

# Thermodynamics

## Semester 4

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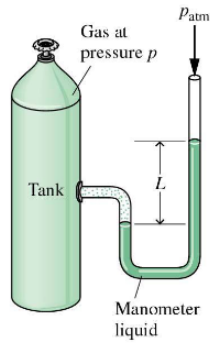
# Chapter 1

## Introductory concepts and definitions

### 1.1 Pressure

When we place an item in a fluid, the fluid exerts a pressure on the item. At a point on the item, the pressure is the same from all directions, however the pressure can differ from one point to another.

A manometer is used to measure pressure, it is a U shaped pipe connected to the medium we are trying to measure pressure in.



$$P_{\text{Tank}} = P_{\text{atm}} + \rho_{\text{Fluid}}gL.$$

**Note:-**

Reminder of the relation between force and pressure

$$P = \frac{F}{A}.$$

### 1.2 Work and Energy

The work of a force is defined to be

$$W = \int_{\ell_1}^{\ell_2} Fd\ell = \int_{V_1}^{V_2} pdV.$$

As well as the relation

$$\frac{1}{2}m(v_2^2 - v_1^2) = \int_{\ell_1}^{\ell_2} Fd\ell.$$

We call a process a polytropic process when it obeys the following relation

$$PV^n = \text{cnst.}$$

The work in a polytropic process is

$$W = \frac{P_2 V_2 - P_1 V_1}{n - 1}.$$

Power is defined to be the change in work

$$\dot{W} = F \cdot v.$$

The change energy in a system  $\Delta E$  is the sum of the changes in kinetic energy, potential energy, and internal energy  $\Delta U$

$$\Delta E = \Delta KE + \Delta PE + \Delta U.$$

### 1.2.1 Heat transfer

The heat energy in the system is denoted as  $Q$

$$Q = \int_{t_1}^{t_2} \dot{Q} dt.$$

the heat flux in the system  $\dot{q}$  can be used to calculate the heat transfer rate in the system

$$\dot{Q} = \int_A \dot{q} dA.$$

$W > 0$  from system

$W < 0$  on system

$Q > 0$  to system

$Q < 0$  from system

#### Conduction (Fourier's law)

Conduction is the transfer of energy from the more energetic particles of a substance to the adjacent less energetic ones as a result of interactions between the particles

$$\dot{Q}_{\text{cond}} = -\kappa A \frac{dT}{dx}.$$

Where  $\kappa$  is constant called thermal conductivity.

#### Convection (Newton's law of cooling)

Convection is the mode of energy transfer between a solid surface and the adjacent liquid or gas that is in motion, and it involves the combined effects of conduction and fluid motion.

$$\dot{Q}_{\text{conv}} = hA(T_0 - T_1).$$

Where  $h$  is constant called the heat transfer coefficient.

#### Radiation

Radiation is the energy emitted by matter in the form of electromagnetic waves (or photons) as a result of the changes in the electronic configurations of the atoms or molecules. Unlike conduction and convection, the transfer of energy by radiation does not require the presence of an intervening medium.

$$\dot{Q}_{\text{emitted}} = \varepsilon \sigma A T_0^4.$$

where  $\varepsilon$  is the emissivity of the surface and  $\sigma$  is the Stefan-Boltzmann constant

$$\sigma = 5.670367 \times 10^{-8}.$$

## 1.3 First Law of Thermodynamics

In a closed system the following relation holds

$$\Delta E = Q - W.$$

In a cycle the energy does not change so

$$W_{\text{cycle}} = Q_{\text{in}} - Q_{\text{out}}.$$

The formula for thermal efficiency is

$$\eta = \frac{W_{\text{cycle}}}{Q_{\text{in}}} = 1 - \frac{Q_{\text{out}}}{Q_{\text{in}}}.$$

The effectiveness or coefficient of performance  $\beta$  of a refrigerator is

$$\beta = \frac{Q_{\text{in}}}{W_{\text{cycle}}} = \frac{Q_{\text{in}}}{Q_{\text{out}} - Q_{\text{in}}}.$$

The effectiveness or coefficient of performance  $\gamma$  of a heat pump is

$$\gamma = \frac{Q_{\text{out}}}{W_{\text{cycle}}} = \frac{Q_{\text{out}}}{Q_{\text{out}} - Q_{\text{in}}}.$$

Under constant  $T$  we can apply Boyle's law

$$P_1 V_1 = P_2 V_2.$$

## Chapter 2

# Evaluating Properties

The quality of a liquid-vapor substance is defined by

$$x = \frac{m_{\text{vapor}}}{m_{\text{liquid}} + m_{\text{vapor}}}.$$

We define the specific volume  $v$  to be the inverse of density  $\rho$

$$v = \frac{1}{\rho} = \frac{V}{m}.$$

The specific volume of a liquid-vapor substance is

$$v = (1 - x)v_{\ell} + xv_g.$$

where  $v_{\ell}$  and  $v_g$  are the specific volume of the liquid portion and the gas portion of the substance respectively.

We define specific internal energy to be

$$u = \frac{U}{m}.$$

Similarly to specific volume

$$u = (1 - x)u_{\ell} + xu_g.$$

We define enthalpy to be

$$H = U + pV.$$

specific enthalpy to be

$$h = u + pv.$$

and similarly to both previous specific quantities

$$h = (1 - x)h_{\ell} + xh_g.$$

We define the specific heats to be

$$C_V = \left. \frac{\partial u}{\partial T} \right|_V$$
$$C_p = \left. \frac{\partial h}{\partial T} \right|_p$$

from these equation we can deduce the formulas

$$\Delta U = mC_V\Delta T$$

$$\Delta H = mC_p\Delta T$$

We also define the specific heat ratio as

$$k = \frac{C_p}{C_v}.$$

In case of incompressible fluids/solids, we have  $C_v = C_p$

$$\begin{aligned}\Delta u &= C\Delta T \\ \Delta h &= C\Delta T + v\Delta p\end{aligned}$$

## 2.1 Ideal gas model

The universal gas constant is

$$\mathcal{R} = 8.314471 \times \text{kJ} \cdot \text{kmol}^{-1} \cdot \text{K}^{-1}.$$

We define the compressibility factor to be

$$Z = \frac{p\bar{v}}{\mathcal{R}T}.$$

such that  $\bar{v} = Mv$  where  $M$  is the atomic/molecular mass

$$Z = \frac{pv}{RT}.$$

where  $R = \frac{M}{\mathcal{R}}$  and it is specific to every gas.

The ideal gas equation is

$$pV = mRT.$$

In an ideal gas, enthalpy only depends on temperature

$$h = u + RT.$$

$$\begin{aligned}\Delta u(T) &= \int_{T_1}^{T_2} C_v(T) dT \\ \Delta h(T) &= \int_{T_1}^{T_2} C_p(T) dT\end{aligned}$$

and

$$\begin{aligned}C_v(T) &= \frac{R}{k-1} \\ C_p(T) &= \frac{kR}{k-1}\end{aligned}$$



## Chapter 3

# Control Volume Analysis

Control Volume is a region in space through which mass and energy can flow. It is defined by a control surface that encloses the control volume.

### 3.1 Mass Conservation

$$\dot{m}_{\text{in}} - \dot{m}_{\text{out}} = \frac{dm_{\text{control}}}{dt}.$$

$$\dot{m} = \int_{A_{\text{in}}} \rho \vec{v} \cdot \vec{n} dA.$$

One dimensional Flow Form of the mass balance rate

$$\dot{m} = \rho A v.$$

$Av$  is the volumetric flow rate

$$\frac{dm_{\text{control}}}{dt} = \sum \frac{A_{\text{in}} v_{\text{in}}}{v_{\text{in}}} - \sum \frac{A_{\text{out}} v_{\text{out}}}{v_{\text{out}}}.$$

When the system reaches a steady state  $\frac{dm_{\text{control}}}{dt} = 0$ , so

$$\sum \dot{m}_{\text{in}} = \sum \dot{m}_{\text{out}}.$$

$$m_{\text{control}}(t) = \int_V \rho dV.$$

### 3.2 Energy Conservation

$$\frac{dE_{\text{control}}}{dt} = \dot{Q} - \dot{W} + \dot{m}_{\text{in}} \left( u_{\text{in}} + \frac{v_{\text{in}}^2}{2} + gz_{\text{in}} \right) - \dot{m}_{\text{out}} \left( u_{\text{out}} + \frac{v_{\text{out}}^2}{2} + gz_{\text{out}} \right).$$

In addition to the work done by the system, there is also work done on the system by the surroundings. This is called flow work and is defined as

$$\dot{W}_{\text{flow}} = \dot{m}(p_{\text{out}} v_{\text{out}} - p_{\text{in}} v_{\text{in}}).$$

$$\dot{W} = \dot{W}_{\text{control}} + \dot{W}_{\text{flow}}.$$

so the equation for the energy of the system becomes

$$\frac{dE_{\text{control}}}{dt} = \dot{Q}_{\text{control}} - \dot{W}_{\text{control}} + \dot{m}_{\text{in}} \left( u_{\text{in}} + p_{\text{in}} v_{\text{in}} + \frac{v_{\text{in}}^2}{2} + gz_{\text{in}} \right) - \dot{m}_{\text{out}} \left( u_{\text{out}} + p_{\text{out}} v_{\text{out}} + \frac{v_{\text{out}}^2}{2} + gz_{\text{out}} \right).$$

Using the definition of the enthalpy, we can write

$$\frac{dE_{\text{control}}}{dt} = \dot{Q}_{\text{control}} - \dot{W}_{\text{control}} + \dot{m}_{\text{in}} \left( h_{\text{in}} + \frac{v_{\text{in}}^2}{2} + gz_{\text{in}} \right) - \dot{m}_{\text{out}} \left( h_{\text{out}} + \frac{v_{\text{out}}^2}{2} + gz_{\text{out}} \right).$$

In practice there may be several locations on the boundary through which mass enters or exists. The energy rate balance in this case becomes

$$\frac{dE_{\text{control}}}{dt} = \dot{Q}_{\text{control}} - \dot{W}_{\text{control}} + \sum \dot{m}_{\text{in}} \left( h_{\text{in}} + \frac{v_{\text{in}}^2}{2} + gz_{\text{in}} \right) - \sum \dot{m}_{\text{out}} \left( h_{\text{out}} + \frac{v_{\text{out}}^2}{2} + gz_{\text{out}} \right).$$

### 3.3 Energy Balance for Steady Flow

$$\dot{m}_{\text{in}} \left( h_{\text{in}} + \frac{v_{\text{in}}^2}{2} + gz_{\text{in}} \right) - \dot{m}_{\text{out}} \left( h_{\text{out}} + \frac{v_{\text{out}}^2}{2} + gz_{\text{out}} \right) = \dot{W}_{\text{control}} - \dot{Q}_{\text{control}}.$$

The mass rate balance reduces simply to  $\dot{m}_{\text{in}} = \dot{m}_{\text{out}}$

$$(h_{\text{in}} - h_{\text{out}}) + \frac{v_{\text{in}}^2 - v_{\text{out}}^2}{2} + g(z_{\text{in}} - z_{\text{out}}) = \frac{\dot{W}_{\text{control}}}{\dot{m}} - \frac{\dot{Q}_{\text{control}}}{\dot{m}}.$$

$\frac{\dot{Q}}{\dot{m}}$  and  $\frac{\dot{W}}{\dot{m}}$  are rates of heat and work transfer per unit mass of fluid flowing through the control volume. They are called specific heat and specific work respectively.

### 3.4 Nozzles and Diffusers

A nozzle is a device that increases the velocity of a fluid at the expense of pressure. A diffuser is a device that increases the pressure of a fluid at the expense of velocity.

The only work is flow work at locations where the fluid enters or exits the control volume. Thus  $\dot{W}_{\text{control}} = 0$ . The energy balance for steady flow becomes

$$0 = \frac{\dot{Q}_{\text{control}}}{\dot{m}} + \frac{v_{\text{in}}^2 - v_{\text{out}}^2}{2} + (h_{\text{in}} - h_{\text{out}}).$$

The term  $\dot{Q}_{\text{control}}$  is often negligible, so the energy balance becomes

$$\frac{v_{\text{in}}^2 - v_{\text{out}}^2}{2} + (h_{\text{in}} - h_{\text{out}}) = 0.$$

### 3.5 Turbines

A turbine is a device that extracts energy from a fluid.

For the turbine the kinetic potential energy of the matter flowing across the boundary is usually small enough to be neglected. The energy balance for steady flow becomes

$$0 = \dot{Q}_{\text{control}} - \dot{W}_{\text{control}} + \dot{m}(h_{\text{in}} - h_{\text{out}}).$$

If  $\dot{Q}_{\text{control}}$  is negligible, then

$$\dot{W}_{\text{control}} = \dot{m}(h_{\text{in}} - h_{\text{out}}).$$

### 3.6 Compressors and Pumps

A compressor is a device that increases the pressure of a gas. A pump is a device that increases the pressure of a liquid. The equations of the model are the same as the equations for the turbines

## 3.7 Heat Exchangers

A heat exchanger is a device that transfers heat from one fluid to another.

The only work is flow work at locations where the fluid enters or exits the control volume. Thus  $W_{\text{control}} = 0$ , and although heat transfer occurs, it is often small enough to be ignored

$$\sum \dot{m}_{\text{in}} h_{\text{in}} = \sum \dot{m}_{\text{out}} h_{\text{out}}.$$

## 3.8 Throttling Devices

A throttling device is a device that reduces the pressure of a fluid without changing its phase.

The only work is flow work at locations where the fluid enters or exits the control volume. Thus  $W_{\text{control}} = 0$ , and although heat transfer occurs, it is often small enough to be ignored, and it is safe to assume that there is no change in potential energy

$$h_1 + \frac{v_1^2}{2} = h_2 + \frac{v_2^2}{2}.$$

Although velocities may be relatively high in the vicinity of the restriction, measurements made upstream and downstream of the reduced flow area show in most cases that the change in the specific kinetic energy of the gas or liquid between these locations can be neglected. With this further simplification, the last equation reduces to

$$h_{\text{in}} = h_{\text{out}}.$$

## Chapter 4

# The Second Law of Thermodynamics

The second law can be used to determine whether a process is possible(spontaneously) and it's direction. It can also be used to determine the equilibrium state of a system, as well as it's maximum efficiency.

### 4.1 Clausius Statement

Clausius statement of the second law of thermodynamics states that it is impossible for any system to operate in such a way that the sole result would be the absorption of heat from a colder body to a warmer one without requiring work.

### 4.2 Kelvin-Planck Statement

Kelvin-Planck statement of the second law of thermodynamics states that it is impossible for any system to operate in a thermodynamic cycle and deliver a net amount of energy by work to its surroundings while receiving energy by heat transfer from a single thermal reservoir. In other words, It is impossible to take a quantity of heat  $Q$  of an energy source and the fully transform into work.

### 4.3 Entropy Statement

Entropy is a property of a system that is a measure of its molecular disorder or randomness. The entropy of a system is a function of state. It is impossible for any system to operate such that it's net entropy decreases.

$$S_{\text{gen}} \geq 0.$$

### Reversibility and Irreversibility

A process is called if the system and all irreversible parts of its surroundings cannot be restored to their initial states without leaving some trace on the surroundings. While a reversible process is one that can be reversed without leaving any trace on the surroundings.

## Chapter 5

# The Second Law of Thermodynamics

The efficiency of a cycle is given by

$$\eta = \frac{\text{Net Work Output}}{\text{Heat Input}} = \frac{W_{\text{net}}}{Q_H} = 1 - \frac{Q_C}{Q_H}.$$

According to the Kelvin-Planck statement,  $\eta_{\text{max}} < 1$  for any cycle.

### Corollary 5.0.1 Carnot Corollary 1

The thermal efficiency of an irreversible power cycle is always less than the thermal efficiency of a reversible power cycle when each operates between the same two thermal reservoirs.

### Corollary 5.0.2 Carnot Corollary 2

All reversible power cycles operating between the same two thermal reservoirs have the same thermal efficiency.

For a refrigeration cycle, the coefficient of performance is given by

$$\beta = \frac{\text{Heat Removed}}{\text{Work Input}} = \frac{Q_C}{W_{\text{net}}} = \frac{Q_C}{Q_H - Q_C}.$$

And for a heat pump cycle, the coefficient of performance is given by

$$\gamma = \frac{\text{Heat Added}}{\text{Work Input}} = \frac{Q_H}{W_{\text{net}}} = \frac{Q_H}{Q_H - Q_C}.$$

According to Clausius statement,  $W_{\text{net}} \neq 0$  so these coefficients are always finite in value.

$$\left( \frac{Q_C}{Q_H} \right)_{\text{rev}} = \frac{T_C}{T_H}.$$

$$\eta_{\text{max}} = 1 - \frac{T_C}{T_H}$$

$$\beta_{\text{max}} = \frac{T_C}{T_H - T_C}$$

$$\gamma_{\text{max}} = \frac{T_H}{T_H - T_C}$$

## 5.1 Carnot Cycle

**Process 1—2** The gas expands isothermally at  $T_C$  while receiving energy  $Q_C$  from the cold reservoir by heat transfer.

**Process 2—3** The gas is compressed adiabatically until its temperature is  $T_H$ .

**Process 3—4** The gas is compressed isothermally at  $T_H$  while it discharges energy  $Q_H$  to the hot reservoir by heat transfer.

**Process 4—1** The gas expands adiabatically until its temperature decreases to  $T_C$ .

## Chapter 6

# Entropy

$$\Delta S = \int \frac{\delta Q_{\text{rev}}}{T}.$$

### 6.1 Using the $T$ - $dS$ Equations

$$T dS = dU + p dV$$

$$T dS = dH - V dp$$

### 6.2 Entropy Change of an Ideal Gas

$$s(T_2, v_2) - s(T_1, v_1) = c_v \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1}$$

$$s(T_2, p_2) - s(T_1, p_1) = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1}$$

### 6.3 Entropy Change of an Incompressible Substance

$$s_2 - s_1 = c \ln \frac{T_2}{T_1}.$$

$$Q_{\text{rev}} = \int_1^2 T dS.$$

### 6.4 Entropy Balance for a Closed System

$$S_2 - S_1 = \int_1^2 \left( \frac{\delta Q}{T} \right)_b + \sigma.$$

where  $\sigma$  is the entropy generated.

The differential form of the entropy balance is

$$dS = \left( \frac{\delta Q}{T} \right)_b + \delta \sigma.$$

Closed system entropy rate balance is given by

$$\frac{dS}{dt} = \sum_j \frac{Q_j}{T_j} + \dot{\sigma}.$$

Control volume entropy rate balance is given by

$$\frac{dS_{\text{control}}}{dt} = \sum_j \frac{\dot{Q}_j}{T_j} + \sum \dot{m}_{\text{in}} s_{\text{in}} - \sum \dot{m}_{\text{out}} s_{\text{out}} + \dot{\sigma}.$$

At steady state, the entropy rate balance becomes

$$\sum_j \frac{\dot{Q}_j}{T_j} + \sum \dot{m}_{\text{in}} s_{\text{in}} - \sum \dot{m}_{\text{out}} s_{\text{out}} + \dot{\sigma} = 0.$$

For one inlet and one outlet, the entropy rate balance becomes

$$s_2 - s_1 = \frac{1}{\dot{m}} \left( \sum_j \frac{\dot{Q}_j}{T_j} \right) + \frac{\dot{\sigma}}{\dot{m}}.$$

## 6.5 Isentropic Processes

An isentropic process is a process during which the entropy of the system remains constant.

$$\Delta S = 0.$$

### 6.5.1 Isentropic Process for an Ideal Gas

$$p_2 = p_1 \exp \left[ \frac{s^\circ(T_2) - s^\circ(T_1)}{R} \right].$$

Assuming constant specific heats,

$$k = \frac{c_p}{c_v}.$$

$$c_p = \frac{kR}{k-1}$$

$$c_v = \frac{R}{k-1}$$

$$0 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1}$$

$$0 = c_v \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1}$$

$$\frac{T_2}{T_1} = \left( \frac{p_2}{p_1} \right)^{\frac{k-1}{k}}$$

$$\frac{T_2}{T_1} = \left( \frac{v_1}{v_2} \right)^{k-1}$$

$$\frac{p_2}{p_1} = \left( \frac{v_1}{v_2} \right)^k$$



## 6.6 Isentropic Efficiency

### 6.6.1 Isentropic Efficiency of a Turbine

$$\eta_t = \frac{h_1 - h_2}{h_1 - h_{2s}}.$$

where  $h_{2s}$  is the enthalpy at state 2 if the turbine were isentropic.

### 6.6.2 Isentropic Efficiency of a Nozzle

$$\eta_{\text{nozzle}} = \frac{h_1 - h_{2s}}{h_1 - h_2}.$$