

Homework Assignment 3

MATH 588 - Introduction to FEM

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MATH 588

Introduction to FEM

Homework assignment 3

Date assigned: March 9, 2025

Due date: **March 24 , 2025**

Problem 1

Consider a one-dimensional problem with a Neumann boundary condition at $x = 1$:

$$-\frac{d^2 p}{dx^2} = f(x), \quad 0 < x < 1$$
$$p(0) = \frac{dp}{dx}(1) = 0$$

Express this problem in a Galerkin variational formulation, formulate the finite element method using piecewise linear functions, and determine the corresponding linear system of algebraic equations for a uniform partition.

Problem 2

Carry out the derivation of the following basis function from the *Global coordinate approach* lecture:

$$\varphi_i(x) = \frac{x - x_{i-1/2}}{h_i} \left(\frac{2}{h_i} (x - x_{i-1/2}) + 1 \right), \quad x \in [x_{i-1}, x_i].$$

Submission

Submit your work as a PDF with handwritten or typed solutions.

Math 588. HW #3.

Problem 1.

One-dim-l- Neumann - boundary pr-m at $x=1$.

$$(1) \quad -\frac{d^2 p}{dx^2} = f(x), \quad 0 < x < 1; \quad x \in (0,1) =: \Omega$$

$$(2) \quad \underbrace{p(0)}_0 = \underbrace{\frac{dp}{dx}(1)}_0 = 0 \quad (3)$$

(i) Galerkin variational form?

(ii) F.E.M with piecewise linear p-m?

(iii) Linear system of equations?

Solution:

(i) Using the given PDE (ODE) in (1), & by letting $v \in V$ s.t.

$$(4) \quad L^2(\Omega) = \{v \in \Omega : \int_{\Omega} v^2 dx = \int_0^1 v^2 dx < \infty\} \text{ \& }$$

$$(5) \quad V = H^1(\Omega) = \{v \in L^2(\Omega) : \frac{\partial v}{\partial x}, \frac{\partial v}{\partial y} \in L^2(\Omega)\} \text{ \& }$$

$$(6) \quad H_0^1(\Omega) = \{v \in H^1(\Omega) : v=0 \text{ on } \partial\Omega\}, \text{ so}$$

for $v \in H_0^1(\Omega)$, let's multiply BS of (1) by $v(x)$ & integrate over the domain $\Omega=(0,1)$.

$$(7) \quad -\frac{d^2 p}{dx^2} \cdot v(x) = f(x) \cdot v(x) \quad \int_0^1 \cdot \text{ of BS.}$$

$$(8) \quad \int_0^1 \frac{d^2 p}{dx^2} \cdot v(x) dx = \int_0^1 f(x) \cdot v(x) dx$$

(1)

Applying IBP to the LHS of (8), we get

$$\begin{aligned}
 (9) \quad \int_0^1 -p''(x) v(x) dx &= - \int_0^1 v(x) d(p'(x)) \\
 &= -v(x) \cdot p'(x) \Big|_0^1 + \int_0^1 p'(x) dv(x) \\
 &= -\underbrace{v(1) \cdot p'(1)}_{\text{by (3)} = 0} + \underbrace{v(0) \cdot p'(0)}_{= 0 \text{ by (2)}} + \int_0^1 p'(x) \cdot v'(x) dx \\
 &= -v(1) \cdot 0 + 0 \cdot p'(0) + \int_0^1 p'(x) \cdot v'(x) dx \\
 &= 0 + \int_0^1 p'(x) \cdot v'(x) dx; \quad \text{so}
 \end{aligned}$$

$$(10) \quad \int_0^1 -p''(x) v(x) dx = \int_0^1 p'(x) \cdot v'(x) dx$$

plugging this back to (8), we have

$$(11) \quad \underbrace{\int_0^1 p'(x) \cdot v'(x) dx}_{:=} = \underbrace{\int_0^1 f(x) \cdot v(x) dx}_{:=} \Rightarrow$$

$$(12) \quad \langle p'(x), v'(x) \rangle = \langle f(x), v(x) \rangle \quad \text{by our}$$

class notation, where

$$\langle f, g \rangle := \int_0^1 f(x) \cdot g(x) dx.$$

This (12) is called Galerkin (weak) formulation of (1)-(3)

$$(12) \quad \boxed{\langle p'(x), v'(x) \rangle = \langle f(x), v(x) \rangle}$$

So we have (by class note) our Galerkin formulation in (12)

$$(12) \quad \langle p'(x), v(x) \rangle = \langle f(x), v(x) \rangle$$

Let K_n be the uniform partition of $\Omega = (0, 1)$ into triangles. But in our case $\Omega = (0, 1)$, a one-dim-l space, the triangles are all collapsed to small subintervals, mesh.

$$(13) \quad \Omega: \quad \begin{array}{ccccccc} & K_1 & K_2 & K_3 & & & K_n \\ | & \bullet & \bullet & \bullet & | & & | \\ 0 & h & 2h & 3h & \dots & & nh & 1 \end{array}$$

and of course by our uniform triangulation,

$$(13) \quad h = \text{diam}(K_i) \quad \forall i = 1, n.$$

Now we introduce the finite element space, V_h , a discrete analog space of continuous space V .

$$(14) \quad V_h = \left\{ v: \begin{array}{l} \text{(i)} \quad v \text{ is continuous f.n on } \Omega = (0, 1) \\ \text{(ii)} \quad v \text{ is linear on each } K_i, \\ \text{(iii)} \quad v = 0 \text{ on } \partial\Omega, \text{ i.e. } \underline{v(0)=0 \text{ \& } v(1)=0.} \end{array} \right\}$$

$$V_h \subset V.$$

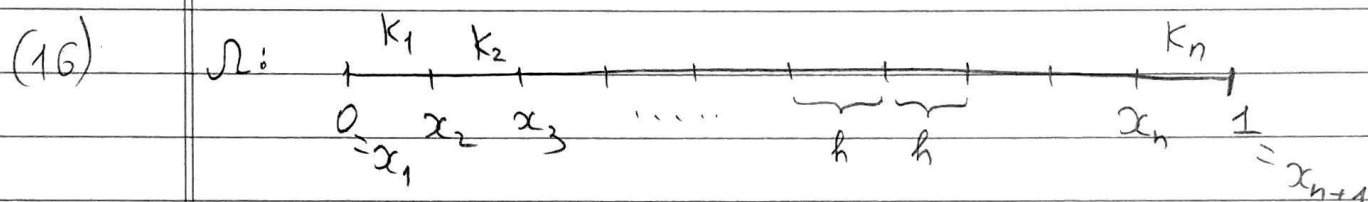
No need indeed, but doesn't hurt to have

The finite element method now can be formulated like in Dirichlet problem in HW#1.

Find $p_h \in V_h$ s.t., using (12)

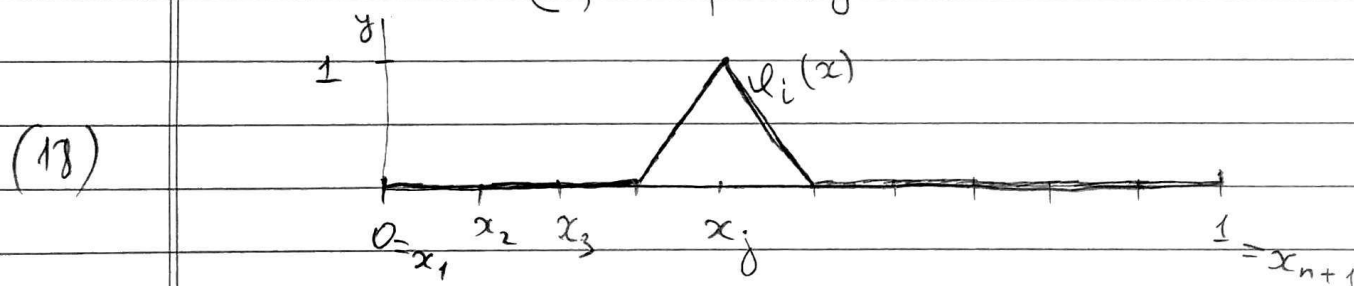
$$(15) \quad \langle p_h'(x), v'(x) \rangle = \langle f(x), v(x) \rangle, \quad \forall v \in V_h$$

Denote the vertices (nodes) of triangles (subintervals) in K by x_1, x_2, \dots, x_n , so (13) is as ↓:



The basis functions $\varphi_i \in V_h$, $i = 1, 2, \dots, n+1$

$$(17) \quad \varphi_i(x_j) = \begin{cases} 1, & \text{if } i=j \\ 0, & \text{if } i \neq j \end{cases}$$



Let x_2, x_3, \dots, x_n be the interior points (nodes) in each K .

Since $p_h(x)$ satisfies (15) $\forall v \in V_h$, then

for choices of v , we can choose easy choice which is basis f-n1 φ_i 's i.e.

$$(19) \quad v = \varphi_1 \in V_h, \quad v = \varphi_2 \in V_h, \quad \dots \quad v = \varphi_{n+1}(x) \in V_h$$

(19) satisfies (15), then we can write the ↓:

(4)

$$\langle p_h'(x), \psi_1'(x) \rangle = \langle f(x), \psi_1(x) \rangle \text{ is valid by (15)}$$

$$\langle p_h'(x), \psi_2'(x) \rangle = \langle f(x), \psi_2(x) \rangle \text{ is valid by (15)}$$

$$\langle p_h'(x), \psi_{n+1}'(x) \rangle = \langle f(x), \psi_{n+1}(x) \rangle \text{ is valid by (15), i.e.,}$$

$$(20) \quad \langle p_h'(x), \psi_j'(x) \rangle = \langle f(x), \psi_j(x) \rangle \quad j = 1, 2, \dots, (n+1).$$

Also $p_h \in V_h = \text{span}\{\psi_1, \psi_2, \dots, \psi_{n+1}\}$, then $p_h(x)$ can be written by L.C. of basis f-ns ψ_i 's:

$$(21) \quad p_h(x) = \sum_{i=1}^{n+1} z_i \cdot \psi_i(x) \quad \text{where}$$

$$(22) \quad \boxed{z_i = p_h(x_i)} \quad \leftarrow p_h(x) \text{ is discrete, hence captures } p(x) \text{ at node points.}$$

Then, using (21) the expression of $p_h(x)$, if we consider the inner product of $p_h'(x)$ with all basis f-ns ψ_j' 's, we come to linear system of equations for unknowns z_1, z_2, \dots, z_{n+1} :

$$(23) \quad \langle p_h'(x), \psi_1'(x) \rangle \stackrel{(21)}{=} \left\langle \left[\sum_{i=1}^{n+1} z_i \psi_i(x) \right]', \psi_1'(x) \right\rangle$$

$$= \left\langle \sum_{i=1}^{n+1} z_i \psi_i'(x), \psi_1'(x) \right\rangle$$

$$= z_1 \langle \psi_1'(x), \psi_1'(x) \rangle + z_2 \langle \psi_2'(x), \psi_1'(x) \rangle + \dots + z_{n+1} \langle \psi_{n+1}'(x), \psi_1'(x) \rangle$$

Similarly

$$(24) \quad \langle p_h'(x), \psi_2'(x) \rangle = z_1 \langle \psi_1'(x), \psi_2'(x) \rangle + z_2 \langle \psi_2'(x), \psi_2'(x) \rangle + \dots + z_{n+1} \langle \psi_{n+1}'(x), \psi_2'(x) \rangle$$

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$$= z_1 \langle \varphi_1'(x), \varphi_{n+1}'(x) \rangle + z_2 \langle \varphi_2'(x), \varphi_{n+1}'(x) \rangle + \dots + z_{n+1} \langle \varphi_{n+1}'(x), \varphi_{n+1}'(x) \rangle$$

Yet on the other hand by (20), we have these left hand sides of (23), (24), (25) :

(26) $\langle p_n'(x), \varphi_j'(x) \rangle = \langle f(x), \varphi_j(x) \rangle$, hence
we have the \downarrow system: $j = 1, 2, \dots, (n+1).$

$$(2A) \begin{cases} z_1 \langle \varphi'_1, \varphi'_1 \rangle + z_2 \langle \varphi'_2, \varphi'_1 \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}, \varphi'_1 \rangle = \langle f, \varphi_1 \rangle \\ z_2 \langle \varphi'_1, \varphi'_2 \rangle + z_2 \langle \varphi'_2, \varphi'_2 \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}, \varphi'_2 \rangle = \langle f, \varphi_2 \rangle \\ \vdots \\ z_1 \langle \varphi'_1, \varphi'_{n+1} \rangle + z_2 \langle \varphi'_2, \varphi'_{n+1} \rangle + \dots + z_{n+1} \langle \varphi'_{n+1}, \varphi'_{n+1} \rangle = \langle f, \varphi_{n+1} \rangle \end{cases} \Rightarrow$$

$$(28) \quad \begin{bmatrix} \langle \varphi_1', \varphi_1' \rangle & \langle \varphi_2', \varphi_1' \rangle & \dots & \langle \varphi_{n+1}', \varphi_1' \rangle \\ \langle \varphi_1', \varphi_2' \rangle & \langle \varphi_2', \varphi_2' \rangle & \dots & \langle \varphi_{n+1}', \varphi_2' \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \varphi_1', \varphi_{n+1}' \rangle & \langle \varphi_2', \varphi_{n+1}' \rangle & \dots & \langle \varphi_{n+1}', \varphi_{n+1}' \rangle \end{bmatrix} \cdot \begin{bmatrix} z_1 \\ z_2 \\ \vdots \\ z_{n+1} \end{bmatrix} = \begin{bmatrix} \langle f, \varphi_1 \rangle \\ \langle f, \varphi_2 \rangle \\ \vdots \\ \langle f, \varphi_{n+1} \rangle \end{bmatrix}$$

(29) $A \cdot \bar{z} = \bar{f}$ ← is the required matrix eq-n for unknown \bar{z} .

where $A = [a_{ij}]_{i,j=1}^{n+1} = \langle \psi'_j, \psi'_i \rangle$;

$$\bar{z} = [z_1, z_2 \dots z_{n+1}]^T \quad \& \quad z_i = p_h(x_i) = p(x_i) ; \quad \boxed{\text{shaded box}}$$

$$\bar{f} = [\langle f, \varphi_1 \rangle, \langle f, \varphi_2 \rangle, \dots, \langle f, \varphi_{n+1} \rangle]^T.$$

⑥

To determine the linear functions using the property in (17)

$$(30) \quad \psi_j(x) = \begin{cases} \frac{x - x_j}{h}, & x \in [x_{j-1}, x_j] \\ \frac{x_{j+1} - x}{h}, & x \in [x_j, x_{j+1}] \\ 0, & \text{otherwise} \end{cases}$$

Then $a_{ij} = \langle \psi_i', \psi_j' \rangle = 0$ if $|i - j| > 1$.

Then the stiffness matrix A in (29) is tridiagonal, symmetric & positive-definite as follows:

$$(31) \quad A = \frac{1}{h} \begin{bmatrix} 2 & -1 & 0 & 0 & \dots & 0 & 0 \\ -1 & 2 & -1 & 0 & \dots & 0 & 0 \\ 0 & -1 & 2 & -1 & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 2 & -1 \\ 0 & 0 & 0 & 0 & \dots & -1 & 2 \end{bmatrix}$$

completing the solution.



Problem 2.

Global coordinate approach

$$(1) \quad \varphi_i(x) = \frac{x - x_{i-\frac{1}{2}}}{h_i} \left[\frac{2}{h_i} (x - x_{i-\frac{1}{2}}) + 1 \right], \quad x \in [x_{i-1}, x_i]$$

Solution:

This approach is about finding the quadratic basis functions φ_{i-1} and φ_i on a subinterval (x_{i-1}, x_i) s.t.

$$(2) \quad \varphi_{i-1}(x_{i-1}) = 1, \quad \varphi_i(x_i) = 0, \quad \varphi_i(x_{i-1}) = 0, \quad \varphi_i(x_i) = 1,$$

a quadratic function $\varphi_{i-1}(x)$ is in the form

$$(3) \quad \varphi_{i-1}(x) = a + bx + cx^2, \quad x \in [x_{i-1}, x_i]$$

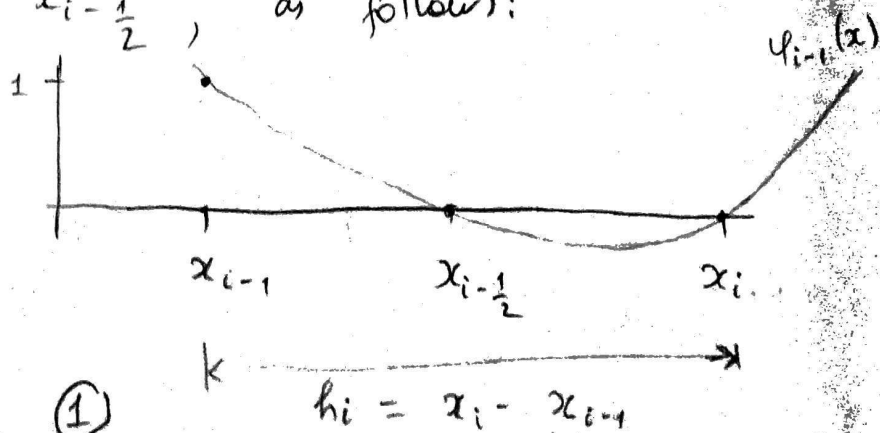
We have now two equations from (2):

$$(4) \quad \varphi_{i-1}(x_{i-1}) = 1 \Leftrightarrow a + bx_{i-1} + c(x_{i-1})^2 = 1$$

$$(5) \quad \varphi_{i-1}(x_i) = 0 \Leftrightarrow a + bx_i + c(x_i)^2 = 0$$

We need one more independent eq-n to find unknowns a, b, c . To do this, let's introduce an additional node $x_{i-\frac{1}{2}}$, as follows:

$$(6) \quad \left. \begin{aligned} h_i &:= x_i - x_{i-1} \\ x_{i-\frac{1}{2}} &:= \frac{x_{i-1} + x_i}{2} \\ x_i &= x_{i-\frac{1}{2}} + \frac{h_i}{2} \\ x_{i-1} &= x_{i-\frac{1}{2}} - \frac{h_i}{2} \end{aligned} \right\} \Rightarrow$$



So this $x_{i-\frac{1}{2}}$ point satisfies

$$(7) \quad \psi_{i-1}(x_{i-\frac{1}{2}}) = 0 \Leftrightarrow a + b x_{i-\frac{1}{2}} + c x_{i-\frac{1}{2}}^2 = 0;$$

Thus we have 3 equations, (4), (5), (7), to find unknowns a, b, c making a 3×3 system of linear equations to solve:

$$(8) \quad \begin{cases} a + b x_{i-1} + c (x_{i-1})^2 = 1 & \text{--- (4)} \\ a + b x_i + c (x_i)^2 = 0 & \text{--- (5)} \\ a + b x_{i-\frac{1}{2}} + c (x_{i-\frac{1}{2}})^2 = 0 & \text{--- (7)} \end{cases}$$

where $x_i, x_{i-1}, x_{i-\frac{1}{2}}$ are all

(4)-(5) subtraction gives:

$$\begin{aligned} b(x_{i-1} - x_i) + c[(x_{i-1})^2 - (x_i)^2] &= 1 \\ -b \cdot h_i + c \cdot (-h_i) \cdot (x_{i-1} + x_i) &= 1 \quad \} \times (-1) \end{aligned}$$

$$b h_i + c h_i \cdot (2 x_{i-\frac{1}{2}}) = -1$$

$$(9) \quad \boxed{b h_i + 2 c h_i x_{i-\frac{1}{2}} = -1}$$

Now similarly subtracting eq-n (7) from eq-n (5) yields;
(5)-(7):

$$b \cdot (x_i - x_{i-\frac{1}{2}}) + c \cdot \left[(x_i - x_{i-\frac{1}{2}}) \cdot \overset{\text{by (6)}}{(x_i + x_{i-\frac{1}{2}})} \right] = 0$$

$$b \cdot \frac{h_i}{2} + c \cdot \frac{h_i}{2} \cdot \left(x_{i-\frac{1}{2}} + \frac{h_i}{2} + x_{i-\frac{1}{2}} \right) = 0 \Rightarrow$$

(2)

$$(10) \Rightarrow \boxed{b \cdot \frac{h_i}{2} + \frac{c h_i}{2} \left(2x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = 0} \text{ by eq. (5)-(7).}$$

So by (9) & (10) we have now 2×2 system to solve for b & c :

$$(11) \begin{cases} b h_i + 2c h_i \cdot x_{i-\frac{1}{2}} = -1 \\ \frac{b h_i}{2} + \frac{c h_i}{2} \left(2x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = 0 \end{cases} \quad \text{+ 2 of BS:}$$

$$\Rightarrow \begin{cases} b h_i + c h_i \cdot (2x_{i-\frac{1}{2}}) = -1 \\ b h_i + c h_i \left(2x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = 0 \end{cases} \quad \text{eq (1) - (2):}$$

$$c h_i \cdot \left(2x_{i-\frac{1}{2}} - 2x_{i-\frac{1}{2}} - \frac{h_i}{2} \right) = -1 \Rightarrow$$

$$c = \frac{2}{h_i \cdot h_i} = \frac{2}{h_i^2};$$

$$(12) \quad \boxed{c = \frac{2}{h_i^2}}$$

Then using (10) & substituting c in (12) back into the eq-n (12), we get

$$\begin{aligned} b &= -c \left(2x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = - \frac{2}{h_i^2} \cdot \left(2x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) \\ &= - \frac{1}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right) \quad ; \quad \Rightarrow \end{aligned}$$

(3)

$$(13) \quad \boxed{b = -\frac{1}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right)} \quad \& \text{ finally for } a,$$

using eq-n (7), we have

$$\begin{aligned} a &= -b x_{i-\frac{1}{2}} - c \cdot (x_{i-\frac{1}{2}})^2 \\ &= \frac{1}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right) \cdot x_{i-\frac{1}{2}} - \frac{2}{h_i^2} \cdot (x_{i-\frac{1}{2}})^2 = \\ &= \frac{x_{i-\frac{1}{2}}}{h_i} \left[\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 - \frac{2x_{i-\frac{1}{2}}}{h_i} \right] \\ &= \frac{x_{i-\frac{1}{2}}}{h_i} \cdot \left[\frac{2x_{i-\frac{1}{2}}}{h_i} + 1 \right] ; \quad \Rightarrow \end{aligned}$$

$$(14) \quad \boxed{a = \frac{x_{i-\frac{1}{2}}}{h_i} \left[\frac{2x_{i-\frac{1}{2}}}{h_i} + 1 \right]}, \quad \text{and putting}$$

these, (12), (13), (14) back to (3), we have

$$\begin{aligned} (15) \quad u_{i-1}(x) &= a + bx + cx^2 \\ &= \frac{x_{i-\frac{1}{2}}}{h_i} \left[\frac{2x_{i-\frac{1}{2}}}{h_i} + 1 \right] - \frac{x}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right) + \frac{2x^2}{h_i^2} \\ &= \frac{x - x_{i-\frac{1}{2}}}{h_i} \cdot \left[\frac{2}{h_i} (x - x_{i-\frac{1}{2}}) - 1 \right], \quad x \in [x_{i-1}, x_i], \end{aligned}$$

as desired. 