

LECTURE 4

INTRODUCTION TO QUANTUM COMPUTING, THE CIRCUIT MODEL

INF587 Quantum computer science and applications

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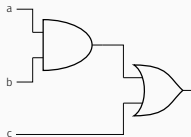
Computer science: art of computing. . .

What do we mean by **quantum computing**?

→ The **quantum circuit model**!

1. Notation and basic circuits
 - Quantum circuits: representation of unitaries and measurement
 - The quantum gate **CNOT**
 - Controlled unitaries
2. The Solovay-Kitaev theorem and the quantum gate model (universal quantum gates)
3. Simulating classical circuits with quantum circuits
4. Quantum parallelism and interference
5. A quantum algorithm: Simon's algorithm

Boolean circuit: finite directed acyclic (no loop) graph with **AND**, **OR** and **NOT** **classical** gates which has input and output nodes.



A circuit computes $f: \{0, 1\}^n \rightarrow \{0, 1\}^m$ if given n input bits x , it outputs m bits given by $f(x)$.

A circuit C_n decides a language $L \subseteq \{0, 1\}^n$ if C_n given $x \in \{0, 1\}^n$ outputs one if and only if $x \in L$.

Two questions:

- What are the **classical** gates that enable to compute any function $f: \{0, 1\}^n \rightarrow \{0, 1\}^m$?
- What class of languages circuits recognize?

Universality:

Logic gates **AND**, **OR** and **NOT** are enough to compute any function $f: \{0, 1\}^n \longrightarrow \{0, 1\}^m$.

→ Is it doable quantumly?

Problem: *any quantum operation is invertible (even unitary) but **AND** is not invertible. . .*

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Toffoli (also CCNOT) gate:

The Toffoli gate takes 3 input bits and outputs 3 bits as follows:

$$\text{Toffoli}(x, y, z) = (x, y, z \text{ XOR } (x \text{ AND } y))$$

Inversability and universality:

- The Toffoli gate is **invertible**,
- Any classical circuit computing a function f consisting of N gates in the set $\{\text{AND}, \text{OR}, \text{NOT}\}$ can be computed using $O(N)$ **Toffoli gates**.

→ In particular: the number of Toffoli gates is **roughly the same**

Many different circuits can compute a function $f : \{0, 1\}^n \rightarrow \{0, 1\}$

How can we distinguish them?

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How can we distinguish them?

→ Some circuits are more efficient than others!

Running time:

We define the running time of a circuit computing f as the number of used gates **AND**, **OR** and **NOT**

Our ideal situation: an efficient circuit

The circuit uses $O(n^k)$ gates for some constant k : we say that it has a cost $\text{poly}(n)$.

In this course: we only care of being $\text{poly}(n)$ (even if the constant k is large. . .)

Exercise:

Is it equivalent to define our running-time model as the number of **Toffoli** gates to compute a function f ? Why?

But is the classical circuit model meaningful?

Complexity theory: uniformly polynomial circuits

Family of circuits $C \stackrel{\text{def}}{=} \{C_n\}_n$ with n input bits and one output bit *such that* there is $\text{polylog}(n)$ -space Turing machine that outputs C_n given n

$$L_C \stackrel{\text{def}}{=} \bigcup_n \{x \in \{0, 1\}^n : C_n(x) = 1\}$$

P: class of languages “for which it exists an efficient **algorithm**” to decide $x \in L$ or not

$L \in P$ if and only if there exists a uniform family of circuits C such that $L = L_C$.

→ Given a uniform family of circuits $C = \{C_n\}$: C_n has at most $\text{poly}(n)$ -gates!

What about quantum computation?

Is the circuit model reasonable? If yes, what is doable quantumly and at which cost?

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Two intuitions:

- ▶ “Quantum circuit” **can simulate classical circuits** because Toffoli gates are universal and invertible. . .
 - Therefore: **quantum circuits** define a “reasonable” model of computation
- ▶ Complexity of computation will be taken into account from **the number of “quantum gates”**
 - Therefore: we expect quantum circuits to **measure the complexity** as in the classical case

NOTATION AND BASIC CIRCUITS

During this course we consider the state space $\mathbb{C}^{2^n} = \underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}}$ of n -qubits register

State space, computational basis and measurement:

We will always write n -qubits registers as

$$\sum_{\mathbf{x} \in \{0,1\}^n} \alpha_{\mathbf{x}} |\mathbf{x}\rangle \quad \text{where } |\mathbf{x}\rangle = |x_1, \dots, x_n\rangle (= |x_1\rangle \otimes \dots \otimes |x_n\rangle) \text{ and } \sum_{\mathbf{x} \in \{0,1\}^n} |\alpha_{\mathbf{x}}|^2 = 1$$

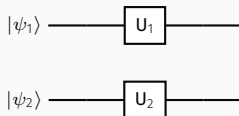
The family $(|\mathbf{x}\rangle)_{\mathbf{x} \in \{0,1\}^n}$ is known as the **computational basis**

→ All the considered measurements (in this course) will be in the computational basis

Given two quantum states $|\psi_1\rangle, |\psi_2\rangle$ and two unitaries U_1, U_2 , the circuit representation of

$$(U_1 \otimes U_2) (|\psi_1\rangle \otimes |\psi_2\rangle)$$

is given by



Exercise:

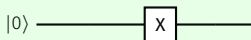
1. What becomes $\frac{|00\rangle + |01\rangle}{\sqrt{2}}$ when feeding to the above circuit?
2. Describe a quantum circuit that transforms $|00\rangle$ into $\frac{|10\rangle - |11\rangle}{\sqrt{2}}$.

Solution:

1. What becomes $\frac{|00\rangle + |01\rangle}{\sqrt{2}}$ when feeding to the above circuit?

It becomes: $U_1 |0\rangle \otimes U_2 \left(\frac{|0\rangle + |1\rangle}{\sqrt{2}} \right) = \frac{1}{\sqrt{2}} U_1 |0\rangle \otimes U_2 |0\rangle + \frac{1}{\sqrt{2}} U_1 |0\rangle \otimes U_2 |1\rangle.$

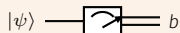
2. Describe a quantum circuit that transforms $|00\rangle$ into $\frac{|10\rangle - |11\rangle}{\sqrt{2}}.$



A measurement converts $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$ into a probabilistic classical bit $b \in \{0, 1\}$ where

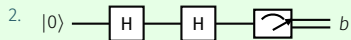
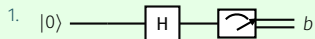
$$\mathbb{P}(b = 0) = |\alpha|^2 \quad \text{and} \quad \mathbb{P}(b = 1) = |\beta|^2$$

The circuit representation of a measurement is:



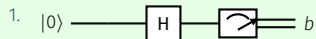
Exercise:

Give the distribution of the following probabilistic bits b :

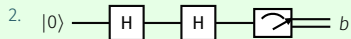


Solution:

Give the distribution of the following probabilistic bits b :



The output bit b is uniform, namely: $\mathbb{P}(b = 0) = \mathbb{P}(b = 1) = \frac{1}{2}$.



As $H^2 = I_2$, the output bit b is always zero.

THE QUANTUM CNOT GATE:

Let us introduce the Controlled-NOT gate (unitary) over 2-qubits:

$$|a, b\rangle \mapsto |a, a \oplus b\rangle$$

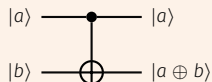
It is a unitary (it maps the computational basis to the computation basis).

Quantum CNOT-gate $|a, b\rangle \mapsto |a, a \oplus b\rangle$

- Matrix representation:

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

- Circuit representation:



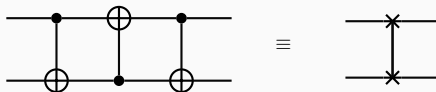
$$|a, b\rangle \mapsto |a, a \oplus b\rangle$$

is the quantum generalization of the **XOR** operation!

Be careful:

The XOR operation $(a, b) \mapsto a \oplus b$ cannot be a quantum operation **because is not invertible**

Given two wires, is it possible to **swap** two qubits?



$$\begin{aligned}
 |a, b\rangle &\longrightarrow |a, a \oplus b\rangle \\
 &\longrightarrow |a \oplus (a \oplus b), a \oplus b\rangle \\
 &\longrightarrow |b, (a \oplus b) \oplus b\rangle \\
 &= |b, a\rangle.
 \end{aligned}$$

Given a qubit $|\psi\rangle$, is it possible to build a quantum circuit that copies it?

→ **No!** Because the no-cloning theorem (see Exercise session 1)

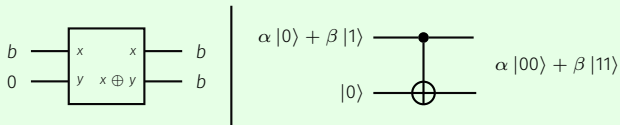
But it is doable for classical bit $(b, 0) \mapsto (b, 0 \oplus b) = (b, b) \dots$

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Take a look at the quantum case:

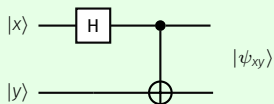


We have built an **entangled** state!

Bell states:

$$|\psi_{xy}\rangle \stackrel{\text{def}}{=} \frac{|0, y\rangle + (-1)^x |1, (1 \oplus y)\rangle}{\sqrt{2}}$$

The quantum circuit building Bell states:



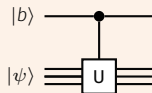
$$|x, y\rangle \xrightarrow{H \otimes I_2} \frac{|0\rangle + (-1)^x |1\rangle}{\sqrt{2}} \otimes |y\rangle = \frac{|0, y\rangle + (-1)^x |1, y\rangle}{\sqrt{2}} \xrightarrow{\text{CNOT}} \frac{|0, y\rangle + (-1)^x |1, (1 \oplus y)\rangle}{\sqrt{2}}$$

Controlled U-gate:

Let U be any unitary over n -qubits. The controlled U -gate has one control qubit $|b\rangle$ and n target qubits $|\psi\rangle$. It is defined as

- If $b = 0$, it outputs $|b\rangle \otimes |\psi\rangle$
- If $b = 1$, it outputs $|b\rangle \otimes U|\psi\rangle$

Circuit representation:



→ Controlled- $U \equiv$ If condition then instruction U

Exercise:

Show that the CNOT gate is the controlled X -gate.

QUANTUM CIRCUITS

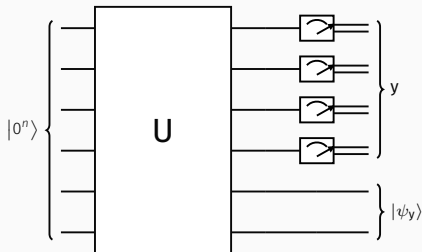
Quantum circuits: starting from n qubits initialized at $|0^n\rangle$ and then successively apply the two admissible operations (unitary and measurements).

Applying U_1 and then U_2 is equivalent to applying $(U_2 U_1)$

→ we can assume the algorithm performs a unitary, then a measurement, then a unitary, then measurement and so on. . .

We will consider only algorithms where **we first perform all the unitary operations and then perform measurements in the computational basis.**

→ As powerful as general algorithms (admitted)



$$U : |\psi\rangle \longrightarrow U |\psi\rangle$$

→ It is often easier to build $U' : |\psi\rangle |0\rangle_{\text{aux}} \longrightarrow U(|\psi\rangle) |0\rangle_{\text{aux}}$

Extra qubits are called **auxiliary qubits**, **ancillary qubits** or **workspace**

→ it is important that they start at $|0\rangle$ and end at $|0\rangle$ (see Exercise session)

SOLOVAY-KITAEV THEOREM AND GATE MODEL

Any classical function can be computed with gates {AND, OR, NOT} (**universal gates**)

What are the quantum universal gates?

The following gate is important (**first time in this course**)

The $\pi/8$ -gate:

It maps $|0\rangle \mapsto |0\rangle$ and $|1\rangle \mapsto e^{i\pi/4} |1\rangle$:

$$\mathbf{T} \stackrel{\text{def}}{=} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{pmatrix}$$

Origin of the terminology:

Up to an unimportant global phase \mathbf{T} is equal to $\mathbf{T} = e^{i\pi/8} \begin{pmatrix} e^{-i\pi/8} & 0 \\ 0 & e^{i\pi/8} \end{pmatrix}$

$\{\text{CNOT}, \text{H}, \text{T}\}$ are universal quantum gates

Solovay-Kitaev theorem (admitted):

Let $\mathcal{G} = \{\text{CNOT}, \text{H}, \text{T}\}$. Any unitary \mathbf{U} over n -qubits can be approximated by applying

$$O\left(2^{2n} \log^4\left(\frac{1}{\varepsilon}\right)\right)$$

gates from \mathcal{G} with accuracy ε .

In other words, from the description of \mathbf{U} , one can construct a sequence $\mathbf{G}_1, \dots, \mathbf{G}_N \in \mathcal{G}$ with $N = O(2^{2n} \log^4(\frac{1}{\varepsilon}))$ and

$$\|\mathbf{G}_N \dots \mathbf{G}_1 - \mathbf{U}\| \leq \varepsilon,$$

where $\|\mathbf{G}_N \dots \mathbf{G}_1 - \mathbf{U}\| \stackrel{\text{def}}{=} \max_{|\psi\rangle} \|\mathbf{G}_N \dots \mathbf{G}_1 |\psi\rangle - \mathbf{U} |\psi\rangle\|$ is the operator norm.

→ The **log** term is important: exponential accuracy with a **polynomial** number of gates

Other universal quantum gates?

Yes! The **CNOT** and qubits gates are also universal

SOLOVAY-KITAEV THEOREM: BE CAREFUL

How many resources are needed to compute a fixed unitary U over n qubits?

► **First definition:** it requires one resource, the unitary U

→ Stupid definition: same thing that saying, to compute classically **any** function f asks one resource, the function f

We want a **the smallest and simplest** set of operations to define the needed resources

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Does any unitary need an exponential number of $\{\text{CNOT}, \text{H}, \text{T}\}$ to be built?

No! As for classical computations there are algorithms/unitaries easy to compute, other not. . .

A reasonable model to define the cost of a quantum computation, i.e. computing a unitary

The number of $\{\text{CNOT}, \text{H}, \text{T}\}$ to approximate well-enough the unitary

But would you be happy to implement X or Y with this set of quantum gates?

→ *A priori no!* The set of operations $\{\text{CNOT}, \text{H}, \text{T}\}$ is not very flexible. . .

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Wouldn't be more reasonable to use as model of cost: the number of unitaries over 1 and 2-qubits?

Yes and by Solovay-Kitaev both models are “poly(λ)-equivalent”

We can approximate any unitary over 1 and 2 qubits with accuracy $2^{-\lambda}$ and

$O(\lambda^4)$ quantum gates {CNOT, H, T}

The quantum gate model:

The quantum running time of a unitary U is the amount of 1 and 2-qubit gates needed to apply U .

The running time of a single-qubit measurement is 1.

One may say that estimating the running time as the number of 1-2 qubits unitaries is an overkill
→ It can be hard to build some 1 or 2 qubits unitary. . .

A more reasonable model:

Running time: number **H**, **T** and **CNOT** gates that are used

→ The “difficulty” to implement quantum circuits reduces to **build** this small set of gates!

By the Solovay-Kitaev theorem:

The running time of the above model is the same than in the quantum gate model, **but up to polynomial factor (in the input length n) if one targets an exponentially close accuracy. . .**

In conclusion: lot of debates to define the running time of quantum circuits. . .

For us: no debates, we don't care of polynomial factors (**even if it is a hard problem to handle in “practice”. . .**) and we will use the quantum gate model

TO TAKE AWAY: YOU SAID ALGORITHM?

- ▶ Algorithm: series of simple and determined in advance instructions
→ Efficient algorithm: small number of instructions!
- ▶ Quantum algorithm: series of 1, 2-qubits unitaries and then measurements
→ Efficient quantum algorithm: small number of 1, 2-qubits unitaries and measurement!

Efficient quantum algorithm: $\text{poly}(n)$ -repetitions of a circuit starting from $|0^n\rangle$ with $\text{poly}(n)$ unitaries and measurements over 1, 2-qubits

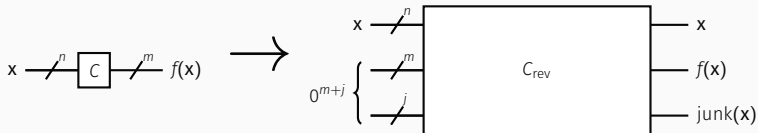
Efficient computing: a difficult task

For many problems, it is (very) hard to find a small number of instructions solving it

Shor's quantum algorithm has been a breakthrough: it solves with "few" quantum-instructions a problem (factoring) such that all known classical algorithms ask a huge number of instructions. . .

CLASSICAL CIRCUITS WITH QUANTUM CIRCUITS

Computing classically a function f with T gates can be transformed into a reversible circuit C_{rev} that only consists of $O(T)$ Toffoli gates, possibly with some junk state $\text{junk}(\mathbf{x})$.



Informally, the junk part keeps a place to perform intermediary computations

Simulating classical circuits with quantum circuits:

Classical Toffoli gates can be interpreted as a quantum unitary acting on three qubits:

$$\text{Toffoli } |x, y, z\rangle \stackrel{\text{def}}{=} |x, y, z \oplus xy\rangle$$

Therefore: C_{rev} can be interpreted as a unitary U :

$$U |x, 0^{m+j}\rangle \stackrel{\text{def}}{=} |x\rangle |f(x)\rangle |j\text{junk}(x)\rangle$$

—→ Quantum computers are at least as powerful as classical computers!

The unitary U_f :

For any function $f : \{0, 1\}^n \rightarrow \{0, 1\}^m$ that can be computed classically with a circuit that runs in time T , there exists a quantum circuit on $n + m$ qubits that runs in time $O(T)$ that can perform the unitary

$$U_f : |x\rangle |y\rangle \rightarrow |x\rangle |y \oplus f(x)\rangle .$$

Be careful:

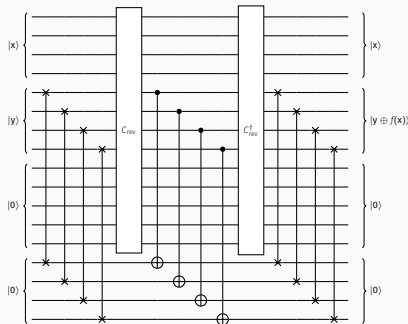
$|x\rangle \mapsto |f(x)\rangle$ may not be a quantum operation (for instance f be the zero function).

→ The auxiliary qubit $|y\rangle$ ensures that U_f is a unitary!

REMOVING THE JUNK PART AND IMPLEMENTING u_f

Proof:

1. On input $|x\rangle |y\rangle |0\rangle |0\rangle$, first swap the second and fourth registers to get $|x\rangle |0\rangle |0\rangle |y\rangle$.
2. Apply C_{rev} on the 3 first registers to get the state $|x\rangle |f(x)\rangle |junk(x)\rangle |y\rangle$.
3. For i from 1 to m , apply a **CNOT** gate between the i^{th} wire of the second register and the i^{th} wire of the forth register. We then have the state $|x\rangle |f(x)\rangle |junk(x)\rangle |y \oplus f(x)\rangle$.
4. Apply C_{rev}^\dagger on the three first registers to get the state $|x\rangle |0\rangle |0\rangle |y \oplus f(x)\rangle$.
5. Swap the second and forth register to get the state $|x\rangle |y \oplus f(x)\rangle |0\rangle |0\rangle$.



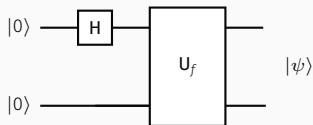
QUANTUM PARALLELISM AND INTERFERENCE

QUANTUM PARALLELISM: ONE BIT FUNCTIONS

Let $f : \{0, 1\} \rightarrow \{0, 1\}$

$$U_f : |x\rangle |y\rangle \rightarrow |x\rangle |y \oplus f(x)\rangle$$

Consider the following quantum circuit:



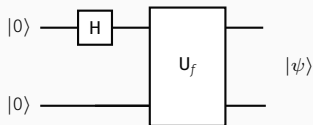
What quantum state is $|\psi\rangle$?

QUANTUM PARALLELISM: ONE BIT FUNCTIONS

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Consider the following quantum circuit:



What quantum state is $|\psi\rangle$?

1. After the first gate we have: $\frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes |0\rangle = \frac{|00\rangle + |10\rangle}{\sqrt{2}}$,
2. Applying U_f leads to (use the linearity):

$$|\psi\rangle = \frac{|0, f(0)\rangle + |1, f(1)\rangle}{\sqrt{2}}$$

→ We have a **superposition of the values of f**

Tensorization of the Hadamard gate:

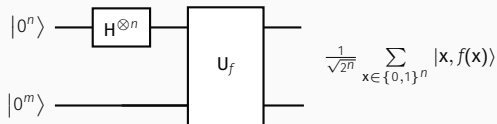
Consider,

$$H^{\otimes n} \stackrel{\text{def}}{=} \underbrace{H \otimes \dots \otimes H}_{n \text{ times}}$$

Then (see Exercise session 1),

$$H^{\otimes n} |0^n\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle$$

The following circuit performs the quantum parallelism (here $f : \{0,1\}^n \rightarrow \{0,1\}^m$)



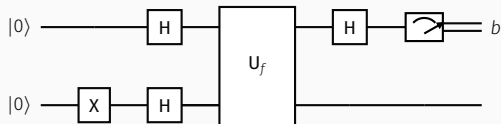
Measurement of $\frac{1}{\sqrt{2^n}} \sum_x |x, f(x)\rangle$ gives $f(x)$ for only one value of x . . .

→ **Interference** is a nice example of how using quantum parallelism!

Remember, the “−1” of the Hadamard gate gives you a huge power. . .

INTERFERENCE (DEUTSCH'S ALGORITHM)

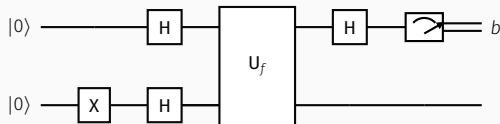
Consider the following circuit (here $f : \{0, 1\} \rightarrow \{0, 1\}$)



What is the value of b ?

INTERFERENCE (DEUTSCH'S ALGORITHM)

Consider the following circuit (here $f : \{0, 1\} \rightarrow \{0, 1\}$)



What is the value of b ?

1. After applying the **X** and **H** gates: $\frac{|0\rangle+|1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle-|1\rangle}{\sqrt{2}} = \frac{|00\rangle - |01\rangle + |10\rangle - |11\rangle}{2},$

2. Applying U_f leads to (use the linearity):

$$\frac{|0, f(0)\rangle - |0, 1 \oplus f(0)\rangle + |1, f(1)\rangle - |1, 1 \oplus f(1)\rangle}{2} = \begin{cases} \pm \frac{|0\rangle+|1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle-|1\rangle}{\sqrt{2}} & \text{if } f(0) = f(1) \\ \pm \frac{|0\rangle-|1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle+|1\rangle}{\sqrt{2}} & \text{if } f(0) \neq f(1) \end{cases}$$

3. Applying the last Hadamard gate leads to (use that $H^2 = I_2$):

$$\begin{cases} \pm |0\rangle \otimes \frac{|0\rangle-|1\rangle}{\sqrt{2}} & \text{if } f(0) = f(1) \\ \pm |1\rangle \otimes \frac{|0\rangle+|1\rangle}{\sqrt{2}} & \text{if } f(0) \neq f(1) \end{cases}$$

4. Measuring the first qubit always leads to $f(0) \oplus f(1)$!

→ We obtained a global property of f (i.e., $f(0) \oplus f(1)$) with **only one evaluation of $f(x)$!**

SIMON'S ALGORITHM

The problem:

- **Input:** A function $f : \{0, 1\}^n \rightarrow \{0, 1\}^n$.
- **Promise:** $\exists s \in \{0, 1\}^n: (f(x) = f(y) \iff (x = y) \text{ or } (x = y \oplus s))$.
- **Goal:** Find s .

1. Start from the $2n$ qubit state, with 2 registers of n qubits.

$$|\psi_0\rangle = |0^n\rangle |0^n\rangle$$

2. Apply $H^{\otimes n}$ on the first n qubits to get

$$|\psi_1\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle |0^n\rangle$$

3. Apply U_f on the state to get

$$|\psi_2\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle |f(x)\rangle = \frac{1}{\sqrt{|\text{Im}(f)|}} \sum_{y \in \text{Im}(f)} \frac{1}{\sqrt{2}} (|x_y\rangle + |x_y \oplus s\rangle) |y\rangle$$

4. Measure the second register and obtain some value $y \in \text{Im}(f)$. The resulting state on the first register is

$$|\psi_4(y)\rangle = \frac{1}{\sqrt{2}} (|x_y\rangle + |x_y \oplus s\rangle)$$

5. Apply $H^{\otimes n}$ on the first register to get

$$|\psi_5(y)\rangle = \frac{1}{\sqrt{2^n}} \sum_{z \in \{0,1\}^n} \left(\frac{1}{\sqrt{2}} (-1)^{x_y \cdot z} + \frac{1}{\sqrt{2}} (-1)^{(x_y \oplus s) \cdot z} \right) |z\rangle$$

5. Apply $H^{\otimes n}$ on the first register to get

$$|\psi_5(\mathbf{y})\rangle = \frac{1}{\sqrt{2^n}} \sum_{\mathbf{z} \in \{0,1\}^n} \left(\frac{1}{\sqrt{2}} (-1)^{\mathbf{x}_y \cdot \mathbf{z}} + \frac{1}{\sqrt{2}} (-1)^{(\mathbf{x}_y \oplus \mathbf{s}) \cdot \mathbf{z}} \right) |\mathbf{z}\rangle.$$

Now, if $\mathbf{s} \cdot \mathbf{z} = 0 \bmod 2$, we have $\left(\frac{1}{\sqrt{2}} (-1)^{\mathbf{x}_y \cdot \mathbf{z}} + \frac{1}{\sqrt{2}} (-1)^{(\mathbf{x}_y \oplus \mathbf{s}) \cdot \mathbf{z}} \right) = \sqrt{2} (-1)^{\mathbf{x}_y \cdot \mathbf{z}}$ and if $\mathbf{s} \cdot \mathbf{z} = 1 \bmod 2$, we have $\left(\frac{1}{\sqrt{2}} (-1)^{\mathbf{x}_y \cdot \mathbf{z}} + \frac{1}{\sqrt{2}} (-1)^{(\mathbf{x}_y \oplus \mathbf{s}) \cdot \mathbf{z}} \right) = 0$. Therefore, we can write

$$|\psi_5(\mathbf{y})\rangle = \sqrt{\frac{2}{2^n}} \sum_{\substack{\mathbf{z} \in \{0,1\}^n \\ \mathbf{s} \cdot \mathbf{z} = 0 \bmod 2}} (-1)^{\mathbf{x}_y \cdot \mathbf{z}} |\mathbf{z}\rangle.$$

6. Measure this state in the computational basis. You get a random \mathbf{z} satisfying $\mathbf{z} \cdot \mathbf{s} = 0$.

\mathbb{F}_2 denotes $\{0, 1\}$ modulo 2. It is a field.

\mathbb{F}_2^n is a n -dimensional \mathbb{F}_2 vector space.

$\{\mathbf{z} \in \mathbb{F}_2^n : \mathbf{z} \cdot \mathbf{s} = \sum_{i=1}^n z_i s_i = 0 \in \mathbb{F}_2\}$ is a subspace of \mathbb{F}_2^n with dimension $n - 1$

The above algorithm gives (z_1, \dots, z_n) s.t. $\sum_{i=1}^n z_i s_i = 0 \pmod 2$.

We **repeat the algorithm** m times to get m random values $\mathbf{z}^{(1)}, \dots, \mathbf{z}^{(m)} \in \mathbb{F}_2^n$ satisfying $\mathbf{z}^{(k)} \cdot \mathbf{s} = 0$

We obtain the following system (\mathbf{s} is the unknown): $\mathbf{Z}\mathbf{s} = \mathbf{0}$ where $\mathbf{Z} \stackrel{\text{def}}{=} \begin{pmatrix} \mathbf{z}_j^{(i)} \end{pmatrix}_{\substack{1 \leq i \leq m \\ 1 \leq j \leq n}}$.

→ If $\mathbf{Z} \in \mathbb{F}_2^{m \times n}$ has rank $n - 1$, we perform a Gaussian elimination to recover \mathbf{s} !

It will be verified with high probability if m large enough, $m = Cn$ for some constant $C > 0$

CONCLUSION: RUNNING TIME OF SIMON'S ALGORITHM

T be the classical running-time of f

Running time in the quantum gate model of one iteration:

- In Step 3 we apply U_f : **it can be done** by using $O(T)$ quantum gates over qubits
- In Steps 2 and 5 we apply $2n$ times H ,
- In Step 4 we perform a measurement on n -registers qubits: n measurements over qubits (in the computational basis)

One iteration:

It costs quantumly $3n + O(T)$

- We repeat $O(n)$ times an iteration: it costs $O(n^2 + nT)$
- We solve a system by a classical Gaussian elimination: it costs $O(n^3)$

Overall cost in the quantum gate model:

Simon's algorithm costs $O(n^2 + n^3 + nT) = O(n^3 + nT)$

- ▶ We have solved Simon's problem in polynomial time with high probability with **only $O(n)$ queries to f** (i.e., $O(n)$ calls to U_f , step 3)

→ Is it doable classically?

- ▶ Simon has proved that any classical randomized algorithm that finds s with high probability needs **to make $\geq C\sqrt{2^n}$ queries to f** where C constant

→ **Quantum computing provides an exponential advantage!**

There are many results about the query complexity of quantum algorithm

- ▶ *Ronald de Wolf's lecture notes, Chapters 11-12.*

<https://arxiv.org/pdf/1907.09415.pdf>

A LAST CONCLUSION

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But one may say that solving Simon's problem is useless. . .

Simon's algorithm has been “the starting point” of Shor's algorithm that quantumly breaks all current deployed public-key cryptography

→ Come at Lecture 6!

EXERCISE SESSION
