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连续性方程 $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$

动量方程 $\begin{cases} x \text{ 方向: } u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) \\ y \text{ 方向: } u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + \nu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right) \end{cases}$

然后分析每个量的量纲。

$$\frac{\partial v}{\partial x} \sim \frac{\Delta u}{\Delta x} \sim \frac{U}{L}, \quad \frac{\partial u}{\partial y} \sim \frac{\Delta u}{\Delta y} \sim \frac{U}{h}$$

$$\frac{\partial^2 u}{\partial x^2} \sim \frac{\Delta u}{(\Delta x)^2} \sim \frac{U}{L^2}, \quad \frac{\partial^2 u}{\partial y^2} \sim \frac{\Delta u}{(\Delta y)^2} \sim \frac{U}{h^2}$$

而 $h \ll L$, 故只考虑 $\frac{\partial^2 u}{\partial y^2}$, 在第一个动量方程中。

$$\frac{1}{\rho} \frac{\partial p}{\partial x} = \nu \frac{\partial^2 u}{\partial y^2}$$

$$\frac{\partial^2 u}{\partial y^2} = \frac{1}{\rho \nu} \frac{\partial p}{\partial x}$$

积分两次可得 $u = \frac{1}{2\mu} \frac{\partial p}{\partial x} y^2 + cy + D$

边界条件 $u(y=0) = 0$,

$$u(y=h(x)) = 0.$$

$$\Rightarrow u = \frac{1}{2\mu} \frac{\partial p}{\partial x} y(y-h) + U \left(1 - \frac{y}{h}\right)$$

把连续性方程改写为 $\frac{\partial u}{\partial x} = -\frac{\partial v}{\partial y}$

$$\int_0^h \frac{\partial u}{\partial x} dy + \int_0^h -\frac{\partial v}{\partial y} dy = -v(h) + v(0) = 0$$

把 u 代入这个式子。

$$\text{得 } \frac{2(h^3 \frac{\partial \bar{p}}{\partial x})}{\partial x} = 6\mu U \frac{\partial h}{\partial x}$$

$$\text{代入 } h = h_L + (h_0 - h_L) \left(1 - \frac{x}{L}\right)^2$$

$$\text{设 } p_1 = \frac{\bar{p} - p_{\infty}}{\mu U L / h_0^2} \Rightarrow \bar{p} = \frac{\mu U L}{h_0^2} p_1 + p_{\infty} \text{ 代入原方程得}$$

$$\frac{\partial \bar{p}}{\partial x} = \frac{\mu U L}{h_0^2} \frac{\partial p_1}{\partial x}$$

$$\frac{2(h^3 \frac{\mu U L}{h_0^2} \frac{\partial p_1}{\partial x})}{\partial x} = 6\mu U \frac{\partial h}{\partial x}$$

$$\text{此微分方程的边界条件是 } \bar{p}(0) = \bar{p}(L) = p_{\infty}$$

$$\Rightarrow p_1(0) = p_1(L) = 0$$

$$\Rightarrow h^3 p_1'' + 3h^2 h' p_1' = \frac{6h_0^2}{L} h'$$

$$\xi \quad x_1 = \frac{x}{L}, \quad \therefore \frac{dp_1}{dx} = \frac{dp_1}{L dx_1}, \quad \frac{d^2 p_1}{dx^2} = \frac{d^2 p_1}{L^2 dx_1^2}$$

$$\Rightarrow \frac{h^3}{L^2} \frac{dp_1}{dx_1^2} + \frac{3h^2 h'}{L} \frac{dp_1}{dx_1} = \frac{6h_0^2}{L} h'$$

对方程进行无量纲化后的新方程的

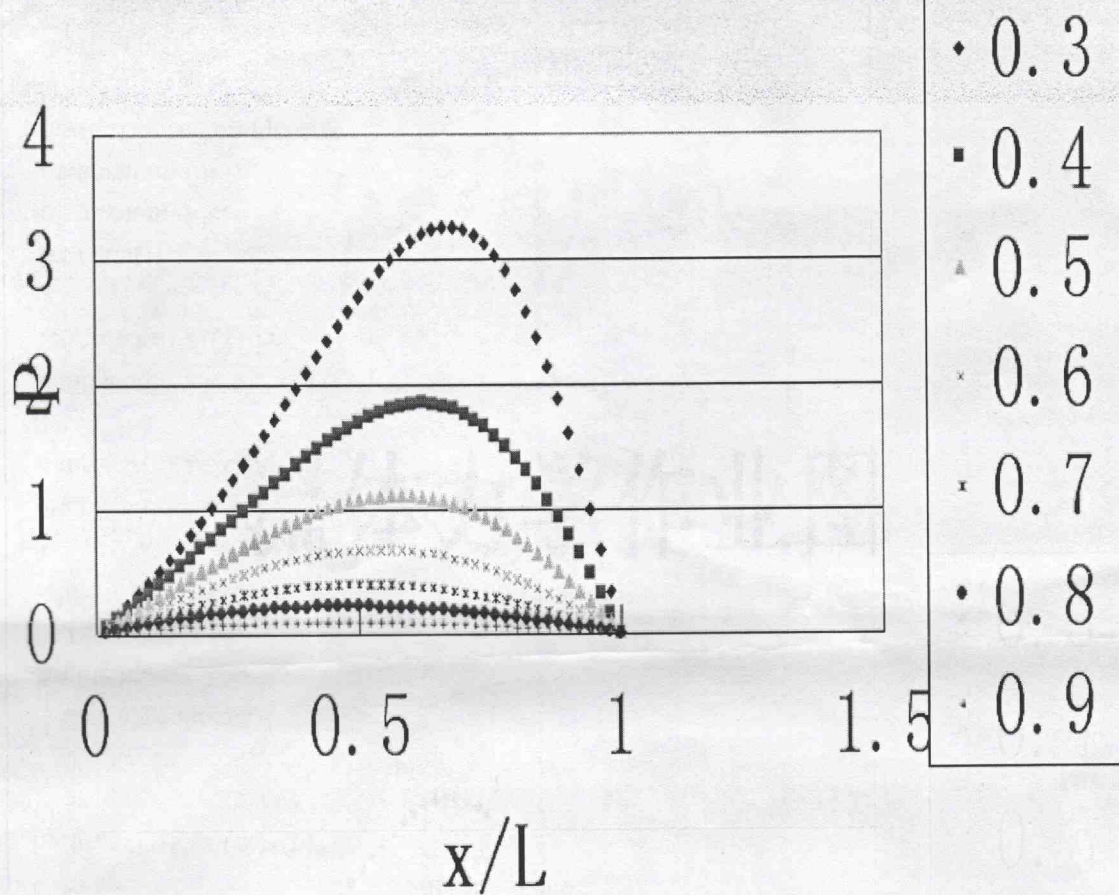
$$p_1 = \frac{\bar{p} - p_{\infty}}{\mu U L / h_0^2}, \quad x_1 = \frac{x}{L}, \quad \dots$$

$$\Rightarrow -\left[\frac{h_L}{h_0} + \left(1 - \frac{h_L}{h_0}\right)(1 - x_1)^2\right]^3 p_1''$$

$$+ 6\left[1 - \frac{h_L}{h_0}\right]\left[\frac{h_L}{h_0} + \left(1 - \frac{h_L}{h_0}\right)(1 - x_1)^2\right]^2 (1 - x_1) p_1'$$

$$= 12\left(1 - \frac{h_L}{h_0}\right)(1 - x_1), \text{ 也需对 } x_1 = 0 \text{ 及 } 1 \text{ 处 } p_1 = 0$$

流体力学作业图



差分法, $\Delta x = 0.02$ 21003079
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