Lecture 2: Differential and Difference Equations

Introduction

Consider a scenario where a thin bar is put on an infinite block of ice at both ends and there a blowtorch heating it along it's length. The task is to find the temperature at several locations along the bar when it reaches steady state (i.e. after a very very long time of heating).



The differential equation (with boundary conditions) explaining this phenomenon is given as

$$-rac{d^2u}{dx^2} = f(x) = 1; \qquad \quad u(0) = 0, u(1) = 0$$
 (1)

where u(x) is the temperature at location $x \in [0,1]$ and the "ice at the ends" mean the temperature at u(0) = u(1) = 0 always! The heat from blowtorch is represented using the function f(x) and it's set to a constant because it's uniform.

Let's look at this differential equation analytically. What we need is some function u(x) that satisfies $-\frac{d^2u}{dx^2}=1$. In other words, the second derivative of u(x) should be -1. A function $u(x)=-\frac{1}{2}x^2$ satisfies this differential equation. So does another function $u(x)=-\frac{1}{2}x^2+Cx+D$, where C and D can be any constants! This is where Boundary

Conditions come in and fix C, D values! This way we get an exact solution $u(x)=-rac{1}{2}x^2+rac{1}{2}x$

Boundary Conditions enables a unique solution to the PDE! Hence they are very important in solving PDEs.

Difference Equations

In order to solve the differential equation numerically, we have to represent it as a difference equation where the derivatives are defined or approximated numerically. We can approximate the second derivative in Equation (1) using Finite Difference and write the difference equation as

$$-rac{u_{i+1}-2u_i+u_{i-1}}{(\Delta x)^2}=f(x_i)=1 \hspace{1.5cm} (2)$$

Let's take a look at a few finite difference formulae approximating the 1st order derivative:

Forward Difference

$$\Delta_F u = rac{u(x+h)-u(x)}{h} pprox u' + \mathcal{O}(h)$$

Backward Difference

$$\Delta_B u = rac{u(x) - u(x-h)}{h} pprox u' + \mathcal{O}(h)$$

Centered Difference

$$\Delta_C u = rac{u(x+h) - u(x-h)}{2h} pprox u' + \mathcal{O}(h^2)$$
 (5)

These finite difference formulae are derived using Taylor Series, which says that

$$u(x+h) = u(x) + hu'(x) + \frac{h^2}{2!}u''(x) + \frac{h^3}{3!}u'''(x) + \cdots$$
 (6)

$$u(x-h) = u(x) - hu'(x) + \frac{h^2}{2!}u''(x) - \frac{h^3}{3!}u'''(x) + \cdots$$
 (7)

Rearranging Equation (6), we can get the Forward Difference approximation to the 1st order derivative:

$$rac{u(x+h)-u(x)}{h} = u'(x) + rac{h}{2!}u''(x) + rac{h^2}{3!}u'''(x) + \cdots \ = u'(x) + \mathcal{O}(h)$$
 (8)

Similarly using Equation (7), we can get the Backward Difference approximation:

$$egin{aligned} rac{u(x) - u(x - h)}{h} &= u'(x) - rac{h}{2!} u''(x) + rac{h^2}{3!} u'''(x) + \cdots \ &= u'(x) + \mathcal{O}(h) \end{aligned}$$

Finally, subtracting (7) from (6), we get Centered Difference approximation:

$$rac{u(x+h)-u(x-h)}{2h} = u'(x) + rac{2h^2}{3!}u'''(x) + \cdots \ = u'(x) + \mathcal{O}(h^2)$$
 (10)

Notice how we get 1st and 2nd order accuracy depending on the terms left after moving them around.

Now we can get 2nd order derivatives in three ways:

- $lacksquare \Delta_F \Delta_B$
- lacksquare $\Delta_B\Delta_F$
- $\Delta_C\Delta_C$: This would be a weird choice because it will stretch too far on the stencil in a way that we get coefficients as $1 \quad 0 \quad -2 \quad 0 \quad 1$.

This might sound unfamiliar, but I'll discuss discretisation and stencil in the next section.

Apart from this, we can also get $\Delta_F \Delta_B$ or $\Delta_B \Delta_F$ by adding (6) and (7) together:

$$rac{u(x+h)-2u(x)+u(x-h)}{h^2} = u''(x) + rac{2h^2}{4!}u''''(x) + \cdots \ = u''(x) + \mathcal{O}(h^2)$$
 (11)

Numerical Discretisation and Solution

The very first step to solving a differential equation numerically is discretisation of the domain. Consider the PDE in Equation (1), where the domain lies between 0 and 1. What discretisation does is instead of using a continuous domain, it uses a set of finite points on the domain to represent it (shown in diagram below). In this example, we will discretise the domain into 5 equispaced points with $\Delta x=0.25$ as the distance between two consecutive points. Also, each point is indexed using the small letter i=0,1,2,3,4.

$$n=5$$
 equispaced discrete points
$$i=0 \qquad 1 \qquad 2 \qquad 3 \qquad 4$$

$$x=0 \qquad \longleftrightarrow \qquad x=1$$

Using 5 points to discretise any practical problem might be a bad idea, but for understanding it's better to keep n small.

The problem now is to find u_0, u_1, u_2, u_3 and u_4 , as these 5 values will essentially represent the function describing u. This is where we turn our attention to Equation (2), and start setting different values to subscript i and form a set of equations.

• When i=1

$$-u_2 + 2u_1 - u_0 = (0.25)^2 (12)$$

• When i=2

$$-u_3 + 2u_2 - u_1 = (0.25)^2 (13)$$

• When i=3

$$-u_4 + 2u_3 - u_2 = (0.25)^2 (14)$$

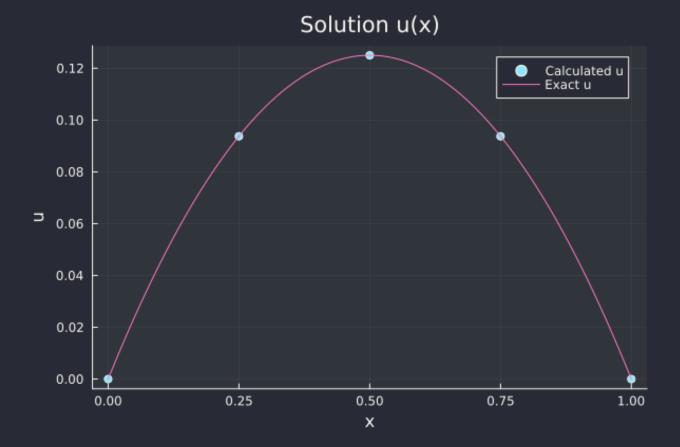
We now have 3 equations and 3 unknowns (u_1, u_2, u_3). Remember that we already know the boundary conditions which states that $u_0 = u_4 = 0$.

The above defined set of equations can also be rewritten in terms of matrix vector multiplication,

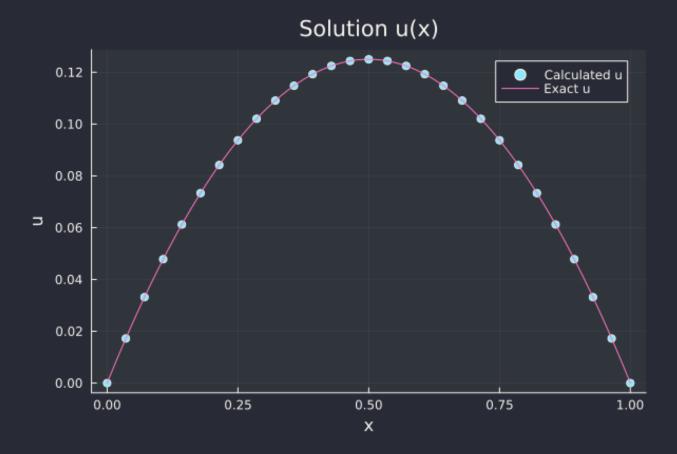
$$\begin{pmatrix} 2 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 2 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ u_3 \end{pmatrix} = \begin{pmatrix} (0.25)^2 \\ (0.25)^2 \\ (0.25)^2 \end{pmatrix}$$
(15)

If you remember from Lecture 1, the left side 3×3 matrix is K_3 and it's invertible! For now, let's just use some computer program to solve this $A\mathbf{u} = b$ and get the vector \mathbf{u} .

```
A = zeros(n, n)
    for i=1:n
        A[i,i] = 2
        if i \neq 1
            A[i, i-1] = -1
        end
        if i ≠ n
            A[i, i+1] = -1
        end
    end
    return A
end
A = coefficient(n)
b = ones(n) * (\Delta x^2)
# Solve
u = A b
println("Solution u: ", u)
pushfirst!(u, 0);
push!(u, 0);
scatter(x, u, label="u", xlabel="x", ylabel="u", title="Solution u(x)")
```



We can now increase the number of points n to get more accurate solution, for example if n=30.



Changing the Boundary Conditions

Let's change something with the problem setup and remove the ice block at x=0! What this means is that we now have a PDE with different boundary conditions

$$-rac{d^2u}{dx^2}=f(x)=1; \qquad \left.rac{\partial u}{\partial x}
ight|_{x=0}=0, u(1)=0 \qquad \qquad (16)$$

Remember that the problem is of steady state, hence the change in temperature at x=0 is set to 0!

Again setting n=5, the problem now is to find u_0,u_1,u_2,u_3 and u_4 . In this case we know that, $u_4=0$ which leaves 4 unknown variables. Again using the approximations to convert differential equation to difference equation and also approximating the boundary condition at x=0 using Equation (8) we get a set of equations as follows

Boundary Condition Approximation

$$u_1 - u_0 = 0 (17)$$

• When i=1, we can use the fact that $u_1=u_0$ and reduce the unknown variables to 3!

$$-u_2 + 2u_1 - u_0 = (0.25)^2$$
 or (18) $-u_2 + u_1 = (0.25)^2$

• When i=2

$$-u_3 + 2u_2 - u_1 = (0.25)^2 (19)$$

• When i=3

$$-u_4 + 2u_3 - u_2 = (0.25)^2 (20)$$

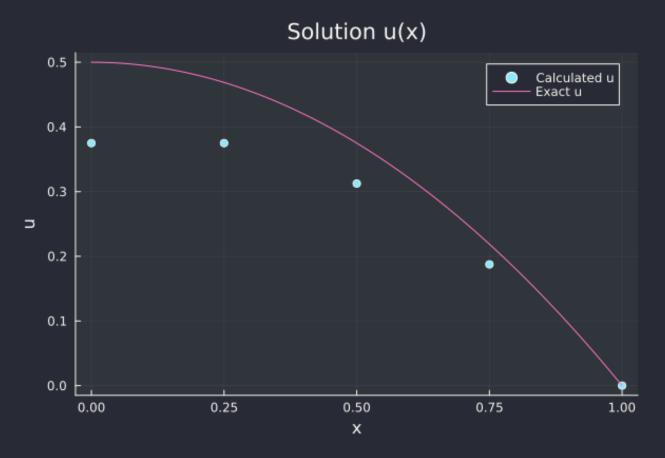
Now, we have 3 equations and 3 variables, which can again be represented as a matrix vector multiplication

$$\begin{pmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 2 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ u_3 \end{pmatrix} = \begin{pmatrix} (0.25)^2 \\ (0.25)^2 \\ (0.25)^2 \end{pmatrix}$$
(21)

Again from Lecture 1, the left side 3×3 matrix is T_3 and it's invertible! Let's again use the inbuilt matrix solver to solve this $A\mathbf{u}=b$ and get the vector \mathbf{u} .

```
function coefficient_T(n::Int)
    # A matrix
A = zeros(n, n)
for i=1:n
    # Diagonal Elements
    if i == 1
        A[i,i] = 1
    else
        A[i,i] = 2
```

```
end
    if i ≠ 1
        # Lower Diagonal
        A[i, i-1] = -1
    end
    if i ≠ n
        # Upper diagonal
        A[i, i+1] = -1
    end
    end
    return A
end
```



Similarly, if we increase the number of points, we get a more accurate solution!

Solution u(x) Calculated u Exact u 0.4 0.3 0.2

An interesting observation is that for the previous case where the boundary conditions were fixed, we got accurate solution even when n=3. It's just a matter of getting lucky where the equispaced points and the constant force term matched up such that it ended up as a perfect scenario. Rest be assured, if you change anything like variable forcing term, it'll give less accurate solution for smaller n values but that will still be $\mathcal{O}(h^2)$ accurate.

0.50

х

0.75

1.00

However, in this case we approximated the left boundary using the backward difference formula which is $\mathcal{O}(h)$ accurate. Hence, we'll get the final solution which is also $\mathcal{O}(h)$ accurate (error propagates!).

For this simple problem, we can actually get a $\mathcal{O}(h^2)$ accuracy! Let's do that...

0.25

0.0

0.00

Now, we have an idea that the solution should be parabolic (this can be confirmed using the exact solution). What we can do is use this fact which will give that $u_{-1}=u_1$. Another way to get to this same conclusion is by using the centered difference formula and approximate the left side boundary condition

$$\left. \frac{\partial u}{\partial x} \right|_{x=0} = \frac{u_1 - u_{-1}}{2\Delta x} = 0 \tag{22}$$

By using this, we again increase the unknown variables to 4, i.e. we need to find x_0, x_1, x_2 and x_3 .

lacktriangle When i=0, we can use the fact that $u_1=u_{-1}$

$$-u_1 + 2u_0 - u_{-1} = (0.25)^2$$
 or (23) $2u_0 - 2u_1 = (0.25)^2$

lacksquare When i=1

$$-u_2 + 2u_1 - u_0 = (0.25)^2 (24)$$

lacksquare When i=2

$$-u_3 + 2u_2 - u_1 = (0.25)^2 (25)$$

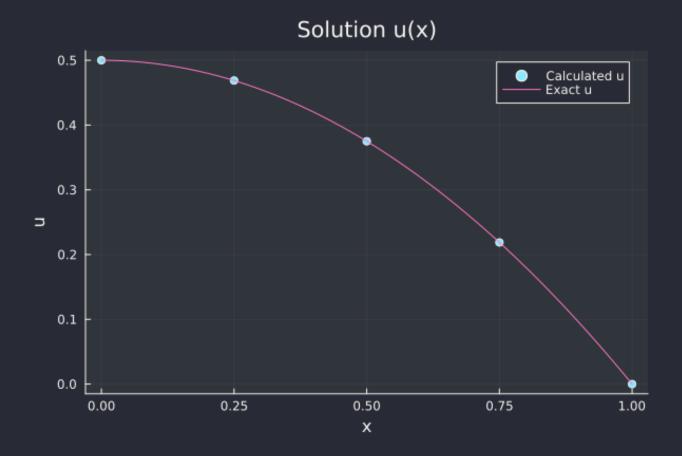
lacksquare When i=3

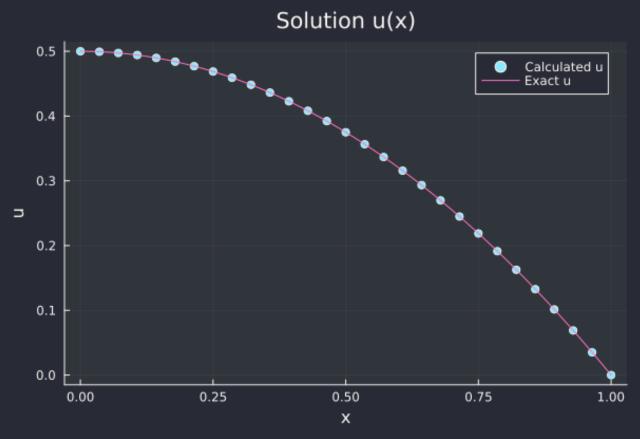
$$-u_4 + 2u_3 - u_2 = (0.25)^2 (26)$$

Now, we have 4 equations and 4 variables, which can again be represented as a matrix vector multiplication

$$\begin{pmatrix} 1 & -1 & 0 & 0 \\ -1 & 2 & -1 & 0 \\ 0 & -1 & 2 & -1 \\ 0 & 0 & -1 & 2 \end{pmatrix} \begin{pmatrix} u_0 \\ u_1 \\ u_2 \\ u_3 \end{pmatrix} = \begin{pmatrix} \frac{(0.25)^2}{2} \\ (0.25)^2 \\ (0.25)^2 \\ (0.25)^2 \end{pmatrix}$$
(21)

Solving this system of equation, we get the following solutions...





This is $\mathcal{O}(h^2)$ accurate!