

Homework Assignment 3

MATH 588 - Introduction to FEM

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MATH 588

Introduction to FEM

Homework assignment 3

Date assigned: March 9, 2025

Due date: March 24 , 2025

Problem 1

Consider a one-dimensional problem with a Neumann boundary condition at $x = 1$:

$$\begin{aligned}-\frac{d^2p}{dx^2} &= f(x), \quad 0 < x < 1 \\ p(0) &= \frac{dp}{dx}(1) = 0\end{aligned}$$

Express this problem in a Galerkin variational formulation, formulate the finite element method using piecewise linear functions, and determine the corresponding linear system of algebraic equations for a uniform partition.

Problem 2

Carry out the derivation of the following basis function from the *Global coordinate approach* lecture:

$$\varphi_i(x) = \frac{x - x_{i-1/2}}{h_i} \left(\frac{2}{h_i} (x - x_{i-1/2}) + 1 \right), \quad x \in [x_{i-1}, x_i].$$

Submission

Submit your work as a PDF with handwritten or typed solutions.

Math 588. HW #3.

Problem 1.

One-dim-l- Neumann-boundary pr-m at $x=1$.

$$(1) \quad -\frac{d^2 p}{dx^2} = f(x), \quad 0 < x < 1; \quad x \in (0,1) =: \mathcal{R}$$

$$(2) \quad \underbrace{p(0)}_0 = \underbrace{\frac{dp}{dx}(1)}_0 = 0 \quad (3)$$

(i) Galerkin variational form?

(ii) F.E.M with piecewise linear p-m?

(iii) Linear system of equations?

Solution:

(i) Using the given PDE (ODE) in (1), & by letting
 $v \in V$ s.t.

$$(4) \quad L^2(\mathcal{R}) = \{v \in \mathcal{R} : \int_{\mathcal{R}} v^2 dx = \int_0^1 v^2 dx < \infty\} \quad \&$$

$$(5) \quad V = H^1(\mathcal{R}) = \{v \in L^2(\mathcal{R}) : \frac{\partial v}{\partial x}, \frac{\partial v}{\partial y} \in L^2(\mathcal{R})\} \quad \&$$

$$(6) \quad H_0^1(\mathcal{R}) = \{v \in H^1(\mathcal{R}) : v = 0 \text{ on } \partial \mathcal{R}\}, \text{ so}$$

for $\psi \in H_0^1(\mathcal{R})$, let's multiply BS of (1) by $\psi(x)$ &
integrate over the domain $\mathcal{R} = (0,1)$.

$$(7) \quad -\frac{d^2 p}{dx^2} \cdot \psi(x) = f(x) \cdot \psi(x) \quad \left\{ \int_0^1 \cdot \text{ of BS.} \right.$$

$$(8) \quad \int_0^1 -\frac{d^2 p}{dx^2} \cdot \psi(x) dx = \int_0^1 f(x) \cdot \psi(x) dx$$

①

Applying IBP to the LHS of (8), we get

$$\begin{aligned}
 (9) \quad \int_0^1 -p''(x) v(x) dx &= - \int_0^1 v(x) d(p'(x)) \\
 &= -v(x) \cdot p'(x) \Big|_0^1 + \int_0^1 p'(x) dv(x) \\
 &= -v(1) \cdot \underbrace{p'(1)}_{\text{by (3)} = 0} + \underbrace{v(0) \cdot p'(0)}_{= 0 \text{ by (2)}} + \int_0^1 p'(x) \cdot v'(x) dx \\
 &= -v(1) \cdot 0 + 0 \cdot p'(0) + \int_0^1 p'(x) \cdot v'(x) dx \\
 &= 0 + \int_0^1 p'(x) \cdot v'(x) dx \quad \text{so}
 \end{aligned}$$

$$(10) \quad \int_0^1 -p''(x) v(x) dx = \int_0^1 p'(x) \cdot v'(x) dx$$

plugging this back to (8), we have

$$(11) \quad \int_0^1 p'(x) \cdot v'(x) dx = \int_0^1 f(x) \cdot v(x) dx \Rightarrow$$

$$(12) \quad : \quad \langle p'(x), v'(x) \rangle = \langle f(x), v(x) \rangle \quad \text{by our}$$

class notation, where

$$\langle f, g \rangle := \int_0^1 f(x) \cdot g(x) dx.$$

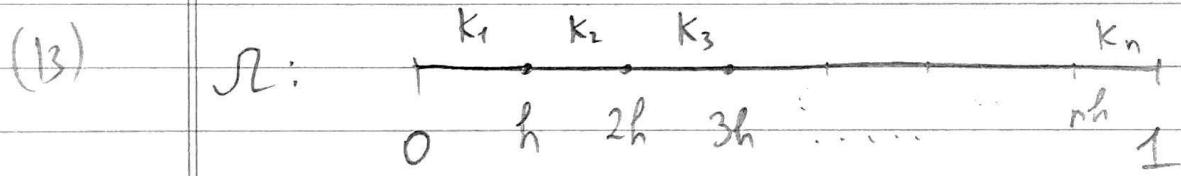
This (12) is called Galerkin (weak) formulation of (1)-(3)

$$(12) \quad \boxed{\langle p'(x), v'(x) \rangle = \langle f(x), v(x) \rangle}$$

So we have (by c Pao's note) our Galerkin formulation in (12)

$$(12) \quad \langle p'(x), v(x) \rangle = \langle f(x), v(x) \rangle$$

Let K_h be the uniform partition of $\Omega = (0, 1)$ into triangles. But in our case $\Omega = (0, 1)$, a one-dim'l space, the triangles are all collapsed to small subintervals, mesh.



and of course by our uniform triangulation,

$$(13) \quad h = \text{diam}(k_i) \quad \forall i = 1, n.$$

Now we introduce the finite element space, V_h , a discrete analog space of continuous space V .

(14)

$$V_h = \{v: v \text{ is } \underline{\text{continuous}} \text{ fn on } \Omega = (0, 1)$$

- (i) v is linear on each k_i ,
- (ii) $v = 0$ on $\partial\Omega$, i.e. $v(0) = 0 \text{ and } v(1) = 0$.

$$V_h \subset V.$$

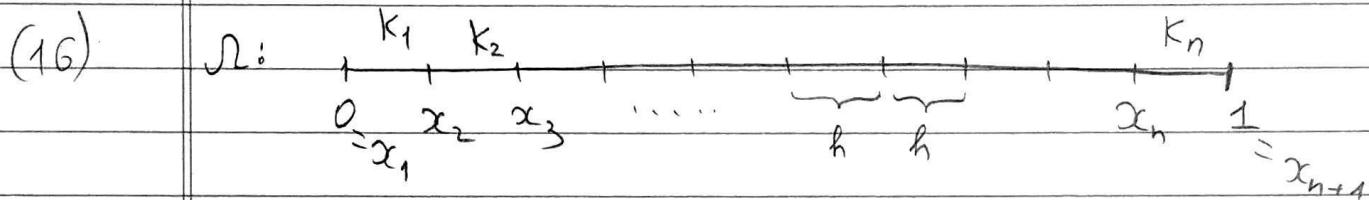
No need indeed, but
doesn't hurt to have

The finite element method now can be formulated like in Dirichlet problem in HW#1.

Find $p_h \in V_h$ s.t., using (12)

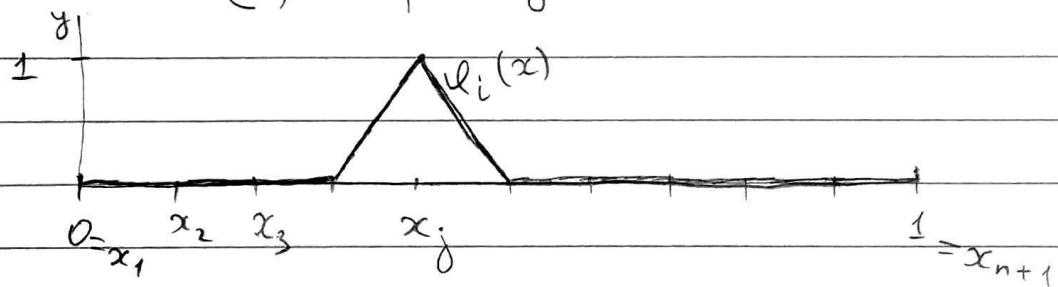
$$(15) \quad \langle p'_h(x), v'(x) \rangle = \langle f(x), v(x) \rangle, \quad \forall v \in V_h$$

Denote the vertices (nodes) of triangles (subintervals) in K by x_1, x_2, \dots, x_n , so (13) is as↓:



The basis functions $\varphi_i \in V_h$, $i = 1, 2, \dots, n+1$

$$(17) \quad \varphi_i(x_j) = \begin{cases} 1, & \text{if } i=j \\ 0, & \text{if } i \neq j \end{cases}$$



Let x_2, x_3, \dots, x_n be the interior points (nodes) in each K .

Since $p_h(x)$ satisfies (15) $\forall v \in V_h$, then

for choices of v , we can choose easy choice which
is basis $f-n$ φ_i 's i.e.

$$(19) \quad v = \varphi_1 \in V_h, \quad v = \varphi_2 \in V_h, \dots, \quad v = \varphi_{n+1}(x) \in V_h$$

(19) satisfies (15), then we can write the ↓:

$$\langle p_h'(x), \varphi_1'(x) \rangle = \langle f(x), \varphi_1(x) \rangle \text{ is valid by (15)}$$

$$\langle p_h'(x), \varphi_2'(x) \rangle = \langle f(x), \varphi_2(x) \rangle \text{ is valid by (15)}$$

$$\langle p_h'(x), \varphi_{n+1}'(x) \rangle = \langle f(x), \varphi_{n+1}(x) \rangle \text{ is valid by (15), i.e.}$$

$$(20) \quad \langle p_h'(x), \varphi_j'(x) \rangle = \langle f(x), \varphi_j(x) \rangle \quad j = 1, 2, \dots, (n+1).$$

Also $p_h \in V_h = \text{span} \{ \varphi_1, \varphi_2, \dots, \varphi_{n+1} \}$, then

$p_h(x)$ can be written by L.C. of basis f-ns φ_i 's:

$$(21) \quad p_h(x) = \sum_{i=1}^{n+1} z_i \cdot \varphi_i(x) \quad \text{where}$$

$$(22) \quad \boxed{z_i = p_h(x_i)} \leftarrow p_h(x) \text{ is discrete, hence captures } p(x) \text{ at node points.}$$

Then, using (21) the expression of $p_h(x)$, if we consider the inner product of $p_h'(x)$ with all basis f-ns φ_j' 's, we come to linear system of equations for unknowns z_1, z_2, \dots, z_{n+1} :

$$(23) \quad \begin{aligned} \langle p_h'(x), \varphi_1'(x) \rangle &\stackrel{(21)}{=} \left\langle \left[\sum_{i=1}^{n+1} z_i \cdot \varphi_i(x) \right]', \varphi_1'(x) \right\rangle \\ &= \left\langle \sum_{i=1}^{n+1} z_i \varphi_i'(x), \varphi_1'(x) \right\rangle \\ &= z_1 \langle \varphi_1'(x), \varphi_1'(x) \rangle + z_2 \langle \varphi_2'(x), \varphi_1'(x) \rangle + \dots + z_{n+1} \langle \varphi_{n+1}'(x), \varphi_1'(x) \rangle \end{aligned}$$

Similarly

$$(24) \quad \begin{aligned} \langle p_h'(x), \varphi_2'(x) \rangle &= z_1 \langle \varphi_1'(x), \varphi_2'(x) \rangle + z_2 \langle \varphi_2'(x), \varphi_2'(x) \rangle + \\ &\quad + z_{n+1} \langle \varphi_{n+1}'(x), \varphi_2'(x) \rangle \end{aligned}$$

And continuing this process gives

$$(25) \quad \langle P_h'(x), \varphi_{n+1}'(x) \rangle = \\ = z_1 \langle \varphi_1'(x), \varphi_{n+1}'(x) \rangle + z_2 \langle \varphi_2'(x), \varphi_{n+1}'(x) \rangle + \dots + z_{n+1} \langle \varphi_{n+1}'(x), \varphi_{n+1}'(x) \rangle$$

Yet on the other hand by (20), we have these left hand sides of (23), (24), (25) :

$$(26) \quad \langle P_h'(x), \varphi_j'(x) \rangle = \langle f(x), \varphi_j'(x) \rangle, \text{ hence} \\ \text{we have the } \downarrow \text{ system:} \quad j = 1, 2, \dots, (n+1).$$

$$(27) \quad \left\{ \begin{array}{l} z_1 \langle \varphi_1', \varphi_1' \rangle + z_2 \langle \varphi_2', \varphi_1' \rangle + \dots + z_{n+1} \langle \varphi_{n+1}', \varphi_1' \rangle = \langle f, \varphi_1 \rangle \\ z_1 \langle \varphi_1', \varphi_2' \rangle + z_2 \langle \varphi_2', \varphi_2' \rangle + \dots + z_{n+1} \langle \varphi_{n+1}', \varphi_2' \rangle = \langle f, \varphi_2 \rangle \\ \vdots \\ z_1 \langle \varphi_1', \varphi_{n+1}' \rangle + z_2 \langle \varphi_2', \varphi_{n+1}' \rangle + \dots + z_{n+1} \langle \varphi_{n+1}', \varphi_{n+1}' \rangle = \langle f, \varphi_{n+1} \rangle \end{array} \right. \Rightarrow$$

$$(28) \quad \underbrace{\begin{bmatrix} \langle \varphi_1', \varphi_1' \rangle & \langle \varphi_2', \varphi_1' \rangle & \dots & \langle \varphi_{n+1}', \varphi_1' \rangle \\ \langle \varphi_1', \varphi_2' \rangle & \langle \varphi_2', \varphi_2' \rangle & \dots & \langle \varphi_{n+1}', \varphi_2' \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \varphi_1', \varphi_{n+1}' \rangle & \langle \varphi_2', \varphi_{n+1}' \rangle & \dots & \langle \varphi_{n+1}', \varphi_{n+1}' \rangle \end{bmatrix}}_{A} \cdot \underbrace{\begin{bmatrix} z_1 \\ z_2 \\ \vdots \\ z_{n+1} \end{bmatrix}}_{\bar{z}} = \underbrace{\begin{bmatrix} \langle f, \varphi_1 \rangle \\ \langle f, \varphi_2 \rangle \\ \vdots \\ \langle f, \varphi_{n+1} \rangle \end{bmatrix}}_{\bar{f}}$$

$$(29) \quad A \cdot \bar{z} = \bar{f} \leftarrow \text{is the required matrix eq-n for unknown } \bar{z}. \\ \text{where } A = [a_{ij}]_{i,j=1}^{n+1} = \langle \varphi_j', \varphi_i' \rangle ;$$

$$\bar{z} = [z_1, z_2, \dots, z_{n+1}]^T \text{ & } z_i = p_h(x_i) = p(x_i) ; \quad \square$$

$$\bar{f} = [\langle f, \varphi_1 \rangle, \langle f, \varphi_2 \rangle, \dots, \langle f, \varphi_{n+1} \rangle]^T.$$

To determine the linear functions using the property in (17)

$$(30) \quad \varphi_j(x) = \begin{cases} \frac{x - x_j}{h}, & x \in [x_{j-1}, x_j] \\ \frac{x_{j+1} - x}{h}, & x \in [x_j, x_{j+1}] \\ 0, & \text{otherwise} \end{cases}$$

Then $a_{ij} = \langle \varphi_i', \varphi_j' \rangle = 0$ if $|i-j| > 0$.

Then the stiffness matrix A in (29) is tridiagonal, symmetric & positive definite as follows:

$$(31) \quad A = \frac{1}{h} \begin{bmatrix} 2 & -1 & 0 & 0 & \dots & 0 & 0 \\ -1 & 2 & -1 & 0 & \dots & 0 & 0 \\ 0 & -1 & 2 & -1 & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 2 & -1 \\ 0 & 0 & 0 & 0 & \dots & -1 & 2 \end{bmatrix}$$

completing the solution.



Problem 2.

Global coordinate approach

$$(1) \quad \varphi_i(x) = \frac{x - x_{i-\frac{1}{2}}}{h_i} \left[\frac{2}{h_i} (x - x_{i-\frac{1}{2}}) + 1 \right], \quad x \in [x_{i-1}, x_i]$$

Solution:

This approach is about finding the quadratic basis functions φ_{i-1} and φ_i on a subinterval (x_{i-1}, x_i) s.t.

$$(2) \quad \varphi_{i-1}(x_{i-1}) = 1, \quad \varphi_i(x_i) = 0, \quad \varphi_i(x_{i-1}) = 0, \quad \varphi_i(x_i) = 1,$$

a quadratic function $\varphi_{i-1}(x)$ is in the form

$$(3) \quad \varphi_{i-1}(x) = a + bx + cx^2, \quad x \in [x_{i-1}, x_i]$$

We have now two equations from (2):

$$(1) \quad \varphi_{i-1}(x_{i-1}) = 1 \Leftrightarrow a + bx_{i-1} + c(x_{i-1})^2 = 1$$

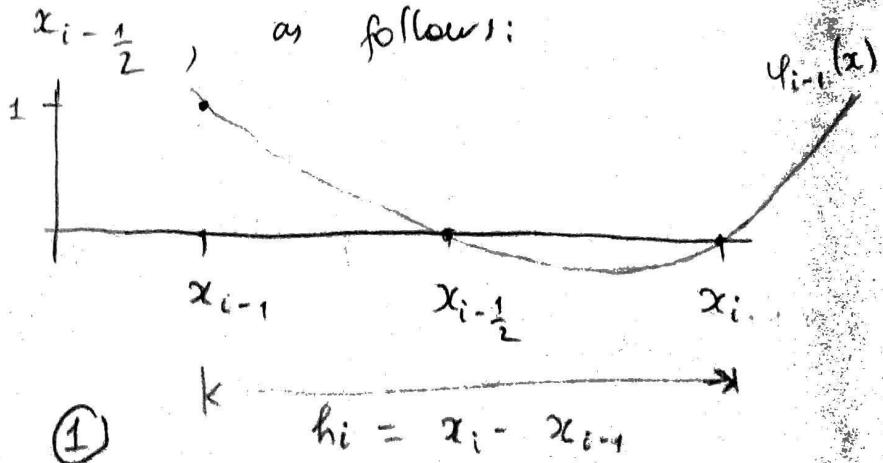
$$(5) \quad \varphi_{i-1}(x_i) = 0 \Leftrightarrow a + bx_i + c(x_i)^2 = 0$$

We need one more independent eq-n to find unknowns a, b, c . To do this, let's introduce an additional node

$x_{i-\frac{1}{2}}$, as follows:

(6)

$$\boxed{\begin{aligned} h_i &:= x_i - x_{i-1} \\ x_{i-\frac{1}{2}} &:= \frac{x_{i-1} + x_i}{2} \Rightarrow \\ x_i &= x_{i-\frac{1}{2}} + \frac{h_i}{2} \quad \& \\ x_{i-1} &= x_{i-\frac{1}{2}} - \frac{h_i}{2} \end{aligned}}$$



So this $x_{i-\frac{1}{2}}$ point satisfies

$$(7) \quad \psi_{i-1}(x_{i-\frac{1}{2}}) = 0 \Leftrightarrow a + bx_{i-\frac{1}{2}} + cx_{i-\frac{1}{2}}^2 = 0;$$

Thus we have 3 equations, (4), (5), (7), to find unknowns a, b, c making a 3×3 system of linear equations to solve:

$$(8) \quad \begin{cases} a + bx_{i-1} + c(x_{i-1})^2 = 1 & \dots \dots \dots (4) \\ a + bx_i + c(x_i)^2 = 0 & \dots \dots \dots (5) \\ a + bx_{i-\frac{1}{2}} + c(x_{i-\frac{1}{2}})^2 = 0 & \dots \dots \dots (7) \end{cases}$$

where $x_i, x_{i-1}, x_{i-\frac{1}{2}}$ are all

(4)-(5) subtraction gives :

$$\begin{aligned} b(\underbrace{x_{i-1} - x_i}_{-h_i}) + c[(x_{i-1})^2 - (x_i)^2] &= 1 \\ -b \cdot h_i + c \cdot (-h_i) \cdot \underbrace{(x_{i-1} + x_i)}_{2x_{i-\frac{1}{2}}} &= 1 \quad \{ +(-1) \end{aligned}$$

$$b h_i + c h_i \cdot (2x_{i-\frac{1}{2}}) = -1$$

$$(9) \quad \boxed{b h_i + 2c h_i x_{i-\frac{1}{2}} = -1}$$

Now similarly subtracting eq-n(7) from eq-n(5) yields;

(5)-(7) :

$$\begin{aligned} b \cdot (\underbrace{x_i - x_{i-\frac{1}{2}}}_{\frac{h_i}{2}}) + c \cdot \underbrace{[(x_i - x_{i-\frac{1}{2}}) \cdot (x_i + x_{i-\frac{1}{2}})]}_{\stackrel{\text{by (6)}}{\downarrow}} &= 0 \\ b \cdot \frac{h_i}{2} + c \cdot \frac{h_i}{2} \cdot \left(x_{i-\frac{1}{2}} + \frac{h_i}{2} + x_{i-\frac{1}{2}} \right) &= 0 \Rightarrow \end{aligned}$$

②

$$(10) \Rightarrow \boxed{b \cdot \frac{h_i}{2} + \frac{ch_i}{2} \left(2x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = 0} \text{ by eq. (5)-(7).}$$

So by (9) & (10) we have now 2×2 system to solve for b & c :

$$(11) \quad \begin{cases} b h_i + 2 c h_i \cdot x_{i-\frac{1}{2}} = -1 \\ \frac{b h_i}{2} + \frac{c h_i}{2} \left(2 x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = 0 \end{cases} \quad \left\{ + 2 \text{ of BS:} \right.$$

$$\Rightarrow \begin{cases} b h_i + c h_i \cdot (2 x_{i-\frac{1}{2}}) = -1 \\ b h_i + c h_i \left(2 x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = 0 \end{cases} \quad \left\{ \text{eq (1) - (2):} \right.$$

$$c h_i \cdot \left(2 x_{i-\frac{1}{2}} - 2 x_{i-\frac{1}{2}} - \frac{h_i}{2} \right) = -1 \Rightarrow$$

$$c = \frac{2}{h_i \cdot h_i} = \frac{2}{h_i^2};$$

$$(12) \quad \boxed{c = \frac{2}{h_i^2}}$$

Then using (10) & substituting c in (12) back into the eqn (12), we get

$$\begin{aligned} b &= -c \left(2 x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) = -\frac{2}{h_i^2} \cdot \left(2 x_{i-\frac{1}{2}} + \frac{h_i}{2} \right) \\ &= -\frac{1}{h_i} \left(\frac{4 x_{i-\frac{1}{2}}}{h_i} + 1 \right); \quad \Rightarrow \end{aligned}$$

(3)

(13) $\left\{ b = -\frac{1}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right) \right\}$ & finally for a,
using eqn (7), we have

$$\begin{aligned} a &= -b x_{i-\frac{1}{2}} - c \cdot (x_{i-\frac{1}{2}})^2 \\ &= \frac{1}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right) \cdot x_{i-\frac{1}{2}} - \frac{2}{h_i^2} \cdot (x_{i-\frac{1}{2}})^2 = \\ &= \frac{x_{i-\frac{1}{2}}}{h_i} \left[\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 - \frac{2x_{i-\frac{1}{2}}}{h_i} \right] \\ &= \frac{x_{i-\frac{1}{2}}}{h_i} \cdot \left[\frac{2x_{i-\frac{1}{2}}}{h_i} + 1 \right]; \quad \Rightarrow \end{aligned}$$

(14) $a = \frac{x_{i-\frac{1}{2}}}{h_i} \left[\frac{2x_{i-\frac{1}{2}}}{h_i} + 1 \right]$, and putting

these (12), (13), (14) back to (3), we have

$$\begin{aligned} (15) \quad u_{i-1}(x) &= a + bx + cx^2 \\ &= \frac{x_{i-\frac{1}{2}}}{h_i} \left[\frac{2x_{i-\frac{1}{2}}}{h_i} + 1 \right] - \frac{x}{h_i} \left(\frac{4x_{i-\frac{1}{2}}}{h_i} + 1 \right) + \frac{2x^2}{h_i^2} \\ &= \frac{x - x_{i-\frac{1}{2}}}{h_i} \cdot \left[\frac{2}{h_i} (x - x_{i-\frac{1}{2}}) - 1 \right], \quad x \in [x_{i-1}, x_i], \end{aligned}$$

as desired.

④