

# Stratiform precipitation, vertical heating profiles, and the Madden-Julian Oscillation

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## ABSTRACT

The observed profile of heating through the troposphere in the Madden-Julian Oscillation (MJO) is found to be very top-heavy: more so than seasonal-mean heating, and systematically more so than all of the seven models for which intraseasonal heating anomaly profiles have been published. Consistently, spaceborne radar shows that stratiform precipitation (known to heat the upper troposphere and cool the lower troposphere) contributes more to intraseasonal rainfall variations than it does to seasonal-mean rainfall. Heating structure in a strong MJO event observed by the TOGA COARE field program agrees with that of a composite MJO from multi-year data sets. Tilted structure is seen in time-height sections of heating, with the altitude of heating rising with time during the passage of an MJO wet phase. Diagnostic partitioning suggests that the trailing-edge heating tilt reflects enhanced stratiform anvil heating, which lags total precipitation by a few days, while leading-edge tilt reflects a change in convective heating profiles from shallower to deeper as the MJO approaches.

## 1. Introduction

Discovered by Madden and Julian (1971, 1972), the Madden-Julian Oscillation (MJO) is the dominant intraseasonal mode of variability in tropical convection and circulation (e.g. Weickmann et al. 1985, Lau and Chan 1985, Salby and Hendon 1994, Wheeler and Kiladis 1999). It affects a wide range of tropical weather such as the onset and breaks of the Indian and Australian summer monsoons (e.g. Yasunari 1979, Hendon and Liebmann 1990), and the formation of tropical cyclones (e.g. Nakazawa 1986, Liebmann et al. 1994). It also drives teleconnections to the extratropics (e.g., Lau and Phillips 1986, Winkler et al. 2001) and impacts some important extratropical weather (e.g. Higgins and Mo 1997, Higgins et al. 2000). On a longer timescale, the MJO is observed to trigger or terminate some El Nino events (e.g. Kessler et al. 1995, Takayabu et al. 1999, Bergman et al. 2001). Therefore, the MJO is important for both extended-range weather forecasting and long-term climate prediction.

The MJO is concentrated in the eastern hemisphere. Figure 1a shows the standard deviation of 30-70 day bandpass filtered precipitation (from CMAP, see section 2 and 3 for details of data and filtering). This intraseasonal rainfall variability has two heating centers, in the Indian ocean and the western Pacific. Eastward propagation is prominent, as seen in lagged correlations between a filtered time series of precipitation at 0N155E and precipitation elsewhere along the equator (Fig. 1b). The propagating part of intraseasonal variability (hereafter called the MJO) moves eastward from the Indian Ocean to the western Pacific, with a phase speed of about 5 m/s (4 degrees/day).

Eastward-propagating intraseasonal signals do occur in some general circulation models (GCMs), but are generally too weak and propagate too fast compared with observations (e.g. Hayashi and Sumi 1986, Hayashi and Golder 1986, 1988, Lau et al. 1988, Slingo et al. 1996). In light of the MJO’s importance, reviewed above, this shortcoming is detrimental to both numerical weather prediction and climate prediction. The overarching goal of this work is to contribute to improved GCM simulations of tropical variability. A first step toward that goal is to better document the observed phenomenon, in order to better define the form and (hopefully) nature of model errors.

Diabatic heating (mainly latent and radiative) is central to the structure and propagation of the MJO (Krishnamurti et al. 1985, Murakami and Nakazawa 1985, Yanai et al. 2000). The heating is strongly associated with a propagating large-scale circulation anomaly (hereafter called a ‘wave’), which in turn feeds back onto the heating (Hendon and Salby 1994, Zhang 1996). This suggests the utility of a “wave-heating feedback” theoretical framework, of which many versions have been explored, with various degrees and types of approximations.

The vertical profile of heating has a large effect on intraseasonal oscillations in such wave-heating theoretical models (discussed further in section 7). In light of this sensitivity to heating profiles, it seems useful to document the observed MJO heating profile, and compare it to the same quantity in models. Might the models’ difficulty in simulating the MJO be related to systematic error in model heating profiles?

Heating profiles over the tropical oceans have been measured as heat budget resid-

uals, using sounding-array data in field experiments such as the Marshall Islands (Yanai et al. 1973), GATE (e.g. Thompson et al. 1979), AMEX (e.g. Frank and McBride 1989), and TOGA COARE (Lin and Johnson 1996, Frank et al. 1996, Yanai et al. 2000, and Zhang and Lin 1999). The profile of total heating in convecting regions is top-heavy, or concentrated high in the upper troposphere. This top-heaviness is contributed by vertical dipole heating profiles observed in areas of stratiform precipitation, with heating in the upper troposphere and cooling in the lower troposphere (Houze 1982, 1989, 1997; Johnson 1984; Mapes and Houze 1995). Stratiform precipitation also lags convective precipitation in time, which might introduce important time-dependence to total heating profiles.

The anomalous vertical heating profile as it varies across the MJO has not been isolated in the published literature. As we shall see, it differs significantly from time-mean profiles. Likewise, the contribution of stratiform precipitation to mean precipitation has been studied (e.g. Schumacher and Houze 2003), but the contribution of anomalous stratiform precipitation to the anomalous precipitation in the MJO has not been reported. This study aims to remedy this lack, as a starting point to identify model errors.

The purpose of this study is to examine the observed anomalous heating profiles and stratiform precipitation in the MJO, consider the consistency of these completely independent measurements, and compare the observed heating profiles with those in some models. Major issues we address are:

- (1) What is the vertical heating profile in the MJO? How does it compared with

the climatological mean profile? How large does stratiform precipitation contribute to the vertical heating profile?

(2) How does the MJO heating anomaly vary with time? Is it simply a reversing sign of a fixed profile, or is there phase tilt? If so, what factors contribute to this tilt?

(3) How well do current GCMs represent the vertical heating profile in the MJO?

The datasets used in this study are described in section 2. The methods are described in section 3. The anomalous vertical heating profiles in the convectively active phase are examined in section 4. The vertical phase tilt of heating anomaly is studied in section 5. Comparisons between the observed heating profiles and the GCM profiles are conducted in section 6. Summary and discussion are given in section 7.

## 2. Data

The datasets used include TOGA COARE data and long-term data.

The TOGA COARE datasets include:

(1) 120 day (November 1, 1992 to February 28, 1993) of the 6-hourly sounding array budgets for the Intensive Flux Array (IFA, Fig. 1) calculated by Ciesielski et al. (2002). The variables we used include diabatic heating  $Q_1$ , moisture sink  $Q_2$ , wind, horizontal divergence, and relative humidity.

(2) 120 day (November 1, 1992 to February 28, 1993) of the daily diabatic heating  $Q_1$  profiles for the Outer Sounding Array (OSA, Fig. 1) calculated by Zhang and Lin (1999) using the method of Zhang and Lin (1997).

(3) 120 day (November 1, 1992 to February 28, 1993) of the hourly radiative heating

$Q_R$  profiles for the IFA Calculated by Qian and Cess (2003).

To diagnose the cumulus convective activity, we calculated the eddy moist static energy flux  $F_h$  for TOGA COARE IFA:

$$F_h(p) = \int_{100mb}^p (Q_1 - Q_2 - Q_R) dp \quad (1)$$

As shown by Yanai and Johnson (1993),  $F_h$  is only contributed by the cumulus convection, not by the stratiform precipitation.

The long term datasets include:

(1) 24 years (1979-2002) of the pentad "chi-corrected" diabatic heating  $Q_1$  profiles calculated by Sardeshmukh et al. (1999) and Winkler et al. (2001). The horizontal resolution is 2.5 degree longitude by 2.5 degree latitude. We average the data between 5N and 5S in longitudes 145-155E. These diabatic heating profiles are determined from an improved iterative solution of the "chi problem" (Sardeshmukh 1993), applied to four times daily NCEP reanalysis wind fields to minimize the nonlinear vorticity budget imbalance at 28 atmospheric levels. The modified divergence is further constrained to satisfy the large-scale mass budget. Diabatic heating rates are finally determined as a residual in the heat budget, using the modified wind circulation to compute advective terms. See Sardeshmukh (1993) for more details of the technique.

(2) 24 years (1979-2002) of NCEP reanalysis upper air wind data. The horizontal resolution is 2.5 degree longitude by 2.5 degree latitude. We average the data along the equator (between 5N and 5S) with a zonal resolution of 10 degree longitude.

(3) 24 years (1979-2002) of the pentad CPC Merged Analysis of Precipitation (CMAP) calculated by Xie and Arkin (1997). The horizontal resolution is 2.5 degree longitude by 2.5 degree latitude. We average the data along the equator (between 5N and 5S) with a zonal resolution of 5 degree longitude.

(4) 5 years (1998-2002) of the half-hourly TRMM precipitation radar (PR) gridded convective/stratiform precipitation (product number 3G68) calculated by Kummerow et al. (2000) and Stocker et al. (2001). The horizontal resolution is 0.5 degree longitude by 0.5 degree latitude. We average the data along the equator (between 5N and 5S) to pentad data with a zonal resolution of 10 degree longitude.

(5) Five years (1998-2002) of twice-daily upper air sounding data at the Atmosphere Radiation Measurement (ARM) Manus site located at 2S147E. We only use the relative humidity in this study. The data are averaged to pentad means. There are only a few missing pentads and they are filled using linear interpolation in time.

### **3. Method**

The MJO is a broad-band phenomenon, with an averaged period of 45 days but a fairly wide spread from 20 to 80 days (see review by Madden and Julian 1994). Its deep convection signal is dominated by wavenumber 1-6, while its circulation signal is dominated by wavenumber 1 (e.g. Salby and Hendon 1994, Wheeler and Kiladis 1999). These characteristics are used to isolate the MJO signal as a master time series from space-time data.

Many different indices have been used in previous observational and modeling



studies. Some are based on circulation signals (e.g. Madden and Julian 1972, Knutson and Weickmann 1987, Lau et al. 1988), others on deep convection signals (e.g. Lau and Chan 1985, Hendon and Salby 1994, among many others), while some use a combination of circulation and deep convection signals (e.g. Hsu 1996). Even in the case of a single variable (often OLR), various methods have been used to construct an MJO index, such as: (1) band-pass filtering and selection of a (usually large) grid box as the reference point (e.g. Lau and Chan 1985, Kiladis and Weickmann 1992, Zhang 1996, Woolnough et al. 2000); (2) band-pass filtering and selection of the leading Empirical Orthogonal Functions (EOFs) of the anomaly (e.g. Weickmann 1983, Lau and Chan 1985); (3) band-pass filtering and selection of strong, eastward propagating events (e.g. Rui and Wang 1990); and (4) space-time filtering to isolate only large-scale eastward propagating signals (e.g. Hendon and Salby 1994). The different methods often give qualitatively similar results in terms of propagation characteristics and phase difference among different variables (e.g. Lau and Chan (1985) compared methods (1) and (2)). This consistency is apparently because strong intraseasonally filtered OLR fluctuations along the equator are dominated, both in number and strength, by coherent eastward propagating events (Wang and Rui 1990a). Linear composite methods used in most of the previous studies (e.g. phase sum, correlation, regression) are apparently dominated by these strong eastward propagating events.

This study uses the above method (1). All datasets are filtered using a 30-70 day Murakami (1976) filter, whose response function is shown in Fig. 2. The central frequency correspond to a period of 45 days, with half amplitude at periods of 30 days

and 70 days. We have also tested the Lanczos filter (Duchan 1979), and the results are not sensitive. For the master MJO index, we use a western Pacific (0N155E) time series of CMAP precipitation filtered in time as above, but also in space, retaining wavenumbers 0-6. Experiments with more complex space-time filtering (to isolate eastward-propagating components more specifically) suggest that our main findings with respect to vertical structure are robustly indicated by this simpler technique.

The MJO structure is constructed differently for the COARE and long-term data sets. For the long term data, composite structure is developed by linear regression with respect to the CMAP MJO index. Regressions have been done for all seasons of the year, as well as for individual seasons (December-February, March-May, June-August, and September-November). The confidence level of linear correlation is estimated following Oort and Yienger (1996). In the TOGA COARE data, there are two MJO events (Fig. 3, Chen et al. 1996). We use filtered data on the December 1992 event, a strong MJO event with its amplitude at the TOGA COARE IFA larger than two standard deviation of the 24 year data (Fig. 1a). It moved eastward with a phase speed of 5 m/s, which is similar to the phase speed of the 24 year composite (Fig. 1b). The maximum of the 30-70 day bandpass filtered precipitation anomaly at the TOGA COARE location occurs on December 21, 1992.

## **4. Heating profiles and stratiform precipitation at the time of maximum precipitation**

### **4.1 Observed heating profile**

The vertical heating profiles at the time of maximum precipitation during the MJO life cycle are shown in Fig. 4 for TOGA COARE IFA, OSA and the composites for the four different seasons (December-February, March-May, June-August, and September-November) using 24 years (1979-2002) of chi-corrected heating data. The two TOGA COARE profiles look quite similar with each another, suggesting that the profile is robust with respect to different array size, and different detailed methods for calculating the sounding array budgets. The two profiles are very top-heavy, i.e., with strong heating in the upper troposphere and weak heating in the lower troposphere. The composite MJO profiles for the four different seasons using chi-corrected heating data are similar to the TOGA COARE profiles, except for being a little less top-heavy, suggesting that the TOGA COARE profiles are statistically representative. The mean chi-corrected heating profiles for the 4-month period of TOGA COARE (not shown) are even closer to the rawinsonde-derived profiles than are the long-term means in Fig.4, suggesting that the modest chi-COARE differences seen there reflect temporal variability, not error in the chi-corrected heating product. The shape of the chi-corrected heating profiles shows little seasonal variation.

The difference between the MJO anomaly profile at the time of maximum precipitation and the seasonal mean profile is illustrated in Fig. 5 for both TOGA COARE and chi-corrected heating data. For the chi-corrected heating data, the difference is calculated between the MJO composite for each season and the mean profile of the corresponding season. In TOGA COARE (Fig. 5a), the MJO anomaly profiles are significantly more top-heavy than the seasonal mean profiles, i.e., with stronger

heating in the upper troposphere and weaker heating in the lower troposphere. This is also the case for the chi-corrected heating data (Fig. 5b), especially in winter and spring, except with the upper troposphere heating differences located at a higher altitude. The difference profiles are similar to the stratiform heating profiles (Houze 1982, 1989, Johnson 1984), suggesting that the MJO may be associated with larger fraction of stratiform precipitation than the seasonal mean value. To pursue this suggestion, the next section examines the contribution of stratiform precipitation using spaceborne radar data.

## 4.2 Contribution of stratiform precipitation to the heating profile

Figure 6 shows the precipitation anomaly during the composite MJO life cycle for 5 years (1998-2002) of TRMM precipitation radar (PR) data at 0N, 155E. The thick solid line is total precipitation, the thin solid line is convective precipitation, and the thin dashed line is stratiform precipitation. At the time of maximum precipitation, the stratiform precipitation contributes about 60% of the anomalous total precipitation. The stratiform precipitation lags the convective precipitation by several days, implying a larger fraction of stratiform precipitation in the later stages (examined further in section 5).

The larger stratiform rain fraction in the MJO wet phase, compared to its climatological mean value, is true across longitude (Fig. 7). In both the western Pacific and the eastern Indian Ocean, the stratiform rain fraction in the MJO is about 0.1 larger than its climatological mean value. This finding is consistent with the observation

in Fig. 5 that the differences between MJO anomaly heating profiles and seasonal mean profiles are similar in shape to stratiform heating profiles. It is also consistent with previous findings that the MJO modulates larger mesoscale convective systems more than small systems (e.g. Mapes and Houze 1993, Chen et al. 1996), in light of the fact that larger systems have a greater stratiform rain fraction (e.g. Nesbitt et al. 2000). Interestingly, the east-west gradient in stratiform rain fraction across the Pacific (c.f. Schumacher and Houze 2003) is seen in both the MJO anomaly and mean curves in Fig. 7.

For a more quantitative consistency check between the stratiform-rain and heating-profile observations, we can formally partition the observed TOGA COARE IFA heating profile into three components: stratiform, radiative and convective, following the method of Johnson (1984):

$$\hat{Q}_1 \approx f\hat{Q}_{1m} + (1 - f)\hat{Q}_{1c} + (Q_R/P_0) \quad (2)$$

where  $Q_{1m}$  is the heating in the stratiform precipitation region,  $Q_{1c}$  the heating in the convective precipitation region,  $Q_R$  the radiative heating,  $P_0$  the area-averaged precipitation rate, and  $f$  the stratiform rain fraction. The caret refers to values of  $Q_1$  normalized by surface precipitation rate. Previous observations show good agreement in rainrate-normalized stratiform heating profiles, but diverse results in convective heating profiles (see comparisons by Houze 1989 and Tao et al. 2001). For this reason,  $Q_{1c}$  instead of  $Q_{1m}$  is derived as a residual.

For the partition, the normalized stratiform heating profile  $Q_{1m}$  is taken from

Houze (1982), while the intraseasonal  $Q_R$  wet-phase anomaly is obtained from the hourly IFA radiative heating profiles calculated by Qian and Cess (2003). The stratiform rain fraction is taken to be 0.6 from Fig. 6. This partition exercise is imperfect, since the stratiform heating profile in Houze (1982) represents the mature phase of strong mesoscale anvils, while the stratiform precipitation reported by TRMM PR includes developing, mature, and decaying phases of mesoscale anvils, as well as the stratiform precipitation not associated with mesoscale convective systems. Nevertheless, the results are adequate for qualitative use.

The partition result is shown in Fig. 8. The convective heating peaks in the middle troposphere near the  $0^\circ\text{C}$  level and only slightly contributes to the top-heaviness of the total profile. The radiative heating is characterized by cloud-top cooling above 250 mb and cloud-base warming below 250 mb, and slightly decreases the top-heaviness of the total profile. The stratiform heating, however, significantly increases the top-heaviness of the total profile. Its upper troposphere heating enhances the convective heating by more than 50%, while its lower troposphere cooling cancels more than 50% of the convective heating. Theoretically, the convective heating can be further divided into two components: net condensation and eddy flux convergence. Considering only the convective heating but not the stratiform heating, Cho and Pendlebury (1997) argued that the only process that might produce a heating maximum in the upper troposphere is the eddy flux convergence, because the net condensation generally has a maximum in the lower troposphere. Fig. 8 shows that the stratiform heating is another major contributor to the upper troposphere heating maximum, and it is more

efficient in increasing the top-heaviness because it simultaneously cools the lower troposphere.

## 5. Vertical phase tilt of heating anomaly

This section extends results from the heating profile at the time of maximum precipitation to the variations of heating profile during the whole MJO life cycle. Figure 9a shows the vertical structure of  $Q_1$  anomaly in TOGA COARE IFA for the December 1992 MJO event, versus time lag (with respect to the time of maximum precipitation, with the axis reversed so that the local time evolution at one longitude visually resembles a spatial cross-section through the eastward-moving MJO (Fig. 3, Fig. 1b)).

A subtle but distinct tilt of the heating contours in this time-height section is evident. Anomalous heating develops first in the lower troposphere near day -15 and then shifts upward as it intensifies. The maximum heating is below 800 mb before day -10, shifts upward to 450 mb at day 0, and shifts further upward to above 400 mb after day +5. Heating for the TOGA COARE OSA (Fig. 9b), shows similar vertical structure as the IFA, indicating that the vertical heating structure is robust with respect to the different array sizes and calculation methods. These COARE results appear to be statistically representative, as the long-term composite (Fig. 9c) exhibits similar tilt. Composites for other seasons have similar vertical structure as this DJF composite (not shown). To the extent that the MJO propagates past the reference longitude (155E) with nearly fixed structure, these tilted time-height

contours correspond to a westward spatial phase tilt of heating with height, as has been found in simulated intraseasonal oscillations in GCMs (Lau et al. 1988, Wang and Schlesinger 1999).

This observed rising of the heating profile with time can be attributed, based on Eq. 2, to one or more of the following three factors: (1) variation of the convective heating  $\hat{Q}_{1c}$  profile (i.e. shallower convective heating in the earlier stage and deeper convective heating in later stage); (2) variation of the stratiform heating fraction  $f$  (i.e. more stratiform heating in later stage); and (3) variation of the radiative heating  $Q_R$  profile.

Factor (2), variation of the stratiform rain fraction  $f$ , is examined in Fig. 10. The stratiform rain fraction lags total precipitation by about five days, consistent with the several-day time lag between stratiform and convective precipitation shown in Fig. 6. The variation of  $f$  is, however, quite small (only about  $\pm 0.02$ ), with  $f$  remaining about 0.6 throughout the composite MJO time series (cf Figs. 6, 7).

The large-scale causes for enhanced stratiform precipitation lagging behind maximum precipitation appear to include wind shear and middle- to upper-troposphere relative humidity. Increasing stratiform rain fraction with vertical wind shear has been found in both observational studies (Saxen and Rutledge 2000) and numerical modeling studies (Shie et al. 2003). This relationship also holds at the MJO time scale in these COARE and long-term data, as shown in Fig. 11. Both low-level (1000-850mb) shear and deep (700-150mb) shear lag precipitation by about 5-10 days, as part of a deep simple baroclinic zonal wind variation (e.g. Madden and Julian 1972).



In addition to this shear effect, middle- to upper-tropospheric relative humidity may be important. Positive correlations between stratiform rain fraction and humidity have been observed by Halverson et al. (1999). A cleaner delineation of cause and effect is found in the cloud modeling work of Tao et al. (1993), who found increases in area coverage and earlier formation of stratiform rain with more humid initial conditions. Similar lags of both stratiform fraction and humidity are seen in these intraseasonal composites (Fig. 12): humidity in the 600-200 hPa layer lags precipitation by 2-5 days in rawinsonde data from both the COARE filtered event data (panel a) and a five year set (1998-2002) of sounding data at the ARM Manus site (2S147E, panel b).

The other factors in the observed tilt of heating contours - time variation in radiative and convective heating profiles - can be separated by the same partitioning exercise used in section 4.2, extended to all the anomalous heating profiles during the MJO life cycle. Assembling the data for this decomposition is a bit tricky, since radiative heating data are only available for COARE while stratiform rain fraction  $f$  is only available from TRMM data unavailable during COARE. Here, the variation of stratiform rain fraction  $f$  is taken from Fig. 10a, added to an assumed total  $f = 0.6$  at the time of maximum precipitation. Radiative heating comes from the computations of Qian and Cess (2003), and convective heating is derived as a residual.

The partitioned heating results are shown in Fig. 13 as a function of lag and height. The tilt of the contours of total heating near day +10 in Fig. 14a is seen to be contributed largely by the extra stratiform heating at positive lag (Fig. 13b),

plus tilted radiative heating (Fig. 13a). The derived convective heating (panel c) has little tilt at the trailing edge of the MJO (positive lags), except in the extreme upper troposphere. However, it retains much of the tilt seen in the total heating contours at the MJO’s leading edge (lags -15 to -10 days).

These partitioned heating tilts are rather subtle, requiring eyeball evaluations of the slope of contours which are somewhat noisy and ragged. Unfortunately, it is not straightforward to make a smoother, more statistically significant picture using long-term data. However, the result about convection growing taller with time at negative lags is bolstered somewhat by an independent estimate of the convective activity from COARE data: eddy moist static energy flux (equation 1, section 2). In the lower free troposphere, this quantity is dominated by eddy moisture flux as saturated convective updrafts rise through unsaturated environmental air. Contours of this quantity also exhibit a slight rearward tilt in the lower-middle free troposphere near day -15 (Fig. 14a). Finally, a tilt in contours of divergence (Fig. 14b) is also seen in the lower-middle free troposphere near days -15 to -10, with less tilt at these levels near day +10.

In summary, a distinct tilt is observed in time-height sections of the total heating field at 155E. The data, while noisy, suggest that this tilt may have different causes at the leading and trailing edges of the COARE MJO event: At the leading edge, convective heating deepens with time, while at the trailing edge, the tilt is more a result of increased latent and radiative heating effects in stratiform anvils. These anvils lag the maximum precipitation by a few days, in step with the lags of wind

shear and middle-upper troposphere humidity.

## 6. Comparison to model heating profiles

Global models tend to simulate the MJO poorly. Theoretical studies within the wave-heating feedback framework suggest that the profile of heating is important to the strength and propagation of disturbances, as reviewed more fully in the discussion below. These prior findings motivated the observational study of MJO heating profiles reported here.

A first step toward improving models with observations is to compare of these observed heating profile against model heating profiles, preferably from multiple models. Obtaining and analyzing model data on all necessary variables with the necessary vertical and time resolution is well beyond the scope of this study. As a smaller step, we surveyed the published literature for all available plots of heating profiles in simulated large-scale equatorial eastward-propagating disturbances. The models include a GFDL GCM (Lau et al. 1988), a Japanese Meteorological Research Institute (MRI) GCM (Tokioka et al. 1988), three GCMs (Goddard Laboratory of Atmospheric Sciences, University of California at Los Angeles, Goddard Laboratory for Atmospheres) compared by Park et al. (1990), the University of Illinois GCM (Wang and Schlesinger 1999), the Seoul National University GCM (Lee et al. 2001), and one idealized model (Sui and Lau 1989). For two papers which show several profiles associated with different GCMs or different experiments with the same GCM (Tokioka et al. 1988, Park et al. 1990), we selected only the most top-heavy profile.

Heating profiles from these seven model studies are displayed in Fig. 15a, along with the observed MJO profile from the COARE IFA data. The mathematical maxima of the model heating profiles vary considerably, as do other, possibly more relevant aspects of profile shape. However, a systematic error is evident: All the model MJO heating profiles are generally middle-heavy, with insufficient heating in the upper troposphere around 450 mb, and excessive heating in the lower troposphere below 600 mb, as highlighted in the difference plot of panel b. In this respect, the differences are similar in shape to the stratiform heating profile (thin solid line in Fig. 8), suggesting that the models systematically lack an intraseasonally-varying stratiform-like heating process.

A model-observation comparison of stratiform rain fraction is not straightforward, or even well-defined. Stratiform precipitation has widely varying definitions even in high-resolution observations and cloud-model data (Lang et al. 2003). It is not clear that any separation of GCM rainfall into “convective” versus “stratiform” parts would correspond meaningfully to any one of the observational separation techniques. In short, we can only say that the models systematically lack a *functional equivalent* of stratiform precipitation.

## 7. Summary and discussion

This study has examined the top-heaviness of anomalous heating profiles in the MJO, and the related question of stratiform precipitation. Intraseasonal heating profiles simulated by models were compared to the observed. The main findings are:

1. The anomalous vertical heating profile at the time of maximum precipitation is very top-heavy, more so than the climatological mean profile. TRMM data show that anomalous stratiform precipitation contributes about 60% to the intraseasonal precipitation anomaly, more than the climatological fraction of stratiform precipitation. These observations are consistent, in that stratiform precipitation is characterized by heating in the upper troposphere and cooling in the lower troposphere.

2. During the passage of the MJO at 155E, anomalous heating has a phase tilt, in the sense of a rise with time in the heating profile. At the trailing edge of the rainy phase, this tilt is contributed largely by latent and radiative heating in stratiform anvils. Heating by these stratiform anvils lags the total precipitation by about five days, associated with the similar lags of vertical wind shear and middle-upper troposphere relative humidity.

3. The anomalous vertical heating profiles in 7 model intraseasonal oscillations differ from observations, in the sense of being middle-heavy rather than top-heavy. Models systematically lack (or under-represent) some functional equivalent of the observed strong modulation of stratiform precipitation by the MJO.

Is this systematic heating-profile error in models related to their poor simulations of the MJO? We will not presume to offer a simple or definitive answer, but theoretical studies offer some guidance for discussion. As mentioned in the introduction, such studies typically assume a wave-heating feedback framework. Although the large-scale circulation ('wave') structure of the observed MJO resembles a combination of forced-damped Kelvin and Rossby wave responses to equatorial heating (e.g. Salby

and Hendon 1994), theories tend to center on the Kelvin component, since it propagates eastward, as does the MJO. Such a Kelvin-wave prototype conveniently limits analysis to a two-dimensional subspace, the zonal-height plane on the equator. Our observations have implicitly assumed a similar equatorially-confined view. It is not known whether this approach is adequate, but it is customary.

The heating in wave-heating feedback theories may be divided into five components, namely the latent heatings associated with free troposphere moisture convergence; boundary layer moisture convergence; surface heat flux (latent plus sensible); and local moisture storage, as well as radiative heating. These five components have been emphasized in different types of theories, including the wave-CISK (Convective Instability of the Second Kind) mechanism (e.g. Hayashi 1970, Lindzen 1974, Lau and Peng 1987, Zhang and Geller 1994), the frictional wave-CISK mechanism (e.g. Hayashi 1971, Wang and Rui 1990b, Salby et al. 1994), the WISHE (Wave Induced Surface Heat Exchange) mechanism (e.g. Emanuel 1987, Neelin et al. 1987), the charge-discharge mechanism (e.g. Blade and Hartmann 1993, Wang and Schlesinger 1999), and the cloud-radiation interaction mechanism (Hu and Randall 1994, Raymond 2001, Lee et al. 2001), respectively.

The role of the vertical heating profile in setting the strength and phase speed of intraseasonal oscillations has been studied in several works. In some theories, such as wave-CISK, the heating profile is specified outright, allowing direct study of the influence of profile shape (Yamasaki 1968a, 1968b, 1969, Lau and Peng 1987, Takahashi 1987, Lau et al. 1988, Chang and Lim 1988, Tokioka et al. 1988, Sui

and Lau 1989, Cho and Pendlebury 1997). Two distinct issues accessible to such a treatment are the instability properties (or growth rate) of oscillations, and their phase speeds.

On the question of instability, Cho and Pendlebury (1997) found that instability requires a sufficient amplitude of a full-wave Fourier component to the heating profile - essentially, the 'stratiform' heating profile discussed in this paper. A 'stratiform instability' mechanism is also active in the model of Mapes (2000). Yamasaki (1968a, 1968b, 1969) found in numerical modeling studies that some tropical systems are unstable only if the cumulus heating profile has a maximum in the upper troposphere. On the basis of these results, a systematic lack of stratiform-like heating in models could be hypothesized to contribute to too-weak intraseasonal variability.

On the question of phase speed, results are available from several wave-CISK studies using linearized heating (Takahashi 1987, Lau et al. 1988) or conditional ("positive only") heating (Lau and Peng 1987, Sui and Lau 1989). In all cases, phase speed is found to increase with increases in the altitude of the maximum heating rate. A consistent result from a full GCM study is that of Tokioka et al. (1988), who introduced a critical (minimum) entrainment rate to the Arakawa-Schubert convection scheme. As this critical parameter was increased, both the altitude of the maximum heating and the phase speed of simulated intraseasonal oscillations decreased. Another GCM result in the same sense is Park et al. (1990)'s comparison of simulated intraseasonal oscillations in three different GCMs.

By these results, the too-fast propagation of intraseasonal oscillation in models

would not seem to be attributable to their systematic heating-profile bias, since the effect appears to be of the wrong sign. However, it is not clear that the theoretical results are as consistent as they appear, since the altitude of peak heating isn't necessarily the salient feature of heating profiles. From the perspective of adiabatic Kelvin wave dynamics, a heating profile should be characterized by what vertical modes it excites. Waves of short vertical wavelength, which propagate slowly, are excited by heating that is either top-heavy or bottom-heavy. If such waves play a role in the slow phase speed of the MJO, a simple peak-altitude characterization of heating may be misleading. This is a big "if," since the spectral power of the MJO does not lie along a Kelvin wave dispersion curve. As shown by Wheeler and Kiladis (1999), convectively-coupled Kelvin waves exist, distinct from the MJO, and propagate faster (10-15 m/s). Instead of "too-fast" MJOs, perhaps many models have reasonable or even too-slow convectively coupled Kelvin waves, and lack the MJO entirely! To summarize this discussion, there is no satisfying basis to conclude that models' systematic errors in heating profiles and intraseasonal oscillation propagation speed are linked.

The difficulties of attempting to link heating and MJO simulation should not be allowed to distract from the striking consistency found among model heating-profile errors (Fig. 15b). This finding suggests there may be a fundamental problem common to many models, which certainly merits attention.

Two possible approaches to increasing the activity of stratiform-like (top-heavy) heating in GCMs suggest themselves: revising the convection parameterization to account for stratiform processes; or changing the partition between parameterized



and explicit rainfall. In the first approach, convection schemes can be modified to include effects of mesoscale downdrafts (e.g. Molinari and Corsetti 1985, Sud and Walker 1993) or both mesoscale updrafts and mesoscale downdrafts (e.g. Donner 1993, Alexander and Cotton 1998, Gray 2000, Donner et al. 2001). In the second approach, a fraction of convective condensate can be passed to the large-scale condensate scheme from the cumulus parameterization, simulating detrainment (e.g. Randall et al. 1989, Tiedtke 1993, Fowler and Randall 1996, Del Genio et al. 1996, Li et al. 2002). None of these studies focused on changes to intraseasonal variability.

It would be naive to think that the approaches listed above will necessarily make intraseasonal heating anomaly profiles more top-heavy. In a GCM, a change to one process tends to induce compensating changes in other processes. For example, Lee et al. (2001) were unable to substantially change intraseasonal heating profiles, even with drastic changes to cloud-radiation processes: cumulus and large-scale heating conspired to cancel the radiative change, such that the anomalous total heating in convective areas was almost identical in all cases (their Fig. 7; LKKM line in Fig.15). A more fruitful approach might be to allow shallow convection to precipitate, heating the lower troposphere and paving the way for a more top-heavy total heating profile in deep convection.

Systematic model errors continue to frustrate MJO research, as it is hard to find a model, much less a suite of models, with variability sufficiently similar to observations to permit comparative studies to lead to incremental improvements. Although we cannot devise a simple definitive model test for the hypothesis that heating-profile

errors and MJO simulation errors are linked, the heating-profile error identified here is an informative diagnostic finding that invites comparative studies in more and different models. Vertical structure seems to contain important clues to the linkages between large-scale horizontal patterns (like MJO development and propagation) and physical processes, and this observational study provides a baseline.

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## FIGURE CAPTIONS

**Fig. 1** (a) Standard deviation of the 30-70 day bandpass filtered anomaly of the CMAP precipitation from 1979-2002. The unit is mm/day. The thick solid polygons are the sounding arrays during TOGA COARE. The inner one is the Intensive Flux Array (IFA). The outer one is the Outer Sounding Array (OSA).

(b) The lag-correlation between the 30-70 day CMAP precipitation anomaly and itself at 0N155E.

**Fig. 2** The response function of the Murakami filter used in this study.

**Fig. 3** The 30-70 day bandpass filtered anomaly of the CMAP precipitation during TOGA COARE (November 1, 1992 to February 28, 1993). The unit is mm/day. The thick dashed line indicates the location of the TOGA COARE sounding arrays.

**Fig. 4** The MJO anomalous vertical heating profile at the time of maximum precipitation for TOGA COARE IFA (thick solid line), TOGA COARE OSA (thick dashed line), and 24 years (1979-2002) chi-corrected heating composite at 0N155E for four different seasons: DJF (thin solid line), MAM (thin dotted line), JJA (thin dashed line), and SON (thin dash-dotted line).

**Fig. 5** Difference between the MJO anomalous heating profile at the time of maximum precipitation and the seasonal mean profile for (a) TOGA COARE (thick solid line for IFA, thick dashed line for OSA), and (b) 24 years (1979-2002) chi-corrected heating composite at 0N155E for four different seasons: DJF (thin solid line), MAM (thin dotted line), JJA (thin dashed line), and SON (thin dash-dotted line). The

scale of the horizontal axis is the same as in Fig. 4.

**Fig. 6** The precipitation anomaly during the composite life cycle of the MJO for 5 years (1998-2002) of TRMM PR data at 0N155E. The thick solid line is the total precipitation. The thin solid line is the convective precipitation. The thin dashed line is the stratiform precipitation.

**Fig. 7** Stratiform rain fraction at the time of maximum precipitation in the composite MJO life cycle (thick line) and the annual mean (thin line) for TRMM PR data along the equator (5N-5S). The composite MJO at each longitude is constructed using linear regression with respect to local temporally filtered (30-70 day) and spatially filtered (wavenumber 0-6) CMAP precipitation.

**Fig. 8** Partition of the observed heating profile (thick solid line) into three components: the stratiform component (thin solid line), the radiative component (thin dotted line), and the convective component (thin dashed line).

**Fig. 9** The vertical structure of the diabatic heating anomaly in the MJO for (a) TOGA COARE IFA. (b) TOGA COARE OSA. (c) DJF composite of 24 years (1979-2002) of chi-corrected heating data at 0N155E. The time lag is with respect to the time of maximum precipitation. The heating anomaly has been normalized by its column-integration at the time of maximum precipitation. The unit is (K/day)/(mm/day). The first contour is 0.03, and the contour interval is 0.15. Negative contours are shaded. In (c) only the anomaly with linear correlation above the 95% confidence level is plotted.

**Fig. 10** As in Fig. 6 except for the stratiform rain fraction anomaly.



**Fig. 11** As in Fig. 6 except for vertical wind shear anomaly for (a) TOGA COARE IFA, and (b) 24 years (1979-2002) of NCEP reanalysis data at 0N155E. The solid line is 1000-850 mb wind shear and the dotted line is 700-150 mb wind shear.

**Fig. 12** As in Fig. 9 except for RH anomaly for (a) TOGA COARE IFA, and (b) 5 years (1998-2002) of sounding data at ARM Manus site.

**Fig. 13** As in Fig. 9 except for TOGA COARE IFA (a)  $Q_R$  anomaly, (b)  $Q_{1m}$  anomaly, and (c)  $Q_{1c}$  anomaly. Contours are the same as in Fig. 9.

**Fig. 14** As in Fig. 9 except for TOGA COARE IFA (a) eddy moist static energy flux anomaly ( $W/m^2$ ), and (b) horizontal divergence anomaly ( $10^{-6}/s$ ). The first contour is 5 in (a) and 0.5 in (b). Contour interval is 5 in (a) and 1.0 in (b).

**Fig. 15** (a) Comparison of the observed heating profile with the model heating profiles. The thick solid line is the observed heating profile. All other lines are the model heating profiles. The models include the GFDL GCM used in Lau et al. 1988 (thick dashed line), the MRI GCM used in Tokioka et al. (1988)'s  $\alpha = 0$  experiment (thin solid line), the GLA GCM used in Park et al. 1990 (thin dotted line), the University of Illinois GCM used in Wang and Schlesinger (1999)'s A3 experiment (thin dashed line) and M2 experiment (thin dot-dashed line), the Seoul National University GCM used in Lee et al. 2001 (thin dot-dot-dashed line), and the theoretical model of Sui and Lau (1989)'s deep convection case (thin long-dashed line). (b) The difference between the observed heating profile and the model heating profiles.

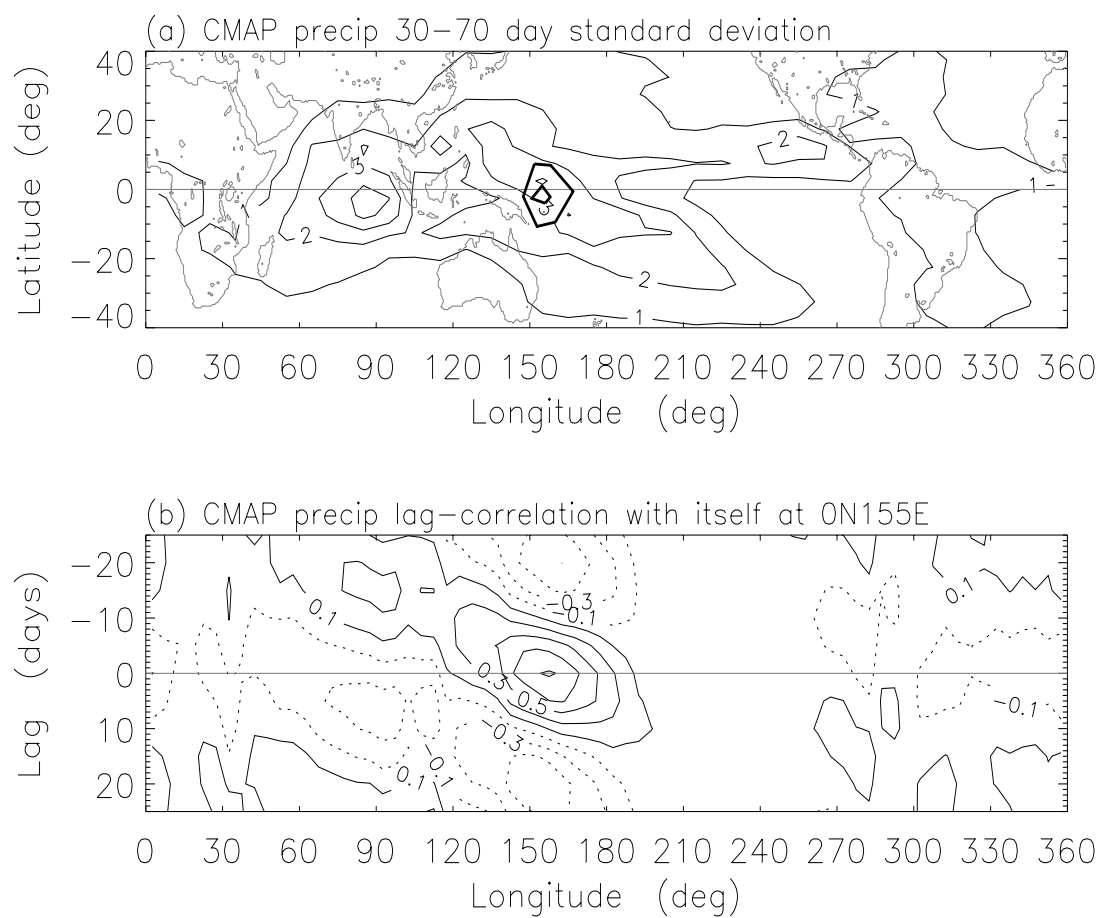


Figure 1:

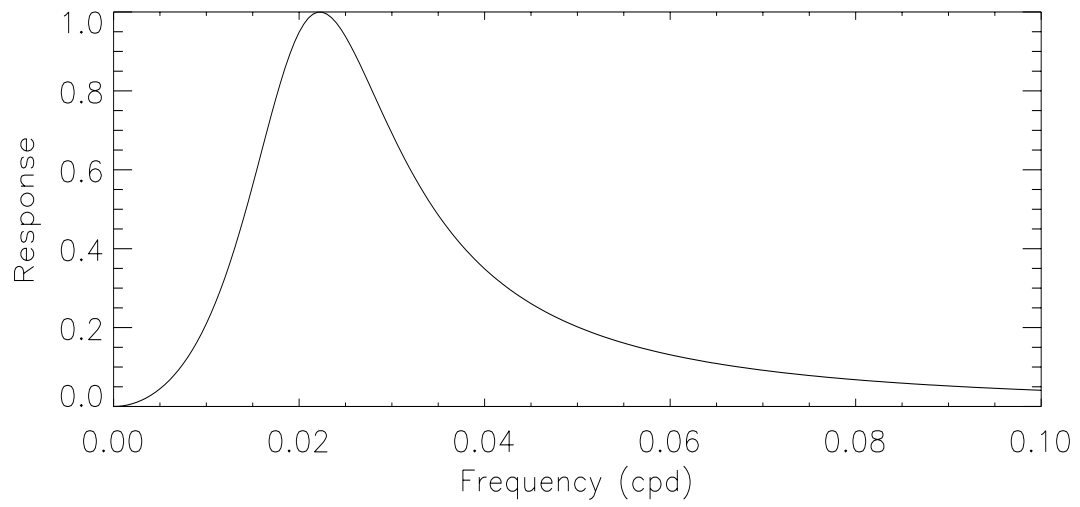


Figure 2:

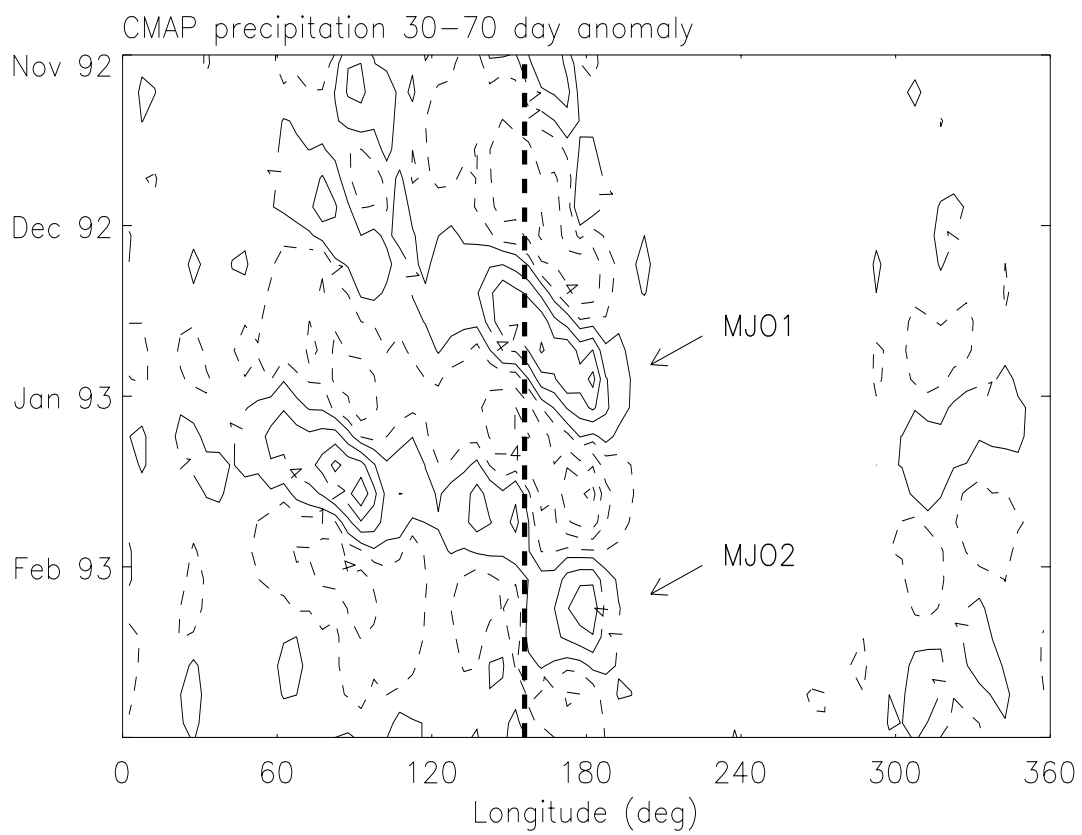


Figure 3:

MJO vertical heating profile  
at time of maximum precipitation

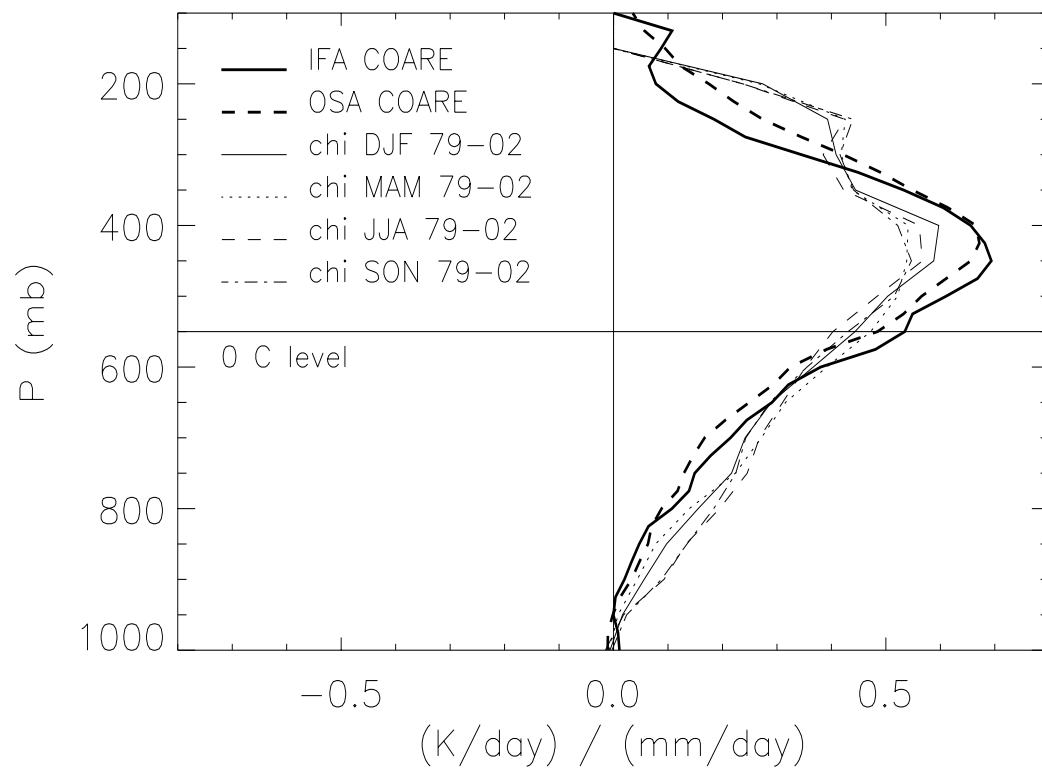


Figure 4:

Difference: MJO – seasonal mean

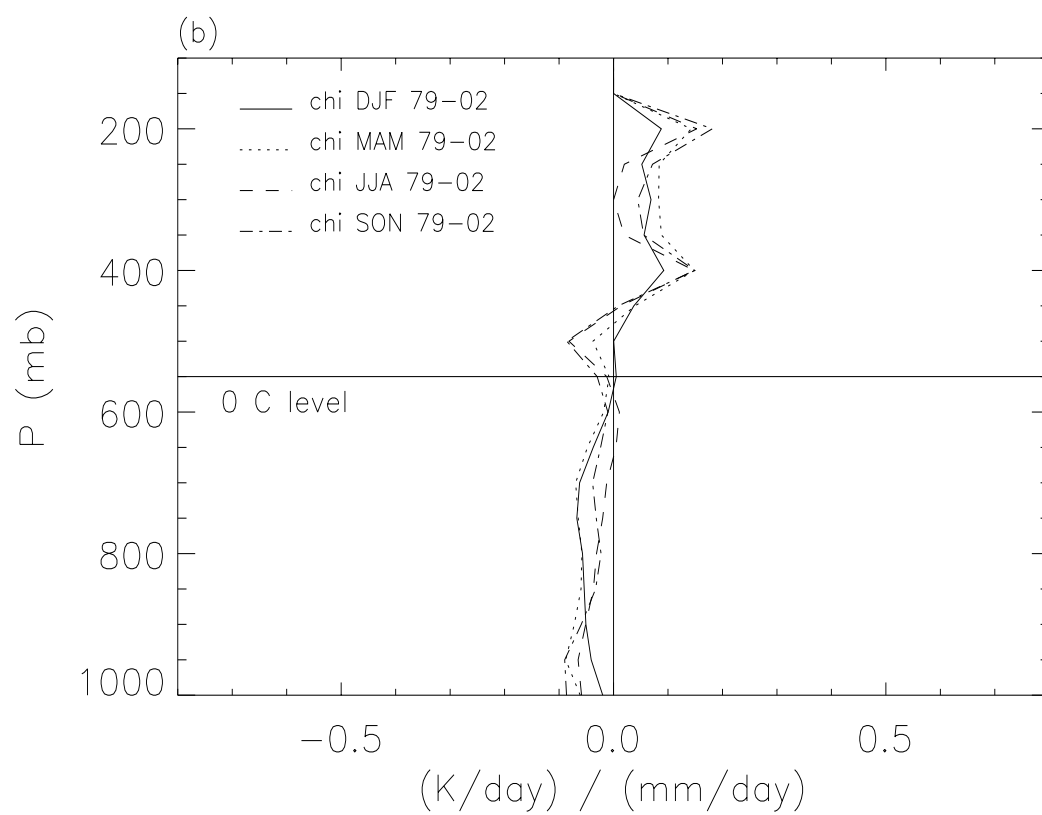
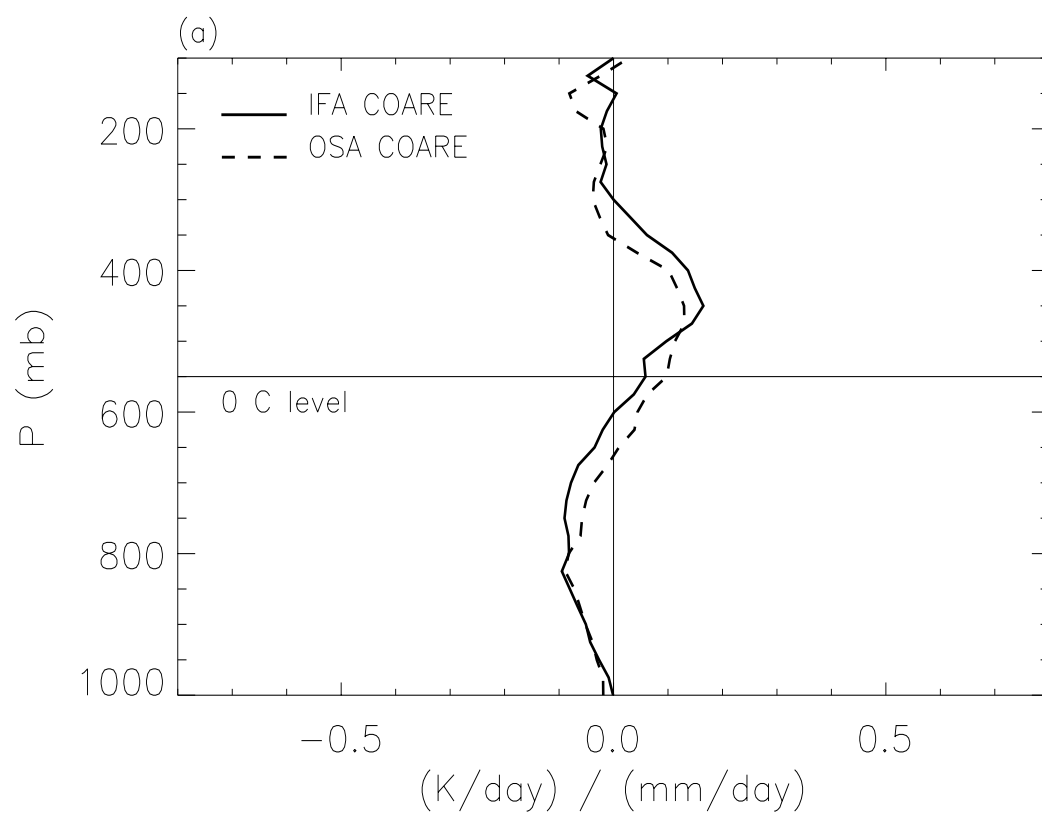


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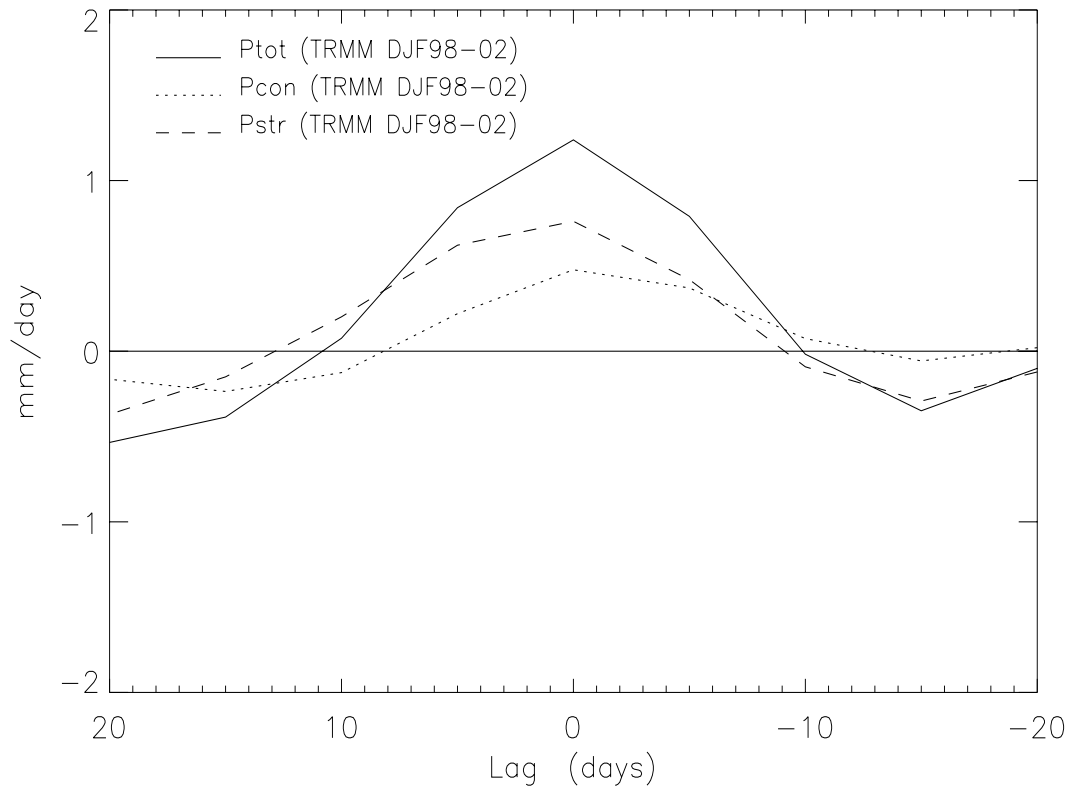


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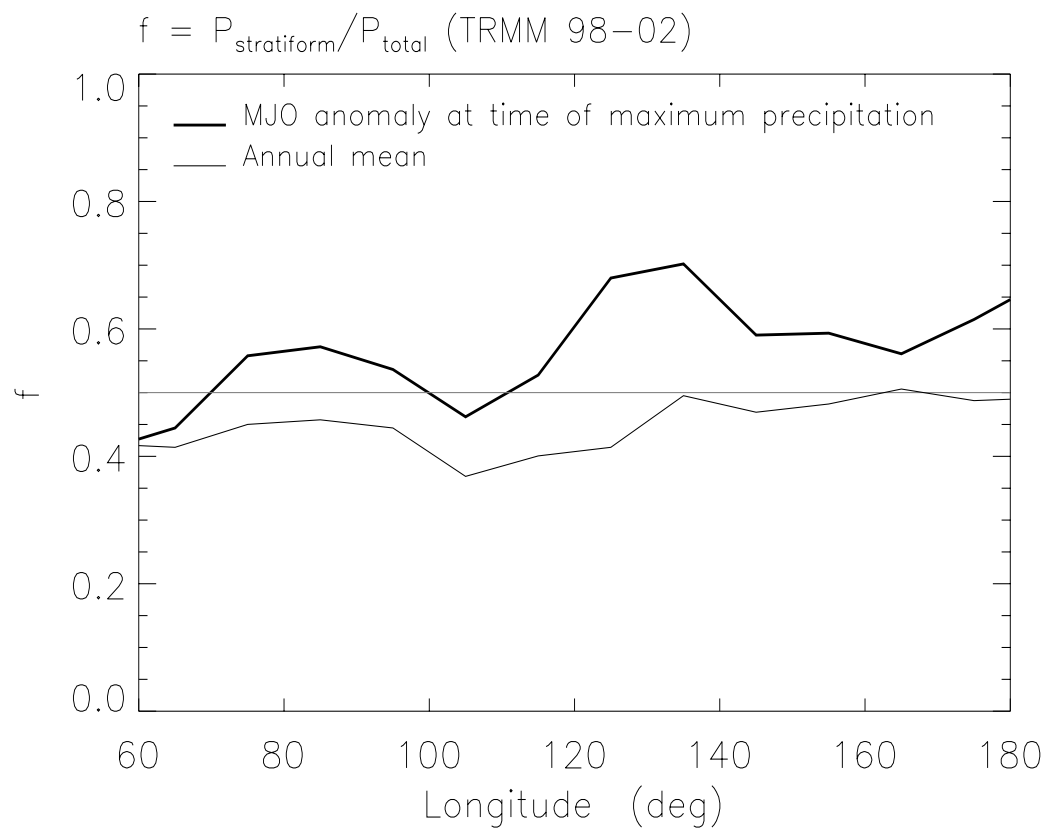


Figure 7:



# Effect of stratiform precipitation on vertical heating profile

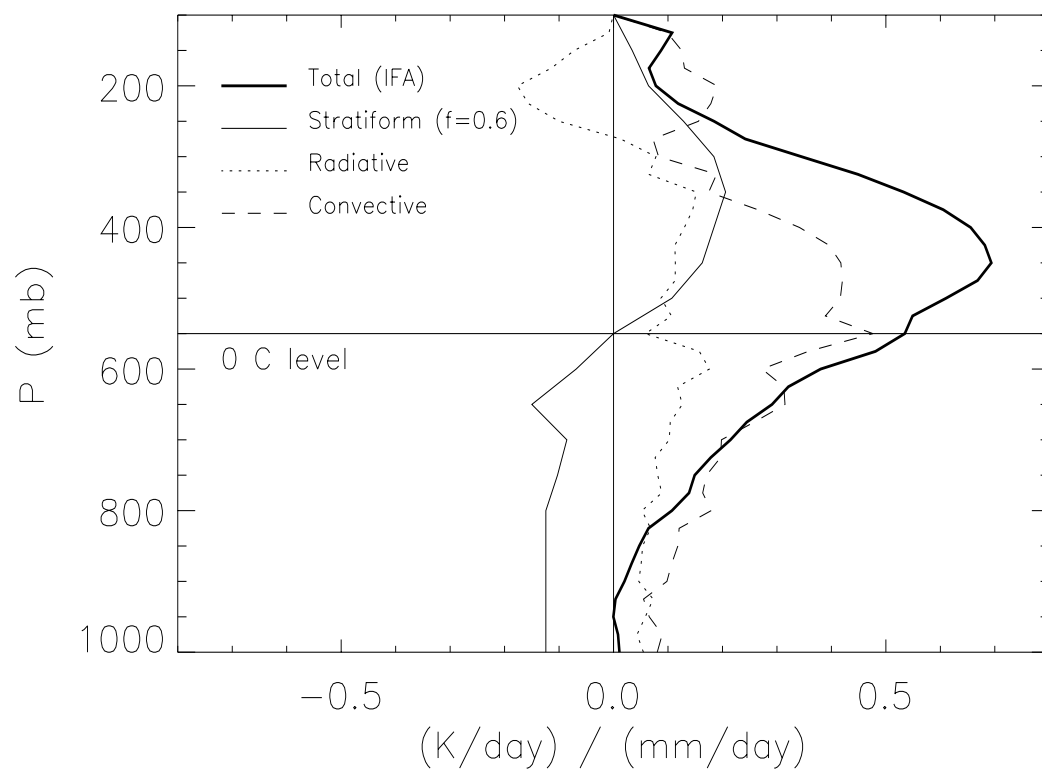


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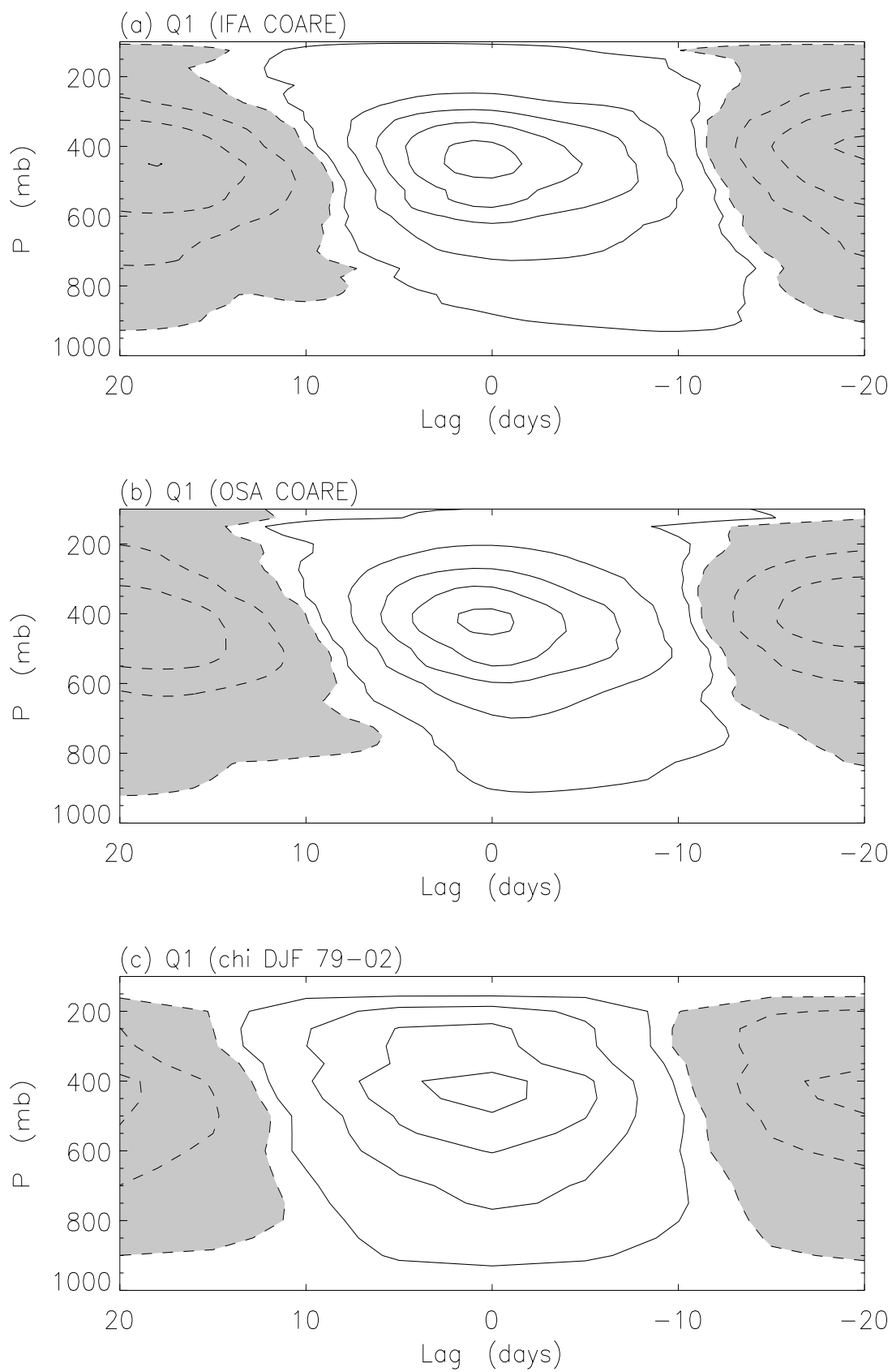


Figure 9:

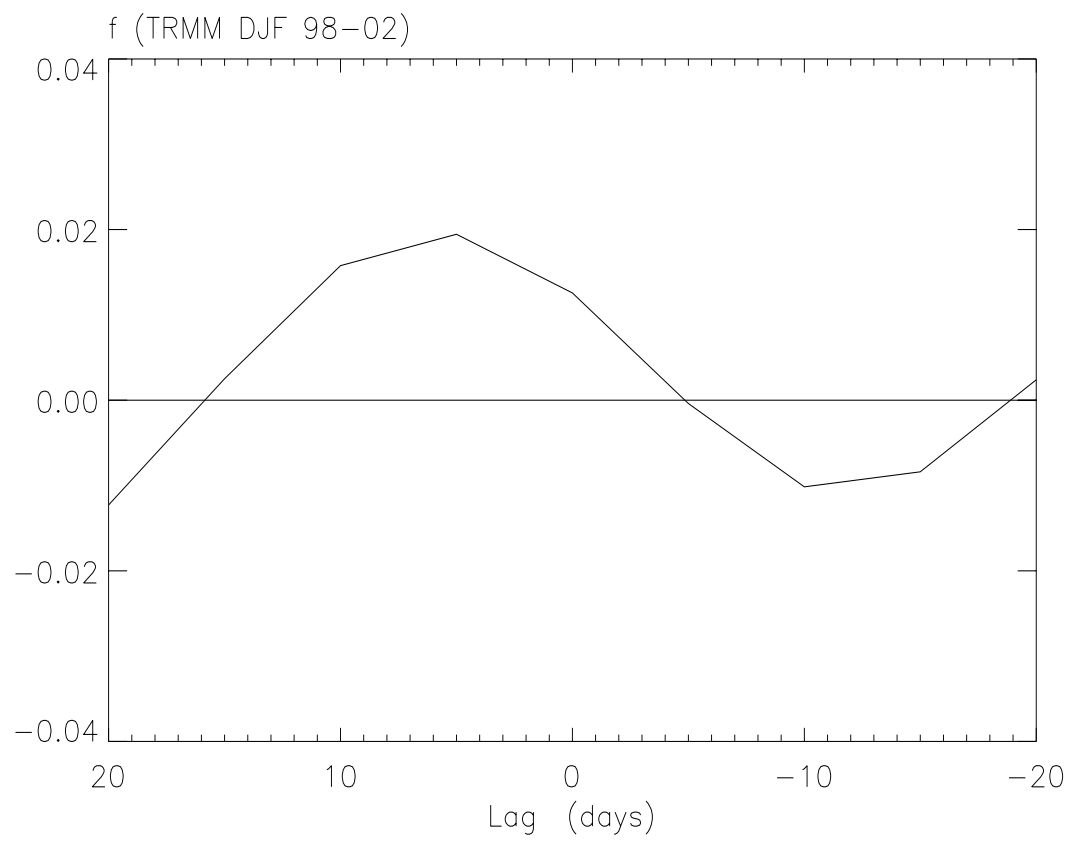


Figure 10:

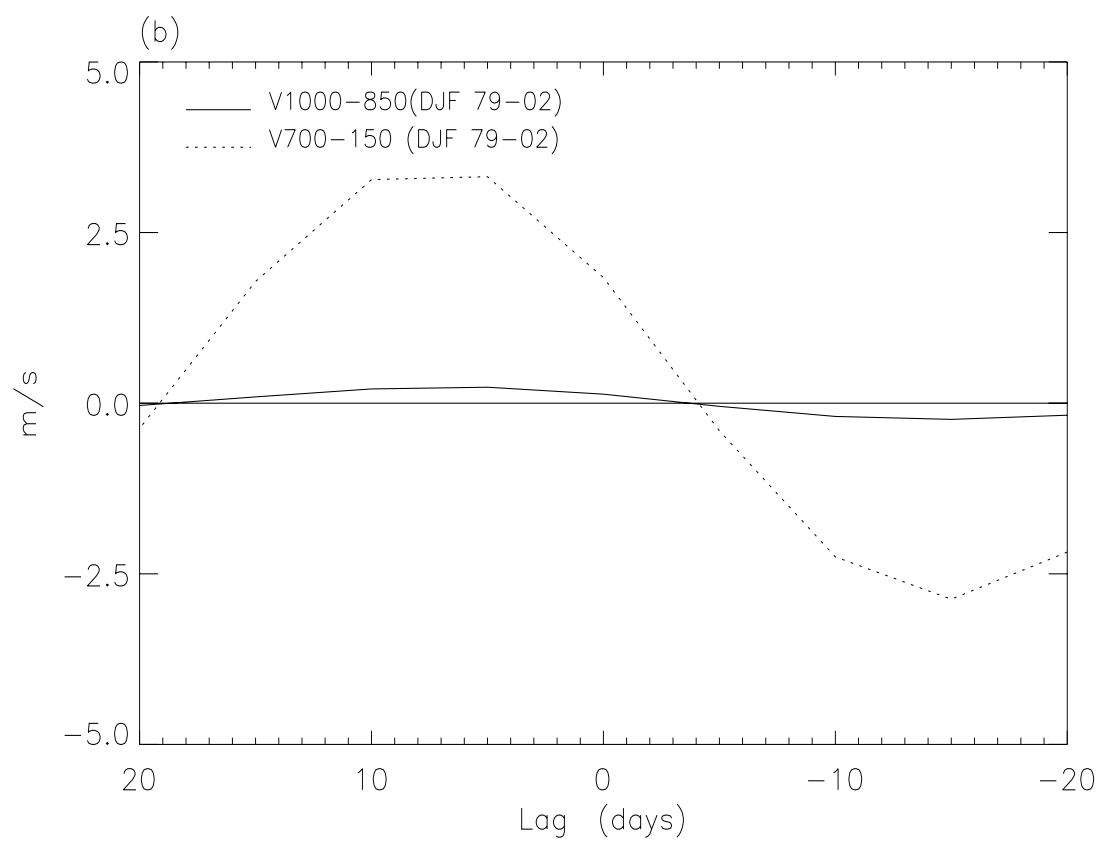
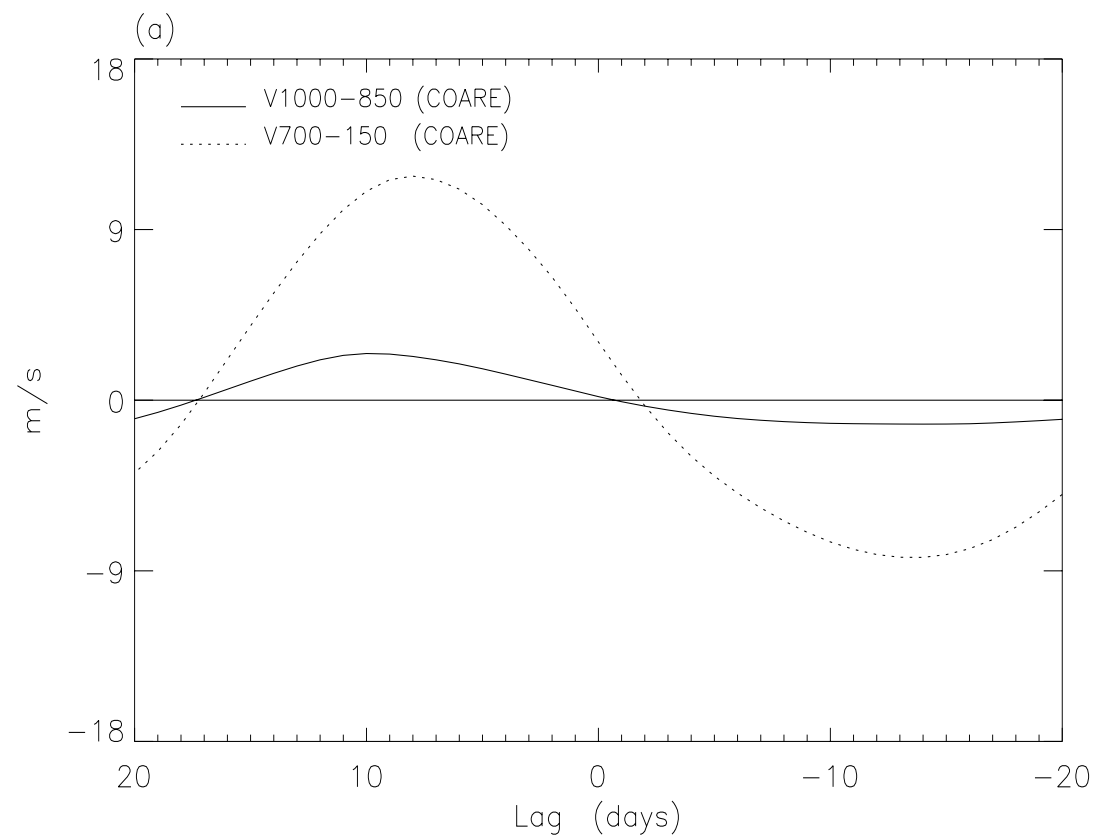


Figure 11:

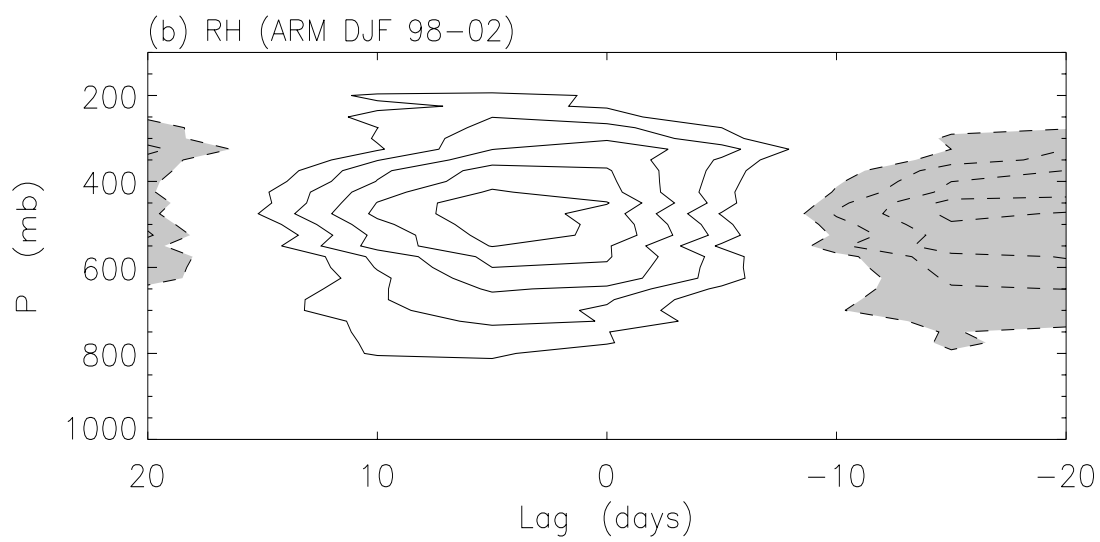
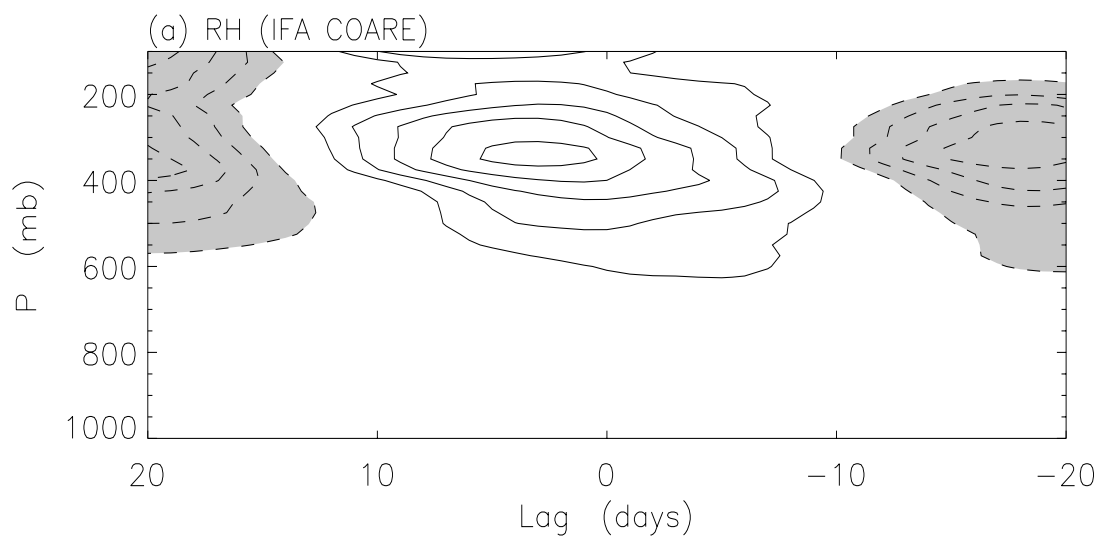


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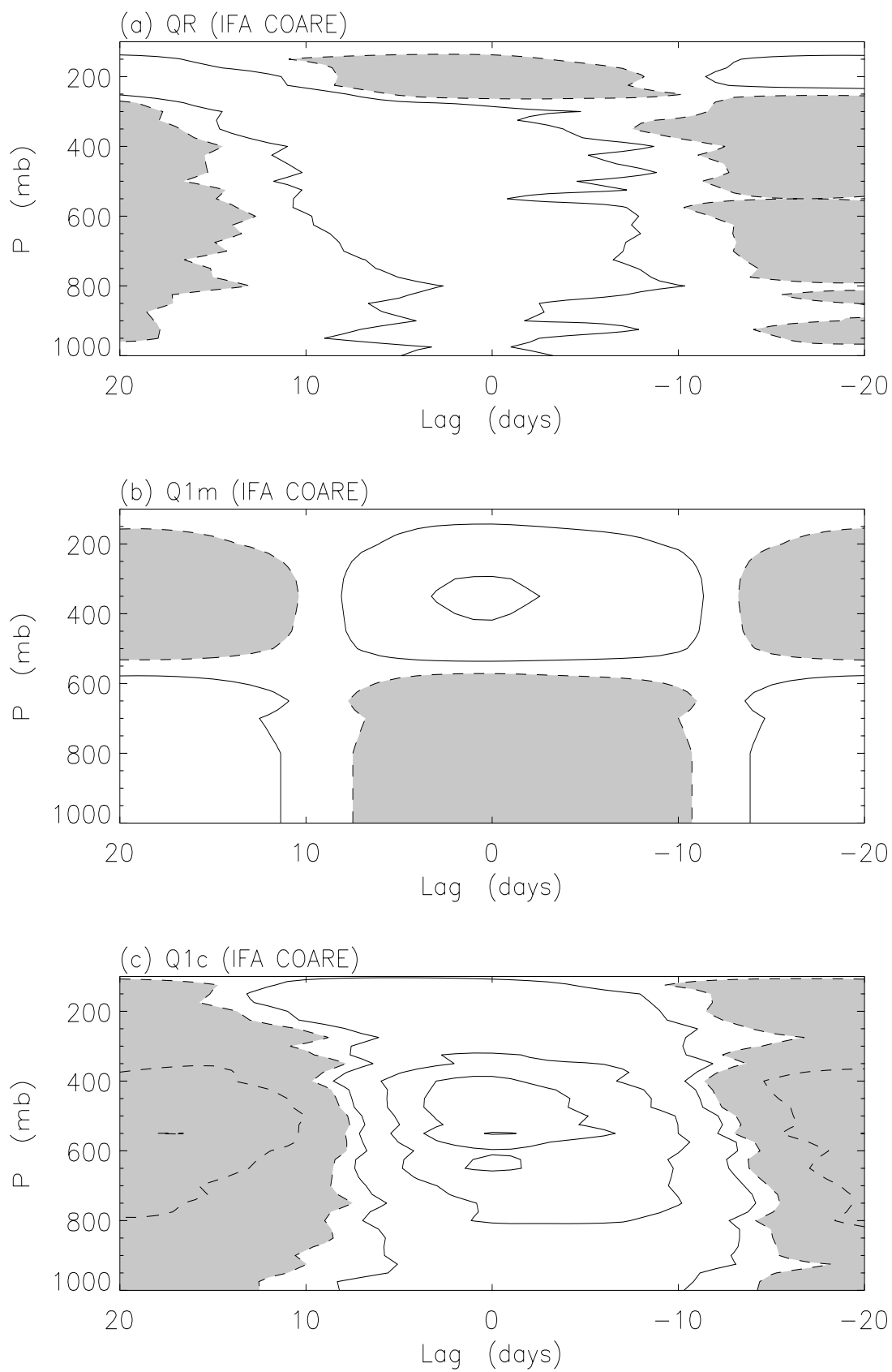


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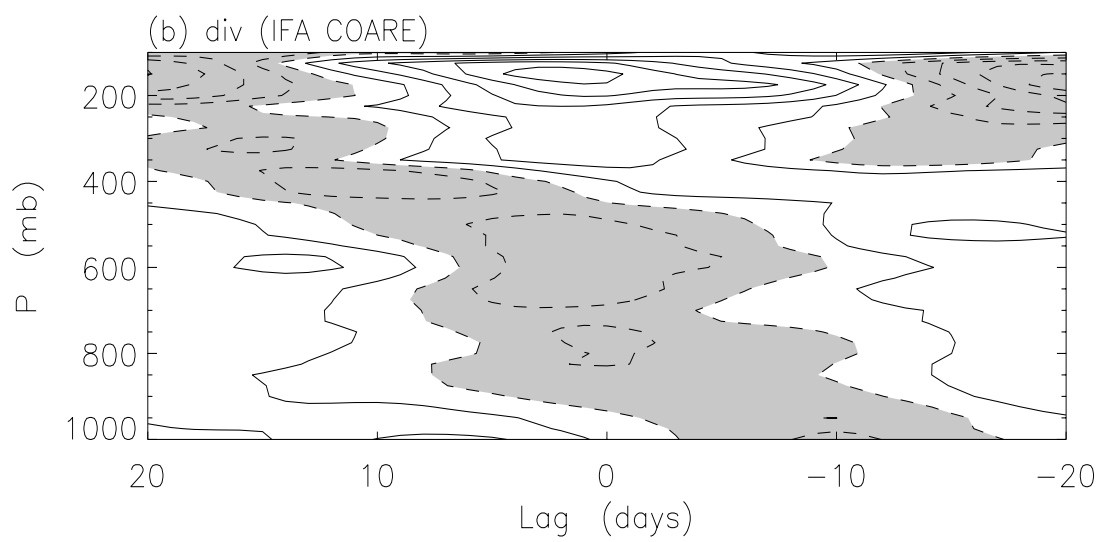
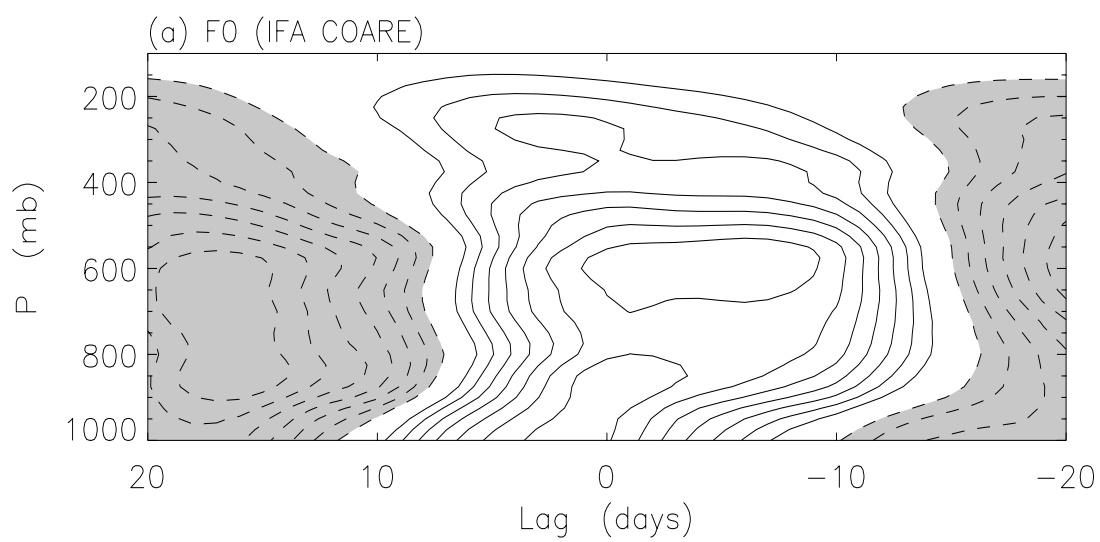


Figure 14:

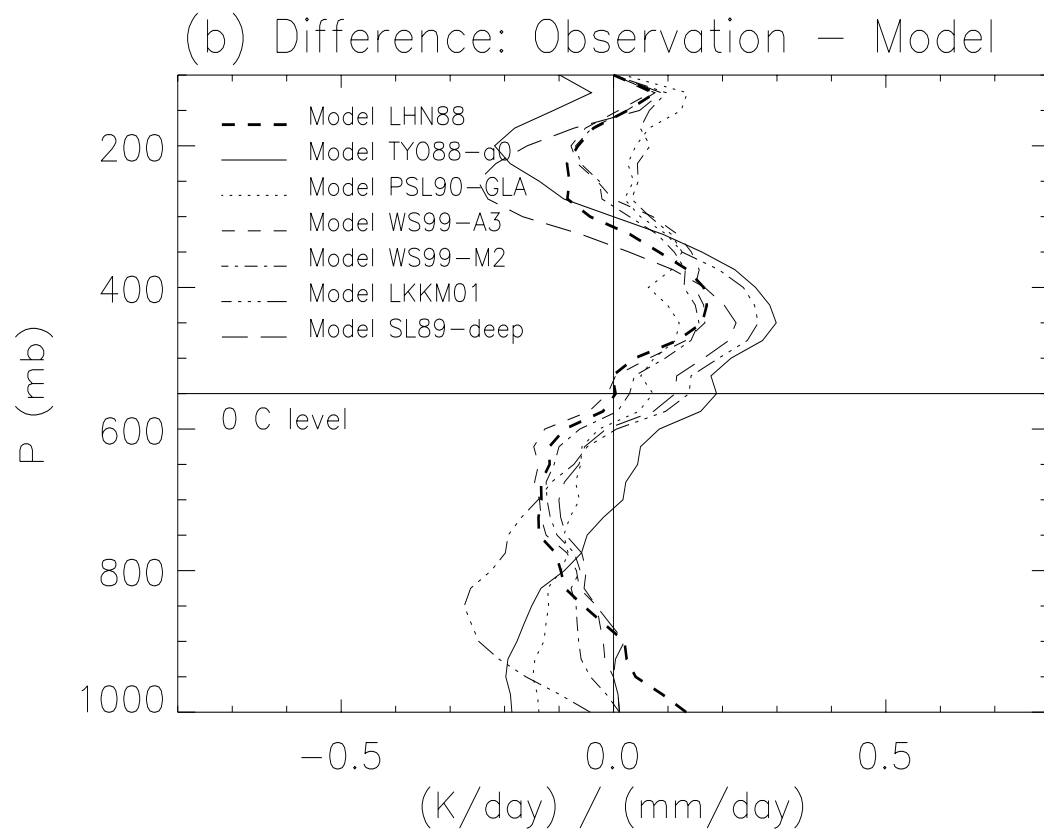
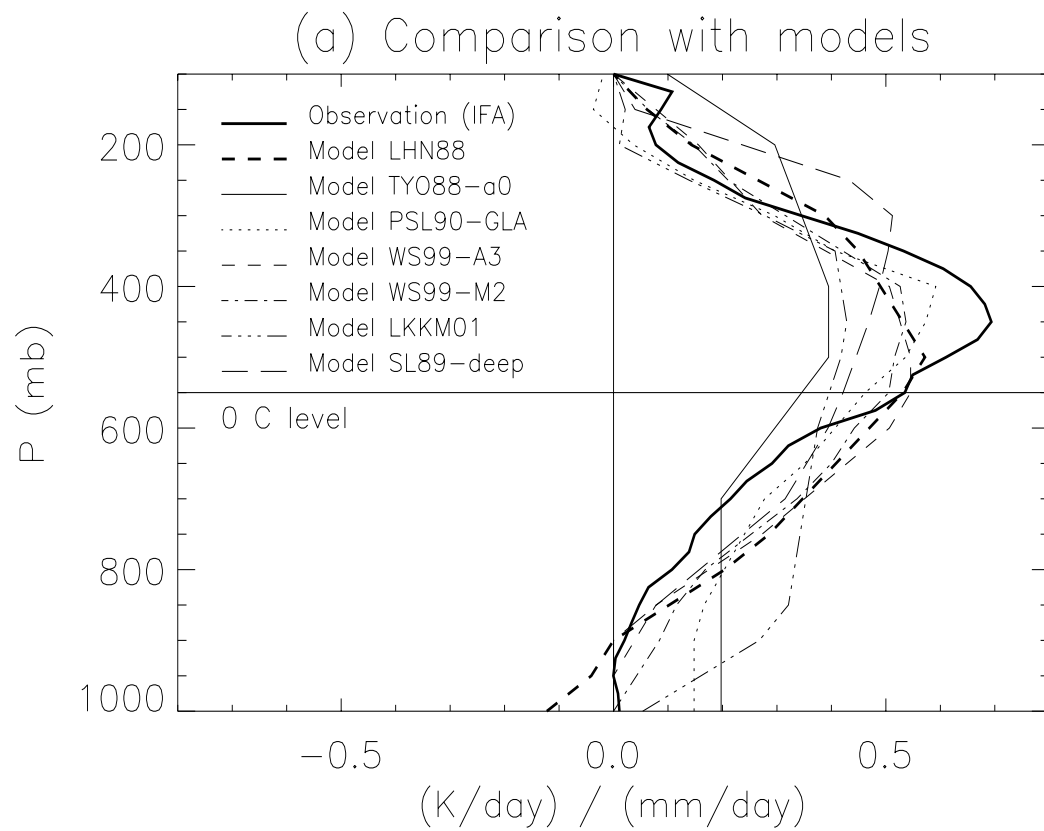


Figure 15: