

# PHYS 512 Problem Set 5

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## 1

Running the provided code with the initial hard-coded parameters gives  $\chi^2 = 15267.937150261654$ . Given 2501 degrees of freedom, the fit for the initial values has a reduced  $\chi_r^2 = \frac{\chi^2}{n_{dof}} = 6.1047329669178945$ . A good fit within uncertainty should produce  $\chi_r^2 \approx 1$ , or  $\chi^2 \approx n_{dof}$ . The initial guesses produce  $\chi^2 \approx 6 \cdot n_{dof}$  which is more than 2 times the variance  $\sigma_{\chi^2}^2 = 2n_{dof}$  away from the mean so I do not believe that the initial guesses make a good fit to the data.

Adjusting the values of the parameters to

$$\begin{aligned}H_0 &= 69 \\ \Omega_b h^2 &= 0.022 \\ \Omega_c h^2 &= 0.12 \\ \tau &= 0.06 \\ A_s &= 2.1e9 \\ n_s &= 0.95\end{aligned}$$

produces a fit with  $\chi^2 = 3272.2053559202204$ . This gives a value of  $\chi_r^2 = 1.3083587988485488$ . With the new parameters,  $\chi^2$  for the fit is approximately equal to  $n_{dof}$  so the new parameters are an acceptable fit.

## 2

I used Newton's method with numerical derivatives to fit the parameters to the data. The model is calculated from the parameter using the `get_spectrum` function defined in the example script "planck\_likelihood.py". The gradients in parameter space were computed using central differences for each parameter. Optimal step size for the central difference derivative is given by

$$dp = \left( \frac{\epsilon f}{f'''} \right)^{1/3} p$$

where  $\epsilon$  is machine precision. My function assumes that  $f \approx \frac{d^3 f}{dx^3}$  and uses step sizes  $dp = p \cdot \epsilon^{1/3}$ . In order for Newton's method to converge, I also had to omit the first line in the dataset. This is an acceptable omission since the first two values of the model (monopole, and multipole 2) are omitted by `get_spectrum`, and skipping the first data line corresponds to multipole 2. Thus, both the model and data in use start at multipole 3. The initial guesses for Newton's method were the poor fit parameters from the example code.

Parameter errors were calculated according to

$$\langle (\Delta m)^2 \rangle = (\nabla A^T N^{-1} \nabla A)^{-1}$$

where  $\nabla A$  is the gradient of the model with respect to the parameters, and  $N$  is a diagonal matrix containing the noise at each point in the data. Best fit parameters and error are saved to "planck\_fit\_params.txt" with the data in the first column (ordered as  $H_0, \Omega_b h^2, \Omega_c h^2, \tau, A_s, n_s$ ), and estimated error in the second column.

The best fit parameters from Newton's method are

$$\begin{aligned}
H_0 &= 68.94 \pm 0.12 \\
\Omega_b h^2 &= 0.022375 \pm 0.000025 \\
\Omega_c h^2 &= 0.17398 \pm 0.00025 \\
\tau &= 0.0522 \pm 0.0041 \\
A_s &= (2.078 \pm 0.016) \cdot 10^{-9} \\
n_s &= 0.96849 \pm 0.00072
\end{aligned}$$

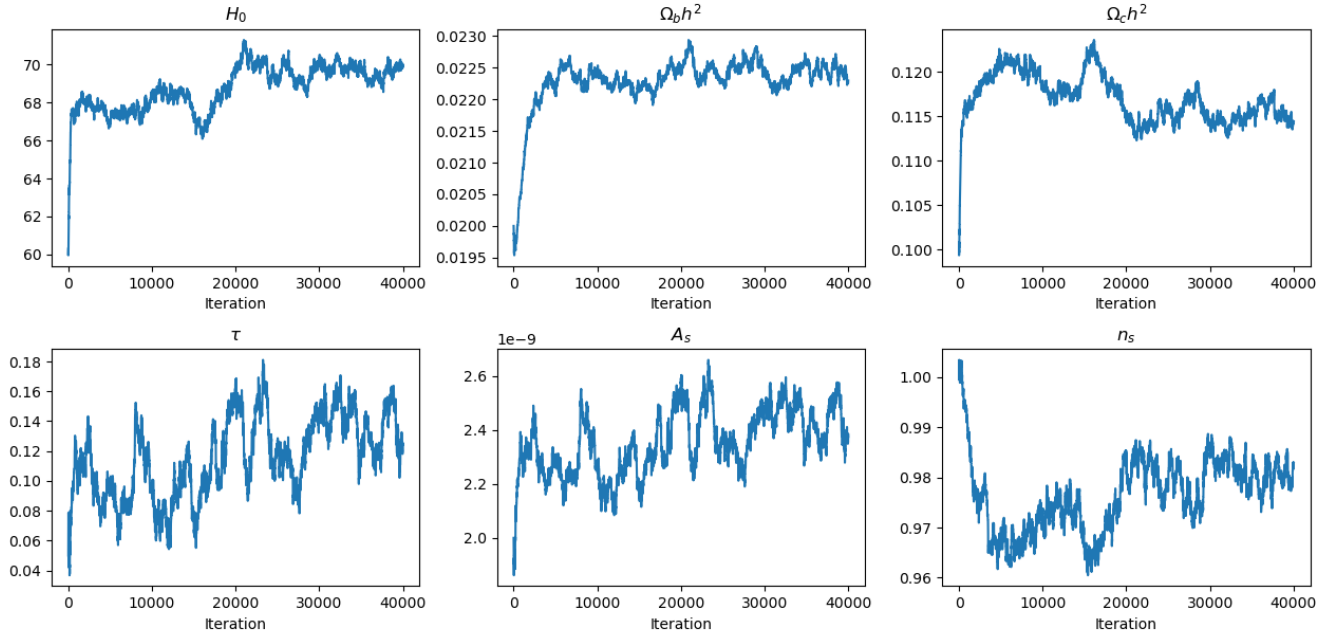
which gives a fit with  $\chi^2 = 2586.62$  for 2500 degrees of freedom. This is a good fit to the data.

time permitting, vary value of dark matter density.

### 3

The starting guess for all MCMC tries was set equal to the poor initial values given in "planck\_likelihood.py". This was chosen as a sufficiently poor initial starting position to show that the MCMC did indeed evolve towards better values. Over multiple different iterations, the best results were from an MCMC running for 40000 steps, with changes for each parameter produced by a random number drawn from a Gaussian distribution centered at 0 with variance 1, multiplied by the estimated error of the Newton's method fit. As shown in figure 1, all parameters appear to stabilize after approximately 20000 steps, so only the second half of the chains was used to estimate parameters. While they remain near a fixed value, none of the parameters appear to converge very well as all the

Figure 1: MCMC chains



chains contain some relatively large scale oscillations. The imperfect convergence of the chains is further shown by the power spectra and corner plot in figure 2. First, the power spectra do contain a flattened left-end, indicating convergence. However, the knee is located near 0.0007 giving an independent sample at approximately every 1400 steps. Over the more stable 20000 steps kept, this gives only 14 independent samples. Furthermore, the pairwise comparisons between parameter spreads shown in the corner plots do not show tight, uncorrelated convergences for

each parameter. For example,  $\tau$  and  $A_s$  appear to have a strong positive correlation, and  $H_0$  and  $\Omega_c h^2$  have an apparent negative correlation. Given the appearances and behaviours of my chains, I do not believe that they are properly converged. This may be due to the fact that  $\tau$  almost immediately jumps to around twice it's expected value and does not explore other regions in parameter space. Thus, with  $\tau$  staying near a poor value, the other parameters are limited in their convergence to a best fit and may be able to vary more freely with less "steep" likelihood controls. Despite  $\tau$  not converging near the expected value, the chains were still able to produce a good fit to the data with  $\chi^2$  values around 2580 for 2501 degrees of freedom throughout most of the fit (figure 3 shows  $\chi^2$  values for the last 20000 steps). Best fit parameters and uncertainty were calculated using the mean and standard deviation of the last 20 000 steps of the MCMC (see print outputs in figure 6).

$$\begin{aligned} H_0 &= 69.69 \pm 0.49 \\ \Omega_b h^2 &= 0.02245 \pm 0.00016 \\ \Omega_c h^2 &= 0.1151 \pm 0.0012 \\ \tau &= 0.132 \pm 0.020 \\ A_s &= (2.422 \pm 0.091) \\ n_s &= 0.9804 \pm 0.0037 \end{aligned}$$

Comparing the CAMB model made using the MCMC parameters to the data gives  $\chi^2 = 2571.45$  for 2501 degrees of freedom which is a slightly better fit than achieved with Newton's method.

Assuming the universe is flat:  $\Omega_b + \Omega_c + \Omega_\Lambda = 1$ . Therefore, the mean value of dark energy  $\Omega_\Lambda$  can be estimated from my MCMC fit:

$$\begin{aligned} \Omega_\Lambda &= 1 - \Omega_b - \Omega_c = 1 - \frac{\Omega_b h^2}{h^2} - \frac{\Omega_c h^2}{h^2} = 1 - \frac{\Omega_b h^2}{(H_0/100)^2} - \frac{\Omega_c h^2}{(H_0/100)^2} \\ \Omega_\Lambda &\approx 0.7168 \pm 0.0043 . \end{aligned}$$

Code for this question is shared between two python files. The MCMC chain was run from "Q3\_mcmc.py", and the chain evaluation, plots, and parameter estimations were performed using "eval\_mcmc.py". Both of these python scripts also rely on the `get_spectrum` function in the provided "planck\_likelihood.py" file.

## 4

A new MCMC chain was run, this time limiting  $\tau$  to values within the range  $[0.0540 - 0.0074; 0.0540 + 0.0074]$ . This constraint was put in place by checking all potential steps to ensure  $\tau$  did not leave its allowed range. If an illegal potential step was found, new random steps were made and verified until an allowed step was found. This constraint was added to the same 40000 step MCMC chain as used in question 3 in order to run the constrained chain (figure 4).

Similarly to the plain MCMC chain, the constrained chains have more stable behaviour after approximately 20000 steps so the first half of the chains were discarded as burn-in. Unfortunately, as shown in the power spectra and corner plot in figure 5, the constrained chain does not fully converge. The power spectra for  $\tau$  and  $A_s$  have the characteristic flat left hand side shape, with an "knee bend" around 0.004. Looking at the chains in figure 4, the nice power spectrum for  $\tau$  is clearly due to the bounds set up for the parameter since  $\tau$  explores it's whole allowed range randomly, but does not converge to any smaller range. The chain for  $A_s$  does look like that of a well converged parameter as suggested by its power spectrum. However, the relationship between  $A_s$  and  $\tau$  shown in the corner plot in figure 5 shows a very strong correlation between the two variables, so the apparent convergence of  $A_s$  can be explained by the limited range of  $\tau$ . The power spectra of all the other parameters do not suggest good convergence as they have much less well defined flat sections on the left hand side with "knee bends" at best near 0.0007, same as for the unconstrained chain. The corner plot also shows strong correlations between multiple other parameters.

Despite poor convergence, calculating parameters and their uncertainty from the mean values and standard deviations of the last 20000 steps in each chain (see print outputs in figure 7) produced another good fit to the data with  $\chi^2 = 2572.73$ . The constrained MCMC best fit parameters are

$$\begin{aligned} H_0 &= 67.74 \pm 0.89 \\ \Omega_b h^2 &= 0.02222 \pm 0.00040 \\ \Omega_c h^2 &= 0.1194 \pm 0.0022 \\ \tau &= 0.0548 \pm 0.0038 \\ A_s &= (2.096 \pm 0.023) \\ n_s &= 0.9695 \pm 0.0063 . \end{aligned}$$

Unfortunately, I am not able to compare these results to the parameters derived by importance sampling of the unconstrained chain. Since  $\tau$  does not stay anywhere near the constrained value for  $\tau$ , evaluating the model using all the unconstrained parameters except switching  $\tau = 0.132 \rightarrow 0.054$  gives a very poor fit of the data with  $\chi^2 = 18229.49$  for 2501 degrees of freedom. This value is too big to be incorporated into a weighted average with weights of the form

$$w_i = e^{-0.5(\chi_i^2 - \chi_\tau^2)}$$

since evaluating any of the weights numerically leads to overflow errors.

Figure 2: MCMC convergence: Power spectra and corner plot

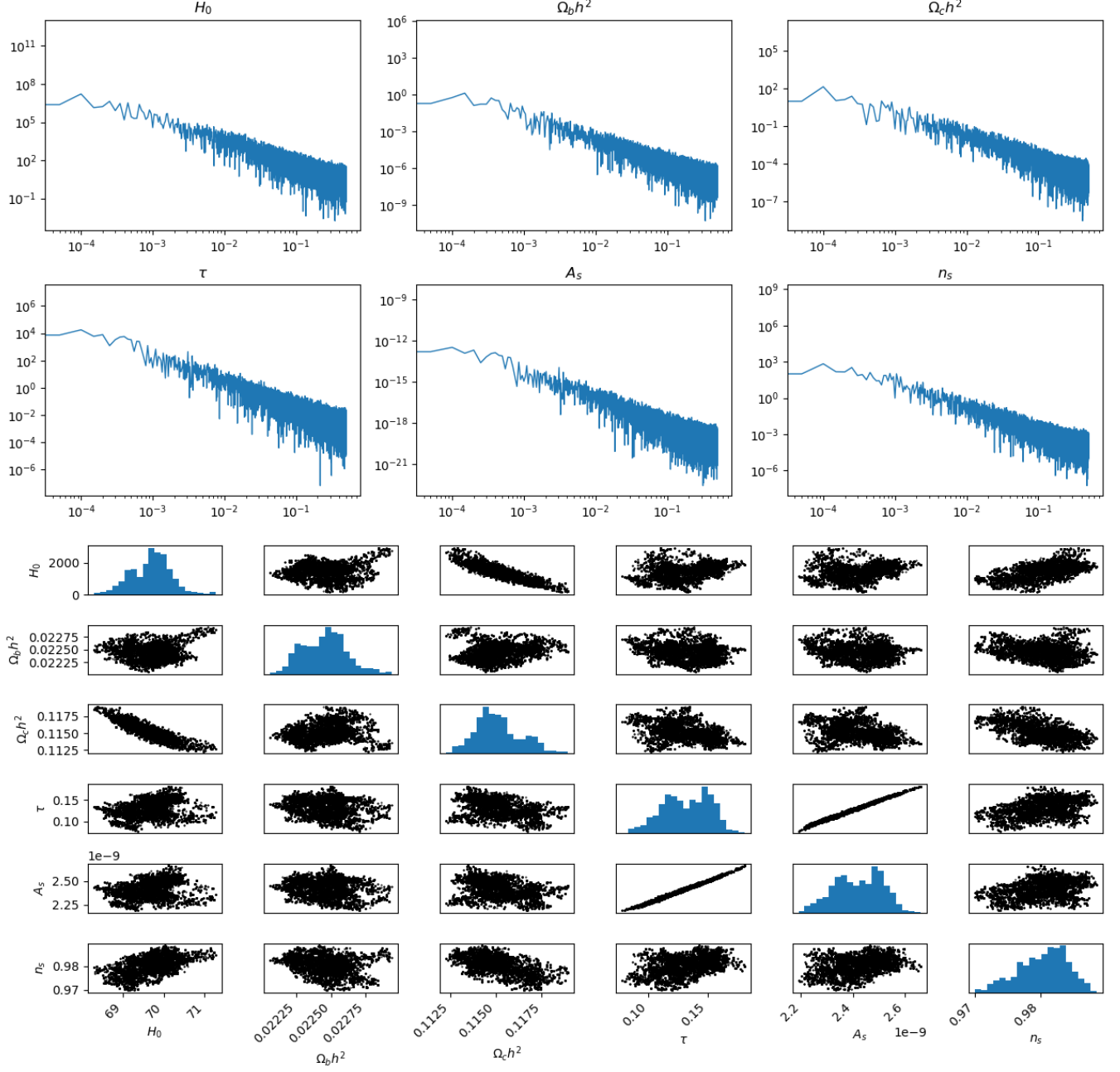


Figure 3: MCMC  $\chi^2$  values over time

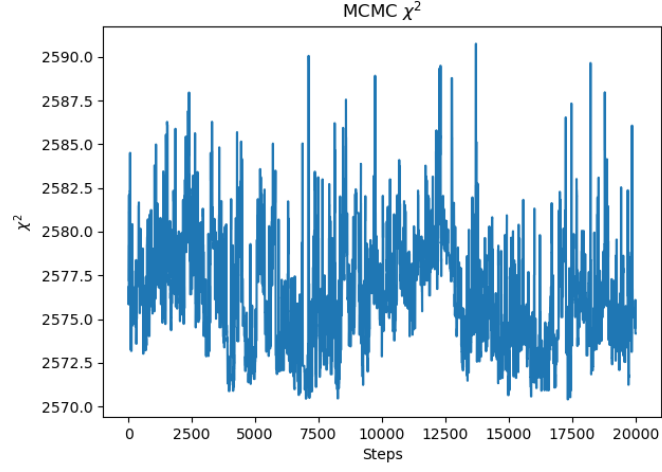


Figure 4: Constrained MCMC chains

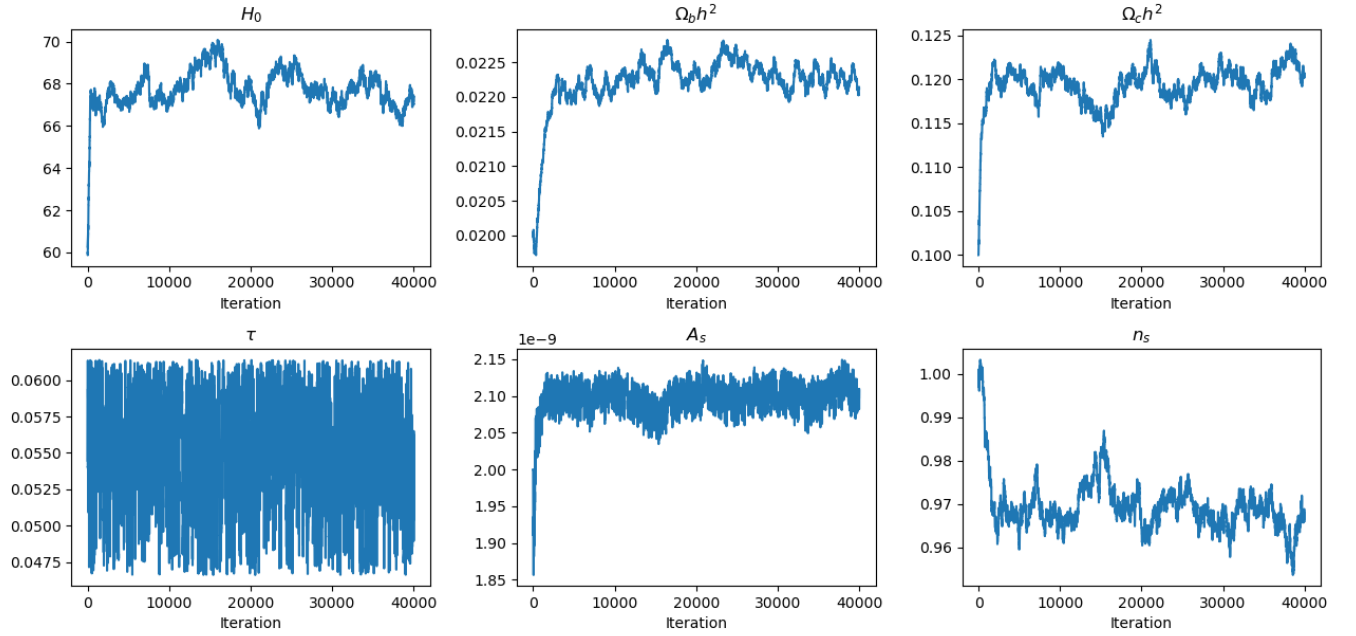


Figure 5: Constrained MCMC convergence: Power spectra and corner plot

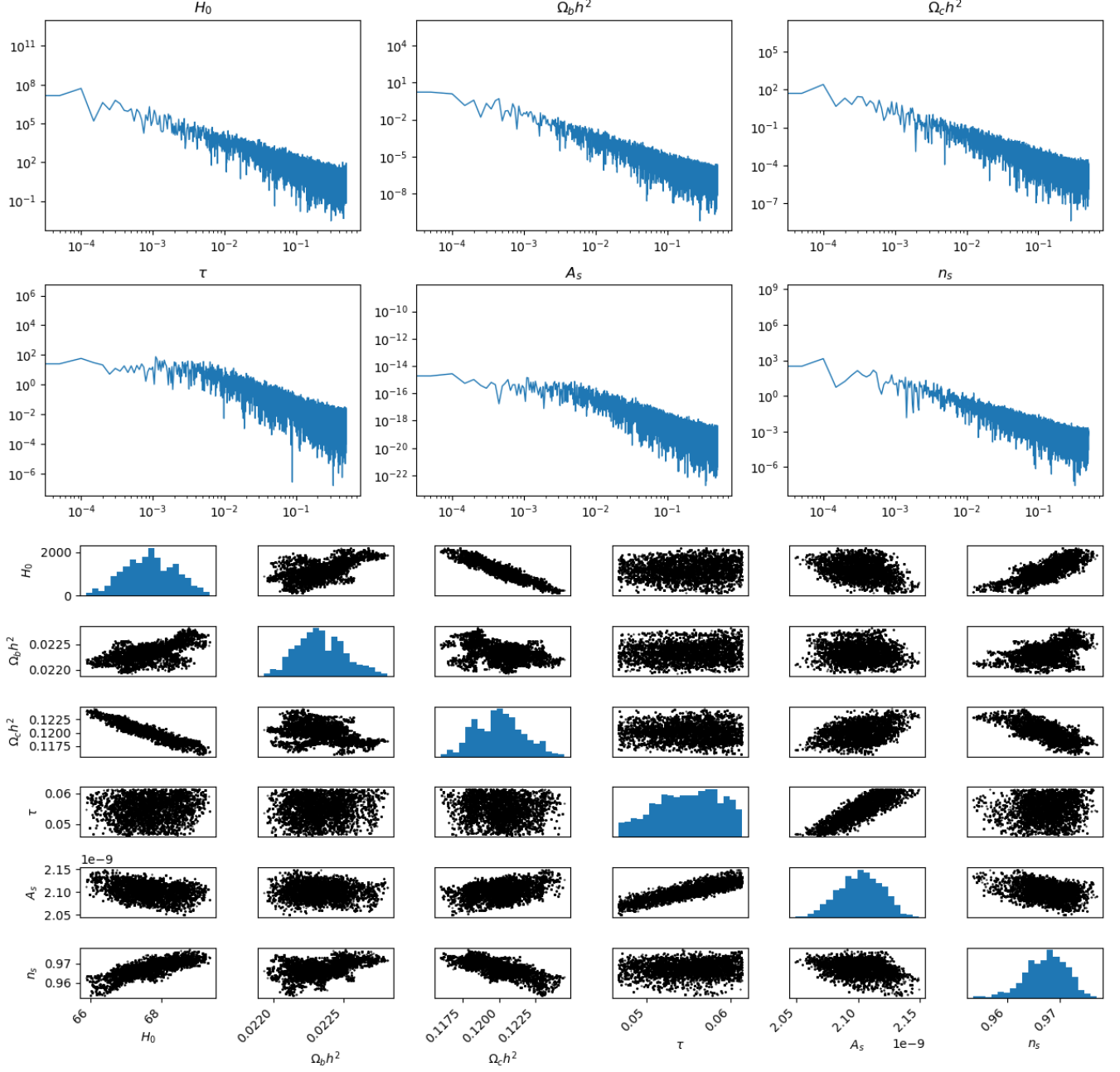


Figure 6: MCMC chain evaluation output

```
parameters:
$H_0$ 69.6979320745904 (+/-) 0.48901054081556317
$\Omega_b h^2$ 0.02245438619206941 (+/-) 0.00015876903947106252
$\Omega_c h^2$ 0.11512093113208911 (+/-) 0.0012398747871251645
$\tau$ 0.13223892049791447 (+/-) 0.019599817158860524
$A_s$ 2.421679970083134e-09 (+/-) 9.067177875394221e-11
$n_s$ 0.9804428829081606 (+/-) 0.0037464566012977625

chisq of mcmc: 2571.4527620705853
chisq for mcmc with swapped tau: 18229.48852743998
```

Figure 7: Constrained MCMC chain evaluation output

```
parameters:
$H_0$ 67.73679931543151 (+/-) 0.888046131425276
$\Omega_b h^2$ 0.022220983870678602 (+/-) 0.00040386359004929357
$\Omega_c h^2$ 0.11935703220457021 (+/-) 0.002173948105569834
$\tau$ 0.05483487657866476 (+/-) 0.0037897900742269513
$A_s$ 2.096495560930782e-09 (+/-) 2.3483633825676578e-11
$n_s$ 0.9694614944737304 (+/-) 0.00625372942360168

chisq of mcmc: 2572.731643295556
chisq for mcmc with swapped tau: 2574.482649738823
```