

Math 588. HW #3.

Problem 1.

One-dim-l- Neumann - boundary pr-m at $x=1$.

$$(1) \quad -\frac{d^2 p}{dx^2} = f(x), \quad 0 < x < 1; \quad x \in (0,1) =: \Omega$$

$$(2) \quad \underbrace{p(0)}_0 = \underbrace{\frac{dp}{dx}(1)}_0 = 0 \quad (3)$$

(i) Galerkin variational form?

(ii) F.E.M with piecewise linear p-m?

(iii) Linear system of equations?

Solution:

(i) Using the given PDE (ODE) in (1), & by letting $v \in V$ s.t.

$$(4) \quad L^2(\Omega) = \{v \in \Omega : \int_{\Omega} v^2 dx = \int_0^1 v^2 dx < \infty\} \text{ \& }$$

$$(5) \quad V = H^1(\Omega) = \{v \in L^2(\Omega) : \frac{\partial v}{\partial x}, \frac{\partial v}{\partial y} \in L^2(\Omega)\} \text{ \& }$$

$$(6) \quad H_0^1(\Omega) = \{v \in H^1(\Omega) : v=0 \text{ on } \partial\Omega\}, \text{ so}$$

for $v \in H_0^1(\Omega)$, let's multiply BS of (1) by $v(x)$ & integrate over the domain $\Omega=(0,1)$.

$$(7) \quad -\frac{d^2 p}{dx^2} \cdot v(x) = f(x) \cdot v(x) \quad \int_0^1 \cdot \text{ of BS.}$$

$$(8) \quad \int_0^1 \frac{d^2 p}{dx^2} \cdot v(x) dx = \int_0^1 f(x) \cdot v(x) dx$$

(1)

Applying IBP to the LHS of (8), we get

$$\begin{aligned}
 (9) \quad \int_0^1 -p''(x) v(x) dx &= - \int_0^1 v(x) d(p'(x)) \\
 &= -v(x) \cdot p'(x) \Big|_0^1 + \int_0^1 p'(x) dv(x) \\
 &= -\underbrace{v(1) \cdot p'(1)}_{\text{by (3)} = 0} + \underbrace{v(0) \cdot p'(0)}_{= 0 \text{ by (2)}} + \int_0^1 p'(x) \cdot v'(x) dx \\
 &= -v(1) \cdot 0 + 0 \cdot p'(0) + \int_0^1 p'(x) \cdot v'(x) dx \\
 &= 0 + \int_0^1 p'(x) \cdot v'(x) dx; \quad \text{so}
 \end{aligned}$$

$$(10) \quad \int_0^1 -p''(x) v(x) dx = \int_0^1 p'(x) \cdot v'(x) dx$$

plugging this back to (8), we have

$$(11) \quad \underbrace{\int_0^1 p'(x) \cdot v'(x) dx}_{:=} = \underbrace{\int_0^1 f(x) \cdot v(x) dx}_{:=} \Rightarrow$$

$$(12) \quad \langle p'(x), v'(x) \rangle = \langle f(x), v(x) \rangle \quad \text{by our}$$

class notation, where

$$\langle f, g \rangle := \int_0^1 f(x) \cdot g(x) dx.$$

This (12) is called Galerkin (weak) formulation of (1)-(3)

$$(12) \quad \boxed{\langle p'(x), v'(x) \rangle = \langle f(x), v(x) \rangle}$$

So we have (by class note) our Galerkin formulation in (12)

$$(12) \quad \langle p'(x), v(x) \rangle = \langle f(x), v(x) \rangle$$

Let K_n be the uniform partition of $\Omega = (0, 1)$ into triangles. But in our case $\Omega = (0, 1)$, a one-dim-l space, the triangles are all collapsed to small subintervals, mesh.

$$(13) \quad \Omega: \quad \begin{array}{ccccccc} & K_1 & K_2 & K_3 & & & K_n \\ | & & & & & & | \\ 0 & h & 2h & 3h & \dots & & nh & 1 \end{array}$$

and of course by our uniform triangulation,

$$(13) \quad h = \text{diam}(K_i) \quad \forall i = 1, n.$$

Now we introduce the finite element space, V_h , a discrete analog space of continuous space V .

$$(14) \quad V_h = \left\{ v: \begin{array}{l} \text{(i)} \quad v \text{ is continuous f.n on } \Omega = (0, 1) \\ \text{(ii)} \quad v \text{ is linear on each } K_i, \\ \text{(iii)} \quad v = 0 \text{ on } \partial\Omega, \text{ i.e. } \underline{v(0)=0 \text{ \& } v(1)=0.} \end{array} \right\}$$

$$V_h \subset V.$$

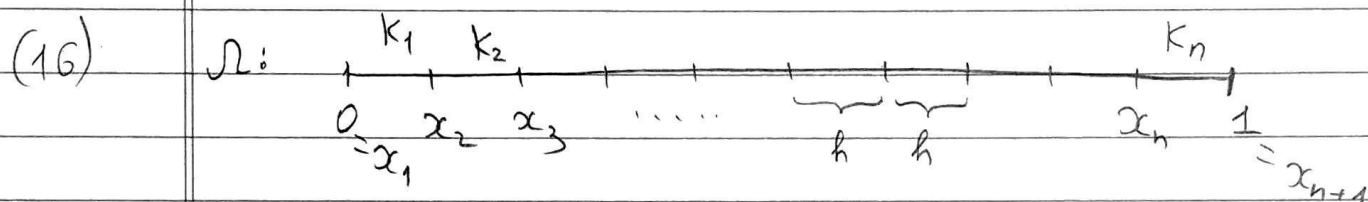
No need indeed, but doesn't hurt to have

The finite element method now can be formulated like in Dirichlet problem in HW#1.

Find $p_h \in V_h$ s.t., using (12)

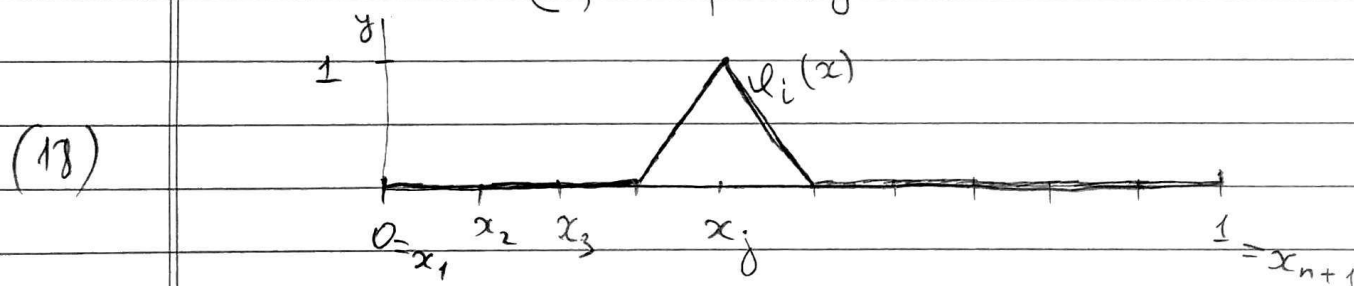
$$(15) \quad \langle p_h'(x), v'(x) \rangle = \langle f(x), v(x) \rangle, \quad \forall v \in V_h$$

Denote the vertices (nodes) of triangles (subintervals) in K by x_1, x_2, \dots, x_n , so (13) is as ↓:



The basis functions $\varphi_i \in V_h$, $i = 1, 2, \dots, n+1$

$$(17) \quad \varphi_i(x_j) = \begin{cases} 1, & \text{if } i=j \\ 0, & \text{if } i \neq j \end{cases}$$



Let x_2, x_3, \dots, x_n be the interior points (nodes) in each K .

Since $p_h(x)$ satisfies (15) $\forall v \in V_h$, then

for choices of v , we can choose easy choice which is basis f-n-1 φ_i 's i.e.

$$(19) \quad v = \varphi_1 \in V_h, \quad v = \varphi_2 \in V_h, \quad \dots \quad v = \varphi_{n+1}(x) \in V_h$$

(19) satisfies (15), then we can write the ↓:

(4)

$$\langle p_h'(x), \psi_1'(x) \rangle = \langle f(x), \psi_1(x) \rangle \text{ is valid by (15)}$$

$$\langle p_h'(x), \psi_2'(x) \rangle = \langle f(x), \psi_2(x) \rangle \text{ is valid by (15)}$$

$$\langle p_h'(x), \psi_{n+1}'(x) \rangle = \langle f(x), \psi_{n+1}(x) \rangle \text{ is valid by (15), i.e.,}$$

$$(20) \quad \langle p_h'(x), \psi_j'(x) \rangle = \langle f(x), \psi_j(x) \rangle \quad j = 1, 2, \dots, (n+1).$$

Also $p_h \in V_h = \text{span}\{\psi_1, \psi_2, \dots, \psi_{n+1}\}$, then $p_h(x)$ can be written by L.C. of basis f-ns ψ_i 's:

$$(21) \quad p_h(x) = \sum_{i=1}^{n+1} z_i \cdot \psi_i(x) \quad \text{where}$$

$$(22) \quad \boxed{z_i = p_h(x_i)} \quad \leftarrow p_h(x) \text{ is discrete, hence captures } p(x) \text{ at node points.}$$

Then, using (21) the expression of $p_h(x)$, if we consider the inner product of $p_h'(x)$ with all basis f-ns ψ_j' 's, we come to linear system of equations for unknowns z_1, z_2, \dots, z_{n+1} :

$$(23) \quad \langle p_h'(x), \psi_1'(x) \rangle \stackrel{(21)}{=} \left\langle \left[\sum_{i=1}^{n+1} z_i \psi_i(x) \right]', \psi_1'(x) \right\rangle$$

$$= \left\langle \sum_{i=1}^{n+1} z_i \psi_i'(x), \psi_1'(x) \right\rangle$$

$$= z_1 \langle \psi_1'(x), \psi_1'(x) \rangle + z_2 \langle \psi_2'(x), \psi_1'(x) \rangle + \dots + z_{n+1} \langle \psi_{n+1}'(x), \psi_1'(x) \rangle$$

Similarly

$$(24) \quad \langle p_h'(x), \psi_2'(x) \rangle = z_1 \langle \psi_1'(x), \psi_2'(x) \rangle + z_2 \langle \psi_2'(x), \psi_2'(x) \rangle + \dots + z_{n+1} \langle \psi_{n+1}'(x), \psi_2'(x) \rangle$$

And continuing this process gives

$$(25) \quad \langle p'_n(x), \varphi'_{n+1}(x) \rangle = \dots \\ = z_1 \langle \varphi'_1(x), \varphi'_{n+1}(x) \rangle + z_2 \langle \varphi'_2(x), \varphi'_{n+1}(x) \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}(x), \varphi'_{n+1}(x) \rangle$$

Yet on the other hand by (20), we have these left hand sides of (23), (24), (25) :

(26) $\langle p_n'(x), \varphi_j'(x) \rangle = \langle f(x), \varphi_j(x) \rangle$, hence
we have the \downarrow system: $j = 1, 2, \dots, (n+1).$

$$(2A) \begin{cases} z_1 \langle \varphi'_1, \varphi'_1 \rangle + z_2 \langle \varphi'_2, \varphi'_1 \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}, \varphi'_1 \rangle = \langle f, \varphi_1 \rangle \\ z_1 \langle \varphi'_1, \varphi'_2 \rangle + z_2 \langle \varphi'_2, \varphi'_2 \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}, \varphi'_2 \rangle = \langle f, \varphi_2 \rangle \\ \vdots \\ z_1 \langle \varphi'_1, \varphi'_{n+1} \rangle + z_2 \langle \varphi'_2, \varphi'_{n+1} \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}, \varphi'_{n+1} \rangle = \langle f, \varphi_{n+1} \rangle \end{cases} \Rightarrow$$

$$(28) \quad \begin{bmatrix} \langle \varphi_1', \varphi_1' \rangle & \langle \varphi_2', \varphi_1' \rangle & \dots & \langle \varphi_{n+1}', \varphi_1' \rangle \\ \langle \varphi_1', \varphi_2' \rangle & \langle \varphi_2', \varphi_2' \rangle & \dots & \langle \varphi_{n+1}', \varphi_2' \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \varphi_1', \varphi_{n+1}' \rangle & \langle \varphi_2', \varphi_{n+1}' \rangle & \dots & \langle \varphi_{n+1}', \varphi_{n+1}' \rangle \end{bmatrix} \cdot \begin{bmatrix} z_1 \\ z_2 \\ \vdots \\ z_{n+1} \end{bmatrix} = \begin{bmatrix} \langle f, \varphi_1 \rangle \\ \langle f, \varphi_2 \rangle \\ \vdots \\ \langle f, \varphi_{n+1} \rangle \end{bmatrix}$$

(29) $A \cdot \bar{z} = \bar{f}$ ← is the required matrix eq-n for unknown \bar{z} .

where $A = [a_{ij}]_{i,j=1}^{n+1} = \langle \psi'_j, \psi'_i \rangle$;

$$\bar{z} = [z_1, z_2 \dots z_{n+1}]^T \quad \& \quad z_i = p_h(x_i) = p(x_i) ; \quad \square$$

$$\bar{f} = [\langle f, \varphi_1 \rangle, \langle f, \varphi_2 \rangle, \dots, \langle f, \varphi_{n+1} \rangle]^T.$$

⑥

To determine the linear functions using the property in (17)

$$(30) \quad \psi_j(x) = \begin{cases} \frac{x - x_j}{h}, & x \in [x_{j-1}, x_j] \\ \frac{x_{j+1} - x}{h}, & x \in [x_j, x_{j+1}] \\ 0, & \text{otherwise} \end{cases}$$

Then $a_{ij} = \langle \psi_i', \psi_j' \rangle = 0$ if $|i - j| > 1$.

Then the stiffness matrix A in (29) is tridiagonal, symmetric & positive-definite as follows:

$$(31) \quad A = \frac{1}{h} \begin{bmatrix} 2 & -1 & 0 & 0 & \dots & 0 & 0 \\ -1 & 2 & -1 & 0 & \dots & 0 & 0 \\ 0 & -1 & 2 & -1 & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 2 & -1 \\ 0 & 0 & 0 & 0 & \dots & -1 & 2 \end{bmatrix}$$

completing the solution.

