\infty} A_k e^{-a^2 \lambda t} \sin \sqrt{\lambda} x = \sum_{k=1}^{\infty} A_k e^{-\frac{k^2 \pi^2}{l^2} a^2 t} \sin \frac{k\pi}{l} x.$$

代入初值条件得

$$v(x,0) = \sum_{k=1}^{\infty} A_k \sin \sqrt{\lambda} x = \sum_{k=1}^{\infty} A_k \sin \frac{k\pi}{l} x = f(x) - u_0.$$

解得

$$A_k = \frac{2}{l} \int_0^l [f(x) - u_0] \sin \sqrt{\lambda} x dx = \frac{2}{l} \int_0^l [f(x) - u_0] \sin \frac{k\pi}{l} x dx = \frac{2}{l} \int_0^l f(x) \sin \frac{k\pi}{l} x dx - \frac{2u_0(1 - \cos k\pi)}{k\pi}.$$

化简得

$$A_k = \frac{2}{I} \int_0^I f(x) \sin \frac{k\pi}{I} x dx + \frac{2u_0[(-1)^k - 1]}{k\pi}$$

故

$$v(x,t) = \sum_{k=1}^{\infty} \left\{ \frac{2}{l} \int_{0}^{l} f(x) \sin \frac{k\pi}{l} x dx + \frac{2u_{0}[(-1)^{k} - 1]}{k\pi} \right\} e^{-\frac{k^{2}\pi^{2}}{l^{2}} a^{2}t} \sin \frac{k\pi}{l} x,$$

$$u(x,t) = u_{0} + e^{-\beta t} \sum_{k=1}^{\infty} \left\{ \frac{2}{l} \int_{0}^{l} f(x) \sin \frac{k\pi}{l} x dx + \frac{2u_{0}[(-1)^{k} - 1]}{k\pi} \right\} e^{-\frac{k^{2}\pi^{2}}{l^{2}} a^{2}t} \sin \frac{k\pi}{l} x.$$

习题 2.2/5

长度为 / 的均匀细杆的初始温度为 $0^{\circ}C$,端点 x = 0 保持常温 u_0 ,而在 x = 1 和侧面上,热量可以发散到周围的介质中去,介质的温度为 $^{\circ}C$,此时杆上的温度分布函数 u(x,t) 满足下述定解问题:

$$\begin{cases} \frac{\partial u}{\partial t} = a^2 \frac{\partial^2 u}{\partial x^2} - b^2 u, \\ u(0, t) = u_0, \quad \left(\frac{\partial u}{\partial x} + Hu \right) \Big|_{x=1} = 0, \\ u(x, 0) = 0, \end{cases}$$

其中 a, b, H 均为常数, 试求出 u(x, t).

设

$$u(x, t) = v(x, t)e^{-b^2t} + f(x).$$

则

$$\begin{split} \frac{\partial u}{\partial t} - a^2 \frac{\partial^2 u}{\partial x^2} + b^2 u &= \frac{\partial v}{\partial t} e^{-b^2 t} - b^2 v e^{-b^2 t} - a^2 \frac{\partial^2 v}{\partial x^2} e^{-b^2 t} - a^2 f''(x) + b^2 v e^{-b^2 t} + b^2 f(x) = 0, \\ v(0, t) e^{-b^2 t} + f(0) &= u_0, \quad \left(\frac{\partial v}{\partial x} + H v \right) \bigg|_{x=I} + f'(I) + H f(I) = 0, \\ v(x, 0) + f(x) &= 0. \end{split}$$

可找到 f(x) 使以下成立

$$\left(\frac{\partial v}{\partial t} - \frac{\partial^2 v}{\partial x^2}\right) e^{-b^2 t} = 0, \quad f''(x) - \frac{b^2}{a^2} f(x) = 0,$$

$$v(0, t) e^{-b^2 t} = 0, \quad f(0) = u_0, \quad \left(\frac{\partial v}{\partial x} + Hv\right) \Big|_{x=I} = 0, \quad f'(I) + Hf(I) = 0.$$

故 f(x) 的通解为

$$f(x) = C_1 e^{\frac{b}{a}x} + C_2 e^{-\frac{b}{a}x}.$$

代入初值条件得

$$f(0) = C_1 + C_2 = u_0,$$

$$f'(I) + Hf(I) = \frac{b}{a} \left(C_1 e^{\frac{bI}{a}} - C_2 e^{-\frac{bI}{a}} \right) + H\left(C_1 e^{\frac{bI}{a}} + C_2 e^{-\frac{bI}{a}} \right) = 0.$$

解得

$$C_{1} = \frac{u_{0}(b - aH)}{aHe^{\frac{2bl}{a}} + be^{\frac{2bl}{a}} - aH + b}, \quad C_{2} = \frac{u_{0}(aH + b)e^{\frac{2bl}{a}}}{aHe^{\frac{2bl}{a}} + be^{\frac{2bl}{a}} - aH + b},$$

$$f(x) = \frac{u_{0}e^{-\frac{bx}{a}} \left[aH\left(e^{\frac{2bl}{a}} - e^{\frac{2bx}{a}}\right) + b\left(e^{\frac{2bl}{a}} + e^{\frac{2bx}{a}}\right) \right]}{aH\left(e^{\frac{2bl}{a}} - 1\right) + b\left(e^{\frac{2bl}{a}} + 1\right)} = \frac{u_{0}\left[aH\sinh\left(\frac{b(l-x)}{a}\right) + b\cosh\left(\frac{b(l-x)}{a}\right) \right]}{aH\sinh\left(\frac{bl}{a}\right) + b\cosh\left(\frac{bl}{a}\right)}.$$

故可以先求解如下定解问题:

$$\begin{cases} \frac{\partial v}{\partial t} = a^2 \frac{\partial^2 v}{\partial x^2}, \\ v(0, t) = 0, \quad \left(\frac{\partial v}{\partial x} + Hv \right) \Big|_{x=1} = 0, \\ v(x, 0) = -f(x). \end{cases}$$

方程的特征值和对应的特征函数为

$$X''(x) + \lambda X(x) = 0,$$

$$X(x) = C_1 \cos \sqrt{\lambda}x + C_2 \sin \sqrt{\lambda}x,$$

代入 X(0) = 0, X'(I) + HX(I) = 0 得

$$C_1=0, \quad C_2(\sqrt{\lambda}\cos\sqrt{\lambda}I+H\sin\sqrt{\lambda}I)=0,$$
 $v=\sqrt{\lambda}I, \quad an v=-rac{\sqrt{\lambda}}{H}.$

存在无数个正根 $v_k > 0$, 满足

$$\left(k-\frac{1}{2}\right)\pi < \nu_k < k\pi, \quad \lambda_k = \left(\frac{\nu_k}{I}\right)^2, \quad X_k(x) = C_k \sin \frac{\nu_k}{I}x, \quad k=1,2,\cdots$$

方程的通解为

$$v(x,t) = \sum_{k=1}^{\infty} A_k e^{-a^2 \lambda t} \sin \sqrt{\lambda} x = \sum_{k=1}^{\infty} A_k e^{-\frac{v_k^2}{l^2} a^2 t} \sin \frac{v_k}{l} x.$$

代入初值条件得

$$v(x,0) = \sum_{k=1}^{\infty} A_k \sin \sqrt{\lambda} x = \sum_{k=1}^{\infty} A_k \sin \frac{v_k}{l} x = -f(x).$$

由固有函数系 $\{X_k\}=\{\sin\sqrt{\lambda_k}x\}$ 的正交性可得

$$M_k = \int_0^I \sin^2 \sqrt{\lambda_k} x dx = \frac{I}{2} + \frac{H}{2(H^2 + \lambda_k)} = \frac{I}{2} + \frac{HI^2}{2(H^2I^2 + v_k^2)},$$
 $A_k = \frac{1}{M_k} \int_0^I -f(x) \sin \sqrt{\lambda_k} x dx = -\frac{1}{M_k} \int_0^I f(x) \sin \frac{v_k}{I} x dx.$

故

$$v(x,t) = -\sum_{k=1}^{\infty} \frac{1}{M_k} \int_0^I f(\xi) \sin \frac{v_k}{I} \xi d\xi \cdot e^{-\frac{v_k^2}{I^2} a^2 t} \sin \frac{v_k}{I} x,$$

$$u(x,t) = -e^{-b^2t} \sum_{k=1}^{\infty} \frac{1}{M_k} \int_0^1 f(\xi) \sin \frac{v_k}{l} \xi d\xi \cdot e^{-\frac{v_k^2}{l^2} a^2 t} \sin \frac{v_k}{l} x + \frac{u_0 \left[aH \sinh \left(\frac{b(l-x)}{a} \right) + b \cosh \left(\frac{b(l-x)}{a} \right) \right]}{aH \sinh \left(\frac{bl}{a} \right) + b \cosh \left(\frac{bl}{a} \right)}.$$

习题 2.3/1

求下列函数的傅里叶变换:

(1)
$$e^{-\eta x^2}$$
 $(\eta > 0)$;
(2) $e^{-a|x|}$ $(a > 0)$.

(2)
$$e^{-a|x|}$$
 $(a > 0)$

(1)

$$g(\lambda) = \int_{-\infty}^{\infty} f(\xi) e^{-i\lambda\xi} d\xi = \int_{-\infty}^{\infty} e^{-\eta\xi^2 - i\lambda\xi} d\xi = \int_{-\infty}^{\infty} e^{-\eta\xi^2 - i\lambda\xi + \frac{\lambda^2}{4\eta} - \frac{\lambda^2}{4\eta}} d\xi = e^{-\frac{\lambda^2}{4\eta}} \int_{-\infty}^{\infty} e^{-\eta\left(\xi + \frac{i\lambda}{2\eta}\right)^2} d\xi.$$

设

$$\int_{-\infty}^{\infty} e^{-\eta x^2} dx = \int_{-\infty}^{\infty} e^{-\eta y^2} dy = A,$$

$$A^2 = \int_{-\infty}^{\infty} e^{-\eta x^2} dx \int_{-\infty}^{\infty} e^{-\eta y^2} dy = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\eta (x^2 + y^2)^2} dx dy.$$

 $r^2 = x^2 + y^2$

$$A^2 = \int_0^{2\pi} \int_0^{\infty} e^{-\eta r^2} r dr d\theta = 2\pi \int_0^{\infty} -\frac{1}{2\eta} e^{-\eta r^2} d(-\eta r^2) = \frac{\pi}{\eta} \int_{-\infty}^0 e^s ds = \frac{\pi}{\eta},$$

$$g(\lambda) = e^{-\frac{\lambda^2}{4\eta}} \int_{-\infty}^{\infty} e^{-\eta \left(\xi + \frac{i\lambda}{2\eta}\right)^2} d\left(\xi + \frac{i\lambda}{2\eta}\right) = e^{-\frac{\lambda^2}{4\eta}} \int_{-\infty}^{\infty} e^{-\eta x^2} dx = Ae^{-\frac{\lambda^2}{4\eta}} = \sqrt{\frac{\pi}{\eta}} e^{-\frac{\lambda^2}{4\eta}}.$$

(2)

$$g(\lambda) = \int_{-\infty}^{\infty} f(\xi) e^{-i\lambda\xi} d\xi = \int_{-\infty}^{\infty} e^{-a|\xi|-i\lambda\xi} d\xi = \int_{-\infty}^{\infty} e^{-a|\xi|} \cos \lambda \xi d\xi + \int_{-\infty}^{\infty} e^{-a|\xi|} i \sin \lambda \xi d\xi.$$

由于 $e^{-a|\xi|}$, $\cos \lambda \xi$ 是偶函数, $i \sin \lambda \xi$ 是奇函数,

$$\int_{-\infty}^{\infty} e^{-a|\xi|} \cos \lambda \xi d\xi = 2 \int_{0}^{\infty} e^{-a\xi} \cos \lambda \xi d\xi, \quad \int_{-\infty}^{\infty} e^{-a|\xi|} i \sin \lambda \xi d\xi = 0,$$

$$g(\lambda) = 2 \int_0^\infty e^{-a\xi} \cos \lambda \xi d\xi = \left. \frac{e^{-a\xi} (\lambda \sin(\lambda \xi) - a \cos(\lambda \xi))}{a^2 + \lambda^2} \right|_0^\infty = \frac{a}{a^2 + \lambda^2}.$$

习题 2.3/2

证明: 当 f(x) 在 $(-\infty, \infty)$ 上绝对可积时, F[f] 为连续函数.

$$F[f](\lambda+h) - F[f](\lambda) = \int_{-\infty}^{\infty} f(\xi)e^{-i\lambda\xi}d\xi - \int_{-\infty}^{\infty} f(\xi)e^{-i(\lambda+h)\xi}d\xi = \int_{-\infty}^{\infty} f(\xi)e^{-i\lambda\xi}(1-e^{-ih\xi})d\xi,$$
$$|F[f](\lambda+h) - F[f](\lambda)| = \left|\int_{-\infty}^{\infty} f(\xi)e^{-i\lambda\xi}(1-e^{-ih\xi})d\xi\right| \leqslant \int_{-\infty}^{\infty} |f(\xi)| \left|e^{-i\lambda\xi}\right| \left|1-e^{-ih\xi}\right|d\xi.$$

易知当 $\xi \in \mathbf{R}$ 时,

$$\begin{split} \left| e^{-i\lambda\xi} \right| &= \sqrt{\cos^2 \lambda \xi + \sin^2 \lambda \xi} = 1, \quad \left| 1 - e^{-ih\xi} \right| &= \sqrt{(1 - \cos h\xi)^2 + \sin^2 h\xi} = \sqrt{2} \cdot \sqrt{1 - \cos h\xi}, \\ \left| F[f](\lambda + h) - F[f](\lambda) \right| &\leq \int_{-\infty}^{\infty} |f(\xi)| \sqrt{2} \cdot \sqrt{1 - \cos h\xi} d\xi. \end{split}$$

现只需证明, 任取 $\varepsilon > 0$, 存在 h > 0 使得 $|F[f](\lambda + h) - F[f](\lambda)| \leq \varepsilon$. 将上式拆分为三个区间的积分

$$|F[f](\lambda+h)-F[f](\lambda)|\leqslant \sqrt{2}\left(\int_{-\infty}^A|f(\xi)|\sqrt{1-\cos h\xi}d\xi+\int_A^B|f(\xi)|\sqrt{1-\cos h\xi}d\xi+\int_B^\infty|f(\xi)|\sqrt{1-\cos h\xi}d\xi\right).$$

由于 f(x) 在 $(-\infty, \infty)$ 上绝对可积, 可找到 A < 0, B > 0 使得

$$\sqrt{2}\int_{-\infty}^{A}|f(\xi)|\sqrt{1-\cos h\xi}d\xi\leqslant\sqrt{2}\int_{-\infty}^{A}|f(\xi)|d\xi\leqslant\frac{\varepsilon}{3},$$

$$\sqrt{2}\int_{B}^{\infty}|f(\xi)|\sqrt{1-\cos h\xi}d\xi\leqslant\sqrt{2}\int_{B}^{\infty}|f(\xi)|d\xi\leqslant\frac{\varepsilon}{3}.$$

设

$$C = \max\{-A, B\}, \quad M = \sup_{x \in [A, B]} |f(x)|, \quad L = B - A,$$

$$\sqrt{2} \int_A^B |f(\xi)| \sqrt{1-\cos h\xi} d\xi \leqslant \sqrt{2} (B-A) \sup_{x \in [A,B]} |f(x)| \sup_{x \in [A,B]} \sqrt{1-\cos hx} = \sqrt{2} ML \sqrt{1-\cos hC}.$$

令

$$h \leqslant rac{1}{C} \arccos \left(1 - rac{arepsilon^2}{18 M^2 L^2}
ight)$$
 ,

则

$$\sqrt{2} \int_{A}^{B} |f(\xi)| \sqrt{1 - \cos h\xi} d\xi \leqslant \frac{\varepsilon}{3},$$
$$|F[f](\lambda + h) - F[f](\lambda)| \leqslant \varepsilon.$$

故得证.