The kaobook class

Use this document as a template

My PhD Thesis

Customise this page according to your needs

Tobias Hangleiter*

March 25, 2025

* A LaTeX lover/hater

The harmony of the world is made manifest in Form and Number, and the heart and soul and all the poetry of Natural Philosophy are embodied in the concept of mathematical beauty.

– D’Arcy Wentworth Thompson

Contents

A flexible Python tool for Fourier-transform noise spectroscopy

1 Introduction

Noise is ubiquitous in condensed matter physics experiments, and in mesoscopic systems in particular it can easily drown out the sought-after signal. Hence, characterizing (and subsequently mitigating) noise is an essential task for the experimentalist. But noise comes in as many different forms as there are types of signal sources and detectors, whether it be a voltage source or a photodetector, and while some instruments have built-in solutions for noise analysis, they vary in functionality and capability. Moreover, the measured signal often does not directly correspond to the noisy physical quantity of interest, making it desirable to be able to manipulate the raw data before processing.

Theory of spectral noise estimation | 2

There exists a multitude of methods for estimating noise properties.

lay out some others

If the noisy process $x(t)$[1] has Gaussian statistics, meaning that the value at a given point in time follows a normal distribution with some mean μ and variance σ^2 over multiple realizations of the process, it can be fully described by the power spectral density (PSD) $S(\omega)$.[2] For the purpose of noise estimation, the assumption of Gaussianity is a rather weak one as the noise typically arises from a large ensemble of individual fluctuators and is therefore well approximated by a Gaussian distribution by the central limit theorem.[3] Even if the process $x(t)$ is not perfectly Gaussian, non-Gaussian contributions can be seen as higher-order contributions if viewed from the perspective of perturbation theory, and therefore the PSD still captures a significant part of the statistical properties. For this reason, the PSD is the central quantity of interest in noise spectroscopy and I will discuss some of its properties in the following.

1: We discuss only classical noise here, meaning $x(t)$ commutes with itself at all times. For descriptions of and spectroscopy protocols for quantum noise refer to Refs. 1 and 2, for example.

2: The term *power spectrum* is often used interchangably. I will do so as well, but emphasize at this point that in digital signal processing in particular, the *spectrum* is a different quantity from the *spectral density*.

maybe a classical signal processing ref?

3: As an example, consider electronic devices, where voltage noise arises from a large number of defects and other charge traps in oxides being populated and depopulated at certain rates γ. The ensemble average over these so-called two-level fluctuators (TLFs) then yields the well-known $1/f$-like noise spectra.

For real signals $x(t) \in \mathbb{R}$, $S(\omega)$ is an even function and one therefore distinguishes the *two-sided* PSD $S^{(2)}(\omega)$ defined over $\mathbb{R}$ from the *one-sided* PSD $S^{(1)}(\omega) = 2S^{(2)}(\omega)$ defined only over $\mathbb{R}^+$. Complex signals $x(t) \in \mathbb{C}$ such as those generated by Lock-in amplifiers after demodulation in turn have asymmetric, two-sided PSDs.

flesh out this sidenote?

for,for,for

2.1 Spectrum estimation from time series

To see how the PSD may be estimated from time-series data, consider a continuous wide-sense stationary[4] signal in the time domain $x(t) \in \mathbb{C}$ that is observed for some time T. We define the windowed Fourier transform of $x(t)$ and its inverse by

$$\hat{x}_T(\omega) = \int_0^T \mathrm{d}t\, x(t)\mathrm{e}^{\mathrm{i}\omega t} \tag{2.1}$$

$$\text{and} \quad x(t) = \int_{-\infty}^{\infty} \frac{\mathrm{d}\omega}{2\pi} \hat{x}_T(\omega)\mathrm{e}^{-\mathrm{i}\omega t}, \tag{2.2}$$

4: For a wide-sense stationary (also called weakly stationary) process $x(t)$, the mean is constant and the auto-correlation function $C(t,t') = \langle x(t)^* x(t')\rangle$ is given by $\langle x(t)^* x(t+\tau)\rangle = \langle x(0)^* x(\tau)\rangle$ with $\tau = t' - t$. That is, it is a function of only the time lag τ and not the absolute point in time. For Gaussian processes as discussed here, this also implies stationarity [3]. The property further implies that $C(\tau)$ is an even function.

sketch of auto-correlation function?

i. e., we assume that outside of the window of observation $x(t)$ is zero. The auto-correlation function of $x(t)$ is given by

$$C(\tau) = \langle x(t)^* x(t+\tau)\rangle \tag{2.3}$$

$$= \lim_{T\to\infty} \frac{1}{T}\int_0^T \mathrm{d}t\, x(t)^* x(t+\tau), \tag{2.4}$$

where $\langle\,\cdot\,\rangle$ is the ensemble average over multiple realizations of the process and the last equality holds true for ergodic processes. Expressing $x(t)$ in terms of its Fourier representation (Equation 2.1) and reordering the integrals, we get[5]

$$C(\tau) = \lim_{T\to\infty} \frac{1}{T}\int_0^T \mathrm{d}t \int_{-\infty}^{\infty} \frac{\mathrm{d}\omega}{2\pi}\hat{x}_T(\omega)^*\mathrm{e}^{\mathrm{i}\omega t}\int_{-\infty}^{\infty}\frac{\mathrm{d}\omega'}{2\pi}\hat{x}_T(\omega')\mathrm{e}^{-\mathrm{i}\omega'(t+\tau)} \tag{2.5}$$

$$= \lim_{T\to\infty} \frac{1}{T}\int_{-\infty}^{\infty} \frac{\mathrm{d}\omega}{2\pi}\int_{-\infty}^{\infty}\frac{\mathrm{d}\omega'}{2\pi}\hat{x}_T(\omega)^*\hat{x}_T(\omega')\mathrm{e}^{-\mathrm{i}\omega'\tau}\int_0^T \mathrm{d}t\,\mathrm{e}^{\mathrm{i}t(\omega-\omega')} \tag{2.6}$$

5: Mathematicians might at this point argue the integrability of $x(t)$, but as we deal with physical processes with finite bandwidth (and have no shame), we do not.

The innermost integral approaches a δ-function for large T,[6] allowing us to further simplify this under the limit as

$$C(\tau) = \lim_{T\to\infty} \frac{1}{T} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \hat{x}_T(\omega)^* \hat{x}_T(\omega') e^{-i\omega'\tau} \delta(\omega - \omega') \tag{2.7}$$

$$= \lim_{T\to\infty} \frac{1}{T} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} |\hat{x}_T(\omega)|^2 e^{-i\omega\tau} \tag{2.8}$$

$$= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S(\omega) e^{-i\omega\tau} \tag{2.9}$$

with the PSD

$$S(\omega) = \lim_{T\to\infty} \frac{1}{T} |\hat{x}_T(\omega)|^2 \tag{2.10}$$

$$= \int_{-\infty}^{\infty} d\tau\, C(\tau) e^{i\omega\tau} \tag{2.11}$$

Equation 2.9 is the Wiener-Khinchin theorem that states that the autocorrelation function $C(\tau)$ and the PSD $S(\omega)$ are Fourier-transform pairs [3]. Furthermore, defining the latter through Equation 2.10 gives us an intuitive picture of the PSD if we recall Parseval's theorem,

$$\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{1}{T} |\hat{x}_T(\omega)|^2 = \frac{1}{T} \int_{-\infty}^{\infty} dt\, |x(t)|^2. \tag{2.12}$$

That is, the total power P contained in the signal $x(t)$ is given by integrating over the PSD. Similarly, the power contained in a band of frequencies $[\omega_1, \omega_2]$ is given by

$$P(\omega_1, \omega_2) = \mathrm{RMS}(\omega_1, \omega_2)^2 \tag{2.13}$$

$$= \int_{\omega_1}^{\omega_2} \frac{d\omega}{2\pi} S(\omega) \tag{2.14}$$

where $\mathrm{RMS}(\omega_1, \omega_2)$ is the root-mean-square within this frequency band. These relations are helpful when analyzing noise PSDs to gauge the relative weight of contributions from different frequency bands to the total noise power.

Equation 2.10 represents the starting point for the experimental spectrum estimation procedure. Instead of a continuous signal $x(t), t \in [0, T]$, consider its discretized version

$$x_n, \quad n \in \{0, 1, \dots, N-1\} \tag{2.15}$$

defined at times $t_n = n\Delta t$ with $T = N\Delta t$ and where $\Delta t = f_s^{-1}$ is the sampling interval (the inverse of the sampling frequency f_s). Invoking the ergodic theorem, we can replace the long-term average in Equation 2.10 by the ensemble average over M realizations of the noisy signal x_n and write

$$S_n = \lim_{M\to\infty} \frac{1}{M} \sum_{i=0}^{M} |\hat{x}_n|^2, \tag{2.16}$$

where $\hat{x}_n$ is the discrete Fourier transform of x_n and S_n the PSD sampled at the discrete frequencies $\omega_n \in 2\pi/T \times \{-N/2, -N/2+1, \dots, N/2\} = 2\pi \times \{-f_s/2, \dots, f_s/2\}$.[7]

6: Note that, because $x(t)$ is wide-sense stationary, we may shift the limits of integration $\int_0^T \to \int_{-T/2}^{+T/2}$.

7: We blithely disregard integer algebra issues occuring here for conciseness and leave it as an exercise for the reader to figure out what the exact bounds of the set of ω_n are.

Characterization and Improvements of a Millikelvin Confocal Microscope

Electrostatic Trapping of Excitons in Semiconductor Membranes

A Filter-Function Formalism For Quantum Operations

Appendix

Bibliography

[1] A. A. Clerk et al. "Introduction to Quantum Noise, Measurement, and Amplification." In: *Rev. Mod. Phys.* 82.2 (Apr. 15, 2010), pp. 1155–1208. DOI: 10.1103/RevModPhys.82.1155. (Visited on 01/19/2022) (cited on page 3).

[2] Gerardo A. Paz-Silva, Leigh M. Norris, and Lorenza Viola. "Multiqubit Spectroscopy of Gaussian Quantum Noise." In: *Phys. Rev. A* 95.2 (Feb. 23, 2017), p. 022121. DOI: 10.1103/PhysRevA.95.022121 (cited on page 3).

[3] Lambert Herman Koopmans. *The Spectral Analysis of Time Series.* 2nd ed. Vol. 22. Probability and Mathematical Statistics. San Diego: Academic Press, 1995 (cited on pages 3, 4).

Special Terms

P

PSD power spectral density. 3, 4

T

TLF two-level fluctuator. 3