

The kaobook class

Use this document as a template

My PhD Thesis

Customise this page according to your needs

Tobias Hangleiter*

December 3, 2024

* A \LaTeX lover/hater

The harmony of the world is made manifest in Form and Number, and the heart and soul and all the poetry of Natural Philosophy are embodied in the concept of mathematical beauty.

– D'Arcy Wentworth Thompson

Contents

Contents	iii
A Python WELCH’S-METHOD SPECTROMETER	1
IMPROVEMENTS AND CHARACTERIZATION OF A MILLIKELVIN CONFOCAL MICROSCOPE	2
ELECTROSTATIC TRAPPING OF EXCITONS IN SEMICONDUCTOR MEMBRANES	3
A FILTER-FUNCTION FORMALISM FOR QUANTUM OPERATIONS	4
1 Monte Carlo and Lindblad master equation simulations	6
1.1 Validation of QFT fidelities	6
2 Reconstruction by frequency-comb time-domain simulation	8
APPENDIX	10
A Test	11
Bibliography	12
List of Terms	13

Publications

- [1] Yaiza Aragonés-Soria, René Otten, Tobias Hangleiter, Pascal Cerfontaine, and David Gross. “Minimising Statistical Errors in Calibration of Quantum-Gate Sets.” June 7, 2022. DOI: [10.48550/arXiv.2206.03417](https://doi.org/10.48550/arXiv.2206.03417). (Visited on 06/08/2022). Pre-published.
- [2] Pascal Cerfontaine, Tobias Hangleiter, and Hendrik Bluhm. “Filter Functions for Quantum Processes under Correlated Noise.” In: *Physical Review Letters* 127.17 (Oct. 18, 2021), p. 170403. DOI: [10.1103/PhysRevLett.127.170403](https://doi.org/10.1103/PhysRevLett.127.170403).
- [3] Thomas Descamps, Feng Liu, Sebastian Kindel, René Otten, Tobias Hangleiter, Chao Zhao, Mihail Ion Lepsa, Julian Ritzmann, Arne Ludwig, Andreas D. Wieck, Beata E. Kardynał, and Hendrik Bluhm. “Semiconductor Membranes for Electrostatic Exciton Trapping in Optically Addressable Quantum Transport Devices.” In: *Phys. Rev. Appl.* 19.4 (Apr. 28, 2023), p. 044095. DOI: [10.1103/PhysRevApplied.19.044095](https://doi.org/10.1103/PhysRevApplied.19.044095). (Visited on 04/28/2023).
- [4] Thomas Descamps, Feng Liu, Tobias Hangleiter, Sebastian Kindel, Beata E. Kardynał, and Hendrik Bluhm. “Millikelvin Confocal Microscope with Free-Space Access and High-Frequency Electrical Control.” In: *Review of Scientific Instruments* 95.8 (Aug. 9, 2024), p. 083706. DOI: [10.1063/5.0200889](https://doi.org/10.1063/5.0200889). (Visited on 08/12/2024).
- [5] Denny Dütz, Sebastian Kock, Tobias Hangleiter, and Hendrik Bluhm. “Distributed Bragg Reflectors for Thermal Isolation of Semiconductor Spin Qubits.”
- [6] Sarah Fleitmann, Fabian Hader, Jan Vogelbruch, Simon Humpohl, Tobias Hangleiter, Stefanie Meyer, and Stefan van Waasen. “Noise Reduction Methods for Charge Stability Diagrams of Double Quantum Dots.” In: *IEEE Transactions on Quantum Engineering* 3 (2022), pp. 1–19. DOI: [10.1109/TQE.2022.3165968](https://doi.org/10.1109/TQE.2022.3165968).
- [7] Fabian Hader, Jan Vogelbruch, Simon Humpohl, Tobias Hangleiter, Chimezie Eguzo, Stefan Heinen, Stefanie Meyer, and Stefan van Waasen. “On Noise-Sensitive Automatic Tuning of Gate-Defined Sensor Dots.” In: *IEEE Transactions on Quantum Engineering* 4 (2023), pp. 1–18. DOI: [10.1109/TQE.2023.3255743](https://doi.org/10.1109/TQE.2023.3255743).
- [8] Tobias Hangleiter, Pascal Cerfontaine, and Hendrik Bluhm. “Filter-Function Formalism and Software Package to Compute Quantum Processes of Gate Sequences for Classical Non-Markovian Noise.” In: *Phys. Rev. Research* 3.4 (Oct. 18, 2021), p. 043047. DOI: [10.1103/PhysRevResearch.3.043047](https://doi.org/10.1103/PhysRevResearch.3.043047). (Visited on 01/19/2022).
- [9] Tobias Hangleiter, Pascal Cerfontaine, and Hendrik Bluhm. “Erratum: Filter-function Formalism and Software Package to Compute Quantum Processes of Gate Sequences for Classical Non-Markovian Noise [Phys. Rev. Research 3, 043047 (2021)].” In: *Phys. Rev. Res.* 6.4 (Oct. 16, 2024), p. 049001. DOI: [10.1103/PhysRevResearch.6.049001](https://doi.org/10.1103/PhysRevResearch.6.049001). (Visited on 10/16/2024).
- [10] Isabel Nha Minh Le, Julian D. Teske, Tobias Hangleiter, Pascal Cerfontaine, and Hendrik Bluhm. “Analytic Filter-Function Derivatives for Quantum Optimal Control.” In: *Phys. Rev. Applied* 17.2 (Feb. 2, 2022), p. 024006. DOI: [10.1103/PhysRevApplied.17.024006](https://doi.org/10.1103/PhysRevApplied.17.024006). (Visited on 02/03/2022).
- [11] Paul Surrey, Julian D Teske, Tobias Hangleiter, Pascal Cerfontaine, and Hendrik Bluhm. “Data-Driven Qubit Characterization and Optimal Control Using Deep Learning.”
- [12] Kui Wu, Sebastian Kindel, Thomas Descamps, Tobias Hangleiter, Jan Christoph Müller, Rebecca Rodrigo, Florian Merget, Beata E. Kardynał, Hendrik Bluhm, and Jeremy Witzens. “Modeling an Efficient Singlet-Triplet-Spin-Qubit-to-Photon Interface Assisted by a Photonic Crystal Cavity.” In: *Phys. Rev. Appl.* 21.5 (May 24, 2024), p. 054052. DOI: [10.1103/PhysRevApplied.21.054052](https://doi.org/10.1103/PhysRevApplied.21.054052). (Visited on 08/21/2024).

Software

The following open-source software packages were developed (at least partially) during the work on this thesis.

- [1] Tobias Hangleiter. *Lindblad_mc_tools*.
- [2] Tobias Hangleiter, Simon Humpohl, Paul Surrey, and Han Na We. *Qutil*. Version 2024.11.1. Zenodo, Nov. 21, 2024. doi: [10.5281/ZENODO.14200303](https://doi.org/10.5281/ZENODO.14200303). (Visited on 11/29/2024).
- [3] Tobias Hangleiter, Isabel Nha Minh Le, and Julian D. Teske. *Filter_functions*. Version v1.1.3. Zenodo, May 14, 2024. doi: [10.5281/ZENODO.4575000](https://doi.org/10.5281/ZENODO.4575000). (Visited on 11/29/2024).
- [4] Tobias Hangleiter, Simon Humpohl, Max Beer, and René Otten. *Python-Spectrometer*. Version 2024.11.1. Zenodo, Nov. 21, 2024. doi: [10.5281/ZENODO.13789861](https://doi.org/10.5281/ZENODO.13789861). (Visited on 11/29/2024).

A Python WELCH'S-METHOD SPECTROMETER

**IMPROVEMENTS AND
CHARACTERIZATION OF A MILLIKELVIN
CONFOCAL MICROSCOPE**

ELECTROSTATIC TRAPPING OF EXCITONS IN SEMICONDUCTOR MEMBRANES

A FILTER-FUNCTION FORMALISM FOR QUANTUM OPERATIONS

Parts of the results presented in this part have been published in Refs. 1 and **Hangleiter2021**. In Ref. **Hangleiter2021**, extensive use was made of Kubo's *cumulant expansion*. Due to an error in Kubo's paper which was only pointed out several years later [3] and the authors were not aware of, those results turned out to be not exact as claimed but approximate [4]. We investigate the accuracy of the expansion in Chapter 2.

[3]: Fox (1976), *Critique of the Generalized Cumulant Expansion Method*

[4]: Hangleiter et al. (2024), *Qutil*

Monte Carlo and Lindblad master equation simulations

1

1.1 Validation of QFT fidelities

In this section, we perform Lindblad master equation and Monte Carlo simulations to verify the fidelities predicted for the QFT circuit in the main text. We focus on noise exclusively on the third qubit, entering through the noise operator $B_\alpha \equiv \sigma_y^{(3)}$.

To validate the fidelity for white noise, we use a Lindblad master equation [Lindblad1976, Gorini1976] in superoperator form. We represent linear maps $\mathcal{A} : \rho \rightarrow \mathcal{A}(\rho)$ by matrices in the Pauli transfer matrix representation as (see Ref. Hangleiter2021 for more details)

$$\mathcal{A}_{ij} := \text{tr}(\sigma_i \mathcal{A}(\sigma_j)) \quad (1.1)$$

and operators as column vectors (*i.e.* generalized Bloch vectors) as

$$\rho_i := \text{tr}(\sigma_i \rho), \quad (1.2)$$

allowing us to write the Lindblad equation

$$\frac{d}{dt}\rho(t) = -i[H(t), \rho(t)] + \sum_{\alpha} \gamma_{\alpha} \left(L_{\alpha} \rho(t) L_{\alpha}^{\dagger} - \frac{1}{2} \{ L_{\alpha}^{\dagger} L_{\alpha}, \rho(t) \} \right) \quad (1.3)$$

as a linear differential equation in matrix form,

$$\frac{d}{dt}\rho_i(t) = \sum_j \left(-i\mathcal{H}_{ij}(t) + \sum_{\alpha} \gamma_{\alpha} \mathcal{D}_{\alpha,ij} \right) \rho_j(t). \quad (1.4)$$

Here, $\mathcal{H}_{ij}(t) = \text{tr}(\sigma_i [H(t), \sigma_j])$ and $\mathcal{D}_{\alpha,ij} = \text{tr}(\sigma_i L_{\alpha} \sigma_j L_{\alpha}^{\dagger} - \frac{1}{2} \sigma_i \{ L_{\alpha}^{\dagger} L_{\alpha}, \sigma_j \})$. By setting $L_{\alpha} \equiv \sigma_y^{(3)}$ as well as $\gamma_{\alpha} \equiv S_0/2$ with S_0 the amplitude of the one-sided noise power spectral density (PSD) so that $S(\omega) = S_0$, we can compare the fidelity obtained from the filter functions to that from the explicit simulation of Equation 1.4. The latter is computed as $F_{\text{avg}} = \text{tr}(\mathcal{Q}^{\dagger} \mathcal{U}) / d^2$, where \mathcal{Q} is the superpropagator due to the Hamiltonian evolution alone (*i.e.* the ideal evolution without noise).

For the Monte Carlo simulation, we explicitly generate time traces of $b_Q(t)$ (c.f. ??) by drawing pseudo-random variables from a distribution whose PSD is $S(f) = S_0/f$. To do this, we draw complex, normally distributed samples in frequency space (*i.e.* white noise), scale it with the square root of the PSD, and finally perform the inverse Fourier transform. We choose an oversampling factor of 16, so that the time discretization of the simulation is $\Delta t_{\text{MC}} = \Delta t/16 = 62.5 \text{ ps}$ ($\Delta t = 1 \text{ ns}$ is the time step of the pulses used in the FF simulation), leading to a highest resolvable frequency of $f_{\text{max}} = 16 \text{ GHz}$. Conversely, we increase the frequency resolution by sampling a time trace longer by a given factor, giving frequencies below f_{min} (16 kHz for pink, 0 Hz for white noise) weight zero, and truncating it to the number of time steps in the algorithm times the oversampling factor. This yields a time trace with frequencies $f \in [f_{\text{min}}, f_{\text{max}}]$ and a given resolution (we choose $\Delta f = 160 \text{ Hz}$).

We then proceed to diagonalize the Hamiltonian $H(t) = H_c(t) + H_n(t)$ and compute the propagator for one noise realization as

$$U(t) = \prod_g V^{(g)} \exp\{-i\Omega^{(g)}\Delta t_{\text{MC}}\} V^{(g)\dagger}, \quad (1.5)$$

where $V^{(g)}$ is the unitary matrix of eigenvectors of $H(t)$ during time segment g and $\Omega^{(g)}$ the diagonal matrix of eigenvalues. Finally, we obtain an estimate for the average gate fidelity F_{avg} from the entanglement fidelity F_e as

$$\langle F_{\text{avg}} \rangle = \left\langle \frac{dF_e + 1}{d + 1} \right\rangle = \left\langle \frac{|\text{tr}(Q^\dagger U(\tau))|^2 + d}{d(d + 1)} \right\rangle. \quad (1.6)$$

Here, $Q \equiv U_c(t = \tau)$ is the noise-free propagator at time τ of completion of the circuit and $\langle \cdot \rangle$ denotes the ensemble average over N Monte Carlo realizations of Equation 1.5, *i.e.*, $\langle A \rangle = \frac{1}{N} \sum_{i=1}^N A_i$. The standard error of the mean can be obtained as $\sigma_{\langle F_{\text{avg}} \rangle} = \sigma_{F_{\text{avg}}} / \sqrt{N}$ with $\sigma_{F_{\text{avg}}}$ the standard deviation over the Monte Carlo traces.

Reconstruction by frequency-comb time-domain simulation

2

For a **given** interaction-picture quantum operation $\tilde{\mathcal{U}}$ resulting from the quantum system's evolution under the noise fully characterized by its one-sided power spectral density (PSD) $S(\omega)$, we define the filter function (FF) $\mathcal{F}(\omega; \tau)$ by

$$\tilde{\mathcal{U}}(\tau) = \exp \left\{ - \int \frac{d\omega}{2\pi} \mathcal{F}(\omega; \tau) S(\omega) \right\}. \quad (2.1)$$

Now, suppose that

$$S_{\omega_i}(\omega) = \sigma_i^2 \delta(\omega - \omega_i), \quad (2.2)$$

that is, the PSD of a monochromatic sinusoid of frequency ω_i and root mean square (RMS) σ_i^2 . Then Equation 2.1 becomes

$$\begin{aligned} \tilde{\mathcal{U}}_{\omega_i}(\tau) &= \exp \left\{ -2\pi\sigma_i^2 \int \frac{d\omega}{2\pi} \mathcal{F}(\omega; \tau) \delta(\omega - \omega_i) \right\} \\ &= \exp \left\{ -\sigma_i^2 \mathcal{F}(\omega_i; \tau) \right\}, \end{aligned} \quad (2.3)$$

where $\tilde{\mathcal{U}}_{\omega_i}(\tau)$ is the noisy quantum operation generated by monochromatic noise with PSD $S_{\omega_i}(\omega)$ according to Equation 2.2. It is now easy to invert Equation 2.3, and we obtain

$$\mathcal{F}(\omega_i; \tau) = -\sigma_i^{-2} \log \tilde{\mathcal{U}}_{\omega_i}(\tau). \quad (2.4)$$

Because we represent quantum operations as matrices in Liouville space, Equation 2.4 is easy to implement on a computer; to sample the exact FF at the set of discrete frequencies $\{\omega_i\}_i$ we simply need to compute $\tilde{\mathcal{U}}_{\omega_i}(\tau)$ using a time-domain simulation method of our choice and take the logarithm!¹

Indeed, we can go a step further and split apart the coherent and incoherent contributions to the noisy evolution. Since (in-)coherent quantum operations are represented by (anti-)symmetric matrices in Liouville space, we may define the incoherent and coherent FFs by

$$\begin{aligned} \mathcal{F}_T(\omega_i; \tau) &= \frac{1}{2} (\mathcal{F}(\omega_i; \tau) + \mathcal{F}(\omega_i; \tau)^T) \\ &= \frac{1}{2\sigma_i^2} (\log \tilde{\mathcal{U}}_{\omega_i}(\tau) + \log \tilde{\mathcal{U}}_{\omega_i}(\tau)^T), \end{aligned} \quad (2.5)$$

and

$$\begin{aligned} \mathcal{F}_\Delta(\omega_i; \tau) &= \frac{1}{2} (\mathcal{F}(\omega_i; \tau) - \mathcal{F}(\omega_i; \tau)^T) \\ &= \frac{1}{2\sigma_i^2} (\log \tilde{\mathcal{U}}_{\omega_i}(\tau) - \log \tilde{\mathcal{U}}_{\omega_i}(\tau)^T), \end{aligned} \quad (2.6)$$

respectively. From a Monte Carlo (MC) point of view, $\tilde{\mathcal{U}}$ is given by

$$\tilde{\mathcal{U}}(\tau) \equiv \langle \tilde{\mathcal{U}}_{\omega_i}(\tau) \rangle = \mathcal{Q}^T \langle \mathcal{U}_{\omega_i}(\tau) \rangle, \quad (2.7)$$

where $\mathcal{U}_{\omega_i}(\tau)$ is the solution of the Schrödinger equation for a single

Equation 2.2 discretizes $S(\omega)$ by sampling it at points ω_i , i.e.,

$$S(\omega) \sim \lim_{n \rightarrow \infty} \sum_{i=1}^n S_{\omega_i}(\omega).$$

1: A similar approach was pursued by Geck in her PhD thesis to compare gate fidelities [5].

realization of the noise and $\langle \cdot \rangle$ denotes the ensemble average. Inserting into Equation 2.3, we find

$$\begin{aligned}\mathcal{F}(\omega_i; \tau) &= \sigma_i^{-2} \log \langle \tilde{\mathcal{U}}_{\omega_i}(\tau) \rangle \\ &= \sigma_i^{-2} \log \mathcal{Q}^T \langle \mathcal{U}_{\omega_i}(\tau) \rangle.\end{aligned}\quad (2.8)$$

If we evaluate Equation 2.8 for a set of frequencies $\{\omega_i\}_i$ sampling the true spectrum $S(\omega)$ sufficiently well, we thus obtain the exact filter function $\mathcal{F}(\omega; \tau)$ (within the accuracy of MC), allowing us to compare the accuracy of the formalism developed in.

So what does a single noise realization of Equation 2.2 in the time domain look like? It's a sinusoid with amplitude $A_i \sim \mathcal{N}(0, \sigma_i^2)$, frequency ω_i , and phase $\phi \sim U(0, 2\pi)$,

$$b(t) = A_i \sin(\omega_i t + \phi). \quad (2.9)$$

Instead of taking the random sampling approach of MC, we can also compute the expectation value of $\mathcal{U}_{\omega_i}(t)$ over A_i and ϕ by integrating over the probability density functions,

$$\mathbb{E}_{A_i, \phi}[\mathcal{U}] = \int dx \rho_{A_i}(x) \int dy \rho_{\phi}(y) \mathcal{U}[x, y], \quad (2.10)$$

where we wrote $\mathcal{U}[x, y]$ for the Liouville representation of the MC propagator for single x and y drawn from their respective distributions. Whether we choose the MC method or direct integration, in practice Equation 2.10 will be evaluated numerically, either by drawing random samples from the distributions $\mathcal{N}(0, \sigma_i^2)$ and $U(0, 2\pi)$ or discretizing Equation 2.10.

For N realizations of the stochastic process underlying $b(t)$, the ensemble average of a quantity $A(t)$ that is a function of $b(t)$ is given by

$$\langle A \rangle(t) = \frac{1}{N} \sum_{i=1}^N A_i(t)$$

where i enumerates the realizations of the stochastic process.

ref

tikz sketch?

We write the expectation value of an observable A with respect to a random variable X with the probability density function ρ_X as $\mathbb{E}_X[A]$.

The probability density function of the zero-mean normal distribution $\mathcal{N}(\mu, \sigma^2)$ is

$$\rho(x) = \frac{1}{\sqrt{2\pi\sigma_i^2}} \exp\left\{-\frac{(x-\mu)^2}{2\sigma^2}\right\}.$$

The probability density function of the uniform distribution $U(a, b)$ is

$$\rho(y) = \begin{cases} (b-a)^{-1} & \text{if } a \leq y < b, \\ 0 & \text{else} \end{cases}$$

APPENDIX

Test

A

Hello, here is some text without a meaning. This text should show what a printed text will look like at this place. If you read this text, you will get no information. Really? Is there no information? Is there a difference between this text and some nonsense like “Huardest gefburn”? Kjift – not at all! A blind text like this gives you information about the selected font, how the letters are written and an impression of the look. This text should contain all letters of the alphabet and it should be written in of the original language. There is no need for special content, but the length of words should match the language.

Bibliography

Here are the references in citation order.

- [1] Pascal Cerfontaine, Tobias Hangleiter, and Hendrik Bluhm. “Filter Functions for Quantum Processes under Correlated Noise.” In: *Physical Review Letters* 127.17 (Oct. 18, 2021), p. 170403. DOI: [10.1103/PhysRevLett.127.170403](https://doi.org/10.1103/PhysRevLett.127.170403) (cited on page 5).
- [2] Ryogo Kubo. “Generalized Cumulant Expansion Method.” In: *J. Phys. Soc. Jpn.* 17.7 (July 1962), pp. 1100–1120. DOI: [10.1143/JPSJ.17.1100](https://doi.org/10.1143/JPSJ.17.1100). (Visited on 01/19/2022) (cited on page 5).
- [3] Ronald Forrest Fox. “Critique of the Generalized Cumulant Expansion Method.” In: *J. Math. Phys.* 17.7 (July 1, 1976), pp. 1148–1153. DOI: [10.1063/1.523041](https://doi.org/10.1063/1.523041). (Visited on 04/23/2024) (cited on page 5).
- [4] Tobias Hangleiter et al. *Qutil*. Version 2024.11.1. Zenodo, Nov. 21, 2024. DOI: [10.5281/ZENODO.14200303](https://doi.org/10.5281/ZENODO.14200303). (Visited on 11/29/2024) (cited on page 5).
- [5] Lotte Geck. “Scalable Control Electronics for a Spin Based Quantum Computer.” PhD thesis. Jülich: Forschungszentrum Jülich GmbH, 2021. (Visited on 11/18/2024) (cited on page 8).

Special Terms

F

FF filter function. 8

M

MC Monte Carlo. 8, 9

P

PSD power spectral density. 8

R

RMS root mean square. 8