

AAE 590: Lie Group Methods for Control and Estimation

SO(2): Lie Theory Fundamentals

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SO(2) — The Simplest Lie Group

- All Lie theory concepts in one dimension
- Lie algebra, exponential/logarithm, adjoint
- Foundation for everything that follows

SO(2) is “too simple” — but perfect for learning the machinery!

The Special Orthogonal Group SO(2) [1]

Definition

$$\text{SO}(2) = \{R \in \mathbb{R}^{2 \times 2} : R^T R = I, \det(R) = 1\}$$

Explicit form:

$$R(\theta) = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

Verify it's a group (CAIN):

- **Closure:** $R(\theta_1)R(\theta_2) = R(\theta_1 + \theta_2)$ ✓
- **Associativity:** matrix multiplication ✓
- **Inverse:** $R(\theta)^{-1} = R(-\theta) = R(\theta)^T$ ✓
- **Neutral:** $R(0) = I$ ✓

$\text{SO}(2)$ is Abelian

Special Property

$\text{SO}(2)$ is **abelian** (commutative): $R(\theta_1)R(\theta_2) = R(\theta_2)R(\theta_1)$

Proof:

$$R(\theta_1)R(\theta_2) = R(\theta_1 + \theta_2) = R(\theta_2 + \theta_1) = R(\theta_2)R(\theta_1)$$

Warning

This is **NOT** true for $\text{SO}(3)$! 3D rotations do not commute.

$\text{SO}(2)$ is “too simple” — many Lie theory concepts become trivial. But it’s perfect for learning!

Why “SO(2)” and not “ S^1

Naming convention: The n in $\text{SO}(n)$ is the dimension of the **space being rotated**, not the group.

Group	Acts on	Group dim.	Topology
$\text{SO}(2)$	\mathbb{R}^2 (plane)	1	$\cong S^1$
$\text{SO}(3)$	\mathbb{R}^3 (3D space)	3	$\cong \mathbb{RP}^3$

$\mathbb{RP}^3 = \text{real projective 3-space}$: S^3 with opposite points identified ($x \sim -x$). This is why q and $-q$ represent the same rotation.

S^1 vs $\text{SO}(2)$:

- S^1 : geometric shape (circle)
- $\text{SO}(2)$: group structure (rotations)

Abstract vs Matrix:

- $\text{SO}(2)$ is an abstract group
- 2×2 matrices = defining representation

$\text{SO}(2)$ as a Manifold

Topology: $\text{SO}(2) \cong S^1$ (the unit circle)

Notation: S^1 and \cong

S^1 denotes the **unit circle** $\{(x, y) \in \mathbb{R}^2 : x^2 + y^2 = 1\}$. More generally, S^n is the n -dimensional sphere.

The symbol \cong means **isomorphic to**. An **isomorphism** is a bijective (one-to-one and onto) map that preserves structure. When we write $A \cong B$, we mean A and B are “essentially the same” — they have identical structure, just different representations.

Key properties:

- 1-dimensional **manifold** (locally looks like \mathbb{R}^1)
- **Compact** (closed and bounded in \mathbb{R}^n)
- **Connected** (one piece)
- **Not simply connected** (has a hole)

Compact — contains all its limit points

Parameterization:

- $\theta \in [0, 2\pi)$ or $(-\pi, \pi]$
- Wrapping: $\theta + 2\pi \equiv \theta$
- No global “flat” coordinates!

Group Actions and Representations

Definition

A **group action** of G on a set X is a map $\cdot : G \times X \rightarrow X$ satisfying:

- ① $e \cdot x = x$ (identity acts trivially)
- ② $(g_1 g_2) \cdot x = g_1 \cdot (g_2 \cdot x)$ (composition rule)

Intuition: Each $g \in G$ “moves” points in X ; composing group elements composes motions.

Definition

A **representation** is a group action on a **vector space** V via **linear** maps: a homomorphism $\rho : G \rightarrow \text{GL}(V)$.

$\text{GL}(V) = \text{invertible linear maps on } V$. For $V = \mathbb{R}^n$: invertible $n \times n$ matrices.

Example: $\text{SO}(2)$ acts on \mathbb{R}^2 by rotation: $R(\theta) \cdot \mathbf{v} = R(\theta)\mathbf{v}$

- Identity: $R(0) \cdot \mathbf{v} = I\mathbf{v} = \mathbf{v}$ ✓
- Composition: $R(\theta_1) \cdot (R(\theta_2) \cdot \mathbf{v}) = R(\theta_1 + \theta_2)\mathbf{v}$ ✓

Representations of SO(2)

Different ways SO(2) can act on vector spaces:

- ① **Standard (defining):** $\rho(\theta) = R(\theta)$ on \mathbb{R}^2 — 2D rotation matrices
- ② **Trivial:** $\rho(\theta) = 1$ on \mathbb{R} — everything maps to identity
- ③ **Complex exponential:** $\rho_n(\theta) = e^{in\theta}$ on \mathbb{C} for any $n \in \mathbb{Z}$

Are They Isomorphic?

No! $\rho_1(\theta) = e^{i\theta}$ and $\rho_2(\theta) = e^{2i\theta}$ are fundamentally different— ρ_2 “winds twice” around the circle as θ goes from 0 to 2π .

The Tangent Space at Identity

Question: What does “infinitesimal rotation” look like?

Taylor expansion:

$$R(\epsilon) = \begin{bmatrix} \cos \epsilon & -\sin \epsilon \\ \sin \epsilon & \cos \epsilon \end{bmatrix} \approx \begin{bmatrix} 1 & -\epsilon \\ \epsilon & 1 \end{bmatrix} = I + \epsilon \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

Derivative at identity:

$$\frac{d}{d\theta} R(\theta) \Big|_{\theta=0} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

*Notation: $\frac{d}{d\theta}$ is the **total derivative** (one variable). Use $\frac{\partial}{\partial \theta}$ for **partial derivatives** (multiple variables).*

This matrix generates all rotations!

Vector Spaces

Definition

A **vector space** over \mathbb{R} is a set V with two operations:

- **Addition:** $+ : V \times V \rightarrow V$
- **Scalar multiplication:** $\cdot : \mathbb{R} \times V \rightarrow V$

satisfying: commutativity, associativity, identity ($\mathbf{0}$), inverses, and distributivity.

Key examples:

- \mathbb{R}^n with standard addition and scalar multiplication
- $\mathbb{R}^{m \times n}$ (matrices) with matrix addition and scalar multiplication
- Polynomials, continuous functions, etc.

Why This Matters

In a vector space we can **add**, **scale**, and do **calculus**. Lie groups are NOT vector spaces (can't add rotations!), but their Lie algebras ARE.

Morphisms (Structure-Preserving Maps) [2]

Definition

A **homomorphism** $\phi : G \rightarrow H$ between groups preserves the group operation:

$$\phi(g_1 \cdot g_2) = \phi(g_1) \cdot \phi(g_2)$$

Definition

An **isomorphism** is a **bijective** homomorphism. If $\phi : G \rightarrow H$ is an isomorphism, we write $G \cong H$ ("G is isomorphic to H").

Definition

An **automorphism** is an isomorphism from a structure to **itself**: $\phi : G \rightarrow G$. The set of all automorphisms of G forms a group $\text{Aut}(G)$.

Example: For $G = \mathbb{R}^2$, the map $\phi_R(\mathbf{v}) = R\mathbf{v}$ (rotation) is an automorphism.

Homeomorphisms (Topology)

Definition

A **homeomorphism** is a **continuous bijection** with a **continuous inverse**.

Homeomorphism \approx “same shape” (topology)

Isomorphism \approx “same algebraic structure”

Examples:

- $SO(2) \cong S^1$ (circle) — homeomorphic *and* isomorphic as Lie groups
- Coffee mug \approx donut — homeomorphic (both have one hole)
- $[0, 1)$ and S^1 are **not** homeomorphic (circle is compact)

For Lie Groups

We care about **Lie group isomorphisms**: smooth maps that preserve both the group structure *and* the manifold structure.

Worked Example: Group Properties of SO(2)

Verify CAIN for $\text{SO}(2)$: show $R(\theta_1)R(\theta_2) = R(\theta_1 + \theta_2)$ and $R(\theta)^{-1} = R(-\theta)$.

Solution: Group Properties of SO(2)

$$R(\theta_1)R(\theta_2) = \begin{bmatrix} c_1 & -s_1 \\ s_1 & c_1 \end{bmatrix} \begin{bmatrix} c_2 & -s_2 \\ s_2 & c_2 \end{bmatrix} = \begin{bmatrix} c_1c_2 - s_1s_2 & -c_1s_2 - s_1c_2 \\ s_1c_2 + c_1s_2 & -s_1s_2 + c_1c_2 \end{bmatrix} = \begin{bmatrix} c_{12} & -s_{12} \\ s_{12} & c_{12} \end{bmatrix} = R(\theta_1 + \theta_2)$$

$$R(\theta)^{-1} = R(\theta)^T = \begin{bmatrix} c & s \\ -s & c \end{bmatrix} = \begin{bmatrix} c_{-\theta} & -s_{-\theta} \\ s_{-\theta} & c_{-\theta} \end{bmatrix} = R(-\theta) \quad (\text{using } \cos(-\theta) = \cos \theta, \sin(-\theta) = -\sin \theta)$$

The Lie Algebra $\mathfrak{so}(2)$ [1, 3]

Definition

The **Lie algebra** $\mathfrak{so}(2)$ is the tangent space at identity:

$$\mathfrak{so}(2) = T_I \mathrm{SO}(2) = \left\{ \begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix} : \omega \in \mathbb{R} \right\}$$

*Notation: $T_I G$ means the **tangent space** to G at the identity I —the set of all “velocity vectors” at that point.*

Properties:

- $\mathfrak{so}(2)$ is a **vector space** (1-dimensional)
- Elements are **skew-symmetric** matrices: $A^T = -A$
- Isomorphic to \mathbb{R} : each element determined by single scalar ω

Key Insight

The Lie algebra is a **flat** vector space where we can do calculus!

Wedge and Vee Operators

Definition

Wedge (\wedge): $\mathbb{R} \rightarrow \mathfrak{so}(2)$ — scalar to matrix

$$\omega^\wedge = \begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}$$

Definition

Vee (\vee): $\mathfrak{so}(2) \rightarrow \mathbb{R}$ — matrix to scalar

$$\begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}^\vee = \omega$$

These are inverses: $(\omega^\wedge)^\vee = \omega$ and $(A^\vee)^\wedge = A$

Notation Convention

We'll use ω for both the scalar and ω^\wedge for the matrix form.

The Exponential Map

Definition

The **exponential map** $\text{Exp} : \mathfrak{so}(2) \rightarrow \text{SO}(2)$:

$$\text{Exp}(\omega^\wedge) = e^{\omega^\wedge} = \sum_{k=0}^{\infty} \frac{(\omega^\wedge)^k}{k!}$$

Key property: $(\omega^\wedge)^2 = -\omega^2 I$

Closed form:

$$\begin{aligned}\text{Exp}(\omega^\wedge) &= I + \omega^\wedge + \frac{(\omega^\wedge)^2}{2!} + \frac{(\omega^\wedge)^3}{3!} + \dots \\ &= I \left(1 - \frac{\omega^2}{2!} + \frac{\omega^4}{4!} - \dots \right) + \omega^\wedge \left(1 - \frac{\omega^2}{3!} + \frac{\omega^4}{5!} - \dots \right) \\ &= I \cos \omega + \omega^\wedge \frac{\sin \omega}{\omega} = R(\omega)\end{aligned}$$

Exponential Map: Geometric Interpretation

Exp maps from the tangent space (flat) to the manifold (curved)

Lie algebra $\mathfrak{so}(2)$:

- Linear space $\cong \mathbb{R}$
- ω is angular velocity
- Can add/subtract freely

Lie group $\text{SO}(2)$:

- Curved manifold $\cong S^1$
- $R(\omega)$ is rotation
- Composition, not addition

$$\boxed{\omega \in \mathbb{R} \xrightarrow{\text{wedge}} \omega^\wedge \in \mathfrak{so}(2) \xrightarrow{\text{Exp}} R(\omega) \in \text{SO}(2)}$$

The Logarithmic Map

Definition

The **logarithmic map** $\text{Log} : \text{SO}(2) \rightarrow \mathfrak{so}(2)$ is the inverse of Exp :

$$\text{Log}(R) = \omega^\wedge \quad \text{where} \quad R = R(\omega)$$

For $\text{SO}(2)$:

$$\text{Log}(R) = \text{atan2}(R_{21}, R_{11})^\wedge$$

Multi-valuedness

Log returns $\omega \in (-\pi, \pi]$. But $R(\omega) = R(\omega + 2\pi k)$ for any $k \in \mathbb{Z}$ (integers).

For $\text{SO}(2)$: Exp is **surjective** (onto: every rotation is reached), Log is only locally defined.

Why Exp/Log Matter

Problem: How to interpolate between rotations R_0 and R_1 ?

Wrong (Euclidean):

$$R(t) = (1 - t)R_0 + tR_1 \quad \text{NOT a rotation matrix!}$$

Correct (Lie group):

$$R(t) = R_0 \cdot \text{Exp} \left(t \cdot \text{Log}(R_0^{-1} R_1) \right)$$

- Compute relative rotation: $\Delta R = R_0^{-1} R_1$
- Map to algebra: $\Delta\omega = \text{Log}(\Delta R)^\vee$
- Scale in algebra: $t \cdot \Delta\omega$
- Map back: $\text{Exp}((t \cdot \Delta\omega)^\wedge)$
- Apply from R_0

Worked Example: SO(2) Exp and Log

Compute $\text{Exp}((\pi/4)^\wedge)$, verify $R^T R = I$, then $\text{Log}(R)$ to recover $\pi/4$.

Solution: SO(2) Exp and Log

$$\text{Exp}((\pi/4)^\wedge) = R(\pi/4) = \begin{bmatrix} \cos(\pi/4) & -\sin(\pi/4) \\ \sin(\pi/4) & \cos(\pi/4) \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix}$$

Verify: $R^T R = \begin{bmatrix} \frac{1}{2} + \frac{1}{2} & 0 \\ 0 & \frac{1}{2} + \frac{1}{2} \end{bmatrix} = I \checkmark$

$$\text{Log}(R)^\vee = \text{atan2}(R_{21}, R_{11}) = \text{atan2}\left(\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2}\right) = \pi/4 \checkmark$$

The Adjoint Representation Ad

Definition

The **adjoint map** $\text{Ad}_R : \mathfrak{so}(2) \rightarrow \mathfrak{so}(2)$ is conjugation:

$$\text{Ad}_R(\omega^\wedge) = R \omega^\wedge R^{-1}$$

For $\text{SO}(2)$: Since $\text{SO}(2)$ is abelian:

$$\text{Ad}_R(\omega^\wedge) = R \omega^\wedge R^T = \omega^\wedge$$

$\text{SO}(2)$ is trivial!

$\text{Ad}_R = I$ for all $R \in \text{SO}(2)$. The adjoint does nothing.

This will be **much more interesting** for $\text{SO}(3)$ and $\text{SE}(2)$!

The Small Adjoint ad

Definition

The **small adjoint** $\text{ad}_\xi : \mathfrak{g} \rightarrow \mathfrak{g}$ is the Lie bracket:

$$\text{ad}_{\omega_1}(\omega_2) = [\omega_1^\wedge, \omega_2^\wedge] = \omega_1^\wedge \omega_2^\wedge - \omega_2^\wedge \omega_1^\wedge$$

For $\text{SO}(2)$:

$$[\omega_1^\wedge, \omega_2^\wedge] = 0 \quad (\text{always zero!})$$

Abelian \Rightarrow Trivial Bracket

The Lie bracket measures “non-commutativity.” For abelian groups, it vanishes.

Relationship: ad is the derivative of Ad at identity.

Baker-Campbell-Hausdorff (BCH) Formula

Question: Given $\text{Exp}(\xi_1^\wedge)$ and $\text{Exp}(\xi_2^\wedge)$, what is $\text{Exp}(\xi_1^\wedge) \cdot \text{Exp}(\xi_2^\wedge)$?

Theorem (BCH Formula)

$$\text{Exp}(\xi_1^\wedge) \cdot \text{Exp}(\xi_2^\wedge) = \text{Exp} \left(\left(\xi_1 + \xi_2 + \frac{1}{2}[\xi_1, \xi_2] + \frac{1}{12}[\xi_1, [\xi_1, \xi_2]] + \dots \right)^\wedge \right)$$

For $\text{SO}(2)$: All brackets vanish, so:

$$\text{Exp}(\omega_1^\wedge) \cdot \text{Exp}(\omega_2^\wedge) = \text{Exp}((\omega_1 + \omega_2)^\wedge)$$

i.e., $R(\omega_1)R(\omega_2) = R(\omega_1 + \omega_2)$ — just addition!

For non-abelian groups: BCH is much more complex.

Kinematics on $\text{SO}(2)$

Angular velocity: $\omega(t) \in \mathbb{R}$

Kinematic equation:

$$\dot{R} = R \omega^\wedge$$

Interpretation:

- $\omega^\wedge \in \mathfrak{so}(2)$ is the **body-frame** angular velocity
- Right multiplication: velocity expressed in body frame
- This is a **left-invariant** vector field on $\text{SO}(2)$

$\text{SO}(2)$ is Abelian: Left = Right Invariant

Since $\text{SO}(2)$ is abelian, $R\omega^\wedge = \omega^\wedge R$, so left and right invariant forms are **equivalent**. This is **not** true for $\text{SE}(2)$ or $\text{SO}(3)$!

Solution (constant ω):

$$R(t) = R(0) \cdot \text{Exp}(t \omega^\wedge) = R(0) \cdot R(t\omega)$$

SO(2) Summary

Concept	SO(2)
Group	2×2 rotation matrices
Lie algebra $\mathfrak{so}(2)$	2×2 skew-symmetric
Dimension	1
Wedge ω^\wedge	$\begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}$
$\text{Exp}(\omega^\wedge)$	$R(\omega)$
$\text{Log}(R)$	$\text{atan2}(R_{21}, R_{11})^\wedge$
Ad_R	Identity (trivial)
$[\cdot, \cdot]$	Zero (abelian)

SO(2) is “too nice” — but now we have all the machinery for harder groups!

Looking Ahead: SE(2)

Next lecture: SE(2) — combining rotation AND translation

- Non-abelian (rotations and translations don't commute)
- Semi-direct product structure
- Non-trivial adjoint representation
- Foundation for rigid body control

References I

- [1] Brian C Hall. *Lie Groups, Lie Algebras, and Representations: An Elementary Introduction.* 2nd. Springer, 2015.
- [2] David S Dummit and Richard M Foote. *Abstract Algebra.* 3rd. Wiley, 2004.
- [3] John Stillwell. *Naive Lie Theory.* Springer, 2008.