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Appendix

1. INTRODUCTION

Proteins are biomolecules that are fundamental to nearly all biological processes. Their diverse roles include transporting nutrients, catalyzing chemical reactions, providing structural support, and more. The function of a protein is determined by the composition of its primary sequence created from amino acids. A single protein sequence can vary greatly in length and order of its amino acids, leading to a very large number of possible configurations.

The size of the protein sequence space makes exhaustive experimental exploration infeasible. However, the rapid growth of biological databases, driven by advances in sequencing technologies, has transformed biology into a data-rich discipline. Large repositories such as UniProt, Pfam, and the Protein Data Bank store millions of sequences and structures, providing crucial resources for computational approaches[1].

This wealth of data has enabled the development of advanced statistical and machine learning models capable of simulating protein sequence evolution, prediciting structural conformations, and generating novel sequences with desired properties. In addition, breakthroughs in protein structure prediction–most notably the Nobel Prize winning AlphaFold[2] –that computational methods can rival experimental accuracy, in a much cheaper manner.

In this work, we explore the implementation of an autoregressive network for protein sequence generation leveraging the Direct Coupling Analysis (DCA) method to model residue-residue dependencies. In Section 2, the biological background will be laid out. Section 3 will introduce the mathematical foundations. Section 4 will detail previous iterations of similar models, while Section 5 will explore our implementation (and improvements).

2. Biological Background

2.1. Proteins and their structure

Proteins are essential biological molecules responsible for a wide range of functions in living organisms. Despite their functional diversity, all proteins are polymers of the same set of standard building blocks: the 20 canonical amino acids arranged in different assortments.

Amino acids share a common core structure consisting of a central carbon atom bonded to a hydrogen atom, an amino group, a carboxyl group, and a variable side chain. This side chain is the defining feature of each amino acid, giving rise to differences in size, shape, chemical reactivity, and polarity. Proteins are formed by peptide bonds of distinct amino acids, where the term polypeptides comes from. In this process, certain atoms are lost, thus, within a protein sequence, individual amino acids are typically referred to as residues. Generally, protein sequences are made up of between 50 and 2000 amino acids. The ordering of the amino acids dictates how a protein folds into its three-dimensional structure, known as its conformation. Although each conformation is unique, two common folding patterns occur in many proteins: the α helix and the β sheet [3].

Protein structure is broken down into four levels of organization. The amino acid sequence is known as the primary structure. The regularly repeating local structures stabilized by chemical bonds, such as the previously mentioned helices and sheets, make up the secondary structure. The overall three-dimensional shape of a single protein molecule constitutes the tertiary structure. This tertiary structure is what defines a protein's function. Additionally, the structure of protein molecules formed as a complex of multiple polypeptide chains is known as the quaternary structure [3].

Each amino acid is chemically distinct and can occur at any position in a protein chain, giving rise to 20^n possible polypeptide sequences of length n. For a typical protein of 300

amino acids, the number of possible sequences is astronomically large. However, only a small fraction of these sequences are capable of folding into a stable three-dimensional conformation. Natural selection has enabled living organisms to explore sequence space, favoring those sequences that reliably fold into stable structures.

2.2. Protein families and evolutionary information

Proteins do not evolve in isolation; they often belong to protein families, groups of proteins that share a common evolutionary origin, therefore exhibit similar sequence features and functional properties [4]. Over evolutionary timescales, mutations accumulate in these families: some are beneficial, altering the protein activity in ways that give rise to new functions, while many others are neutral and have no effect on stability or activity. Harmful changes, by contrast, disrupt folding or function and are eliminated by natural selection. The result is a collection of homologous proteins that retain overall structural and functional characteristics, but also display sequence variability that encodes the evolutionary history of the family [3].

The study of evolutionary history and relationships among biological entities is referred to as phylogenetics. Protein families provide the domain for phylogenetic analysis, as examining families provides insight that cannot be obtained from a single sequence. Analyzing homologous proteins across diverse organisms enables the detection of important correlated mutations between amino acid positions, representing constraints enforced by evolution. These statistical patterns are exploited by computational approaches to predict three-dimensional structure and understand protein function.

2.3. Multiple sequence alignments

To extract important evolutionary clues from protein families, the homologous sequences need to be organized in a systematic way. This is done through a multiple sequence alignment (MSA), in which three or more sequences are arranged so that homologous sites are placed in the same column [5]. To maximise the positional

correspondence of sequences with varied length, alignment gaps are introduced when necessary.

Multiple sequence alignments reveal patterns of conservation and variation across the family. Conserved positions, those which are unchanged in multiple sequences, represent sites that are critical for maintaining structure or function, while variable positions indicate sites that can tolerate mutations without disruption to the protein. Beyond conservation, MSAs also capture covariation: pairs of positions that mutate in a correlated way across sequences. These covariation signals reflect couplings, where a mutation at one site requires compensation at another to maintain protein integrity [6].

2.4. PROTEIN CONTACTS

One of the earliest challenges for researchers in computational biology–historically encouraged by the CASP (Critical Assessment of Structure Prediction) competition–has been to understand how the linear sequence of amino acids folds into its conformation. Researchers explored spatially close pairs of residues in a protein's folded three-dimensional structure, called contacts. As the structures can be represented in Cartesian coordinates (x, y, z), contacts are defined using distance thresholds. Two residues are generally considered to be in contact if the distance between the selected atoms is below a set threshold, usually 8 Ångströms (Å).

Residues that are far apart in the linear sequence, may be close together in the 3D structure. The contact distances are further categorized into short, medium, and long range predictions. Most computational approaches evaluate the long range contacts separately, as they are the most important for accurate predictions and unsurprisingly, the hardest to predict [7].

2.5. THE PROBLEM OF INFERENCE

The central inference problem in computational biology, particularly in the context of protein sequence analysis, is to disentangle the true structural and functional constraints

embedded in protein families from the noisy correlations introduced by evolutionary processes.

Correlations can arise indirectly, for example if A and B were correlated and B and C as well, A and C could appear to be correlated even with a direct interaction. Similarly, shared evolutionary history (phylogeny) and biases in sequence databases can create apparent patterns that obscure the true couplings that govern protein folding and function [8]. Distinguishing direct from indirect correlations is therefore a fundamental challenge in computational biology.

3. MATHEMATICAL FOUNDATIONS

3.1. PROTEINS AS STATISTICAL SYSTEMS

Protein sequences can be thought of as random variables produced by a certain distribution. Each sequence of length L can be written as:

$$\boldsymbol{\sigma} = (\sigma_1, \sigma_2, ..., \sigma_L), \quad \sigma_i \in \mathcal{A}, \tag{1}$$

where $\mathcal{A}$ is the alphabet of size $q = 21$ (20 amino acids and the alignment gap) and σ_i are the residue sites. We can organize these sequences into multiple sequence alignments, a table $\left\{\boldsymbol{\sigma}^{(m)}\right\}_{m=1}^{M}$ of M empirical samples. These samples are aligned to have a common length L. Each row in the MSA represents a protein, and each column a position in the sequence.

From the alignments, we can define single and pairwise frequency counts for the columns. The single-site frequency for MSA column i can be computed as:

$$f_i(k) = \frac{1}{M} \sum_{m=1}^{M} \delta\left(\sigma_i^{(m)}, A\right), \quad A \in \mathcal{A}, \quad \text{where } \delta \text{ is the Kronecker delta}^1 \tag{2}$$

while the pairwise frequency of MSA columns i, j is computed as:

$^1\delta(x, y) := \begin{cases} 0 \text{ if } x \neq y \\ 1 \text{ if } x = y \end{cases}$

$$f_{ij}(A, B) = \frac{1}{M} \sum_{m=1}^{M} \delta\left(\sigma_i^{(m)}, A\right) \delta\left(\sigma_j^{(m)}, B\right) \tag{3}$$

The empirical frequencies will serve as constraints for the distribution we want to infer.

3.2. MAXIMUM ENTROPY PRINCIPLE

To find the probability distribution $P(\boldsymbol{\sigma})$ that will satisfy constraints, in other words reproduce the empirical marginals $f_i(A)$ and $f_{ij}(A, B)$, we can use the maximum entropy principle[9], [10].

The first step to set up the MEP, is extracting information from the system. Usually, this information is given in the form of averages of functions $\langle f_k(x) \rangle$. For example, in a physical system, one could compute average magnetization or energy of the observed system. We also need to define a probability of occupancy of states, $p(x)$, which runs over all the possible states. This distribution has the usual property of mapping each state to a value within 0 and 1 and adding up to 1 when considering all states [11].

Our uncertainty on the system is expressed quantitatively through Shannon entropy S[12]:

$$S = -\sum_x p(x) \ln p(x) \tag{4}$$

Generally, the distribution which maximizes entropy is the uniform distribution. In this situation, constraints affect the probabilities of states, so the uniform is not suitable. The constraints on our system are:

$$\sum_x p(x) = 1, \quad \text{and} \quad \sum_x p(x) f_k(x) = \langle f_k \rangle \quad (k = 1, ..., m). \tag{5}$$

Selecting Shannon entropy as our measure of information, we can introduce Lagrange Multipliers $\lambda_0, \lambda_1, ... \lambda_m$ to maximize Eq. 4 subject to the constraints Eq. 5, yielding:

$$p(x) = \exp\left(-\lambda_0 - \sum_{k=1}^{m} \lambda_k f_k(x)\right). \tag{6}$$

Define the partition function Z as:

$$Z(\lambda) = \sum_x \exp\left(-\sum_{k=1}^{m} \lambda_k f_k(x)\right). \tag{7}$$

With normalization $\lambda_0 = \ln Z$, we can write the moments as:

$$\langle f_k \rangle = -\frac{\partial}{\partial \lambda_k} \ln Z(\lambda) \tag{8}$$

The entropy of the distribution then reduces to[2]:

$$S_{\max} = \lambda_0 + \sum_{k=1}^{m} \lambda_k \langle f_k \rangle \tag{9}$$

The maximum-entropy (MaxEnt) distribution for this system is exactly the Potts model.

3.3. CONNECTION TO STATISTICAL MECHANICS

The Potts model is defined in the context of statistical mechanics, where it was introduced as a generalization of the Ising spin model. Both models were created to explore ferromagnetism and phase transitions.

3.3.1. ISING MODEL

The Ising model [13] describes a system of spins σ_i arranged on a lattice. Each spin can take one of two possible values:

$$\sigma_i \in \{+1, -1\} \tag{10}$$

The Hamiltonian function, which represents the energy, of a spin configuration $\{\sigma_i\}$ is

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[2]The full derivation can be found in Appendix A.

$$H_{\text{Ising}}(\{\sigma\}) = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i, \tag{11}$$

where J is the coupling constant calculated over nearest-neighbor pairs (represented by $\langle i,j \rangle$) that defines the strength between paired interactions. The h term is the external magnetic field acting on the spins.

At thermal equilibrium, the probability of a configuration is given by the Boltzmann distribution:

$$P(\{\sigma\}) = \frac{1}{Z} \exp\left(-\beta H_{\text{Ising}}(\{\sigma\})\right), \tag{12}$$

with $\beta = 1/(k_b T)$ and partition function

$$Z = \sum_{\{\sigma\}} \exp\left(-\beta H_{\text{Ising}}(\{\sigma\})\right). \tag{13}$$

3.3.2. Potts Model

The Potts model [14] generalizes the Ising model by allowing each spin to take on q possible states. The q-state Potts model states take values in $\{1, 2, ...q\}$. The Hamiltonian is written as

$$H_{\text{Potts}}(\{\sigma\}) = -J \sum_{\langle i,j \rangle} \delta\left(\sigma_i, \sigma_j\right). \tag{14}$$

J again represents the ferromagnetic coupling that encourages neighboring spins to align in the same state. The partition function becomes

$$Z = \sum_{\{\sigma\}} \exp\left(-\beta H_{\text{Potts}}(\{\sigma\})\right). \tag{15}$$

The Potts model simplifies to an Ising model when $q = 2$.

Putting this together, the probability distribution from the Potts model is written as:

$$P(\boldsymbol{\sigma}) = \frac{1}{Z}\exp(-\beta H_{\text{Potts}}(\{\sigma\})). \tag{16}$$

It is important to note–anticipating further exploration–that the Potts model is over-parametrized: many different (h, J) sets define exactly the same distribution. Without a gauge choice, the parameters will not be uniquely identifiable and the norms of J can be misleading, hindering the process of optimization. Some common gauges, which will be explored in detail further, are reference states, zero-sum, and defining it implicitly via regularization.

3.4. DIRECT COUPLING ANALYSIS

The statistical dependencies observed in MSAs encode important information about protein structure and function. However, naive measures such as covariance cannot distinguish between direct correlations, which arise from true physical or functional contacts, and indirect correlations, propagated via other residues. Direct Coupling Analysis (DCA) addresses this problem by introducing an explicit generative model for the sequences.

DCA assumes that sequences in the MSA are independent realizations of a Potts model distribution:

$$P(\boldsymbol{\sigma}) = \frac{1}{Z}\exp\left(\sum_i h_i(\sigma_i) + \sum_{i<j} J_{ij}(\sigma_i, \sigma_j)\right) \tag{17}$$

where h_i are site-specific fields and J_{ij} are coupling parameters between residue pairs. The central goal of DCA is to infer the interaction parameters J_{ij} that best explain the observed single- and pairwise frequencies.

In practice, direct inference of these parameters is computationally challenging. Evaluating the partition function Z is intractable for realistic proteins, as it requires summing

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over q^L possible sequences. Thus, distinct implementations of DCA introduce approximations to bypass this computation. The original message passing implementation [15], directly attempted to solve for couplings through a slowly converging iterative scheme, making it inapplicable in practice. In this paper, we will explore some important algorithmic innovations in DCA following its conception.

4. Evolution of DCA Methods

4.1. Mean-field DCA (2011)

The mean-field Direct Coupling Analysis (mfDCA) algorithm, introduced by Morcos et al. [16], provided the first computationally feasible approximation of the Potts model. The idea is to approximate weak correlations via a small-coupling expansion, which reduces the inference problem to the inversion of a correlation matrix.

4.1.1. Method

To begin, the raw frequency counts suffer from sampling bias, so the weight of highly similar sequences is reduced. Each sequence A^a is assigned a weight

$$m^a = |\{b \in \{1, ..., M\} | \text{sim}(A^a, A^b) > x\}|, \quad x \approx 0.8 \tag{18}$$

where M is the number of sequences in the MSA and sim is their similarity[3]. The effective weight of a sequence a is $1/m^a$, and the total effective number of sequences is

$$M_{\text{eff}} = \sum_{a=1}^{M} 1/m^a. \tag{19}$$

From these weights, the new regularized empirical single-site frequency is computed as

$$f_{i(A)} = \frac{1}{M_{\text{eff}} + \lambda} \left(\frac{\lambda}{q} + \sum_{a=1}^{M} \frac{1}{m^a} \delta(A, A_i^a) \right) \tag{20}$$

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[3]The original paper used "seqid" to represent percentage identity. We chose "sim" as the future methods adopted this notation.

and the pairwise frequency as

$$f_{ij}(A, B) = \frac{1}{M_{\text{eff}} + \lambda} \left(\frac{\lambda}{q^2} + \sum_{a=1}^{M} \frac{1}{m^a} \delta(A, A_i^a) \delta\left(B, A_j^a\right) \right) \tag{21}$$

where $q = 21$ is the amino acid alphabet with the gap symbol, and λ is a pseudocount parameter used for regularization.

The gauge choice employed in mfDCA is defining a reference state. Typically, it is set as the last amino acid $A = q$. This gives us

$$\forall i, j : J_{ij}(a, q) = J_{ij}(q, a) = h_i(q) = 0 \tag{22}$$

To circumvent the intractable Z computation, mfDCA assumes weak correlations between sites, expanding the exponential in Eq. 17 by the Taylor series to first order[17], [18]. This results in the relation

$$C_{ij}(A, B) = f_{ij}(A, B) - f_i(A) f_j(B). \tag{23}$$

between the couplings and the connected correlation matrix. The couplings are then approximated as through naive mean-field inversion

$$e_{ij}(A, B) = -\left(C^{-1}\right)_{ij}(A, B), \tag{24}$$

where C is treated as a $((q-1)L) \times ((q-1)L)$ matrix, and the pair (i, A) is regarded as a single index. The full derivation can be found in Appendix B.

In practice, the correlation matrix is often singular without regularization. mfDCA opts to use a strong pseudocount ($\lambda \approx M_{\text{eff}}$) to stabilize the inversion and prevent spurious large couplings.

Once the estimate of the pair couplings $e_{ij}(A, B)$ is inferred, we need a way to rank residue pairs by their interaction strength. The $(q-1){\times}(q-1)$-dimensional coupling matrices need to map to a single scalar parameter. We can do this through the direct information (DI) [15]. To begin, for each pair (i, j) a two-site model is constructed:

$$P_{ij}^{(\text{dir})}(A, B) = \frac{1}{Z_{ij}} \exp\big(e_{ij}(A, B) + \tilde{h}_i(A) + \tilde{h}_j(B)\big), \tag{25}$$

with auxiliary fields $\tilde{h}_i, \tilde{h}_j$ chosen such that

$$\sum_{B=1}^{q} P_{ij}^{(\text{dir})}(A, B) = f_i(A), \quad \sum_{A=1}^{q} P_{ij}^{(\text{dir})}(A, B) = f_j(B). \tag{26}$$

The direct information is the mutual information associated to this distribution:

$$\mathcal{S}_{ij}^{\text{DI}} = \sum_{A,B=1}^{q} P_{ij}^{(\text{dir})}(A, B) \ln \frac{P_{ij}^{(\text{dir})}(A, B)}{f_i(A) f_j(B)} \tag{27}$$

The top-scoring pairs are predicted to be structural contacts.

4.1.2. LIMITATIONS

While mfDCA represented a breakthrough in usability of DCA, the simplifying approximations impose limitations on the model:

- Weak coupling assumptions: the small-coupling expasion assumes nearly linear correlations, which can underestimate strong epistatic effects in proteins. {add in-text}
- Computational scaling: the inversion of the correlation matrix scales as $\mathcal{O}(L)^3$, which is costly for very large MSAs.
- Pseudocount dependence: due to the algorithm's vast parameter size (around $400N^2$) strong regularization is required. This makes the choice of the pseudocount λ significantly affect performance.

4.2. Pseudo-likelihood Maximization DCA (2013)

The pseudolikelihood maximization DCA (plmDCA) algorithm replaces the intractable full-likelihood fit of the Potts model (which earlier methods approximated) with a tractable product of conditional likelihoods. Concretely, the inverse Potts problem on an alignment of length L becomes L coupled multinomial logistic regressions rather than a single optimization over the global partition function.

4.2.1. Method

We retain the Potts parametrization of fields and couplings introduced in Eq. 17. To mitigate redundacy, sequences are reweighted using an empirically found identity threshold $x = 0.9$. For sequence m, let

$$w_m = \frac{1}{m_m}, \quad \text{with } m_m = |\{a : \text{sim}(\sigma^{(a)}, \sigma^{(b)}) \geq x\}|, \tag{28}$$

and define the effective sample size $M_{\text{eff}} = \sum_{m=1}^{M} w_m$.

Method Setup

The pivotal idea is to optimize pseudolikelihoods. For site r, the conditional distribution given all other sites $\sigma_{\backslash r}$ is

$$P\left(\sigma_r = l \mid \sigma_{\backslash r}\right) = \frac{\exp\left(h_r(l) + \sum_{i \neq r} J_{ri}(l, \sigma_i)\right)}{\sum_{k=1}^{q} \exp\left(h_r(k) + \sum_{i \neq r} J_{ri}(k, \sigma_i)\right)}). \tag{29}$$

The weighted sitewise negative log-pseudolikelihood is

$$g_r(h_r, J_r) = -\frac{1}{M_{\text{eff}}} \sum_{m=1}^{M} w_m \log P\left(\sigma_r = \sigma_r^{(m)} | \sigma_{\backslash r} = \sigma_{\backslash r}^{(m)}\right). \tag{30}$$

and the global objective aggregates these points:

$$\mathcal{L}_{\text{pseudo}(h,J)} = \sum_{r=1}^{L} g_{r(h_r, J_r)}. \qquad (31)$$

To curtail overfitting, we add convex l_2 penalties,

$$R_{\ell_2} = \lambda_h \sum_{r=1}^{L} \|h_r\|_2^2 + \lambda_J \sum_{1 \leq i < j \leq L} \|J_{ij}\|_2^2 \qquad (32)$$

and minimize

$$\{h^{\text{PLM}}, J^{\text{PLM}}\} = \arg \min_{h,J} \{\ell_{\text{pseudo}}(h, J) + R_{\ell_2}(h, J)\} \qquad (33)$$

The ℓ_2 penalty fixes the gauge implicitly by selecting a unique representative among gauge-equivalent parameters.

Each g_r is precisely the loss of a multinomial logistic regression (softmax): classes are the q states of σ_r; features are one-hot encodings of $\{\sigma_i\}_{i \neq r}$ with $(L-1)(q-1)$ degrees of freedom after dropping a reference state. Hence (asymmetric) plmDCA is implemented as L independent, weighted softmax problems with $\mathcal{L}_2$ penalty (e.g. solved with L-BFGS or mini-batch SGD). Two variants are used in practice: asymmetric PLM, which fits each independently and then symmetrizes averaging:

$$\hat{J}_{ij} \leftarrow \frac{1}{2}\left(J_{ij}^{(i)} + J_{ij}^{(j)}\right) \qquad (34)$$

and symmetric (joint) PLM, which minimizes $\mathcal{L}_{\text{pseudo}}$ over all parameters at once. Results are typically similar, but optimization in the symmetric variant can be heavier.

Pair Scoring

For pair scoring, plmDCA avoids the previously defined Direct Information (DI) as it would introduce a regularization parameter for the pseudocounts, on top of the existing λ_h and λ_J. Instead it considers the Frobenius Norm (FN)

$$\left\| J_{ij} \right\|_2 = \sqrt{\sum_{k,l=1}^{q} J_{ij}(k,l)^2}. \tag{35}$$

Unlike the DI score, the FN is not independent of gauge choice so a direct decision must be made. As noted in [15], the zero-sum gauge minimizes the FN, making it the most appropriate gauge choice available. The procedure is hence (i) convert to zero-sum gauge J'_{ij}:

$$J'_{ij}(k,l) = J_{ij}(k,l) - J_{ij}(\cdot,l) - J_{ij}(k,\cdot) + J_{ij}(\cdot,\cdot), \tag{36}$$

where "·" denotes a simple average over the q states at that position; (ii) compute the Frobenius norm

$$\mathcal{S}_{ij}^{\mathrm{FN}} = \left\| J'_{ij} \right\|_2 = \sqrt{\sum_{k,l=1}^{q} \left(J'_{ij}(k,l) \right)^2}; \tag{37}$$

(iii) apply Average Product Correction (APC) adjusted from its use in [19] to reduce phylogenetic bias,

$$\mathcal{S}_{ij}^{\mathrm{CN}} = \mathcal{S}_{ij}^{\mathrm{FN}} - \frac{\mathcal{S}_{i\cdot}^{\mathrm{FN}} \mathcal{S}_{\cdot j}^{\mathrm{FN}}}{\mathcal{S}_{\cdot\cdot}^{\mathrm{FN}}}, \tag{38}$$

where $\mathcal{S}_{i\cdot}^{\mathrm{FN}}$ and $\mathcal{S}_{\cdot j}^{\mathrm{FN}}$ are row/column means and $\mathcal{S}_{\cdot\cdot}^{\mathrm{FN}}$ is the grand mean. Residue pairs (i,j) are ranked by $\mathcal{S}_{ij}^{\mathrm{CN}}$ to predict structural contacts.

4.2.2. LIMITATIONS

Despite its practical impact, plmDCA remains sensitive to sampling: accurate contact recovery still requires large M_{eff}, and sparse or biased MSAs degrade estimates. Phylogenetic and positional bias persist (reweighting and APC help but do not eliminate them), which can inflate false positives.

4.2.3. FAST PLMDCA (2014)

To make plmDCA deployable at scale, two of the original authors and a third collaborator revisited the optimization choices of the method [20]. In this second version, the asymmetric variant is used: L independent, weighted softmax regressions are computed, one per site, and the final couplings are symmetrized by averaging,

$$J_{ij} = \frac{1}{2}\left(J_{ij}^{(i)} + J_{ij}^{(j)}\right). \tag{39}$$

This decomposition reduces per-solve dimensionality and, crucially, enables trivial parallelization across CPU cores or nodes, which is the primary source of the runtime gain.

Because each J_{ij} is regularized twice (once in the regression for i and once for j), the coupling penalty parameter must be halved relative to the symmetric formulation. Typical values are $\lambda_h = 0.01$ and $\lambda_{J'} = 0.005$. Before averaging, the two independently inferred coupling blocks are each shifted into the zero-sum gauge to ensure consistency,

$$\hat{J}_{ij}(k,l) = J_{ij}(k,l) - J_{ij}(:,l) - J_{ij}(k,:) + J_{ij}(:,:). \tag{40}$$

For scoring, the method adopts the Corrected Frobenius Norm (CFN): first compute the Frobenius norm of $\hat{J}_{ij}$ excluding the gap state,

$$\mathrm{FN}_{ij} = \sqrt{\sum_{k,l \neq \text{ gap}} \hat{J}_{ij}(k,l)^2}, \tag{41}$$

then apply the Average Product Correction

$$\mathcal{S}_{ij}^{\mathrm{CFN}} = \mathrm{FN}_{ij} - \frac{\mathrm{FN}_{i:}\mathrm{FN}_{:j}}{\mathrm{FN}_{::}} \tag{42}$$

The combination of asymmetric regression and symmetrization, regularization scaling, gauge alignment, and CFN scoring yields essentially the same contact accuracy as the

original plmDCA, but at a fraction of the runtime, making large protein families and long sequences tractable in practice.

4.3. BOLTZMANN MACHINE DCA (2018)

In mfDCA accuracy was traded for speed via the small-coupling inversion; in plmDCA, a product of sitewise conditionals was optimized, avoiding the global partition function. bmDCA[21] instead proposes a solution closer to the statistical ideal: fitting the full Potts model by maximizing the true likelihood, using Monte Carlo Markov Chains to estimate intractable expectations. Improving on this, adabmDCA[22] introduces an adaptive procedure to keep sampling reliable and efficient. The result is a generative model that (by construction) matches the empirical one- and two-site statistics of the reweighted MSA.

4.3.1. METHOD

bmDCA aims to infer the full Potts model defined as Eq. 17 that reproduces the single- and pair-wise frequencies of the MSA.

The training procedure follows the classical Boltzmann machine learning algorithm:

1. Sampling step. For current parameters $\{J, h\}$, estimate model frequencies p_i, p_{ij} by MCMC

2. Update step. Adjust parameters when the estimated frequencies deviate from empirical ones: $\Delta h_i \propto f_i - p_i, \quad \Delta J_{ij} \propto f_{ij} - p_{ij}$

3. Iterate until empirical and model moments match within tolerance.

Regularization is imposed to stabilize the fit and avoid overfitting.

4.3.2. LIMITATIONS

- Computational cost: Direct Boltzmann machine learning with MCMC is extremely slow for realistic protein lengths ($10^7 - 10^9$ parameters).
- Sampling difficulty: Accurate moment estimation requires long chains and careful equilibration, making the approach impractical without approximations.

- Scaling: While bmDCA is theoretically optimal (full likelihood optimization), it is limited in practice to small systems where large-scale computation is feasible.

4.3.3. ADABMDCA (2021)

adabmDCA revisits the same goal of fitting the full Potts by likelihood maximization, but introduces an adaptive MCMC procedure to make the approach more practical. Instead of relying on fixed sampling parameters, it dynamically tunes equilibration, waiting times, and chain persistance. adabmDCA pushes the bmDCA approach closer to practical usability, at the cost of a more complex implementation and still significant compute requirements.

Method

Let the MSA have length L, and M sequences $\boldsymbol{\sigma} = (\sigma_1, ..., \sigma_L)$. The maximum-entropy distribution that reproduces chosen moments is the Potts model defined in Eq. 17.

To reduce phylogenetic redundancy, assign each sequence a weight

$$w_\mu = \frac{1}{|\{a : \text{sim}(s^{(a)}, s^{(\mu)}) \geq x\}|}, \quad x \approx 0.8, \quad M_{\text{eff}} = \sum_{\mu=1}^{M} w_\mu \tag{43}$$

Empirical frequencies are then reweighted and smoothed with pseudocount λ as

$$f_i(a) = (1-\lambda)f_i^{\text{data}(a)} + \frac{\lambda}{q}, \quad f_{ij}(a,b) = (1-\lambda)f_{ij}^{\text{data}(a,b)} + \frac{\lambda}{q^2} \tag{44}$$

$$f_i^{\text{data}}(a) = \frac{1}{M_{\text{eff}}} \sum_\mu w_\mu \delta\left(\sigma_i^{(\mu)}, a\right), \quad f_{ij}^{\text{data}}(a,b) = \frac{1}{M_{\text{eff}}} \sum_\mu w_\mu \delta\left(\sigma_i^{(\mu)}, a\right)\left(\sigma_j^{(\mu)}, b\right) \tag{45}$$

A practical starting point for the system is the profile model, an independent-site Potts model where the first empirical moments are matched by means of the fields,

$$h_i^{\text{prof}}(a) = \log f_i(a) + \text{const}, \tag{46}$$

but all couplings are set to zero $J \equiv 0$. In addition, zero or user-provided initial parameters can also work.

Likelihood and moment matching

The average log-likelihood of the MSA under $P(\cdot \, | J, h)$ is

$$\mathcal{L}(J, h) = \frac{1}{M} \sum_{\mu=1}^{M} \left[\sum_i h_i\Big(\sigma_i^{(\mu)}\Big) + \sum_{i<j} J_{ij}\Big(\sigma_i^{(\mu)}, \sigma_j^{(\mu)}\Big) \right] - \log Z(J, h). \qquad (47)$$

As an exponential-family model, $\mathcal{L}$ is concave in the natural parameters, so gradient ascent converges to the unique optimum. The gradients are moment gaps:

$$\frac{\partial \mathcal{L}}{\partial h_i(a)} = f_i(a) - p_i(a), \quad \frac{\partial \mathcal{L}}{\partial J_{ij}(a,b)} = f_{ij}(a,b) - p_{ij}(a,b), \qquad (48)$$

where p_i and p_{ij} are model marginals under current (J, h). Hence the update

$$h_i^{t+1}(a) \leftarrow h_i^t(a) + \eta_h \Big[f_i(a) - p_i^{(t)}(a) \Big],$$

$$J_{ij}^{t+1}(a,b) \leftarrow J_{ij}^t(a,b) + \eta_J \Big[f_{ij}(a,b) - p_{ij}^{(t)}(a,b) \Big]. \qquad (49)$$

drives the model toward exact moment matching f = p. The obstacle is that p_i and p_{ij} are not analytically computable at scale; adabmDCA estimates them by MCMC at each epoch.

Estimating model expectations via adaptive MCMC

At training epoch t, run N_s independent Markov chains using Metropolis-Hastings [23], [24] (Gibbs sampling strategy may also be used [25]), each producing N_c samples after an equilibration period T_{eq} and with an inter-sample waiting time T_{wait}. The Monte Carlo estimators are

$$p_i^{(t)}(a) = \frac{1}{N_s N_c} \sum_{\mu=1}^{N_s N_c} \delta\left(\sigma_i^{(\mu)}(t), a\right),$$

$$p_{ij}^{(t)}(a, b) = \frac{1}{N_s N_c} \sum_{\nu=1}^{N_s N_c} \delta\left(\sigma_i^{(\mu)}(t), a\right)\delta\left(\sigma_j^{(\mu)}(t), b\right).$$

(50)

Chains may be transient, reinitialized every epoch, or persistent, initialized only at the first epoch. Equilibration is often sped up through persistence of chains. For more details on the adaptive scheme, refer to Appendix C.

Convergence and quality control

A practical convergence proxy is the differnece between the empriical and the model two-site connected correlations:

$$\varepsilon_c = \max_{i,j,a,b}\left|c_{ij}^{\text{model}}(a, b) - c_{ij}^{\text{emp}}(a, b)\right| \quad \text{where}$$

$$c_{ij}^{\text{model}} = p_{ij} - p_i p_j, \quad c_{ij}^{\text{emp}} = f_{ij} - f_i f_j$$

(51)

with a target $\varepsilon_c \approx 10^{-}2$ In addition to this, some other commonly used diagnostics is the Pearson correlation between c^{model} and c^{emp}, one- and two-site fitting errors, and optionally, a third connected correlation on a subset of triples is used to assess generative fidelity beyond pairwise constraints.

Priors and sparsity

A fully connected Potts model has $\sim \frac{L(L-1)}{2}q^2 + Lq$, meaning a number in the order of 10^7-10^9 parameters for a realistic L [22]. Due to the finite sample size of MSAs, they rarely contain enough independent information to estimate all of them robustly. Not controlling this uncertainty could lead to overfitting, high variance or instability, and bad conditioning in the model. To address these issues, adabmDCA employs two complementary strategies. First, we can place a prior on $P(J, h)$ and maximize the

posterior, equivalent to adding penalties to the objective. The two standard choices are the ℓ_1 and ℓ_2 priors:

$$R_{\ell_1}(J, h) = \theta_{1,h} \sum_i \|h_i\|_1 + \theta_{1,J} \sum_{i<j} \|J_{ij}\|_1$$

$$R_{\ell_2}(J, h) = \theta_{2,h} \sum_i \|h_i\|_2^2 + \theta_{2,J} \sum_{i<j} \|J_{ij}\|_2^2$$

(52)

Under ℓ_2, the gradients are shrunk toward zero:

$$\frac{\partial}{\partial h_i(a)} : f_i(a) - p_i(a) - \theta_{2,h} h_i(a),$$

$$\frac{\partial}{\partial J_{ij}(a,b)} : f_{ij}(a,b) - p_{ij}(a,b) - \theta_{2,h} J_{ij}(a,b).$$

(53)

Under ℓ_1, they include subgradient terms that promote exact zeros.

$$\frac{\partial}{\partial h_i(a)} : f_i(a) - p_i(a) - \theta_{1,h} \ \text{sign}(h_i(a)),$$

$$\frac{\partial}{\partial J_{ij}(a,b)} : f_{ij}(a,b) - p_{ij}(a,b) - \theta_{1,h} \ \text{sign}(J_{ij}(a,b)).$$

(54)

ℓ_2 reduces variance and improves conditioning by smoothly shrinking all parameters. It selects a unique gauge and tends to preserve the relative ordering of strong couplings. On the other hand, ℓ_1 induces sparsity by zeroing weak parameters, also reducing overfitting, though at a cost of biasing small effects downward. Generally, in stochastic settings, elastic-net is used (a combination of both parameters), that stabilizes training near zero. In addition, a separate parameter is used for fields and couplings, typically regularizing J more strongly than h.

The second method is introducing sparsity via pruning or decimation. The reason for this is that true contact maps in nature are indeed sparse; most residue pairs are not in

direct physical contact. Encoding this structural prior can reduces variance and speed up learning. There are two approaches:

1. A priori topology. Reduce the number of parameters by starting from a restricted edge set, for example pairs with high MI, and learn only those J_{ij}, omitting the rest.

2. Information-based decimation. In this approach, start dense and iteratively remove the least informative couplings until target sparsity. This can be done by comparing the KL divergence for a candidate element. This directly controls overfitting by only keeping the paramets that actually affect the model's predictions.

In short, pruning and decimation prevent overfitting by removing parameters that don't materially alter the model, and have the added benefit of aligning with the biological prior of contact sparsity.

4.4. Autoregressive Network DCA

ArDCA was built to explore the ability of generative models for protein design coming from sequence-data. It was not the first generative model, DeepSequence, … ArDCA emerges as a capable model for the extraction of structural and functional protein information which encoded in rapidly growing protein databases.

4.4.1. Method

In arDCA, the exponential-family MaxEnt distribution is replaced with a conditional probability model, where each residue is predicted from the previous ones. This arises from the chain rule decomposition of the joint probability distributions:

$$P(\boldsymbol{\sigma}) = \prod_{i=1}^{L} P(\sigma_i | \sigma_1, ..., \sigma_{i-1}) \tag{55}$$

Parameter inference

The inference of the parameters is done through likelihood maximization. Following a Bayesian setting with a uniform prior, the optimal parameters are those that maximize the probability of the data:

$$\{J^*, h^*\} = \arg\max_{\{J,h\}} P(\mathcal{M}|\{J,h\})$$

$$= \arg\max_{\{J,h\}} \log P(\mathcal{M}|\{J,h\})$$

$$= \arg\max_{\{J,h\}} \sum_{m=1}^{M} \log \prod_{i=1}^{L} P(a_i^m|a_{i-1}^m, ..., a_1^m)$$

$$= \arg\max_{\{J,h\}} \sum_{m=1}^{M} \sum_{i=1}^{L} \log P(a_i^m|a_{i-1}^m, ..., a_1^m)$$

(56)

Each $h_{i(a)}$ and $J_{ij}(a,b)$ is present in only one conditional probability $P(a_i|a_{i-1}, ..., a_1)$, thus we can maximize each conditional probability indepdently in Eq. 56:

$$\{J_{ij}^*, h_i^*\} = \arg\max_{\{J_{ij}, h_i\}} \sum_{m=1}^{M} \left[h_i(a_i^m) + \sum_{j=1}^{i-1} J_{ij}\Big(a_i^m, a_j^m\Big) - \log z_i(a_{i-1}^m, ..., a_1^m) \right] \quad (57)$$

where

$$z_i(a_{i-1}, .., a_1) = \sum_{a_i} \exp\left(h_i(a_i) + \sum_{j=1}^{i-1} J_{ij}\Big(a_i, a_j\Big) \right) \quad (58)$$

is the normalization factor of the conditional probability of a_i.

Taking the derivative with respect to $h_{i(a)}$ or $J_{ij}(a,b)$, with $j = 1, ..., i-1$, we get:

$$0 = \frac{1}{M} \sum_{m=1}^{M} \left[\delta(a, a_i^m) - \frac{\partial \log z_i(a_{i-1}^m, .., a_1^m)}{\partial h_i(a)} \right]$$

$$0 = \frac{1}{M} \sum_{m=1}^{M} \left[\delta(a, a_i^m)\delta\Big(b, a_j^m\Big) - \frac{\partial \log z_i(a_{i-1}^m, .., a_1^m)}{\partial J_{ij}(a,b)} \right].$$

(59)

Using Eq. 58, we find

$$\frac{\partial \log z_i(a_{i-1}^m, .., a_1^m)}{\partial h_i(a)} = P(a_i = a \mid a_{i-1}^m, ..., a_1^m)$$

$$\frac{\partial \log z_i(a_{i-1}^m, .., a_1^m)}{\partial J_{ij}(a, b)} = P(a_i = a \mid a_{i-1}^m, ..., a_1^m)\delta\left(a_j^m, b\right). \tag{60}$$

The set of equations reduces to a simple form:

$$f_i(a) = \langle P(a_i = a \mid a_{i-1}^m, ..., a_1^m)\rangle_{\mathcal{M}},$$

$$f_{ij}(a, b) = \langle P(a_i = a \mid a_{i-1}^m, ..., a_1^m)\delta\left(a_j^m, b\right)\rangle_{\mathcal{M}}, \tag{61}$$

where $\langle \cdot \rangle_{\mathcal{M}} = \frac{1}{M}\sum_{m=1}^{M} \cdot^m$ denotes the empirical average. The first variable, $i = 1$, is unconditioned, therefore the coupling equation is $J \equiv 0$ and the equation for the field is the profile model also used in bmDCA, $h_1(a) = \log f_1(a) + \text{const}$.

Implementation details

Unlike bmDCA, the equations do not enforce exact matching between model marginals and empirical frequencies. The ability to reproduce the frequencies is thus a good proxy for the generative properties of the model, on top of the fitting quality of current parameters.

In practice, the parameters are updated using gradient descent on the likelihood. As we are able to take the expectations directly over the MSA, not needed MCMC, the gradients are exact. A commonly used optimizer for this task is the quasi-Newton optimizer (low-storage BFGS).

For regularization, arDCA makes use of ℓ_2 penalties, with different strengths depending on whether the task is generating sequences or just contact prediction. It was noted that small regularization values improved the generative quality while larger were beneficial for contact prediction. Similarly to the previous methods, sequence reweighting was also used in this algorithm.

4.4.2. Key Contributions

This lightweight approach performs at a similar accuracy of previous iterations, but at a substantially lower computational cost (a factor between 10^2 and 10^3) [1]. It presents an important innovation also due to its mathematical advantages, which will be explored further, leading to improved applicability in sequence generation and evaluation.

5. Implementation and Extension

The original ArDCA model was developed in Julia, a language designed for high-performance numerical and scientific computing. Julia's strengths lie in just-in-time (JIT) compilation, seamless handling of linear algebra, and built-in support for parallelism, making it well-suited for implementing large-scale statistical models.

In contrast, the re-implementation was carried out in Python, a widely adopted language for machine learning and data science. While Python itself is generally slower due to its interpreted nature, performance-critical operations are offloaded to highly optimized libraries (e.g., NumPy, SciPy, PyTorch), which use underlying C/C++ or Fortran code. This ecosystem makes Python particularly attractive for model development, as it provides a wide range of frameworks, pre-built optimization routines, and strong community support.

Important Differences Between Julia and Python

There are some features of the languages that makes the porting not a simple one-to-one translation. To begin, Julia uses column-major order, which akes column-wise operations very fast and easily cached. This is in complete contrast to Python which is strucutred with row-major order where row-wise operations are more optimized. The impact of this difference is that loops and reshaping operations need to be reconsidered for memory efficiency. In addition, Julia is 1-indexed whereas Python is 0-indexed. This has no effect on the performance, but it changes the bounds on loops, the indexing of objects, and the range functions.

Furthermore, loops have different behaviors in the languages. Julia loops are compiled down to efficient machine code with JIT. There exist a macro that removes the bounds checking, making loops fast without the need for vectorization. On the other hand, pure Python loops are slow. To reach the same level of efficiency vectorization is needed in Python, implemented via NumPy or broadcasting.

The final important difference is in compilation versus interpretation. Julia JIT compiles functions to machine code, which has an initial compilation cost, but allows future calls to run much faster. On the other hand, Python is interpreted, thus heavy computations require external libraries which are precompiled in other, more efficient, languages.

5.1. IMPLEMENTATION DETAILS

We implement the autoregressive network within the PyTorch ecosystem, using the `nn.Module` interface to leverage its training and optimization capabilities. The central operation of the model is the computation of the autoregressive logits, which define the conditional distribution at each sequence position. For site i, the logit vector is given by

$$z[m, i, a] = h[i, a] + \sum_{j<i} \sum_{b} J[i, j, a, b] X[m, j, b] \tag{62}$$

where h and J are the previously discussed local biases and pairwise couplings. X is a tensor holding the one-hot encoding of the symbol b observed at site j in sequence m.

To compute the autoregressive logits, our initial implementation employed a masked matrix multiplication, performed in a single step using the `einsum` operator. While correct, and more efficient than naive loops, this aproach proved computationally inefficient: the full interaction matrix was calculated, even though the contributions from the upper-triangular part were masked to zero. To address this, we devised a more efficient formulation that explicitly exploits the block-sparse structure of the coupling matrix. Specifically, we collect all lower-triangular index pairs (i, j) with $j < i$, extract the corresponding coupling blocks extracted $J[i, j]$, and accumulate their contributions

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through an `einsum` operation. The resulting implementation only computes terms required by the autoregressive factorization, removing unnecessary computation.

```python
def compute_ar_logits(self, X_oh: torch.Tensor):
    """z[m,i,a] = h[i,a] + sum_{j<i} sum_b J[i,j,a,b] * X[m,j,b]"""
    X_oh = X_oh.to(self.J.dtype)
    M, L, q = X_oh.shape
    logits = self.h_pos.unsqueeze(0).expand(M, -1, -1).clone()  # (M,L,q)


    # collect lower-triangular index pairs
    i_idx, j_idx = torch.tril_indices(L, L, offset=-1)
    J_blocks = self.J[i_idx, j_idx]    # (n_pairs, q, q)
    X_blocks = X_oh[:, j_idx]          # (M, n_pairs, q)


    # compute pairwaise contributions and accumulate into logits
    contrib = torch.einsum("mpq,pqr->mpr", X_blocks, J_blocks)
    logits = logits.index_add(1, i_idx, contrib)
    return logits
```

5.2. Computational challenges

- Benchmarks

Evaluation with Boltz-2 / AlphaFold3

6. Results and Discussion

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A Full Langrage Multipliers Calculation for Maximum Entropy Principle

$$p_i = e^{-\lambda - \mu f(x_i)}$$

$$\sum_i p_i = 1$$

$$\sum_i e^{-\lambda - \mu f(x_i)} = 1, \text{ factor out } e^{-\lambda}$$

$$e^{-\lambda} \sum_i e^{-\mu f(x_i)} = 1,$$

(63)

Define the partition function $Z(\mu)$:

$$Z(\mu) = \sum_i e^{-\mu f(x_i)}, \text{ thus}$$

$$e^{-\lambda} Z(\mu) = 1 \Rightarrow \lambda = \ln Z(\mu)$$

(64)

B Small Coupling Mean-Field Derivation

Start with the perturbed Hamiltonian

$$\mathcal{H}(\alpha) = -\alpha \sum_{1 \le i < j \le L} e_{ij}(A_i, A_j) - \sum_{i=1}^{L} h_i(A_i) \tag{1}$$

which allows interpolation between independent couplings, $\alpha = 0$, and the original model, $\alpha = 1$. We also define the Gibbs potential

$$-\mathcal{G}(\alpha) = \ln \left[\sum_{\{A_i | i=1,\dots,L\}} e^{-\mathcal{H}(\alpha)} \right] - \sum_{i=1}^{L} \sum_{B=1}^{q-1} h_i(B) P_i(B) \tag{2}$$

as the Legendre transform of the free energy $\mathcal{F} = -\ln Z$. The fields can be found via

$$h_i(A) = \frac{\partial \mathcal{G}(\alpha)}{\partial P_i(A)}, \tag{3}$$

and

$$(C^{-1})_{ij}(A, B) = \frac{\partial h_i(A)}{\partial P_j(B)} = \frac{\partial^2 \mathcal{G}(\alpha)}{\partial P_i(A) \partial P_j(B)}. \tag{4}$$

Our aim is to expand the Gibbs potential up to first order around the independent-site case $\alpha = 0$,

$$\mathcal{G}(\alpha) = \mathcal{G}(0) + \frac{d\mathcal{G}(\alpha)}{d\alpha}\bigg|_{\alpha=0} \alpha + \mathcal{O}(\alpha^2) \tag{5}$$

Independent-site approximation

To start, let us consider the Gibbs potential in $\alpha = 0$. In this case, the Gibbs potential equals the negative entropy of an ensemble of L uncoupled Potts spins,

$$\begin{aligned} \mathcal{G}(0) &= \sum_{i=1}^{L} \sum_{A=1}^{q} P_i(A) \ln P_i(a) \\ &= \sum_{i=1}^{L} \sum_{A=1}^{q-1} P_i(A) \ln P_i(a) + \sum_{i=1}^{L} \left[1 - \sum_{A=1}^{q-1} P_i(A) \right] \ln \left[1 - \sum_{A=1}^{q-1} P_i(A) \right] \end{aligned} \tag{6}$$

Mean-field approximation

To get the first order in Eq. 5, we have to determine the derivative at $\alpha = 0$. Recalling the definition of Gibbs potential in Eq. 2,

$$\begin{aligned} \frac{d\mathcal{G}(\alpha)}{d\alpha} &= \frac{-d}{d\alpha} \ln Z(\alpha) - \sum_{i=1}^{L} \sum_{A=1}^{q-1} \frac{dh_i(A)}{d\alpha} P_i(A) \\ &= -\sum_{\{A_i\}} \left[\sum_{i<j} e_{ij}(A_i, A_j) + \sum_i \frac{dh_i(A)}{d\alpha} \right] \frac{e^{-\mathcal{H}(\alpha)}}{Z(\alpha)} - \sum_{i=1}^{L} \sum_{A=1}^{q-1} \frac{dh_i(A)}{d\alpha} P_i(A) \\ &= -\left\langle \sum_{i<j} e_{ij}(A_i, A_j) \right\rangle_{\alpha} \end{aligned} \tag{7}$$

The first derivative of the Gibbs potential with respect to α this is the average of the coupling term in the Hamiltonian. At $\alpha = 0$, this average can be done easily due to the joint distribution of all variables becoming factorized over the single sites,

$$\left. \frac{d\mathcal{G}(\alpha)}{d\alpha} \right|_{\alpha=0} = -\sum_{i<j} \sum_{A,B} e_{ij}(A_i, A_j) P_i(A) P_j(B). \tag{8}$$

Plugging this and Eq. 6 into Eq. 5, we find the first-order approximation of the Gibbs potential. The first and second partial derivatives with respect to $P_i(A)$ provide self-consistent equations for the local fields,

$$\frac{P_i(A)}{P_i(q)} = \exp\left(h_i(A) + \sum_{\{j|j\neq i\}} \sum_{B=1}^{q-1} e_{ij}(A, B) P_j(B) \right) \tag{9}$$

and the inverse of the connected correlation matrix,

$$\left. (C^{-1})_{ij}(A, B) \right|_{\alpha=0} = \begin{cases} -e_{ij}(A, B) \text{ for } i \neq j \\ \frac{\delta(A,B)}{P_i(A)} + \frac{1}{P_i(q)} \text{ for } i = j \end{cases} . \tag{10}$$

This equation allows us to solve the original inference in the mean-field approximation in a single step. To determine the marginals of the empirical freqiencies, we just need to determine the empirical connected correlation matrix and invert it to get the couplings e_{ij}.

C Adaptive MCMC Sampling in adabmDCA

To ensure accurate estimation of the marginals, adabmDCA monitors both equilibration and decorrelation of the Markov chains through sequence overlaps. For two sampled sequences $s_n^i, s_m^k \in \mathcal{A}^L$, the normalized overlap is defined as

$$O(s_n^i, s_m^k) = \frac{1}{L} \sum_{j=1}^{L} \delta(s_n^i(j), s_m^k(j)) \tag{11}$$

Three types of overlaps are used:

- External overlap (O_{ext}): between samples from different chains at the same time n.

- Internal-1 overlap (O_{int1}): between consecutive samples of the same chain at times n and $n + 1$.

- Internal-2 overlap (O_{int2}): between samples of the same chain at times n and $n + 2$.

At each epoch, adabmDCA computes averages μ_α and standard errors σ_α ($\alpha \in \{\text{ext}, \text{int1}, \text{int2}\}$). In equilibrium, all three overlaps should agree within statistical error.

Adaptive update of T_{wait}

To control correlation between successive samples, T_{wait} is adjusted dynamically:

Increase T_{wait} (double it) if:

$$|\mu_{\text{ext}} - \mu_{\text{int2}}| > 5\sqrt{\sigma_{\text{ext}}^2 + \sigma_{\text{int2}}^2} \tag{12}$$

i.e. samples at lag $2T_{\text{wait}}$ are still too correlated.

Reduce T_{wait} (average with last pre-increase value) if:

$$|\mu_{\text{ext}} - \mu_{\text{int1}}| < 5\sqrt{\sigma_{\text{ext}}^2 + \sigma_{\text{int1}}^2} \tag{13}$$

i.e. decorrelation already occurs at lag T_{wait}. Equilibration time is then set to $T_{\text{eq}} = 2T_{\text{wait}}$.

This adaptive scheme ensures that chain samples are equilibrated, initial bias is removed after T_{eq} steps, and decorrelated, independence between samples at lag $T_{\text{wait}} - 2T_{\text{wait}}$. As a result, the estimated marginals match those of the equilibrium Potts distribution, guaranteeing stable gradient estimates and convergence of the Boltzmann machine.