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# **My PHD Thesis**

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The harmony of the world is made manifest in Form and Number, and the heart and soul and all the poetry of Natural Philosophy are embodied in the concept of mathematical beauty.

– D'Arcy Wentworth Thompson

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## Publications

This might come in handy for PhD theses: some ideas and figures have appeared previously in the following publications:

- [1] Yaiza Aragonés-Soria et al. “Minimising Statistical Errors in Calibration of Quantum-Gate Sets.” June 7, 2022. doi: [10.48550/arXiv.2206.03417](https://doi.org/10.48550/arXiv.2206.03417). (Visited on 06/08/2022). Pre-published.
- [2] Pascal Cerfontaine, Tobias Hangleiter, and Hendrik Bluhm. “Filter Functions for Quantum Processes under Correlated Noise.” In: *Physical Review Letters* 127.17 (Oct. 18, 2021), p. 170403. doi: [10.1103/PhysRevLett.127.170403](https://doi.org/10.1103/PhysRevLett.127.170403).
- [3] Thomas Descamps et al. “Semiconductor Membranes for Electrostatic Exciton Trapping in Optically Addressable Quantum Transport Devices.” In: *Phys. Rev. Appl.* 19.4 (Apr. 28, 2023), p. 044095. doi: [10.1103/PhysRevApplied.19.044095](https://doi.org/10.1103/PhysRevApplied.19.044095). (Visited on 04/28/2023).
- [4] Thomas Descamps et al. “Millikelvin Confocal Microscope with Free-Space Access and High-Frequency Electrical Control.” In: *Review of Scientific Instruments* 95.8 (Aug. 9, 2024), p. 083706. doi: [10.1063/5.0200889](https://doi.org/10.1063/5.0200889). (Visited on 08/12/2024).
- [5] Sarah Fleitmann et al. “Noise Reduction Methods for Charge Stability Diagrams of Double Quantum Dots.” In: *IEEE Transactions on Quantum Engineering* 3 (2022), pp. 1–19. doi: [10.1109/TQE.2022.3165968](https://doi.org/10.1109/TQE.2022.3165968).
- [6] Fabian Hader et al. “On Noise-Sensitive Automatic Tuning of Gate-Defined Sensor Dots.” In: *IEEE Transactions on Quantum Engineering* 4 (2023), pp. 1–18. doi: [10.1109/TQE.2023.3255743](https://doi.org/10.1109/TQE.2023.3255743).
- [7] Tobias Hangleiter, Pascal Cerfontaine, and Hendrik Bluhm. “Filter-Function Formalism and Software Package to Compute Quantum Processes of Gate Sequences for Classical Non-Markovian Noise.” In: *Phys. Rev. Research* 3.4 (Oct. 18, 2021), p. 043047. doi: [10.1103/PhysRevResearch.3.043047](https://doi.org/10.1103/PhysRevResearch.3.043047). (Visited on 01/19/2022).
- [8] Tobias Hangleiter, Pascal Cerfontaine, and Hendrik Bluhm. “Erratum: Filter-function Formalism and Software Package to Compute Quantum Processes of Gate Sequences for Classical Non-Markovian Noise [Phys. Rev. Research 3, 043047 (2021)].” In: *Phys. Rev. Res.* 6.4 (Oct. 16, 2024), p. 049001. doi: [10.1103/PhysRevResearch.6.049001](https://doi.org/10.1103/PhysRevResearch.6.049001). (Visited on 10/16/2024).
- [9] Isabel Nha Minh Le et al. “Analytic Filter-Function Derivatives for Quantum Optimal Control.” In: *Phys. Rev. Applied* 17.2 (Feb. 2, 2022), p. 024006. doi: [10.1103/PhysRevApplied.17.024006](https://doi.org/10.1103/PhysRevApplied.17.024006). (Visited on 02/03/2022).
- [10] Kui Wu et al. “Modeling an Efficient Singlet-Triplet-Spin-Qubit-to-Photon Interface Assisted by a Photonic Crystal Cavity.” In: *Phys. Rev. Appl.* 21.5 (May 24, 2024), p. 054052. doi: [10.1103/PhysRevApplied.21.054052](https://doi.org/10.1103/PhysRevApplied.21.054052). (Visited on 08/21/2024).

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# Software

The following open-source software packages were developed (at least partially) during the work on this thesis.

- [1] Tobias Hangleiter, Isabel Nha Minh Le, and Julian D. Teske. *Filter\_functions*. Version v1.1.3. Zenodo, May 14, 2024. doi: [10.5281/ZENODO.4575000](https://doi.org/10.5281/ZENODO.4575000). (Visited on 09/26/2024).
- [2] Tobias Hangleiter et al. *Python-Spectrometer*. Version 2024.11.1. Zenodo, Nov. 21, 2024. doi: [10.5281/zenodo.14198682](https://doi.org/10.5281/zenodo.14198682). (Visited on 11/21/2024).
- [3] Tobias Hangleiter et al. *Qutil*. Version 2024.11.1. Zenodo, Nov. 21, 2024. doi: [10.5281/zenodo.14200666](https://doi.org/10.5281/zenodo.14200666). (Visited on 11/21/2024).

# **A FILTER-FUNCTION FORMALISM FOR QUANTUM OPERATIONS**



# Filter Functions

# 1

For a **given** quantum operation  $\tilde{\mathcal{U}}$  resulting from the quantum system's evolution under the noise fully characterized by its one-sided power spectral density (PSD)  $S(\omega)$ , we define the filter function (FF)  $\mathcal{F}(\omega; \tau)$  by

$$\tilde{\mathcal{U}}(\tau) = \exp \int \frac{d\omega}{2\pi} \mathcal{F}(\omega; \tau) S(\omega). \quad (1.1)$$

Now, suppose that

$$S_{\omega_i}(\omega) = \sigma_i^2 \delta(\omega - \omega_i), \quad (1.2)$$

that is, the PSD of a monochromatic sinusoid of frequency  $\omega_i$  and root mean square (RMS)  $\sigma_i^2$ . Then Equation 1.1 becomes

$$\begin{aligned} \tilde{\mathcal{U}}_{\omega_i}(\tau) &= \exp \left\{ \sigma_i^2 \int \frac{d\omega}{2\pi} \mathcal{F}(\omega; \tau) \delta(\omega - \omega_i) \right\} \\ &= \exp \left\{ \frac{\sigma_i^2}{2\pi} \mathcal{F}(\omega_i; \tau) \right\}, \end{aligned} \quad (1.3)$$

where  $\tilde{\mathcal{U}}_{\omega_i}(\tau)$  is the noisy quantum operation generated by monochromatic noise with PSD  $S_{\omega_i}(\omega)$  according to Equation 1.2. It is now easy to invert Equation 1.3, and we obtain

$$\mathcal{F}(\omega_i; \tau) = \frac{2\pi}{\sigma_i^2} \log \tilde{\mathcal{U}}_{\omega_i}(\tau). \quad (1.4)$$

Because we represent quantum operations as matrices in Liouville space, Equation 1.4 is easy to implement on a computer; we simply need to evaluate  $\tilde{\mathcal{U}}_{\omega_i}(\tau)$  for a set of discrete frequencies  $\{\omega_i\}_i$  using Monte Carlo (MC) and take the logarithm! [1].

Indeed, we can go a step further and split apart the coherent and incoherent contributions to the noisy evolution. Since (in-)coherent quantum operations are represented by (anti-)symmetric matrices in Liouville space, we may define the incoherent and coherent FFs by

$$\begin{aligned} \mathcal{F}_T(\omega; \tau) &= \frac{1}{2} (\mathcal{F}(\omega; \tau) + \mathcal{F}(\omega; \tau)^T) \\ &= \frac{\pi}{\sigma_i^2} (\log \tilde{\mathcal{U}}_{\omega_i}(\tau) + \log \tilde{\mathcal{U}}_{\omega_i}(\tau)^T), \end{aligned} \quad (1.5)$$

and

$$\begin{aligned} \mathcal{F}_\Delta(\omega; \tau) &= \frac{1}{2} (\mathcal{F}(\omega; \tau) - \mathcal{F}(\omega; \tau)^T) \\ &= \frac{\pi}{\sigma_i^2} (\log \tilde{\mathcal{U}}_{\omega_i}(\tau) - \log \tilde{\mathcal{U}}_{\omega_i}(\tau)^T), \end{aligned} \quad (1.6)$$

respectively. From a MC point of view,  $\tilde{\mathcal{U}}$  is given by

$$\tilde{\mathcal{U}} \equiv \langle \tilde{\mathcal{U}}_{\omega_i} \rangle(\tau) = \mathcal{Q}^T \langle \mathcal{U}_{\omega_i}(\tau) \rangle, \quad (1.7)$$

where  $\mathcal{U}_{\omega_i}(\tau)$  is the solution of the Schrödinger equation for a single realization of the noise and  $\langle \cdot \rangle$  denotes the ensemble average. Solving Equa-

Equation 1.2 discretizes  $S(\omega)$  by sampling it at points  $\omega_i$ , i.e.,

$$S(\omega) = \lim_{n \rightarrow \infty} \sum_{i=1}^n S_{\omega_i}(\omega).$$

[1]: Geck (2021), *Scalable Control Electronics for a Spin Based Quantum Computer*

For  $N$  realizations of the stochastic process underlying  $b(t)$ , the ensemble average of a quantity  $A(t)$  that is a function of  $b(t)$  is given by

$$\langle A \rangle(t) = \frac{1}{N} \sum_{i=1}^N A_i(t)$$

where  $i$  enumerates the realizations of

tion 1.1 for  $\mathcal{F}(\omega; \tau)$ , we find

$$\begin{aligned}\mathcal{F}(\omega; \tau) &= \int \frac{d\omega'}{2\pi} \delta(\omega - \omega') \log \langle \tilde{\mathcal{U}}_{\omega_i}(\tau) \rangle \\ &= \mathcal{Q}^\top + \int \frac{d\omega'}{2\pi} \delta(\omega - \omega') \log \langle \mathcal{U}_{\omega_i}(\tau) \rangle.\end{aligned}\quad (1.8)$$

So what does a single noise realization of Equation 1.2 in the time domain look like? It's a sinusoid with amplitude  $A_i \sim \mathcal{N}(0, \sigma_i^2)$ , frequency  $\omega_i$ , and phase  $\phi \sim \mathcal{U}(0, 2\pi)$ ,

$$b(t) = A \sin(\omega_i t + \phi). \quad (1.9)$$

We can compute the expectation value of  $\mathcal{U}_{\omega_i}(t)$  over  $A_i$  and  $\phi$  by integrating over the probability distributions,

$$\mathbb{E}_{\sigma_i, \phi}[\mathcal{U}](t) = \int dx \rho_{\sigma_i}(x) \int_0^{2\pi} d\phi \mathcal{U}[x, \phi]. \quad (1.10)$$

In practice, Equation 1.10 will be evaluated numerically by dividing the integration intervals into sub-intervals of equal size and weighing the integrand with the probability density, approximated to be constant on the sub-interval.

We write the expectation value of an observable  $A$  with respect to a random variable  $X$  with the probability density function  $\rho_X$  as  $\mathbb{E}_X[A]$ .

The probability density function of the zero-mean normal distribution is

$$\rho_{\sigma_i}(x) = \frac{1}{\sqrt{2\pi\sigma_i^2}} \exp\left(-\frac{x^2}{2\sigma_i^2}\right).$$

# **APPENDIX**

Test

A

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# Bibliography

Here are the references in citation order.

- [1] Lotte Geck. “Scalable Control Electronics for a Spin Based Quantum Computer.” PhD thesis. Jülich: Forschungszentrum Jülich GmbH, 2021. (Visited on 11/18/2024) (cited on page 2).