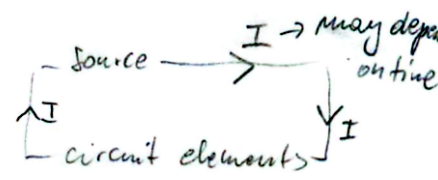


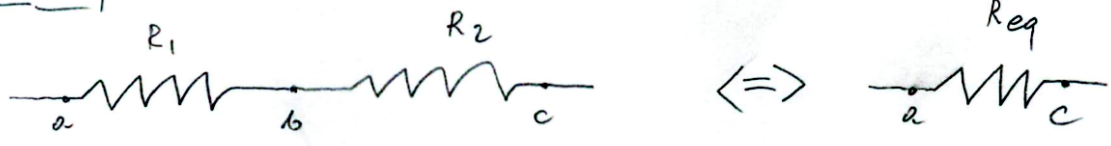
DC circuits

direct-current (no change in direction of the current)



Systems of resistors

SERIES CONNECTION



$$V_{ab} + V_{bc} = V_{ac} \quad (V = \int \vec{E} \cdot d\vec{r} \downarrow \text{additive})$$

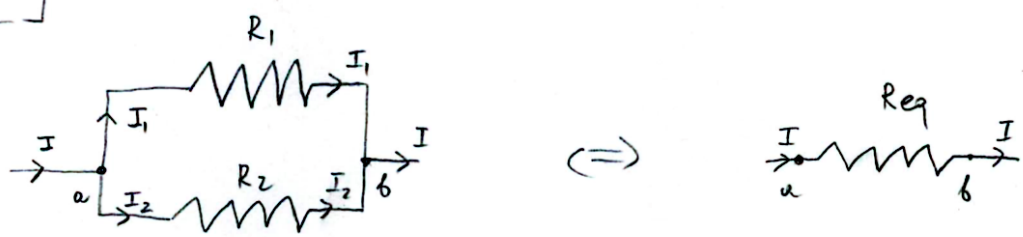
Current through system is I ; use $V = IR$ (Ohm's law)

$$IR_1 + IR_2 = I R_{eq}$$

$R_{eq} = R_1 + R_2$

can generalize to n resistors
 $R_{eq} = \sum_{i=1}^n R_i$

PARALLEL CONNECTION



$$I_1 + I_2 = I \quad (\text{Conservation of charge})$$

Potential change on both resistors is equal (V_{ab}); use $I = \frac{V}{R}$

$$\frac{V_{ab}}{R_1} + \frac{V_{ab}}{R_2} = \frac{V_{ab}}{R_{eq}}$$

$\frac{1}{R_{eq}} = \frac{1}{R_1} + \frac{1}{R_2}$

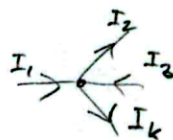
can generalize to n resistors
 $\frac{1}{R_{eq}} = \sum_{i=1}^n \frac{1}{R_i}$

Note. $V_{ab} = I_1 R_1$
 $V_{ab} = I_2 R_2 \Rightarrow \frac{R_1}{R_2} = \frac{I_2}{I_1}$ most current through resistor of least resistance

Kirchhoff's Rules (mid 19th century)

The algebraic sum of the currents into any junction is zero

$$\sum_k I_k = 0$$

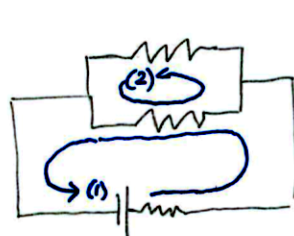


junction rule
(any junction)

The algebraic sum of the potential differences in any loop (including potential differences across emfs and resistors) is zero.

$$\sum_k V_k = 0$$

loop rule
(any loop)



possible loops



Justification

① junction rule - conservation of charge

$$0 = \sum_{i=1}^n I_i = \sum_{i=1}^n \frac{dQ_i}{dt} = \frac{d}{dt} \left(\sum_{i=1}^n Q_i \right) \Rightarrow \sum_{i=1}^n Q_i = \text{const}$$

② loop rule - electric force is conservative

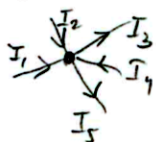
$$\int_{A \rightarrow B} \vec{E} d\vec{r} \text{ does not depend on path or } \oint \vec{E} d\vec{r} = 0$$

$$\begin{aligned} 0 &= \oint_{\text{loop}} \vec{E} d\vec{r} = \int_{A_1 \rightarrow A_2} \vec{E} d\vec{r} + \int_{A_2 \rightarrow A_3} \vec{E} d\vec{r} + \dots + \int_{A_N \rightarrow A_1} \vec{E} d\vec{r} = \\ &= V_{A_1 A_2} + V_{A_2 A_3} + \dots + V_{A_N A_1} \end{aligned}$$

~~~~~ 0 ~~~~~

## SIGN CONVENTION

①



$$I_1 + I_2 - I_3 + I_4 - I_5 = 0$$



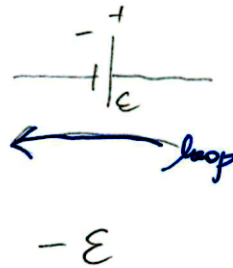
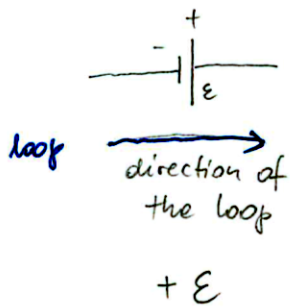
"plus"



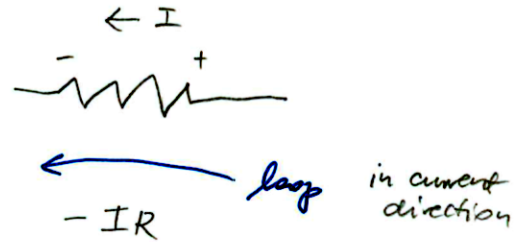
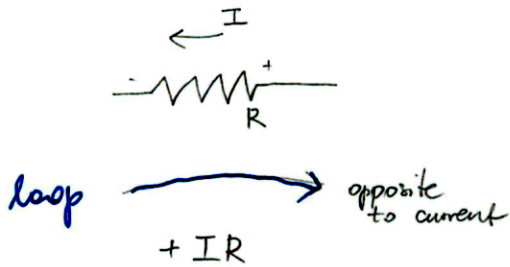
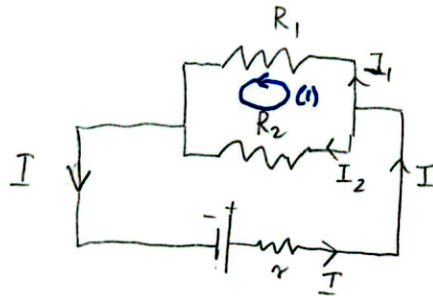
"minus"

(2)

emfs



resistors

Example

$$\text{loop (1): } -I_1 R_1 + I_2 R_2 = 0$$

$$\text{loop (2): } -I r - I_2 R_2 + \varepsilon = 0$$

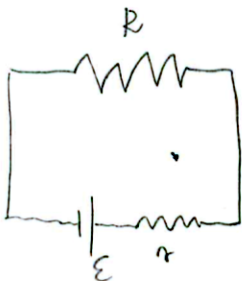
$$\text{junction rule: } I - I_1 - I_2 = 0$$

$$\Downarrow$$

number of independent equations

$$''$$

number of unknowns

Example (1)

Find  $R$  s.t. the power on the resistor  $R$  is maximum ( $\varepsilon, r$  - given)

$$\text{loop rule: } -I r - I R + \varepsilon = 0$$

$$I = \frac{\varepsilon}{R+r}$$

Power

$$P = I^2 R = \frac{\varepsilon^2}{(R+r)^2} R = P(R)$$

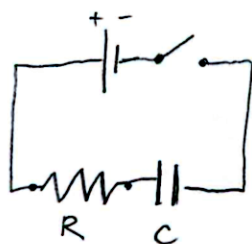
$$\frac{dP}{dR} = \varepsilon^2 \left[ \frac{1}{(R+r)^2} - \frac{2R}{(R+r)^3} \right] = 0 \Rightarrow 2R = R+r$$

$$\boxed{R = r}$$

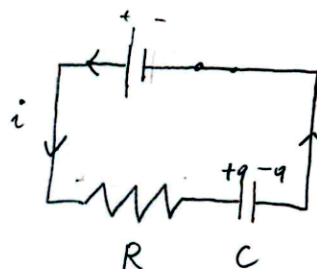
(2) Wheatstone bridge  
(see problem set)

## R-C circuits

Current dependent on time (but still in one direction)



no current



Assume ideal  
emf;  
 $\mathcal{E} = \text{const}$   
 $r = 0$

Notation: lowercase indicates time dependence  $i = i(t)$

$$q = q(t)$$

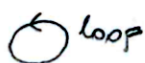
$$v_{ab} = v_{ab}(t)$$

### (1) Charging

Initial state ( $t=0$ )  
 $q(0)=0; i(0)=I_0$

K. loop law: 
$$-\underbrace{I_0 R}_{v_{ab}} + \mathcal{E} = 0 \Rightarrow \boxed{I_0 = \frac{\mathcal{E}}{R}} \quad \text{initial condition}$$

For  $t > 0$

K. loop law:  loop

$$-iR - \frac{q}{C} + \mathcal{E} = 0 \Rightarrow i = \frac{\mathcal{E}}{R} - \frac{q}{RC}$$

1<sup>st</sup> order Differential eq. with separable variables  $i = \frac{dq}{dt}$

$$\frac{dq}{dt} = \frac{\mathcal{E}}{R} - \frac{q}{RC}$$
$$\int_0^{q(t)} \frac{dq}{q - C\mathcal{E}} = - \int_0^t \frac{dt}{RC}$$

$$\ln \left| \frac{q(t) - C\mathcal{E}}{-C\mathcal{E}} \right| = - \frac{t}{RC}$$

Note  $q(t) < C\mathcal{E} = Q_{\text{max}}$

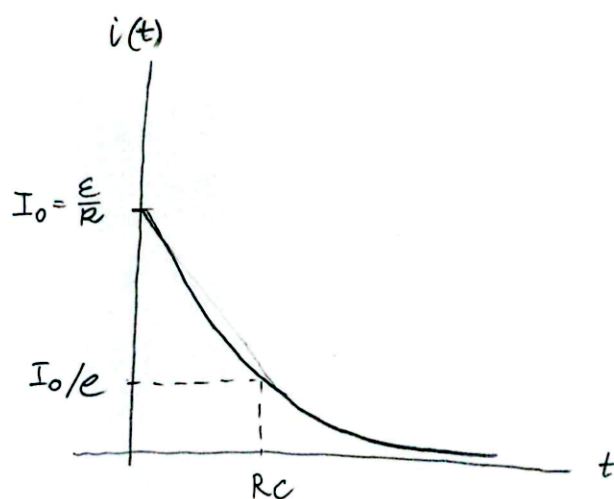
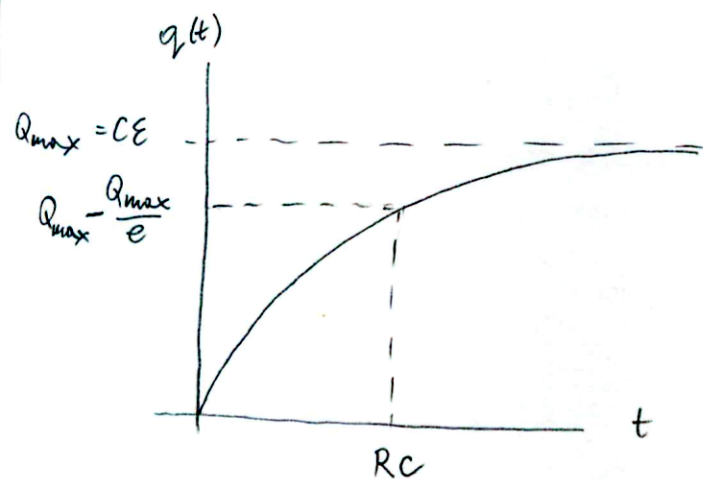
$$\ln \frac{C\mathcal{E} - q(t)}{C\mathcal{E}} = - \frac{t}{RC}$$

$$C\mathcal{E} - q(t) = C\mathcal{E} e^{-t/RC}$$

$$\boxed{q(t) = C\mathcal{E} (1 - e^{-t/RC})}$$

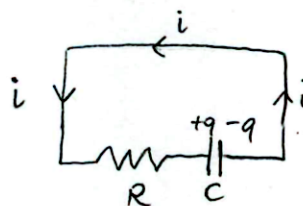
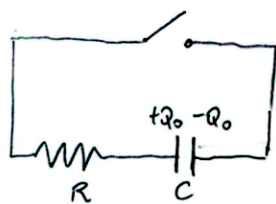
Current

$$i = \frac{dq(t)}{dt} = \frac{\mathcal{E}}{R} e^{-t/RC} = I_0 e^{-t/RC}$$



$RC = \tau$  time constant for R-C circuit [unit is]

(2) Discharging



loop rule

$$-iR - \frac{q}{C} = 0$$
$$i = -\frac{q}{RC}$$

Initial state ( $t=0$ )

$$q(0) = Q_0; \quad i(0) = -\frac{Q_0}{RC} = I_0$$

For  $t > 0$

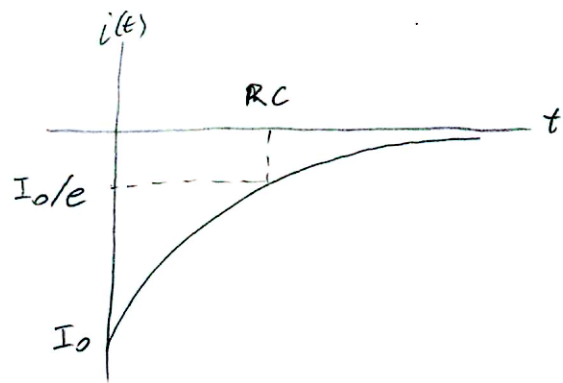
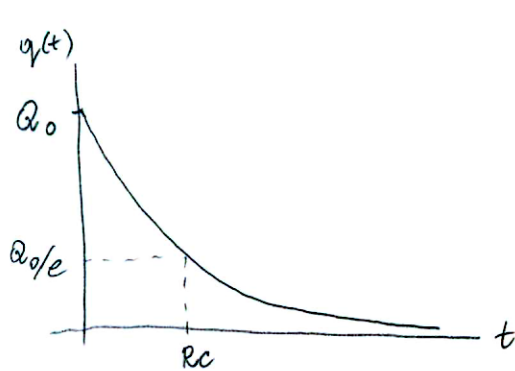
$$i = -\frac{q}{RC} \Rightarrow \frac{dq}{dt} = -\frac{q}{RC} \Rightarrow \int_{Q_0}^{q(t)} \frac{dq}{q} = -\frac{dt}{RC}$$

$$\ln \frac{q(t)}{Q_0} = -\frac{t}{RC} \Rightarrow \boxed{q(t) = Q_0 e^{-t/RC}}$$

Current

$$i = -\frac{q}{RC} = -\frac{Q_0}{RC} e^{-t/RC} = I_0 e^{-t/RC}$$





Energy in the charging process

$$-iR - \frac{q}{C} + E = 0 \quad /i$$

$$-i^2 R - \frac{i q}{C} + E i = 0$$

power dissipated  
in resistor

part stored  
in capacitor

power supplied by battery

Total energy:

\* supplied by battery

$$E_{\text{batt}} = \int_0^{\infty} E i \, dt = I_0 E \int_0^{\infty} e^{-t/RC} \, dt = I_0 R C E = \frac{I_0^2}{R} C E^2$$

\* dissipated in resistor

$$E_{\text{res}} = \int_0^{\infty} i^2 R \, dt = I_0^2 R \int_0^{\infty} e^{-2t/RC} \, dt = \frac{C E^2}{2} \quad (50\%)$$

\* stored in capacitor

$$E_{\text{cap}} = \int_0^{\infty} \frac{i q}{C} = I_0 C E \int_0^{\infty} (e^{-t/RC} - e^{-2t/RC}) \, dt = \frac{C E^2}{2} \quad (50\%)$$