

The evolution of structures in the universe

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Abstract. I compute the density perturbations, and velocities of dark matter and baryons. As well as the gravitational potential Φ and curvature Ψ . More importantly I also find the temperature multi poles Θ_l .

0.1 Introduction

In this project I am following the algorithm presented in Callin (2005)[1] for simulating the cosmic microwave background. This is part three of four for this project.

In the first part I set up the background cosmology of the universe, and made a function that could find the conformal time as a function of x . In the second part I computed the electron fraction, electron density, optical depth and visibility function for times around and during recombination.

In this part I will use some of these functions along with the Einstein-Boltzmann equations without polarization and neutrinos to compute the density perturbations, and velocities of dark matter and baryons. I will also compute the temperature multi poles Θ_l .

As previously done I will continue building on the skeleton code provided.

0.2 Equations

The full set of Einstein-Boltzmann equations without polarization and neutrinos read

$$\Theta'_0 = -\frac{ck}{\mathcal{H}}\Theta_1 - \Phi' \quad (1)$$

$$\Theta'_1 = \frac{ck}{3\mathcal{H}}\Theta_0 - \frac{2ck}{3\mathcal{H}}\Psi + \tau' \left[\Theta_1 + \frac{1}{3}v_b \right] \quad (2)$$

$$\Theta'_l = \frac{lck}{(2l+1)\mathcal{H}}\Theta_{l-1} - \frac{(l+1)ck}{(2l+1)\mathcal{H}}\Theta_{l+1} + \tau' \left[\Theta_l - \frac{1}{10}\Theta_l\delta_{l,2} \right], 2 \leq l < l_{max} \quad (3)$$

$$\Theta'_l = \frac{ck}{\mathcal{H}}\Theta_{l-1} - c\frac{l+1}{\mathcal{H}\eta(x)}\Theta_l + \tau'\Theta_l, l = l_{max} \quad (4)$$

$$\delta' = \frac{ck}{\mathcal{H}}v - 3\Phi' \quad (5)$$

$$v' = -v - \frac{ck}{\mathcal{H}}\Psi \quad (6)$$

$$\delta'_b = \frac{ck}{\mathcal{H}}v_b - 3\Phi' \quad (7)$$

$$v'_b = -v_b - \frac{ck}{\mathcal{H}}\Psi + \tau'R(3\Theta_1 + v_b) \quad (8)$$

$$\Phi' = \Psi - \frac{c^2k^2}{3\mathcal{H}^2}\Phi + \frac{H_0^2}{2\mathcal{H}^2} \left[\Omega_m a^{-1}\delta + \Omega_b a^{-1}\delta_b + 4\Omega_r a^{-2}\Theta_0 \right] \quad (9)$$

$$R = \frac{4\Omega_r}{3\Omega_b a}. \quad (10)$$

These equations are the ones we will use when we are not in the tight coupling regime. When in the tight coupling regime the factor $(3\Theta_1 + v_b)$ is very close to zero. In the equation for v'_b this is multiplied by τ' which is very large in, making this equation terribly unstable. The same thing makes the equation for Θ'_1 unstable. Because of this one can expand $(3\Theta_1 + v_b)$ in powers of $1/\tau'$. This results in a slight change in the set of equations for v'_b and Θ'_1 .

$$q = \frac{-[(1-2R)\tau' + (1+R)\tau''](3\Theta_1 + v_b) - \frac{ck}{\mathcal{H}}\Psi + (1 - \frac{\mathcal{H}'}{\mathcal{H}})\frac{ck}{\mathcal{H}}(-\Theta_0 + 2\Theta_2) - \frac{ck}{\mathcal{H}}\Theta'_0}{(1+R)\tau' + \frac{\mathcal{H}'}{\mathcal{H}} - 1} \quad (11)$$

$$v'_b = \frac{1}{1+R} \left[-v_b - \frac{ck}{\mathcal{H}}\Psi + R(q + \frac{ck}{\mathcal{H}}(-\Theta_0 + 2\Theta_2) - \frac{ck}{\mathcal{H}}\Psi) \right] \quad (12)$$

$$\Theta'_1 = \frac{1}{3}(q - v'_b). \quad (13)$$

So far I have not stated what tight coupling means. Basically because of the mathematical operations we have done, like expanding in powers of $1/\tau'$, we end up with some conditions on when these equations above can be used. The conditions are $|ck/(\mathcal{H}\tau')| < 1/10$, that the time is before recombination, and $|\tau'| > 10$.

All of these equations refer to their respective quantities in Fourier space. This means that the k 's everywhere refer to Fourier modes. This is done to separate the quantities into the different scales at which they take place, with low k 's referring to large scales, and high k 's to small scales.

Because of this it is expected that quantities relating to large k 's start changing earlier than those relating to low k 's, simply because no causal physics has been able to work on such large scales.

0.3 Implementation

The way to solve all these equations is to first make a function that finds the time where tight coupling ends. Note that this function clearly depends on which k mode we are working on.

For each k we insert the initial conditions, and then run through every value of x from some start value early in the universe after inflation is done. In this case I have chosen to use the x value corresponding to $a_{init} = 10^{-8}$. At some point through this the x value becomes larger than the x at the end of tight coupling. At that point we change the equations for the relevant quantities, and continue on until we reach today. This has to be done for all k values. I have chosen to set the limit at $k = 100$ so far. With this low k value the program completes in about 3 minutes.

It should also be said that we limit our number of l 's to 6. This can be done because we are using line of sight integration. Historically people used to include thousands of variables to trace multi poles. If we had to use this the program would use days or weeks instead of minutes.

0.4 Results

To get a good distribution of k modes we use a quadratic distribution in k such that

$$k_i = k_{min} + (k_{max} - k_{min})(i/100)^2, \quad (14)$$

where $k_{max} = 1000H_0/c$, and $k_{min} = 0.1H_0/c$. The results show the various quantities for six different k values. These k values are $k_1, k_5, k_{10}, k_{40}, k_{60}$, and k_{100} . See figures 1a and 1b for Φ and Ψ . Figures 2a, and 2b for δ , and δ_b . Figures 3a, and 3b for v and v_b . And figure 4 for Θ_0 . Most of what is happening is explained in the captions for each of them.

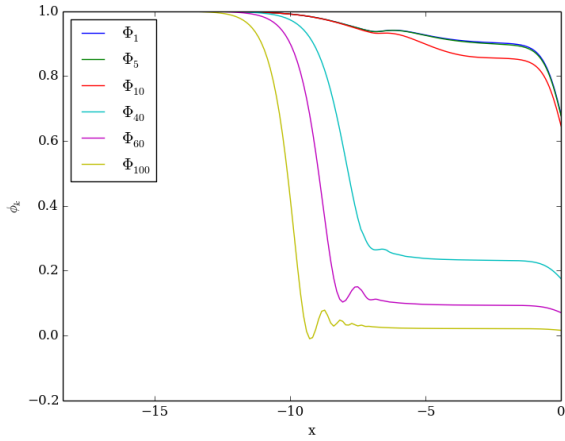
0.5 Conclusions

In this project we have calculated the density perturbations and velocities of dark matter and baryons. We also found the gravitational potential Φ and curvature Ψ . And most importantly we found the evolution of the temperature multi poles Θ_l which will enable us to make a CMB power spectrum in the next and final part of the project. We saw that dark matter and regular matter behaved as one would expect considering their respective properties with regards to electromagnetic radiation. This caused the gravitational potential to decrease when the dark matter density increased. This also meant that the curvature increased accordingly. All of these contributed to the calculation of the temperature multi poles, of which we only plotted Θ_0 at this point. We saw that the different modes stabilized at different temperatures. Exactly why the mean temperature of different Fourier modes should do this is not obvious at this point.

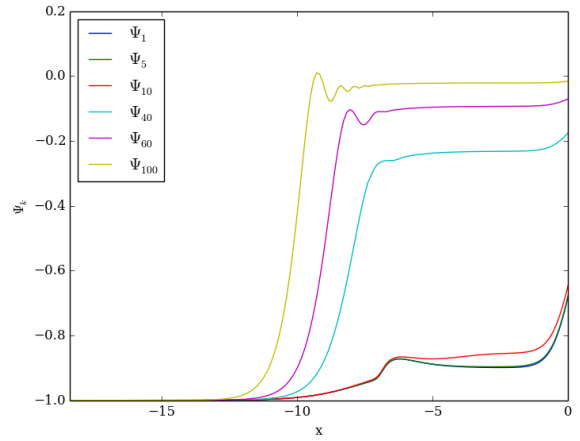
I was also able to appreciate the power of precomputing certain quantities, this enabled me to speed up the program by a factor of 10. There is still potential for improvement on this part, but it is good enough for its purpose at this point. If at a later point there is a need to increase the number of Fourier modes one should look at more quantities to precompute to reduce the runtime further. There is also some indication that the x grid used may benefit from a bit finer resolution when looking at the higher modes in Θ_0 .

0.6 References

- [1] P. Callin, astro-ph/0606683

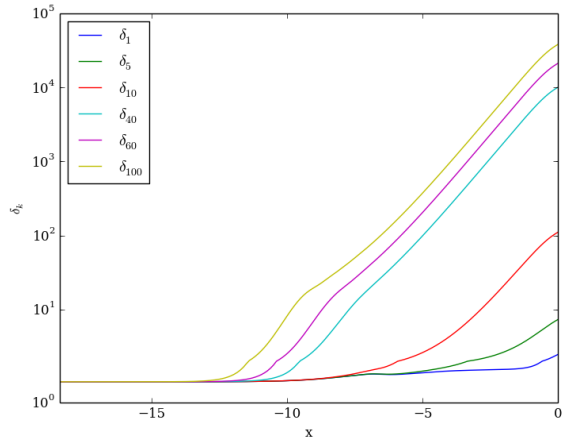


(a) The plot for Φ for six different k 's.

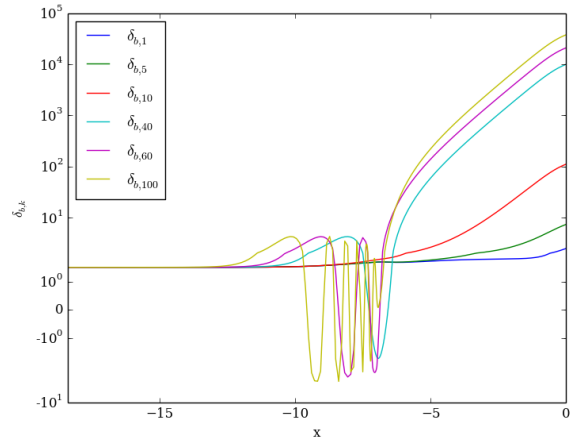


(b) The plot for Ψ for six different k 's.

Figure 1: The first thing to note is that the plots seem to be the inverse of each other. Also all the modes of Φ start at 1, and decrease as they leave tight coupling. The higher the mode, the closer to 0 they stabilize after recombination. The higher modes oscillate a bit before stabilizing around recombination. This is not seen in the lower modes. The smallest modes do not show this oscillating behavior at all, instead they slowly decrease after tight coupling before doing a nose dive near the end. The behavior of Ψ is the complete inverse of this. This behavior is as expected since the gravitational potential should be high with things far apart, and small when they are close. Thus higher modes fall first since matter and dark matter for high modes can clump together earlier, shrinking the potential. (Dark matter is the dominant factor here since matter is prevented from clumping by the photons until later.) This of course creates curvature, and that is why we see the opposite effect for Ψ . The lowest modes simply have not had time to collapse as much as they can yet.

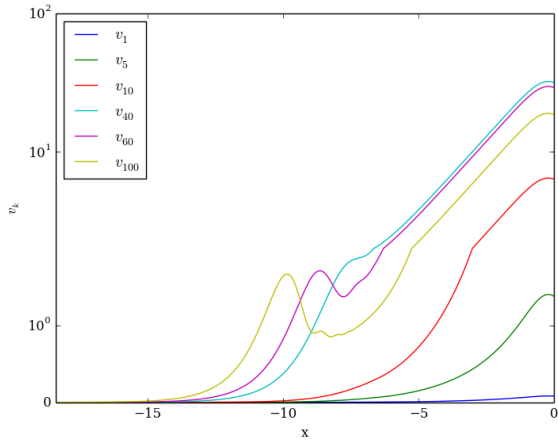


(a) The plot for δ for six different k 's.

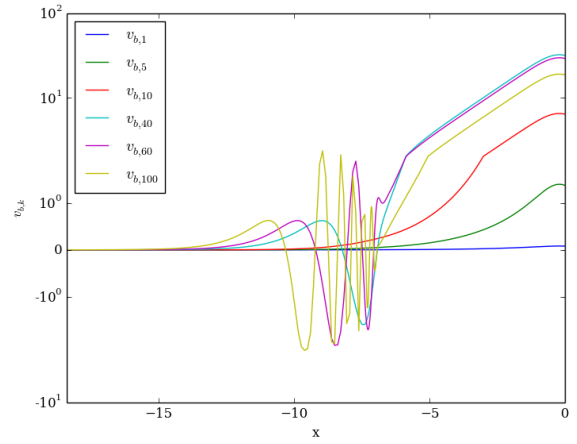


(b) The plot for δ_b for six different k 's.

Figure 2: Note the special y axis on the plot for δ_b . The dark matter perturbations behave nicely. The higher the mode the earlier it can start growing, just as expected due to the fact that forces don't propagate instantaneously. Note also that all the dark matter perturbations remain positive. This is not the case for baryons. When we start approaching recombination all the higher order modes start oscillating, and it is only after recombination is done that they are allowed to start growing again, and now they can behave like dark matter. The reason for this is of course that dark matter does not interact with electromagnetic radiation while regular matter does. Thus dark matter has no other force than gravity acting on it and thus it can collapse even though the universe is dominated by radiation. This is not the case for regular matter, as we see in the plot. The radiation rips everything apart before it can grow large, causing it to oscillate until recombination lets the photons loose.



(a) The plot for v for six different k 's.



(b) The plot for v_b for six different k 's.

Figure 3: As before there is a clear difference between the velocities of the two components. The highest modes start growing first, and since they are larger when approaching recombination they slow down more. Note the yellow and purple curves in the v plot. The lower modes don't slow down, but they don't increase as fast, while the smallest modes don't notice anything special happening. The velocity of the baryons is very different. As the dark matter, the highest modes start growing first, but as they exit tight coupling they start to oscillate, expanding and contracting, until after recombination where they speed up again to catch up with the dark matter. Note that the lowest modes don't notice anything special happening here either. In fact, they are fairly equal for dark matter and baryons. All of this has the same explanation as for the density perturbations. This is after all the velocity of those perturbations.

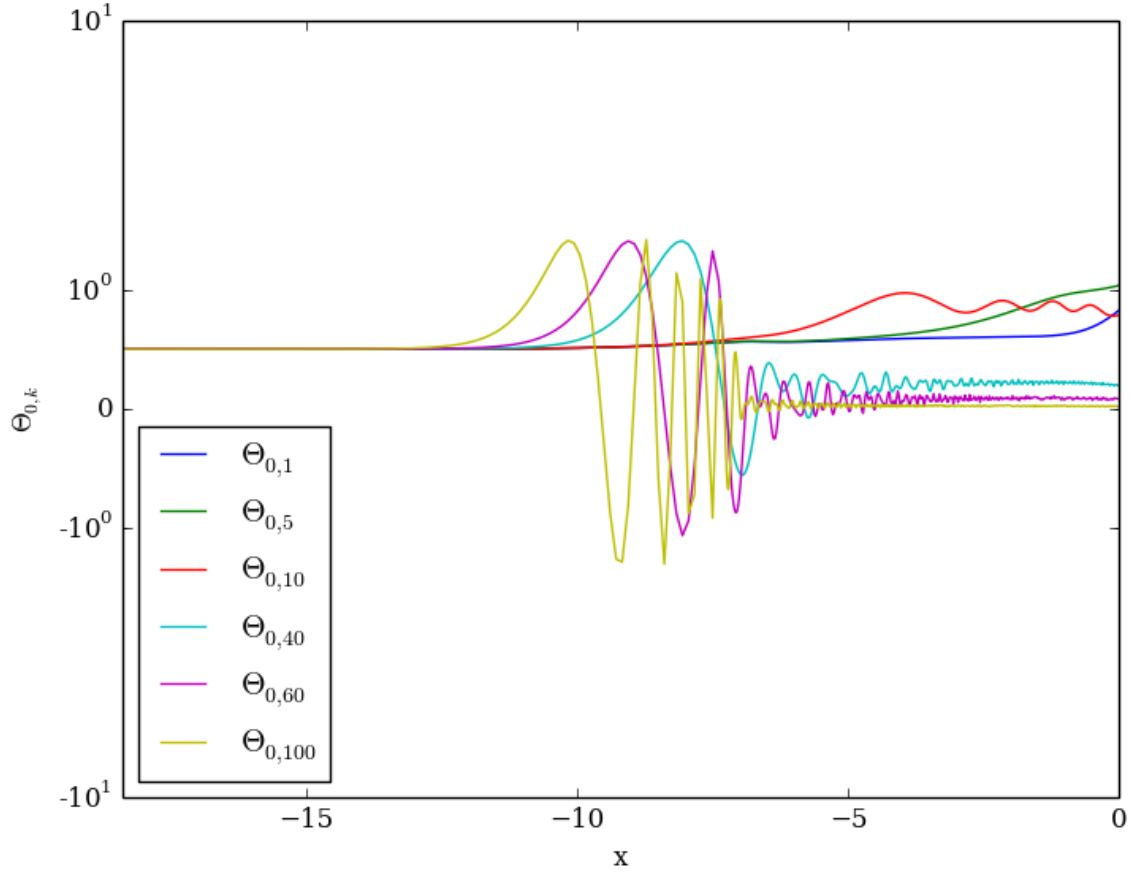


Figure 4: The plot shows Θ_0 for six different k 's. The low modes stay constant until the time at which their corresponding baryon perturbations start growing. This is as expected since Θ_0 measures the mean temperature, and this should increase as baryons clump together because of gravity. The high modes oscillate like their corresponding perturbation modes do. However, they do not grow after recombination, instead they stabilize around values fairly close to zero. The intermediate modes like k_{10} oscillate around some value between 0.5 and 1.

0.7 Source code

The source code for the `evolution_mod` file is included for inspection. This file depends on all files previously used in the two earlier parts of the project.

```
module evolution_mod
  use healpix_types
  use params
  use time_mod
  use ode_solver
  use rec_mod
  implicit none

  !Use j,k,l as global variable
  integer(i4b) :: j,k,l

  ! Accuracy parameters
  real(dp), parameter, private :: k_min = 0.1d0 * H_0 / c
  real(dp), parameter, private :: k_max = 1.d3 * H_0 / c
  integer(i4b), parameter :: n_k = 100
  integer(i4b), parameter, private :: lmax_int = 6

  ! Perturbation quantities
  real(dp), allocatable, dimension(:,:,:): Theta
  real(dp), allocatable, dimension(:,:) : delta
  real(dp), allocatable, dimension(:,:) : delta_b
  real(dp), allocatable, dimension(:,:) : Phi
  real(dp), allocatable, dimension(:,:) : Psi
  real(dp), allocatable, dimension(:,:) : v
  real(dp), allocatable, dimension(:,:) : v_b
  real(dp), allocatable, dimension(:,:) : dPhi
  real(dp), allocatable, dimension(:,:) : dPsi
  real(dp), allocatable, dimension(:,:) : dv_b
  real(dp), allocatable, dimension(:,:,:): dTheta

  !Precomputing variables
  real(dp), allocatable, dimension(:) :: dtau
  real(dp), allocatable, dimension(:) :: ddtau
  real(dp), allocatable, dimension(:) :: H_p
  real(dp), allocatable, dimension(:) :: dH_p
  real(dp), allocatable, dimension(:) :: eta_precomp

  ! Fourier mode list
  real(dp), allocatable, dimension(:) :: ks

  ! Book-keeping variables
  real(dp), private :: k_current, ck_current, ckH_p
  integer(i4b), private :: npar = 6+lmax_int
```


contains

```
! NB!!! New routine for 4th milestone only; disregard until then!!!
! subroutine get_hires_source_function(k, x, S)
!   implicit none

!   real(dp), pointer, dimension(:), intent(out) :: k, x
!   real(dp), pointer, dimension(:,:), intent(out) :: S

!   integer(i4b) :: i, j
!   real(dp)      :: g, dg, ddg, tau, dt, ddt, H_p, dH_p, ddHH_p, Pi, dPi, ddPi
!   real(dp), allocatable, dimension(:,:) :: S_lores

! Task: Output a pre-computed 2D array (over k and x) for the
!       source function, S(k,x). Remember to set up (and allocate) output
!       k and x arrays too.
!
! Substeps:
! 1) First compute the source function over the existing k and x
!    grids
! 2) Then spline this function with a 2D spline
! 3) Finally, resample the source function on a high-resolution uniform
!    5000 x 5000 grid and return this, together with corresponding
!    high-resolution k and x arrays

! end subroutine get_hires_source_function

! Routine for initializing and solving the Boltzmann and Einstein equations
subroutine initialize_perturbation_eqns
  implicit none
  integer(i4b) :: i
  real(dp)      :: k_min = 0.1d0*H_0/c
  real(dp)      :: k_max = 1000.d0*H_0/c

! Initialize k-grid, ks; quadratic between k_min and k_max
  allocate(ks(n_k))
  do k=1,n_k
    ks(k) = k_min + (k_max - k_min)*((k-1)/100.d0)**2
  end do

! Allocate arrays for perturbation quantities
  allocate(delta(1:n_t, n_k))
  allocate(delta_b(1:n_t, n_k))
  allocate(v(1:n_t, n_k))
  allocate(v_b(1:n_t, n_k))
  allocate(Phi(1:n_t, n_k))
```

```

allocate(Theta(1:n_t, 0:lmax_int, n_k))
allocate(Psi(1:n_t, n_k))

allocate(dPhi(1:n_t, n_k))
allocate(dPsi(1:n_t, n_k))
allocate(dv_b(1:n_t, n_k))
allocate(dTheta(1:n_t, 0:lmax_int, n_k))

!Allocate arrays for precomputed variables
allocate(dtau(n_t),H_p(n_t),dH_p(n_t))
allocate(ddtau(n_t),eta_precomp(n_t))

!Precompute useful variables
do i=1,n_t
    dtau(i) = get_dtau(x_t(i))
    ddtau(i) = get_ddtau(x_t(i))
    H_p(i) = get_H_p(x_t(i))
    dH_p(i) = get_dH_p(x_t(i))
    eta_precomp(i) = get_eta(x_t(i))
end do
write(*, '(2X, ES14.6)') H_p(1),dH_p(1),ddtau(1),dtau(1),ks(1)
!Set up initial conditions for the Boltzmann and Einstein equations
Phi(1,:) = 1.d0
delta(1,:) = 1.5d0*Phi(1,:)
delta_b(1,:) = delta(1,:)
Theta(1,0,:) = 0.5d0*Phi(1,:)
do k = 1, n_k
    v(1,k) = c*ks(k)/(2.d0*H_p(1))*Phi(1,k)
    v_b(1,k) = v(1,k)
    Theta(1,1,k) = -c*ks(k)/(6.d0*H_p(1))*Phi(1,k)
    Theta(1,2,k) = -20.d0*c*ks(k)/(45.d0*H_p(1)*dtau(1))*Theta(1,1,k) !without polarization
    do l = 3, lmax_int
        Theta(1,l,k) = -1/(2.d0*l+1.d0)*c*ks(k)/(H_p(1)*dtau(1))*Theta(1,l-1,k)
    end do
    Psi(1,k) = -Phi(1,k) - 12.d0*H_0**2/(ks(k)*c*a_t(1))**2*Omega_r*Theta(1,2,k)
end do
end subroutine initialize_perturbation_eqns

subroutine integrate_perturbation_eqns
implicit none
real(dp) :: x1, x2, x_init
real(dp) :: eps, hmin, h1, x_tc, j_tc, dt, t1, t2
real(dp) :: R,d_v,d_v_b,q
real(dp), allocatable, dimension(:) :: y, y_tight_coupling, dydx
eps = 1.d-8
hmin = 0.d0
h1 = 1.d-5
allocate(y(npar))
allocate(dydx(npar))

```

```

allocate(y_tight_coupling(7))

dydx(:) = 0

! Propagate each k-mode independently
do k = 1, n_k
  write(*,*) 'Current k', k
  k_current = ks(k) ! Store k_current as a global module variable
  ck_current = c*ks(k) !store c*k

  ! Initialize equation set for tight coupling
  y_tight_coupling(1) = delta(1,k)
  y_tight_coupling(2) = delta_b(1,k)
  y_tight_coupling(3) = v(1,k)
  y_tight_coupling(4) = v_b(1,k)
  y_tight_coupling(5) = Phi(1,k)
  y_tight_coupling(6) = Theta(1,0,k)
  y_tight_coupling(7) = Theta(1,1,k)

  ! Find the time to which tight coupling is assumed,
  ! and integrate equations to that time
  x_tc = get_tight_coupling_time(k_current)
  !write(*,*) 'x_tc =',x_tc
  !write(*,*) 'under x_tc'

  !Integrate from x_init until the end of tight coupling, using
  !the tight coupling equations

  !write(*,*) 'Start of tight coupling'
  !write (*, '(*(2X, ES14.6))') delta(1,k), delta_b(1,k), &
  !v(1,k), v_b(1,k), Phi(1,k), Theta(1,0,k), Theta(1,1,k),Psi(1,k)
  !write (*, '(*(2X, ES14.6))') x_t(1),dv_b(1,k),dPsi(1,k), &
  !dPhi(1,k),dTheta(1,0,k),dTheta(1,1,k),dTheta(1,2,k)

  do j=2,n_t
    if (x_t(j)< x_tc) then
      !precompute some variables
      ckH_p = ck_current/H_p(j)

      !Solve next step
      call odeint(y_tight_coupling,x_t(j-1),x_t(j),eps,h1,hmin,derivs_tc, bsstep, output3)

      !Save variables
      delta(j,k) = y_tight_coupling(1)
      delta_b(j,k) = y_tight_coupling(2)
      v(j,k) = y_tight_coupling(3)
      v_b(j,k) = y_tight_coupling(4)
      Phi(j,k) = y_tight_coupling(5)
      Theta(j,0,k) = y_tight_coupling(6)

```

```

Theta(j,1,k) = y_tight_coupling(7)
Theta(j,2,k) = -(20.d0*ckH_p)/(45.d0*dtau(j))*Theta(j,1,k)
do l = 3, lmax_int
    Theta(j,l,k) = -1/(2.d0*1+1.d0)*ckH_p/dtau(j)*Theta(j,l-1,k)
end do
Psi(j,k)      = -Phi(j,k) - 12.d0*H_0**2/(ck_current*a_t(j))**2*Omega_r*Theta(j,2,k)

!Store derivatives that are required for C_l estimation
dPhi(j,k)     = Psi(j,k) - (ckH_p)**2/3.d0*Phi(j,k) + (H_0**2/H_p(j))**2/2.d0 &
    *(Omega_m/a_t(j)*delta(j,k) + Omega_b/a_t(j)*delta_b(j,k) &
    + 4.d0*Omega_r/a_t(j)**2 *Theta(j,0,k))

dPsi(j,k)     = -dPhi(j,k) - 12.d0*H_0**2/(ck_current*a_t(j))**2&
    *Omega_r*(-2.d0*Theta(j,2,k)+dTheta(j,2,k))

dTheta(j,0,k) = -ckH_p*Theta(j,1,k) - dPhi(j,k)

R              = 4.d0*Omega_r/(3.d0*Omega_b*a_t(j))

q              = (-(1.d0-2.d0*R)*dtau(j) + &
    (1.d0+R)*ddtau(j))*(3.d0*Theta(j,1,k)+v_b(j,k)) - &
    ckH_p*Psi(j,k) +&
    (1.d0-dH_p(j)/H_p(j))*ckH_p*(-Theta(j,0,k) + 2.d0*Theta(j,2,k))-&
    ckH_p*dTheta(j,0,k))/((1.d0+R)*dtau(j)+dH_p(j)/H_p(j) -1.d0)

dv_b(j,k)     = 1.d0/(1.d0+R)*(-v_b(j,k)-ckH_p*Psi(j,k)+&
    R*(q+ckH_p*(2.d0*Theta(j,2,k)-Theta(j,0,k)))-&
    ckH_p*Psi(j,k))

dTheta(j,1,k) = 1.d0/3.d0*(q-dv_b(j,k))
dTheta(j,2,k) = 0
do l = 3, lmax_int
    dTheta(j,l,k) = 0
end do

!write (*,'*(2X, ES14.6)') delta(j,k), delta_b(j,k), &
!v(j,k), v_b(j,k), Phi(j,k), Theta(j,0,k), Theta(j,1,k),Psi(j,k)
!write (*,'*(2X, ES14.6)') x_t(j),dPsi(j,k),dPhi(j,k),dv_b(j,k),&
!dTheta(j,0,k),dTheta(j,1,k),dTheta(j,2,k)

else
    j_tc = j
    exit
end if
end do
!write(*,*) 'End of tight coupling'

!Set up variables for integration from the end of tight coupling

```

```

!until today
y(1:7) = y_tight_coupling(1:7)
y(8) = Theta(1,2,k)
do l = 3, lmax_int
    y(6+l) = Theta(1,l,k)
end do

!Continue after tight coupling
!write(*,*) 'start of rec'
do j = j_tc, n_t

    !Precompute some variables
    ckH_p = ck_current/H_p(j)

    !Integrate equations from tight coupling to today
    !write(*,*) 'running odeint with j =', j
    call odeint(y, x_t(j-1), x_t(j), eps, h1, hmin, derivs, bsstep, output3)

    !Store variables at time step j in global variables
    delta(j,k) = y(1)
    delta_b(j,k) = y(2)
    v(j,k) = y(3)
    v_b(j,k) = y(4)
    Phi(j,k) = y(5)

    do l = 0, lmax_int
        Theta(j,l,k) = y(6+l)
    end do
    Psi(j,k) = - Phi(j,k) - 12.d0*H_0**2/(ck_current*a_t(j))**2*Omega_r*Theta(j,2,k)

    !Store derivatives that are required for C_l estimation
    dPhi(j,k) = Psi(j,k) - c**2*k_current**2/(3.d0*H_p(j)**2)*&
        Phi(j,k) + H_0**2/(2.d0*H_p(j))*(&
        Omega_m/a_t(j)*&
        delta(j,k) + Omega_b/a_t(j)*delta_b(j,k) + 4.d0*&
        Omega_r/a_t(j)**2*Theta(j,0,k))

    dv_b(j,k) = -v_b(j,k) - ckH_p*Psi(j,k) + dtau(j)*R*(3.d0*Theta(j,1,k)+ v_b(j,k))

    dTheta(j,0,k) = -ckH_p*Theta(j,1,k) - dPhi(j,k)
    dTheta(j,1,k) = ckH_p/3.d0*Theta(j,0,k) - &
        2.d0*ckH_p/3.d0*Theta(j,2,k)+&
        ckH_p/3.d0*Psi(j,k) + &
        dtau(j)*(Theta(j,1,k)+ 1.d0/3.d0*v_b(j,k))
    dTheta(j,2,k) = 2.d0*ckH_p/5.d0*Theta(j,1,k) -&
        3.d0*ckH_p/5.d0*Theta(j,3,k)+&
        dtau(j)*0.9d0*Theta(j,2,k)
    do l=3,lmax_int-1
        dTheta(j,l,k) = 1*ckH_p/(2.d0*l+1.d0)*Theta(j,l-1,k) -&

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                                (1+1.d0)*ckH_p/(2.d0*1+1.d0)*Theta(j,l+1,k)+&
                                dtau(j)*Theta(j,l,k)
end do
dTheta(j,lmax_int,k) = ckH_p*Theta(j,l-1,k) -&
                        c*(1+1.d0)/(H_p(j)*eta_precomp(j))*&
                        Theta(j,l,k) + dtau(j)*Theta(j,l,k)
dPsi(j,k) = -dPhi(j,k) - 12.d0*H_0**2/(ck_current*a_t(j))**2*&
            Omega_r*(-2.d0*Theta(j,2,k)+dTheta(j,2,k))
end do
!write(*,*) 'today'
end do
deallocate(y_tight_coupling)
deallocate(y)
deallocate(dydx)
end subroutine integrate_perturbation_eqns

subroutine derivs_tc(x,y_tc, dydx)
  use healpix_types
  implicit none
  real(dp),          intent(in)  :: x
  real(dp), dimension(:), intent(in)  :: y_tc
  real(dp), dimension(:), intent(out) :: dydx

  real(dp) :: d_delta
  real(dp) :: d_delta_b
  real(dp) :: d_v
  real(dp) :: q,R

  real(dp) :: delta,delta_b,v,v_b,Phi,Theta0,Theta1,Theta2
  real(dp) :: Psi,dPhi,dTheta0,dv_b,dTheta1

  delta = y_tc(1)
  delta_b = y_tc(2)
  v = y_tc(3)
  v_b = y_tc(4)
  Phi = y_tc(5)
  Theta0 = y_tc(6)
  Theta1 = y_tc(7)

  Theta2 = -20.d0*ckH_p/(45.d0*dtau(j))*Theta1

  R = (4.d0*Omega_r)/(3.d0*Omega_b*a_t(j))

  Psi = -Phi - 12.d0*(H_0/ck_current/a_t(j))**2.d0*Omega_r*Theta2

  dPhi = Psi - ckH_p**2/3.d0*Phi + (H_0/H_p(j))**2/2.d0*&
        (Omega_m/a_t(j)*delta + Omega_b/a_t(j)*delta_b + &
        4.d0*Omega_r/a_t(j)**2*Theta0)

```

```

dTheta0   = -ckH_p*Theta1 - dPhi

d_delta   = ckH_p*v    - 3.d0*dPhi

d_delta_b = ckH_p*v_b - 3.d0*dPhi

d_v       = -v -ckH_p*Psi

q          = ( -(1.d0-2.d0*R)*dtau(j) + (1.d0+R)*ddtau(j)) *&
              (3.d0*Theta1+v_b) - ckH_p*Psi +(1.d0-dH_p(j)/H_p(j))*&
              ckH_p*(-Theta0 + 2.d0*Theta2) - ckH_p*dTheta0) / &
              ((1.d0+R)*dtau(j)+dH_p(j)/H_p(j) -1.d0)

dv_b      = (1.d0/(1.d0+R)) *(-v_b - ckH_p*Psi + &
              R*(q+ckH_p*(-Theta0 + 2.d0*Theta2)-ckH_p*Psi))

dTheta1   = (1.d0/3.d0)*(q-dv_b)

dydx(1) = d_delta
dydx(2) = d_delta_b
dydx(3) = d_v
dydx(4) = dv_b
dydx(5) = dPhi
dydx(6) = dTheta0
dydx(7) = dTheta1
!write(*,*) 'dydx(1) =',dydx(1)
!write(*,*) 'dydx(2) =',dydx(2)
end subroutine derivs_tc

subroutine derivs(x,y, dydx)
  use healpix_types
  implicit none
  real(dp),          intent(in)  :: x
  real(dp), dimension(:), intent(in)  :: y
  real(dp), dimension(:), intent(out) :: dydx

  real(dp) :: d_delta
  real(dp) :: d_delta_b
  real(dp) :: d_v
  real(dp) :: q,R
  integer(i4b) :: i
  real(dp) :: delta,delta_b,v,v_b,Phi,Theta0,Theta1,Theta2,Theta3,Theta4,Theta5,Theta6
  real(dp) :: Psi,dPhi,dTheta0,dv_b,dTheta1,dTheta2

  delta   = y(1)
  delta_b = y(2)
  v       = y(3)
  v_b     = y(4)
  Phi     = y(5)

```

```

Theta0 = y(6)
Theta1 = y(7)
Theta2 = y(8)
Theta3 = y(9)
Theta4 = y(10)
Theta5 = y(11)
Theta6 = y(12)

R      = (4.d0*Omega_r)/(3.d0*Omega_b*a_t(j))

Psi     = -Phi - 12.d0*(H_0/ck_current/a_t(j))**2.d0*Omega_r*Theta2

dPhi    = Psi - ckH_p**2/3.d0*Phi + (H_0/H_p(j))**2/ &
          2.d0*(Omega_m/a_t(j)*delta + Omega_b/a_t(j)* &
          delta_b + 4.d0*Omega_r/a_t(j)**2*Theta0)

dTheta0 = -ckH_p*Theta1 - dPhi

d_delta = ckH_p*v      - 3.d0*dPhi

d_delta_b = ckH_p*v_b - 3.d0*dPhi

d_v      = -v -ckH_p*Psi

dv_b     = -v_b -ckH_p*Psi +dtau(j)*R*(3.d0*Theta1+v_b)

dTheta1  = ckH_p/3.d0*Theta0 -2.d0/3.d0*ckH_p*Theta2 + &
          ckH_p/3.d0*Psi +dtau(j)*(Theta1+v_b/3.d0)
dTheta2  = 1/(2.d0*1+1)*ckH_p*Theta1 - (1+1.d0)/&
          (2.d0*1+1.d0)*ckH_p*Theta3+dtau(j)*0.9d0*Theta2

do i=3,lmax_int-1
    dydx(6+i) = 1/(2.d0*1+1)*ckH_p*y(5+i) - &
              (1+1.d0)/(2.d0*1+1.d0)*ckH_p*y(7+i) +dtau(j)*y(6+i)
end do

dydx(12) = ckH_p*Theta5 -c*(1+1.d0)/H_p(j)/eta_precomp(j)*Theta6 +dtau(j)*Theta6

dydx(1) = d_delta
dydx(2) = d_delta_b
dydx(3) = d_v
dydx(4) = dv_b
dydx(5) = dPhi
dydx(6) = dTheta0
dydx(7) = dTheta1
dydx(8) = dTheta2
!write(*,*) 'dydx(1) =',dydx(1)
!write(*,*) 'dydx(2) =',dydx(2)
end subroutine derivs

```



```

subroutine output3(x, y)
  use healpix_types
  implicit none
  real(dp),          intent(in)  :: x
  real(dp), dimension(:), intent(in) :: y
end subroutine output3

function get_tight_coupling_time(k)
  implicit none

  real(dp), intent(in)  :: k
  real(dp)          :: get_tight_coupling_time
  integer(i4b)      :: i,n
  real(dp)          :: x
  n = 1d4
  do i=0,n
    x = x_init + i*(0.d0-x_init)/n
    !write(*,*) x,x_start_rec
    if (x < x_start_rec .and. &
        abs(c*k/(get_H_p(x)*get_dtau(x))) <= 0.1d0 .and.&
        abs(get_dtau(x)) > 10.d0) then
      get_tight_coupling_time = x
    end if
  end do
end function get_tight_coupling_time

end module evolution_mod

```