

# Home exam FYS3140

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## Contents

<b>1</b>	<b>Problem 1: Fourier transforms and special functions</b>	<b>2</b>
1.a	.....	2
1.b	.....	3
<b>2</b>	<b>Problem 2: Contour integral</b>	<b>5</b>
2.a	.....	5
2.b	.....	6
2.c	.....	6
2.d	.....	7
2.e	.....	8
<b>3</b>	<b>Problem 3: Fröbenius method</b>	<b>9</b>
3.a	.....	9
3.b	.....	10
3.c	.....	11

# 1 Problem 1: Fourier transforms and special functions

Some special functions are defined in terms of infinite sums. For example, the Riemann zeta function,

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s}, (\Re(s) > 1) \quad (1)$$

and similarly the Dirichlet eta function

$$\zeta(s) = \sum_{n=1}^{\infty} (-1)^{n-1} \frac{1}{n^s}, (\Re(s) > 0). \quad (2)$$

One application of Fourier series is the possibility to compute certain values of this kind of functions.

## 1.a

Compute the Fourier series of  $f(x) = x^2$  with basic interval  $-\pi < x < \pi$ .

### Solution:

Since  $f(x)$  is an even function I start by setting  $b_n = 0$  for all  $n$ . The Fourier series can then be written

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx) \quad (3)$$

and the coefficients can be found with

$$a_n = \frac{2}{\pi} \int_0^{\pi} x^2 \cos(nx) dx. \quad (4)$$

First I find  $a_0$ :

$$\begin{aligned} a_0 &= \frac{2}{\pi} \int_0^{\pi} x^2 dx \\ a_0 &= \frac{2}{\pi} \left[ \frac{x^3}{3} \right]_0^{\pi} \\ a_0 &= \frac{2}{\pi} \frac{\pi^3}{3} \\ a_0 &= \frac{2\pi^2}{3} \end{aligned}$$

The rest of the coefficients can be found by using integration by parts on the integral in equation (4), and setting  $u = x^2$ ,  $u' = 2x$ ,  $v' = \cos(nx)$ ,  $v = \frac{1}{n}\sin(nx)$ :

$$\int x^2 \cos(nx) dx = \frac{x^2}{n} \sin(nx) - \frac{2}{n} \int x \sin(nx) dx$$

Where the integral on the right hand side can be solved with integration by parts. Setting  $u = x$ ,  $u' = 1$ ,  $v' = \sin(nx)$ ,  $v = \frac{-1}{n}\cos(nx)$  gives:

$$\begin{aligned} \int x \sin(nx) dx &= -\frac{x}{n} \cos(nx) + \frac{1}{n} \int \cos(nx) dx \\ &= -\frac{x}{n} \cos(nx) + \frac{1}{n^2} \sin(nx) \end{aligned}$$

Inserting this back into (4) yields

$$\begin{aligned} a_n &= \frac{2}{\pi} \left[ \frac{x^2}{n} \sin(nx) - \frac{2}{n} \left( -\frac{x}{n} \cos(nx) + \frac{1}{n^2} \sin(nx) \right) \right]_0^\pi \\ &= \frac{2}{\pi} \left( \frac{2\pi}{n^2} \cos(n\pi) \right) \\ &= \frac{4}{n^2} \cos(n\pi) \\ &= \frac{4(-1)^n}{n^2} \end{aligned}$$

and thus we have found that

$$f(x) = x^2 = \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4(-1)^n}{n^2} \cos(nx) \quad (5)$$

## 1.b

Use your result to show that

$$\zeta(2) = \frac{\pi^2}{6} \quad (6)$$

and

$$\eta(2) = \frac{\pi^2}{12} \quad (7)$$

### Solution:

We start with  $\zeta(2)$ . Using the definition of the function:

$$\zeta(2) = \sum_{n=1}^{\infty} \frac{1}{n^2}$$

If one now takes equation (6) and sets  $x = \pi$  one finds that

$$\begin{aligned}\pi^2 &= \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4(-1)^n}{n^2} \cos(n\pi) \\ \pi^2 &= \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4(-1)^n(-1)^n}{n^2} \\ \pi^2 &= \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4}{n^2} \\ \pi^2 - \frac{\pi^2}{3} &= 4 \sum_{n=1}^{\infty} \frac{1}{n^2} \\ \sum_{n=1}^{\infty} \frac{1}{n^2} &= \frac{3\pi^2 - \pi^2}{12} \\ \sum_{n=1}^{\infty} \frac{1}{n^2} &= \frac{\pi^2}{6} \\ \zeta(2) &= \frac{\pi^2}{6}\end{aligned}$$

Which was what I wanted to show.

From the definition of  $\eta$  we have

$$\eta(2) = \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n^2}$$

Again I use equation (6) but this time I set  $x = 0$  and get the following:

$$\begin{aligned}0 &= \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4(-1)^n}{n^2} \\ -\frac{\pi^2}{3} &= 4 \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} \\ -\frac{\pi^2}{3} &= -4 \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n^2} \\ \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n^2} &= \frac{\pi^2}{12} \\ \eta(2) &= \frac{\pi^2}{12}\end{aligned}$$

Which was what I wanted to show.

## 2 Problem 2: Contour integral

In this problem we will explore the complex integral

$$I = \oint_{\gamma} \frac{dz}{1+z^n}; n \geq 2 \quad (8)$$

along the closed contour  $\Gamma = \gamma_1 + \gamma_2 + \gamma_3$  where

$\gamma_1$  : Straight line along the x-axis from the origin to  $x = R$

$\gamma_2$  : Arc of a circle of radius  $R$  from  $z = x = R$  to  $z = Re^{2\pi i/n}$

$\gamma_3$  : Straight line from  $z = Re^{2\pi i/n}$  to the origin,

where we will let  $R \rightarrow \infty$  eventually.

### 2.a

Sketch the integration contour.

**Solution:**

From  $\gamma_2$  it is clear that one starts with a half circle, and that this circle is made smaller by increasing  $n$ . Since it is clear that the integration contour is dependent on  $n$  I've chosen to make a figure for  $n = 4$ .

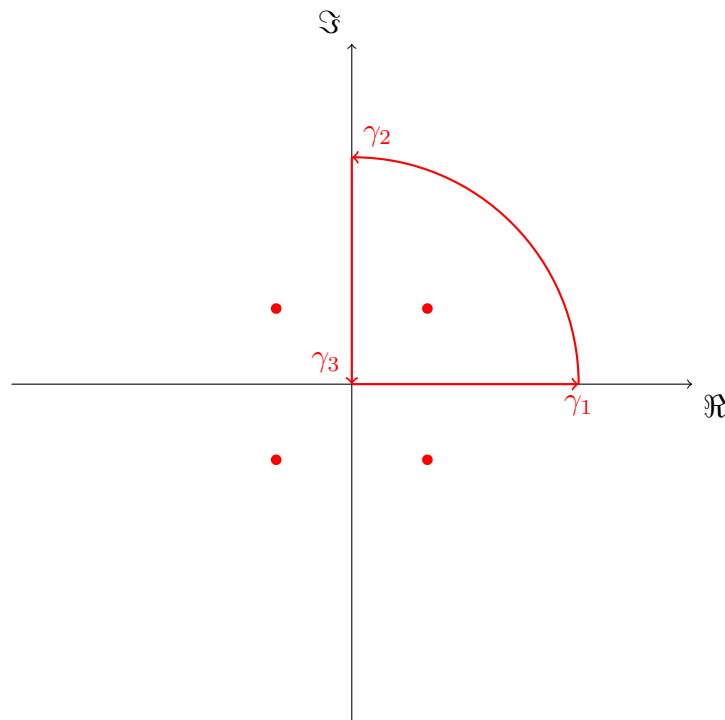


Figure 1: The integration path when  $n = 4$ .

**2.b**

Find all singularities of the integrand and indicate them in your figure from a).

**Solution:**

Since the integrand is  $\frac{dz}{1+z^n}$ , one gets a singularity when

$$z^n = -1 = e^{i\pi(1+2k)}$$

$$z = e^{\frac{i\pi(1+2k)}{n}}; k = 0, 1, \dots, n-1$$

From this it is clear that the number of, and the position of, the singularities is dependent on  $n$ . As an example I get for  $n = 2$ :

$$z_1 = e^{\frac{i\pi}{2}} = i$$

$$z_2 = e^{\frac{3i\pi}{2}} = -i$$

while for  $n = 3$  I get:

$$z_1 = e^{\frac{i\pi}{3}}$$

$$z_2 = e^{\frac{2i\pi}{3}} = -1$$

$$z_3 = e^{\frac{4i\pi}{3}}$$

In fact, there are  $n$  singularities, and if one makes a straight line from one to the next they will form an  $n$ -sided polygon. Since it would be a complete mess to mark of singularities for all  $n$ , I've chosen to mark the ones for  $n = 4$ , since that is the  $n$  I chose for the figure earlier.

**2.c**

Compute the integral  $I$  in the limit  $R \rightarrow \infty$ .

**Solution:**

From the residue theorem we have that

$$\int_{\Gamma} f(z) dz = 2\pi i \sum_N \text{Res}[f(z), z_n]. \quad (9)$$

So instead of calculating the integral  $I$  I can just find all the residues inside  $\Gamma$ . Now I need to find how many are inside  $\Gamma$  depending on what  $n$  is. For a

singularity to be inside  $\Gamma$  the following must be true:

$$\begin{aligned}\frac{2\pi}{n} &\geq \frac{\pi(1+2k)}{n} \\ 2\pi &\geq \pi(1+2k) \\ 2 &\geq 1+2k \\ \frac{1}{2} &\geq k \\ k &\leq \frac{1}{2}\end{aligned}$$

and since  $k$  increases in integer steps, this only happens when  $k = 0$ . This means that

$$\sum_N \text{Res}[f(z), z_n] = \text{Res}[f(z), z_0].$$

And this means that I can conclude that no matter what  $n$  is, there will always be a singularity inside the contour  $\Gamma$  at  $z_0 = e^{i\pi/n}$ . To find this residue I use that for rational functions  $f(z) = \frac{P(z)}{Q(z)}$ , we have

$$\text{Res}[f(z), z_0] = \frac{P(z_0)}{Q'(z_0)}. \quad (10)$$

In our case  $P(z) = 1$ , and  $Q(z) = 1 + z^n$ . Thus  $Q'(z) = nz^{n-1}$ , and

$$\begin{aligned}I &= 2\pi i \text{Res}[f(z), z_0] = 2\pi i \frac{1}{nz_0^{n-1}} = \frac{2\pi i}{n} e^{-\frac{i\pi(n-1)}{n}} \\ I &= -\frac{2\pi i}{n} e^{\frac{i\pi}{n}}\end{aligned}$$

## 2.d

Examine the contribution to the integral from the circle arc  $\gamma_2$  as  $R \rightarrow \infty$ . (Hint: perform an upper bound estimate)

**Solution:**

I want to estimate an upper bound for  $\int_{\gamma_2} \frac{1}{1+z^n} dz$ . This can be calculated by

$$\left| \int_{\gamma_2} f(z) dz \right| \leq ML \quad (11)$$

Where  $L$  is the length of the contour, and  $M$  is the maximum value of  $|f(z)|$  on the contour  $\Gamma$ .

$L$  is straightforward to find:

$$dL = R d\theta$$

$$L = \frac{R2\pi}{n}$$

Finding  $M$ :

$$\left| \frac{1}{1+z^n} \right| = \frac{|1|}{|1+z^n|} \leq \frac{1}{|z^n| - 1} = \frac{1}{R^n - 1} = M$$

$$M = \frac{1}{R^n - 1}$$

And thus

$$\left| \int_{\gamma_2} f(z) dz \right| \leq \lim_{R \rightarrow \infty} \frac{R2\pi}{n} \frac{1}{R^n - 1}$$

$$\leq \lim_{R \rightarrow \infty} \frac{2\pi}{n^2 R^{n-1}}$$

$$\leq 0$$

$$\left| \int_{\gamma_2} f(z) dz \right| \leq 0$$

$$\int_{\gamma_2} f(z) dz = 0.$$

Where I've used L'Hopital's rule to calculate the limit. So the conclusion is that  $\gamma_2$  contributes nothing to the integral as long as  $R \rightarrow \infty$ .

## 2.e

Using your results from c) and d), show that the real integral

$$\int_0^\infty \frac{dx}{1+x^n} = \frac{\pi/n}{\sin(\pi/n)} \quad (12)$$

(Hint: Find a suitable parametrization along  $\gamma_3$ ).

**Solution:**

I found earlier that

$$I = \frac{-2\pi i}{n} e^{i\pi/n}$$



but one can also write the integral for  $I$  as

$$\begin{aligned}
 I &= \int_{\gamma_1} \frac{dx}{1+x^n} + \int_{\gamma_2} \frac{dz}{1+z^n} + \int_{\gamma_3} \frac{dz}{1+z^n} \\
 &= \lim_{R \rightarrow \infty} \int_0^R \frac{dx}{1+x^n} + 0 + \lim_{R \rightarrow \infty} e^{2i\pi/n} \int_R^0 \frac{dx}{1+x^n} \\
 &= \int_0^\infty \frac{dx}{1+x^n} - e^{2i\pi/n} \int_0^\infty \frac{dx}{1+x^n} \\
 &= (1 - e^{2i\pi/n}) \int_0^\infty \frac{dx}{1+x^n}.
 \end{aligned}$$

Since I actually want to find the integral in this last expression, I set this expression equal to the one I found earlier for  $I$ :

$$\begin{aligned}
 \frac{-2\pi i}{n} e^{i\pi/n} &= (1 - e^{2i\pi/n}) \int_0^\infty \frac{dx}{1+x^n} \\
 \int_0^\infty \frac{dx}{1+x^n} &= \frac{\frac{-2\pi i}{n} e^{i\pi/n}}{1 - e^{2i\pi/n}} \\
 &= \frac{\pi}{n} \frac{2i}{e^{\frac{i\pi}{n}} - e^{\frac{-i\pi}{n}}} \\
 \int_0^\infty \frac{dx}{1+x^n} &= \frac{\pi/n}{\sin(\pi/n)}.
 \end{aligned}$$

Which was what I wanted to show.

### 3 Problem 3: Fröbenius method

Here you will use the Fröbenius method to address the differential equation

$$x^2 y'' + xy' + (x^2 - \frac{1}{4})y = 0 \quad (13)$$

#### 3.a

Set up and solve the indicial equation

**Solution:**

Start by assuming a solution of the form

$$y(x) = \sum_{n=0}^{\infty} a_n x^{n+s}$$

which has first derivative

$$y'(x) = \sum_{n=0}^{\infty} (n+s) a_n x^{n+s-1}$$

and second derivative

$$y''(x) = \sum_{n=0}^{\infty} (n+s)(n+s-1) a_n x^{n+s-2}.$$

Inserting these into the DE (14):

$$\sum_{n=0}^{\infty} (n+s)(n+s-1) a_n x^{n+s} + \sum_{n=0}^{\infty} (n+s) a_n x^{n+s} + \left(x^2 - \frac{1}{4}\right) \sum_{n=0}^{\infty} a_n x^{n+s} = 0.$$

Since this has to hold for every power of  $x$  we start with the first one,  $x^s$ :

$$\begin{aligned} s(s-1)a_0 + sa_0 + 0 - \frac{1}{4}a_0 &= 0; a_0 \neq 0 \\ s^2 - s + s - \frac{1}{4} &= 0 \\ s^2 &= \frac{1}{4} \\ s &= \pm \frac{1}{2} \end{aligned}$$

### 3.b

Find the general solution for  $y(x)$ .

**Solution:**

Since  $s_1 - s_2 = \text{integer}$ , there is a possibility that the smallest of them will generate a complete solution. Because of this I want to try using  $s = -1/2$ , and see if I can find the coefficients  $a_n$ . To do this I start with the lowest power of  $x$ , e.g.  $x^{-1/2}$ :

$$\begin{aligned} \left(-\frac{1}{2}\right)\left(-\frac{1}{2}-1\right)a_0 - \frac{1}{2}a_0 - \frac{1}{4}a_0 &= 0 \\ \left(\frac{3}{4} - \frac{3}{4}\right)a_0 &= 0 \end{aligned}$$

Which means  $a_0$  is undetermined.

Next is  $x^{1/2}$ :

$$\begin{aligned}\left(1 - \frac{1}{2}\right)\left(1 - \frac{1}{2} - 1\right)a_1 + \left(1 - \frac{1}{2}\right)a_1 - \frac{1}{4}a_1 &= 0 \\ -\frac{1}{4}a_1 + \frac{1}{2}a_1 - \frac{1}{4}a_1 &= 0 \\ \left(\frac{1}{2} - \frac{1}{2}\right)a_1 &= 0\end{aligned}$$

Which means  $a_1$  is also undetermined. And since I now have two undetermined coefficients, this has produced the complete solution.

Next thing then is to find an equation for  $a_n$  with  $n \geq 2$ :

$$\begin{aligned}\left(n - \frac{1}{2}\right)\left(n - \frac{1}{2} - 1\right)a_n + \left(n - \frac{1}{2}\right)a_n + a_{n-2} - \frac{1}{4}a_n &= 0 \\ \left(n - \frac{1}{2}\right)\left(n - \frac{3}{2}\right)a_n + \left(n - \frac{1}{2}\right)a_n - \frac{1}{4}a_n &= -a_{n-2} \\ \left(n^2 - \frac{3n}{2} - \frac{n}{2} + \frac{3}{4} + n - \frac{1}{2} - \frac{1}{4}\right)a_n &= -a_{n-2} \\ \left(n^2 - n\right)a_n &= -a_{n-2} \\ a_n &= -\frac{a_{n-2}}{n(n-1)}\end{aligned}$$

This means that the complete solution is

$$y(x) = a_0x^{-1/2} + a_1x^{1/2} - \sum_{n=2}^{\infty} \frac{a_{n-2}}{n(n-1)}x^{n-1/2} \quad (14)$$

or if one doesn't like infinite sums one can write out the start of the sum

$$y(x) = a_0x^{-1/2} + a_1x^{1/2} - \frac{a_0}{2!}x^{3/2} - \frac{a_1}{3!}x^{5/2} + \frac{a_0}{4!}x^{7/2} + \frac{a_1}{5!}x^{9/2} + \dots$$

and notice that this can be written

$$y(x) = x^{-1/2}[a_0\cos(x) + a_1\sin(x)] \quad (15)$$

### 3.c

Show that the solution can be written as

$$y(x) = AJ_{1/2}(x) + BJ_{-1/2}(x) \quad (16)$$

where A and B are constants, and  $J_{\pm 1/2}(x)$  are Bessel functions of the first kind, of order  $\pm \frac{1}{2}$ . Bessel functions are important in a number of physics applications, and information about them can be found in mathematical tables.

**Solution:**

To start with I need to know what these Bessel functions look like. It turns out that Bessel functions of the first kind of order p are defined as

$$J_p(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{\Gamma(n+1)\Gamma(n+1+p)} \left(\frac{x}{2}\right)^{2n+p} \quad (17)$$

which means that

$$J_{1/2}(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{\Gamma(n+1)\Gamma(n+3/2)} \left(\frac{x}{2}\right)^{2n+1/2} \quad (18)$$

and

$$J_{-1/2}(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{\Gamma(n+1)\Gamma(n+1/2)} \left(\frac{x}{2}\right)^{2n-1/2}. \quad (19)$$

I will also need that

$$\cos(x) = \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n}}{(2n)!} \quad (20)$$

and

$$\sin(x) = \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n+1}}{(2n+1)!}. \quad (21)$$

First I write  $y(x)$  with these other forms for  $\sin(x)$  and  $\cos(x)$ :

$$\begin{aligned} y(x) &= x^{-1/2} \left[ a_0 \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n}}{(2n)!} + a_1 \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n+1}}{(2n+1)!} \right] \\ y(x) &= a_0 \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n-1/2}}{(2n)!} + a_1 \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n+1/2}}{(2n+1)!} \\ y(x) &= a_0 \sum_{n=0}^{\infty} \frac{2^{2n-1/2} (-1)^n}{(2n)!} \left(\frac{x}{2}\right)^{2n-1/2} + a_1 \sum_{n=0}^{\infty} \frac{2^{2n+1/2} (-1)^n}{(2n+1)!} \left(\frac{x}{2}\right)^{2n+1/2} \end{aligned}$$

For this to be the same as equation (16), the following must be true:

$$\frac{2^{2n-1/2} a_0}{(2n)!} = \frac{B}{\Gamma(n+1)\Gamma(n+1/2)} \quad (22)$$

and

$$\frac{2^{2n+1/2}a_1}{(2n+1)!} = \frac{A}{\Gamma(n+1)\Gamma(n+3/2)} \quad (23)$$

Start with (22):

$$\begin{aligned} \frac{2^{2n-1/2}a_0}{(2n)!} &= 2^{-1/2}a_0 \frac{2^{2n}}{(2n)!} \\ &= 2^{-1/2}a_0 \frac{2^{2n}}{(2n)(2n-1)(2n-2)(2n-3)\dots 3 \times 2 \times 1} \\ &= 2^{-1/2}a_0 \frac{1}{(n)(n-1/2)(n-1)(n-3/2)\dots 3/2 \times 1 \times 1/2} \\ &= 2^{-1/2}a_0 \frac{\Gamma(n)\Gamma(n-\frac{1}{2})\Gamma(n-1)\dots\Gamma(1)\Gamma(\frac{1}{2})}{(n)\Gamma(n)(n-\frac{1}{2})\Gamma(n-\frac{1}{2})(n-1)\Gamma(n-1)\dots\Gamma(1)\frac{1}{2}\Gamma(\frac{1}{2})} \\ &= 2^{-1/2}a_0 \frac{\Gamma(n)\Gamma(n-\frac{1}{2})\Gamma(n-1)\dots\Gamma(1)\Gamma(\frac{1}{2})}{\Gamma(n+1)\Gamma(n+\frac{1}{2})\Gamma(n)\dots\Gamma(1)1/2\Gamma(\frac{1}{2})} \\ &= 2^{-1/2}a_0 \frac{\Gamma(1)\Gamma(\frac{1}{2})}{\Gamma(n+1)\Gamma(n+\frac{1}{2})} \\ &= 2^{-1/2}a_0 \frac{\pi^{1/2}}{\Gamma(n+1)\Gamma(n+\frac{1}{2})} \end{aligned}$$

which works if I define

$$B = a_0 \left( \frac{\pi}{2} \right)^{1/2}. \quad (24)$$

This is okay, since all I've actually done to get  $B$  is multiply  $a_0$  by some constant. Similarly for (23):

$$\begin{aligned} \frac{2^{2n+1/2}a_1}{(2n+1)!} &= 2^{1/2}a_1 \frac{2^{2n}}{(2n+1)!} \\ &= 2^{1/2}a_1 \frac{2^{2n}}{(2n+1)(2n)(2n-1)(2n-2)(2n-3)\dots 3 \times 2 \times 1} \\ &= 2^{1/2}a_1 \frac{1}{(n+1/2)(n)(n-1/2)(n-1)(n-3/2)\dots 3/2 \times 1 \times 1} \\ &= 2^{1/2}a_1 \frac{\Gamma(n+\frac{1}{2})\Gamma(n)\Gamma(n-\frac{1}{2})\dots\Gamma(1)}{(n+1/2)\Gamma(n+\frac{1}{2})(n)\Gamma(n)(n-\frac{1}{2})\Gamma(n-\frac{1}{2})\dots\Gamma(1)} \\ &= 2^{1/2}a_1 \frac{\Gamma(n+\frac{1}{2})\Gamma(n)\Gamma(n-\frac{1}{2})\Gamma(n-1)\dots\Gamma(1)}{\Gamma(n+\frac{3}{2})\Gamma(n+1)\Gamma(n+\frac{1}{2})\Gamma(n)\dots\Gamma(1)} \\ &= 2^{1/2}a_1 \frac{1}{\Gamma(n+1)\Gamma(n+\frac{3}{2})} \end{aligned}$$

which works if I define

$$A = a_1 2^{1/2}. \quad (25)$$

Notice that  $A$  is only  $a_1$  multiplied by some constant, just like  $B$ . Then finally

$$y(x) = x^{-1/2}[a_0 \cos(x) + a_1 \sin(x)] = AJ_{1/2}(x) + BJ_{-1/2}(x) \quad (26)$$

where

$$A = a_1 2^{1/2}, B = a_0 \left(\frac{\pi}{2}\right)^{1/2}.$$

Which is what I wanted to show.