

Force Based Model Analysis

Jerry Wu

June 25, 2025

Helpful Link:

- Geogebra Demo
- Canva Graphics

Assumption 0.1. *The robot body would never contact with the terrain*

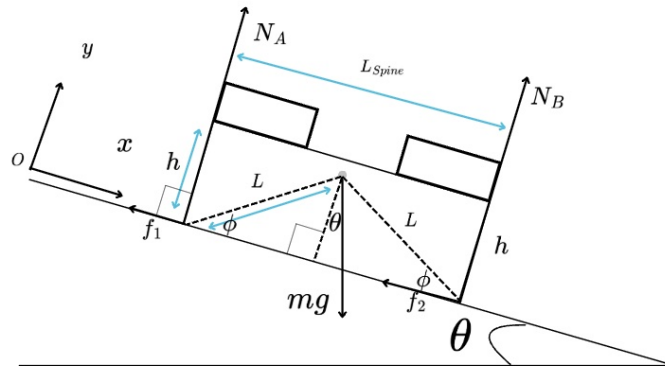
Assumption 0.2. *The CoM of the robot is always at local geometric center in direction parrallel to Spine*

1 Statics

Assumption 1.1. *Static means the robot would not move and achieve static equilibrium in this state*

Slope

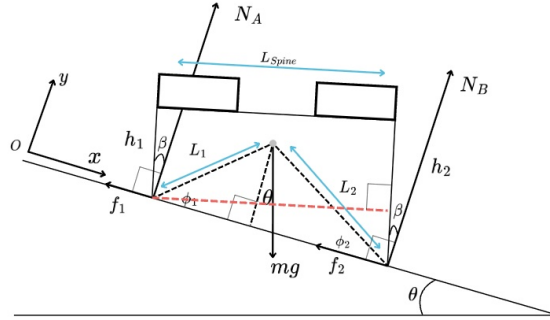
Diagram 1.2. *Simple case where two legs have same height*



To start with, we put robot on a slope with angle θ . To make the analysis easier, assume that two leg have same height h in this case. Therefore, the moment arm are symetric for R_A and R_B . Therefore, we can get follow equations:

$$\begin{aligned}\sum F_x &= (f_1 + f_2) - \sin(\theta)mg = 0 \\ \sum F_y &= (N_A + N_B) - \cos(\theta)mg = 0 \\ \sum M_A &= L_{\text{spine}} \cdot N_B - L \cdot mg \cdot \sin\left(\frac{\pi}{2} - \phi + \theta\right) = 0 \\ \sum M_B &= L_{\text{spine}} \cdot N_A - L \cdot mg \cdot \sin\left(\frac{\pi}{2} - \phi + \theta\right) = 0\end{aligned}$$

Diagram 1.3. *Two legs have different height*



Now extend the case to where the leg length would change. Here the robot extended the right leg to h_2 , which makes the robot rotate angle β around the z-Axis. Now both moment arms on two legs are shifted. We can also form the static equilibrium equation as follows:

$$\begin{aligned}\sum F_x &= (f_1 + f_2) - \sin(\theta)mg = 0 \\ \sum F_y &= (N_A + N_B) - \cos(\theta)mg = 0 \\ \sum M_A &= \frac{L_{\text{spine}}}{\sin\left(\frac{\pi}{2} - \beta\right)} \cdot N_B \cdot \sin\left(\frac{\pi}{2}\right) - L_1 \cdot mg \cdot \sin\left(\frac{\pi}{2} - \phi_1 + \theta\right) = 0 \\ \sum M_B &= \frac{L_{\text{spine}}}{\sin\left(\frac{\pi}{2} - \beta\right)} \cdot N_A \cdot \sin\left(\frac{\pi}{2}\right) - L_2 \cdot mg \cdot \sin\left(\frac{\pi}{2} - \phi_2 - \theta\right) = 0\end{aligned}$$

based on the geometry of our robot, we can also formulate:

$$L_{\text{spine}} \cdot \cos\left(\frac{\pi}{2} - \beta\right) = h_2 - h_1$$

Note 1.4. h_1, h_2 are known values, so β is also known. Therefore, ϕ_1 and ϕ_2 can easily be solved using Theorem 1.5.

Theorem 1.5. *Law of Cosines*

$$\cos A = \frac{b^2 + c^2 - a^2}{2bc}$$

$$\cos B = \frac{a^2 + c^2 - b^2}{2ac}$$

$$\cos C = \frac{a^2 + b^2 - c^2}{2ab}$$

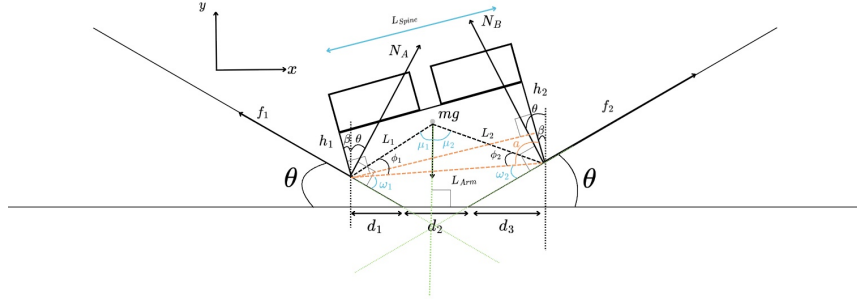
Zig-Zag Terrain

Assumption 1.6. *The terrain is uniform that all triangle obstacle and gap between them are same*

A natural extension for the slope is triangle terrain. There are essentially five cases when robot is static on the triangle terrain: 1. two legs on the same slope 2. two legs on two slopes that face each other 3. two legs on two slopes that does not face each other 4. one leg is on the slope and the other is on the gap. 5. two legs are all on the gap.

For case 1, we have already discussed in the previous section. And for case 5, it is essentially the same like slope when θ is 0.

Diagram 1.7. *Case2: two legs on the slope that face each other*



To make it clearer, I changed the coordinate that aligns with the ground rather slope from now on. We can also formulate the equilibrium equations:

$$\sum F_x = (-f_1 + f_2) \cos(\theta) + (N_A - N_B) \sin(\theta) = 0$$

$$\sum F_y = (f_1 + f_2) \sin(\theta) + (N_A + N_B) \cos(\theta) - mg = 0$$

$$\sum M_A = L_{Arm} \cdot (N_B \cdot \sin\left(\frac{\pi}{2} - \omega_2\right) + f_2 \cdot \sin(\pi - \omega_2)) - L_1 \cdot mg \cdot \sin(\mu_1) = 0$$

$$\sum M_B = L_{Arm} \cdot (N_A \cdot \sin\left(\frac{\pi}{2} - \omega_1\right) + f_1 \cdot \sin(\pi - \omega_1)) - L_2 \cdot mg \cdot \sin(\mu_2) = 0$$

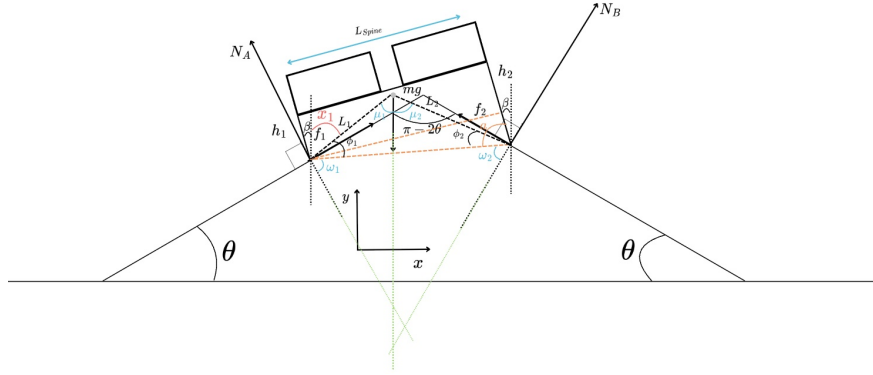
where by geometry of the shape, we can find:

$$\omega_1 = (\theta + \beta) - \left(\frac{\pi}{2} - \alpha\right)$$

$$\begin{aligned}\omega_2 &= (\theta - \beta) + \left(\frac{\pi}{2} - \alpha\right) \\ \mu_1 &= \pi - (\phi_1 + \omega_1) - \left(\frac{\pi}{2} - \theta\right) \\ \mu_2 &= \pi - (\phi_2 + \omega_2) - \left(\frac{\pi}{2} - \theta\right) \\ \alpha &= \arctan\left(\frac{L_{\text{Spine}}}{h_2 - h_1}\right) \\ L_{\text{Arm}} &= \sin(\alpha) \cdot L_{\text{Spine}}\end{aligned}$$

similarly to Note 1.4, this equation can be solved.

Diagram 1.8. *Case3: two legs on two slopes that do not face each other*



$$\begin{aligned}\sum F_x &= (f_1 - f_2) \cos(\theta) + (-N_A + N_B) \sin(\theta) = 0 \\ \sum F_y &= (f_1 - f_2) \sin(\theta) + (-N_A + N_B) \cos(\theta) - mg = 0 \\ \sum M_A &= L_{\text{Arm}} \cdot (N_B \cdot \sin(\pi - \omega_2) + f_2 \cdot \sin\left(\frac{\pi}{2} - \omega_2\right)) - L_1 \cdot mg \cdot \sin(\mu_1) = 0 \\ \sum M_B &= L_{\text{Arm}} \cdot (N_A \cdot \sin(\pi - \omega_1) + f_1 \cdot \sin\left(\frac{\pi}{2} - \omega_1\right)) - L_2 \cdot mg \cdot \sin(\mu_2) = 0\end{aligned}$$

where by geometry of the shape, we can find:

$$\begin{aligned}\omega_1 &= \beta + \alpha - \theta \\ \omega_2 &= \pi - \beta - \alpha - \theta \\ \mu_1 &= \pi - (\phi_1 + \omega_1) - \theta \\ \mu_2 &= \pi - (\phi_2 + \omega_2) - \theta \\ \alpha &= \arctan\left(\frac{L_{\text{Spine}}}{h_2 - h_1}\right)\end{aligned}$$

$$L_{\text{Arm}} = \sin(\alpha) \cdot L_{\text{Spine}}$$

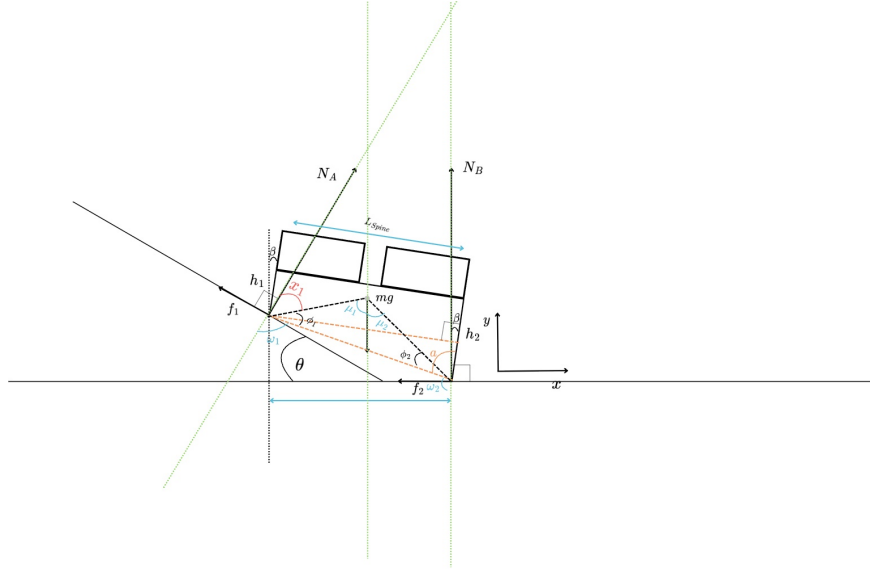
Note 1.9. The equation for ω_1 can be found specifically by following equation:

$$x_1 + \phi_1 + \omega_1 + \theta = \pi \quad (1)$$

$$\beta + x_1 + \phi_1 - \left(\frac{\pi}{2} - \alpha\right) = \frac{\pi}{2} \quad (2)$$

$$\omega_1 = \beta + \alpha - \theta \quad \text{From (1) and (2)}$$

Diagram 1.10. Case4: one leg is on the slope and the other is on the gap



$$\sum F_x = -f_1 \cos(\theta) - f_2 + N_A \sin(\theta) = 0$$

$$\sum F_y = f_1 \sin(\theta) + N_A \cos(\theta) + N_B - mg = 0$$

$$\sum M_A = L_{\text{Arm}} \cdot (N_B \cdot \sin(\pi - \omega_2) + f_2 \cdot \sin\left(\frac{\pi}{2} - \omega_2\right)) - L_1 \cdot mg \cdot \sin(\mu_1) = 0$$

$$\sum M_B = L_{\text{Arm}} \cdot (N_B \cdot \sin(\pi - \omega_1) + f_1 \cdot \sin\left(\frac{3\pi}{2} - \omega_1\right)) - L_2 \cdot mg \cdot \sin(\mu_2) = 0$$

where by geometry of the shape, we can find:

$$\omega_1 = \theta - \beta + \alpha$$

$$\omega_2 = \pi + \beta - \alpha$$

$$\mu_1 = \pi - (\phi_1 + \omega_1) - \theta$$

$$\mu_2 = \pi - (\phi_2 + \omega_2)$$

$$\alpha = \arctan\left(\frac{L_{\text{Spine}}}{h_2 - h_1}\right)$$

$$L_{\text{Arm}} = \sin(\alpha) \cdot L_{\text{Spine}}$$

Note 1.11. The equation for ω_1 can be found specifically by following equation:

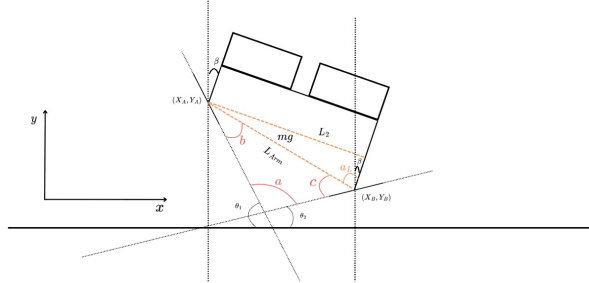
$$x_1 + \phi_1 + \omega_1 = \pi \quad (3)$$

$$\theta - \beta + x_1 + \phi_1 - \left(\frac{\pi}{2} - \alpha\right) = \frac{\pi}{2} \quad (4)$$

$$\omega_1 = \theta - \beta + \alpha \quad \text{From (1) and (2)}$$

Now, by observation, this case is very similar to previous case. In fact, all of the above cases are geometrically similar, which means we can find a general expression for this force based model.

Diagram 1.12. Determine the geometry shape by robot body and terrain



After simplifying the model, it can be seen that given the terrain tangent angle θ_1 , θ_2 , robot pitch angle β and length of the virtual link between end point of two leg L_{Arm} we can determine the triangle with inner angle a, b, c .

Here is how we determine angle a, b, c :

$$a = \pi - \theta_1 - \theta_2$$

$$b = \pi - \beta - \frac{\pi}{2} - \left(\frac{\pi}{2} - \alpha_L\right) - (\pi - \theta_1)$$

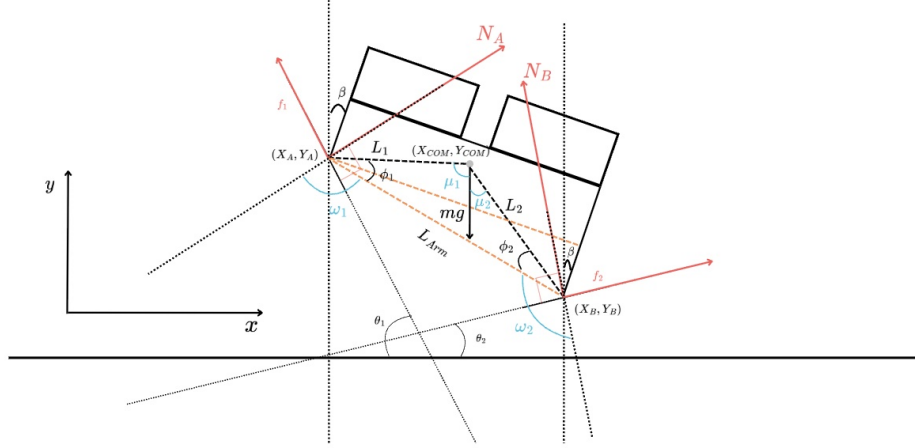
$$c = \pi - (\pi - \theta_2) + \beta - \alpha_L$$

by verification:

$$a + b + c = \pi$$

As for different tangent angles, this property should hold as well. This means there would be **only one solution** by given condition.

Diagram 1.13. *Generailized Free Body Diagram*



Given the determined geometric shape of the robot and terrain, we can easily find the position of the COM through the given of relative position L_1 and L_2 , which can be solved by the geometry as well. Which means, if we put the robot into coordinate, the points (X_A, Y_A) , (X_B, Y_B) , (X_{COM}, Y_{COM}) are given.

Therefore, given (X_A, Y_A) , (X_B, Y_B) , (X_{COM}, Y_{COM}) , θ_1, θ_2 , L_{Spine} we can formulate

$$\sum \vec{F} = \vec{f}_1 + \vec{f}_2 + \vec{N}_A + \vec{N}_B + \vec{W} = \vec{0}$$

$$\sum \vec{M}_A = (\vec{r}_B - \vec{r}_A) \times (\vec{f}_2 + \vec{N}_B) + (\vec{r}_{COM} - \vec{r}_A) \times \vec{W} = 0$$

$$\sum \vec{M}_B = (\vec{r}_A - \vec{r}_B) \times (\vec{f}_1 + \vec{N}_A) + (\vec{r}_{COM} - \vec{r}_B) \times \vec{W} = 0$$

where

$$\hat{t}_1 = (\cos(\pi - \theta_1), \sin(\pi - \theta_1)) \quad (\text{contact tangent 1})$$

$$\hat{n}_1 = (-\sin(\pi - \theta_1), \cos(\pi - \theta_1)) \quad (\text{contact normal 1})$$

$$\hat{t}_2 = (\cos \theta_2, \sin \theta_2)$$

$$\hat{n}_2 = (-\sin \theta_2, \cos \theta_2)$$

$$\begin{aligned}
\vec{f}_1 &= f_1 \cdot \hat{t}_1 = f_1 \cdot (\cos(\pi - \theta_1), \sin(\pi - \theta_1)) \\
\vec{N}_A &= N_A \cdot \hat{n}_1 = N_A \cdot (-\sin(\pi - \theta_1), \cos(\pi - \theta_1)) \\
\vec{f}_2 &= f_2 \cdot \hat{t}_2 = f_2 \cdot (\cos \theta_2, \sin \theta_2) \\
\vec{N}_B &= N_B \cdot \hat{n}_2 = N_B \cdot (-\sin \theta_2, \cos \theta_2) \\
\vec{W} &= -mg \cdot (0, 1)
\end{aligned}$$

Moment arm of \vec{f}_2, \vec{N}_B about A : $\vec{r}_{f_2/A} = \vec{r}_B - \vec{r}_A = (X_B - X_A, Y_B - Y_A)$

Moment arm of \vec{W} about A : $\vec{r}_{W/A} = \vec{r}_{COM} - \vec{r}_A = (X_{COM} - X_A, Y_{COM} - Y_A)$

Moment arm of \vec{f}_1, \vec{N}_A about B : $\vec{r}_{f_1/B} = \vec{r}_A - \vec{r}_B = (X_A - X_B, Y_A - Y_B)$

Moment arm of \vec{W} about B : $\vec{r}_{W/B} = \vec{r}_{COM} - \vec{r}_B = (X_{COM} - X_B, Y_{COM} - Y_B)$

Therefore,

$$\sum \vec{F}_x = -f_1 \cos \theta_1 + f_2 \cos \theta_2 - N_A \sin \theta_1 - N_B \sin \theta_2 = 0$$

$$\sum \vec{F}_y = f_1 \sin \theta_1 + f_2 \sin \theta_2 - N_A \cos \theta_1 + N_B \cos \theta_2 - mg = 0$$

$$\begin{aligned}
\sum \vec{M}_A &= (X_B - X_A)(f_2 \sin \theta_2 + N_B \cos \theta_2) - (Y_B - Y_A)(f_2 \cos \theta_2 - N_B \sin \theta_2) - mg(X_{COM} - X_A) = 0 \\
\sum \vec{M}_B &= (X_A - X_B)(f_1 \sin \theta_1 - N_A \cos \theta_1) - (Y_A - Y_B)(-f_1 \cos \theta_1 - N_A \sin \theta_1) - mg(X_{COM} - X_B) = 0
\end{aligned}$$

solve,

$$\begin{aligned}
N_A &= \frac{gm(X_B - X_{COM}) \cos \theta_1}{X_A - X_B} \\
N_B &= \frac{gm(X_A - X_{COM}) \cos \theta_2}{X_A - X_B} \\
f_1 &= \frac{gm(X_{COM} - X_B) \sin \theta_1}{X_A - X_B} \\
f_2 &= \frac{gm(X_A - X_{COM}) \sin \theta_2}{X_A - X_B}
\end{aligned}$$

By friction function $|f| \leq \mu N$, we can solve this equation and find

$$|f_1| = \left| \frac{gm(X_{\text{COM}} - X_B) \sin \theta_1}{X_A - X_B} \right| \leq \mu \cdot \frac{gm(X_B - X_{\text{COM}}) \cos \theta_1}{X_A - X_B} = \mu N_A$$

$$|f_2| = \left| \frac{gm(X_A - X_{\text{COM}}) \sin \theta_2}{X_A - X_B} \right| \leq \mu \cdot \frac{gm(X_A - X_{\text{COM}}) \cos \theta_2}{X_A - X_B} = \mu N_B$$

The above expression is derived in case where the θ_1 and θ_2 are both acute angles from graphic definition, but it is clear that in other cases, they are the same, a even more genral expression is actually:

$$\begin{aligned}\vec{f}_1 &= f_1 \cdot \hat{t}_1 \\ \vec{N}_A &= N_A \cdot \hat{n}_1 \\ \vec{f}_2 &= f_2 \cdot \hat{t}_2 \\ \vec{N}_B &= N_B \cdot \hat{n}_2 \\ \vec{W} &= -mg \cdot (0, 1)\end{aligned}$$

$$\begin{aligned}\sum F_x &= f_1 t_{1x} + f_2 t_{2x} + N_A n_{1x} + N_B n_{2x} = 0 \\ \sum F_y &= f_1 t_{1y} + f_2 t_{2y} + N_A n_{1y} + N_B n_{2y} - mg = 0 \\ \sum M_A &= (x_B - x_A)(f_2 t_{2y} + N_B n_{2y}) - (y_B - y_A)(f_2 t_{2x} + N_B n_{2x}) \\ &\quad - mg(x_{\text{COM}} - x_A) = 0 \\ \sum M_B &= (x_A - x_B)(f_1 t_{1y} + N_A n_{1y}) - (y_A - y_B)(f_1 t_{1x} + N_A n_{1x}) \\ &\quad - mg(x_{\text{COM}} - x_B) = 0\end{aligned}$$

In matrix form:

$$\begin{bmatrix} t_{1x} & t_{2x} & n_{1x} & n_{2x} \\ t_{1y} & t_{2y} & n_{1y} & n_{2y} \\ 0 & \Delta x_{BA} t_{2y} - \Delta y_{BA} t_{2x} & 0 & \Delta x_{BA} n_{2y} - \Delta y_{BA} n_{2x} \\ \Delta x_{AB} t_{1y} - \Delta y_{AB} t_{1x} & 0 & \Delta x_{AB} n_{1y} - \Delta y_{AB} n_{1x} & 0 \end{bmatrix} \cdot \begin{bmatrix} f_1 \\ f_2 \\ N_A \\ N_B \end{bmatrix} = \begin{bmatrix} 0 \\ mg \\ mg(x_{\text{COM}} - x_A) \\ mg(x_{\text{COM}} - x_B) \end{bmatrix}$$

Friction Cone Constraints

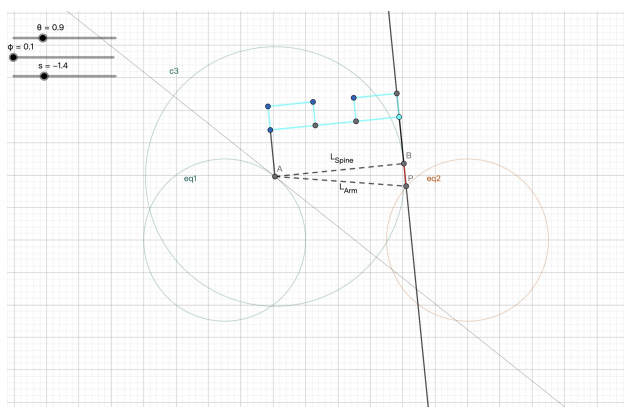
$$|\vec{f}_1| \leq \mu |N_A|, \quad |\vec{f}_2| \leq \mu |N_B|$$

Cylinder Terrain

Diagram 1.14. *Generailized expression on Cylinder Terrain*

I made a simple demo for robot geometric shape can be found in this [GeoGebra link](#). Here A represents the contact point of left leg on the left circle and AB represents of L_{Spine} . The moving P point represents the right leg extension and AP represents the L_{Arm} . For both following cases, we form equations as follows in coordinates:

Diagram 1.17. *When Point P contacts the terrain*



Obstacle c_1	$(x - c_1)^2 + y^2 = r^2$
----------------	---------------------------

Obstacle c_2 $(x - c_2)^2 + y^2 = r^2$

Control Parameters

$\theta \in [0, \pi]$	(defines point A on the upper half of the circle centered at c_1)
$\varphi \in [0, 2\pi]$	(controls the direction of the robot spine extended from point A)
$s \in [-L_{\text{leg}}, 0]$	(represents the extension length of the right leg from point B)

Point Definitions

$$A = (x_0, y_0) = (c_1 + r \cos \theta, r \sin \theta)$$

$$B = (x_1, y_1) = (x_0 + L_{\text{spine}} \cos \varphi, y_0 + L_{\text{spine}} \sin \varphi)$$

Body Reachability Circle

$$c_A : (x - x_0)^2 + (y - y_0)^2 = L_{\text{Spine}}^2$$

Direction Vector and Point P

$$\vec{t} = (t_x, t_y) = (-\sin \varphi, \cos \varphi)$$

$$P = (x_P, y_P) = (x_1 + st_x, y_1 + st_y)$$

Collision Condition Generally speaking, segment \overline{BP} must intersect the boundary of Obstacle 2 exactly once:

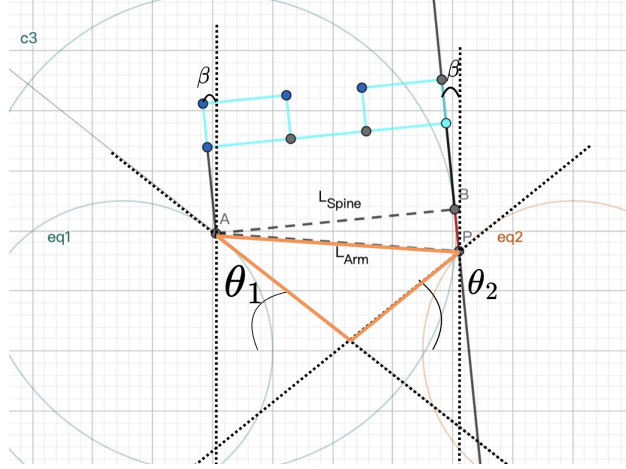
$$|\overline{BP} \cap \partial c_2| = 1$$

$$\overline{BP} \cap c_2^\circ = \emptyset$$

Case1

In this case, it is very similar to what we did in the previous sections that we would consider pitch angle β in consideration. With that, we can determine triangle with knowing that A and P are on c_1 and c_2 and their angle.

Diagram 1.19. *Determine the geometric shape*



- Two circles c_1 and c_2 with centers $C_1 = (x_1, y_1)$ and $C_2 = (x_2, y_2)$ respectively, both of radius r .
- Unknown points $A \in \partial c_1$, $P \in \partial c_2$.
- Length $L = \|A - P\|$ is known.
- Direction angle of \vec{AP} is known: $\psi \in [0, 2\pi]$ from pitch angle β .

Note 1.20. To be more specific, direction angle is derived from both the geometric angle formed by triangle ABP and β .

Then the direction vector is:

$$\vec{d} = \vec{AP} = L(\cos \psi, \sin \psi)$$

We parameterize:

$$A = (x_A, y_A)$$

$$P = A + \vec{d} = (x_A + L \cos \psi, y_A + L \sin \psi)$$

Apply the circle constraints:

$$\|A - C_1\|^2 = r^2 \tag{1}$$

$$\|P - C_2\|^2 = r^2 \tag{2}$$

Substituting into (2):

$$(x_A + L \cos \psi - x_2)^2 + (y_A + L \sin \psi - y_2)^2 = r^2 \tag{2'}$$

Now, equations (1) and (2') form a nonlinear system in x_A, y_A .

Once A is solved, compute:

$$P = A + L(\cos \psi, \sin \psi)$$

Then, compute the angle between the **tangent direction** at the contact points and the ground (i.e., the horizontal x-axis). The tangent direction at a point on a circle is perpendicular to the radial vector from the circle's center to that point.

Let the centers of the circles be $C_1 = (x_1, y_1)$ and $C_2 = (x_2, y_2)$, and the contact points be $A = (x_A, y_A)$ and $P = (x_P, y_P)$. The radial vectors are:

$$\vec{n}_1 = A - C_1 = (x_A - x_1, y_A - y_1), \quad \vec{n}_2 = P - C_2 = (x_P - x_2, y_P - y_2)$$

The corresponding tangent vectors (obtained by rotating the radial vectors 90° clockwise) are:

$$\vec{t}_1 = (-(y_A - y_1), x_A - x_1), \quad \vec{t}_2 = (-(y_P - y_2), x_P - x_2)$$

The angles between the tangent vectors at the contact points and the ground (horizontal axis) are given by:

$$\theta_1 = \text{mod}(\arctan 2(t_{1,y}, t_{1,x}), 2\pi)$$

$$\theta_2 = \text{mod}(\arctan 2(t_{2,y}, t_{2,x}), 2\pi)$$

These angles represent the **oriented direction** of the tangents relative to the ground, measured counterclockwise from the positive x -axis and wrapped into the interval $[0, 2\pi)$.

Case2

For case 2, we can express that tangent line for robot body reachability circle at point B also must tangent with c_2 . Therefore we can formulate:

Let $\vec{n} := B - A$

Define a line ℓ such that:

$$B \in \ell, \quad \vec{d}_\ell \cdot \vec{n} = 0 \quad \text{i.e., } \ell \text{ is tangent to } c_A \text{ at } B$$

$$\exists P' \in \ell \cap \partial c_2, \quad \vec{d}_\ell \cdot (P' - C_2) = 0 \quad \text{i.e., } \ell \text{ is tangent to } c_2 \text{ at } P'$$

$$|BP'| \leq |BP|$$

Note 1.21. For Case2, the pitch angle β does not matter and we can still find at most one solution.

Let \vec{d}_ℓ be the direction vector of ℓ , defined by:

$$\vec{d}_\ell = (-n_y, n_x) \quad (\text{perpendicular to } \vec{n} = B - A)$$

Then, the parametric equation of ℓ is:

$$\ell(\lambda) = B + \lambda \vec{d}_\ell$$

To ensure tangency with c_2 , we require:

$$\|\ell(\lambda) - C_2\|^2 = r^2$$

which leads to a quadratic in λ :

$$\|B + \lambda \vec{d}_\ell - C_2\|^2 = r^2$$

Solving this yields candidate point(s) $P' = B + \lambda^* \vec{d}_\ell$. Select the point that satisfies:

$$\|BP'\| \leq \|BP\|$$

Finally, the center A of c_A is obtained by inverting \vec{n} :

$$\begin{aligned} \vec{n}_{\text{unit}} &= \frac{\vec{n}}{\|\vec{n}\|} \\ A &= B - R \cdot \vec{n}_{\text{unit}} \end{aligned}$$

This determines A and P' geometrically under the tangency constraints.

2 From d'Alembert principle to Dynamics

The d'Alembert principle provides a natural bridge from statics to dynamics.

Theorem 2.1.

$$\sum_i (\mathbf{F}_i - m_i \mathbf{a}_i) \cdot \delta \mathbf{r}_i = 0$$

where \mathbf{F}_i is the applied force on the i -th particle, $m_i \mathbf{a}_i$ is its inertial force, and $\delta \mathbf{r}_i$ denotes any virtual displacement compatible with the system's constraints.

Let the position of the robot's center of mass be denoted by $\mathbf{r}_G^{\mathcal{I}} \in \mathbb{R}^2$. And two leg contact position as $\mathbf{p}_A^{\mathcal{I}}$ and $\mathbf{p}_B^{\mathcal{I}}$.

2.1 Holonomic constraints:

Using the identity $\sin(\arctan(x)) = \frac{x}{\sqrt{1+x^2}}$, the geometric constraint imposed by the spine length can be rewritten as:

$$L_{\text{Arm}} = \|\mathbf{p}_B^{\mathcal{I}} - \mathbf{p}_A^{\mathcal{I}}\| = \frac{L_{\text{Spine}}^2}{\sqrt{L_{\text{Spine}}^2 + (L_{\text{Leg2}} - L_{\text{Leg1}})^2}}$$

The terrain contact constraints are then written as:

$$f_A := p_{A,y}^{\mathcal{I}} - h(p_{A,x}^{\mathcal{I}}) = 0, \quad f_B := p_{B,y}^{\mathcal{I}} - h(p_{B,x}^{\mathcal{I}}) = 0$$

In general, the center of mass can be written as a function of the two foot positions $\mathbf{p}_A^{\mathcal{I}}$ and $\mathbf{p}_B^{\mathcal{I}}$, depending on the mass distribution and geometry:

$$\mathbf{r}_G^{\mathcal{I}} = \mathbf{f}_G(\mathbf{p}_A^{\mathcal{I}}, \mathbf{p}_B^{\mathcal{I}}, \beta)$$

where β is the pitch angle of the spine, measured from horizontal and it can be derived through $\mathbf{f}_\beta(\mathbf{p}_A^{\mathcal{I}}, \mathbf{p}_B^{\mathcal{I}}, L_{\text{Spine}}, L_{\text{Leg1}}, L_{\text{Leg2}})$ by geometric constrains following:

1. $\|\mathbf{q}_1 - \mathbf{p}_A\| = L_{\text{Leg1}}$
2. $\|\mathbf{q}_2 - \mathbf{p}_B\| = L_{\text{Leg2}}$
3. $\|\mathbf{q}_2 - \mathbf{q}_1\| = L_{\text{Spine}}$
4. Directional constraint:

$$\mathbf{q}_1 - \mathbf{p}_A \parallel \mathbf{q}_2 - \mathbf{p}_B \Rightarrow (\mathbf{q}_1 - \mathbf{p}_A) = \lambda(\mathbf{q}_2 - \mathbf{p}_B)$$

Therefore, we get:

$$\begin{aligned}(x_1 - x_A)^2 + (y_1 - y_A)^2 &= L_{\text{Leg1}}^2 \\(x_2 - x_B)^2 + (y_2 - y_B)^2 &= L_{\text{Leg2}}^2 \\(x_2 - x_1)^2 + (y_2 - y_1)^2 &= L_{\text{Spine}}^2 \\(x_1 - x_A)(y_2 - y_B) - (y_1 - y_A)(x_2 - x_B) &= 0\end{aligned}$$

And therefore, we can get the solution of x_1, x_2, y_1, y_2 through numerical method. And then, we can find the β by:

$$\beta = \text{atan2}(y_1 - y_A, x_1 - x_A)$$

Furthure more, we can even determine the **CoM** by:

$$\mathbf{r}_G^{\mathcal{I}} = \begin{bmatrix} x_m \\ y_m \end{bmatrix} + d \begin{bmatrix} \cos \beta \\ \sin \beta \end{bmatrix}, \quad \text{where} \quad \begin{bmatrix} x_m \\ y_m \end{bmatrix} = \frac{1}{2} \begin{bmatrix} x_1 + x_2 \\ y_1 + y_2 \end{bmatrix}, \text{ and } d \text{ is offset}$$

2.2 Non-Holonomic constraints:

There are none.

Therefore, the total **DoF** is $3N - 5 = 1$. This indicates we can desribe the whole system through one generalized coordinate. And we can choose $x_A^{\mathcal{I}}$ for simplicity. We can describe as $q := x_A^{\mathcal{I}}$. All other quantities can be written as functions of q due to the holonomic constraints.

$$\mathbf{p}_A^{\mathcal{I}}(q) = \begin{bmatrix} q \\ h(q) \end{bmatrix}, \quad \mathbf{p}_B^{\mathcal{I}}(q) = \begin{bmatrix} \mathbf{f}_B(q) \\ h(\mathbf{f}_B(q)) \end{bmatrix}, \quad \mathbf{r}_G^{\mathcal{I}} = \begin{bmatrix} \mathbf{f}_{Gx}(q) \\ \mathbf{f}_{Gy}(q) \end{bmatrix}$$

Then, the velocity and acceleration of the **CoM** can be expressed as:

$$\dot{\mathbf{r}}_G = \frac{d\mathbf{f}_G}{dq} \cdot \dot{q}, \quad \ddot{\mathbf{r}}_G = \frac{d\mathbf{f}_G}{dq} \cdot \ddot{q} + \frac{d^2\mathbf{f}_G}{dq^2} \cdot \dot{q}^2$$

Therefore, through d'Alembert principle, we write:

$$(\mathbf{f}_A + \mathbf{f}_B + \mathbf{F}_g - m\ddot{\mathbf{r}}_G) \cdot \delta \mathbf{r}_G = 0$$

2.3 Sinusoidal Terrain

Lets start with the sinusoidal terrain where the terrain is easy to express with one function:

$$h(x^{\mathcal{I}}) = A \cdot \sin\left(\frac{2\pi}{T}x^{\mathcal{I}} + \phi\right) + B$$

We can set B as $\frac{A}{2}$ and θ as 0 and gives us:

$$h(x^{\mathcal{I}}) = A \cdot \sin\left(\frac{2\pi}{T} x^{\mathcal{I}}\right) + \frac{A}{2}$$

Therefore, we can express our three key points explicitly as:

$\mathbf{p}_A^{\mathcal{I}}(q) = \begin{bmatrix} q \\ A \sin\left(\frac{2\pi}{T} q\right) + \frac{A}{2} \end{bmatrix}$	directly defined by q
$\mathbf{p}_B^{\mathcal{I}}(q) = \begin{bmatrix} x_B(q) \\ A \sin\left(\frac{2\pi}{T} x_B(q)\right) + \frac{A}{2} \end{bmatrix}$	
$x_B(q)$ solves: $(x_B - q)^2 + \left[A \sin\left(\frac{2\pi}{T} x_B\right) - A \sin\left(\frac{2\pi}{T} q\right)\right]^2 = L_{\text{Arm}}^2$	
$x_m = \frac{1}{2}(x_1 + x_2), \quad y_m = \frac{1}{2}(y_1 + y_2)$	
$\beta = \text{atan2}(y_1 - y_A, x_1 - x_A)$	
$\mathbf{r}_G^{\mathcal{I}}(q) = \begin{bmatrix} x_m \\ y_m \end{bmatrix} + d \begin{bmatrix} \cos \beta \\ \sin \beta \end{bmatrix}$	x_1, x_2, y_1, y_2 solves: $(x_1 - x_A)^2 + (y_1 - y_A)^2 = L_{\text{Leg1}}^2$ $(x_2 - x_B)^2 + (y_2 - y_B)^2 = L_{\text{Leg2}}^2$ $(x_2 - x_1)^2 + (y_2 - y_1)^2 = L_{\text{Spine}}^2$ $(x_1 - x_A)(y_2 - y_B) - (y_1 - y_A)(x_2 - x_B) = 0$

Therefore, we can now expand out the equation for virtual work as follows:

$$\left(\mathbf{f}_A \cdot \frac{d\mathbf{p}_A}{dq} + \mathbf{f}_B \cdot \frac{d\mathbf{p}_B}{dq} + \mathbf{F}_g \cdot \frac{d\mathbf{r}_G}{dq} - m\ddot{\mathbf{r}}_G \cdot \frac{d\mathbf{r}_G}{dq} \right) \delta q = 0$$

$$\mathbf{f}_A = \lambda_N^A \hat{\mathbf{n}}_A + \lambda_f^A \hat{\mathbf{t}}_A + \lambda_{\parallel}^A \hat{\mathbf{e}}_{\beta} + \lambda_{\perp}^A \hat{\mathbf{e}}_{\beta}^{\perp}$$

$$\mathbf{f}_B = \lambda_N^B \hat{\mathbf{n}}_B + \lambda_f^B \hat{\mathbf{t}}_B + \lambda_{\parallel}^B \hat{\mathbf{e}}_{\beta} + \lambda_{\perp}^B \hat{\mathbf{e}}_{\beta}^{\perp}$$

where we can expand out as:

$$\begin{aligned} & \underbrace{\left(\lambda_N^A \hat{\mathbf{n}}_A \cdot \frac{d\mathbf{p}_A}{dq} + \lambda_f^A \hat{\mathbf{t}}_A \cdot \frac{d\mathbf{p}_A}{dq} + \lambda_{\parallel}^A \hat{\mathbf{e}}_{\beta} \cdot \frac{d\mathbf{p}_A}{dq} + \lambda_{\perp}^A \hat{\mathbf{e}}_{\beta}^{\perp} \cdot \frac{d\mathbf{p}_A}{dq} \right)}_{\mathbf{f}_A \cdot \frac{d\mathbf{p}_A}{dq}} \\ & + \underbrace{\left(\lambda_N^B \hat{\mathbf{n}}_B \cdot \frac{d\mathbf{p}_B}{dq} + \lambda_f^B \hat{\mathbf{t}}_B \cdot \frac{d\mathbf{p}_B}{dq} + \lambda_{\parallel}^B \hat{\mathbf{e}}_{\beta} \cdot \frac{d\mathbf{p}_B}{dq} + \lambda_{\perp}^B \hat{\mathbf{e}}_{\beta}^{\perp} \cdot \frac{d\mathbf{p}_B}{dq} \right)}_{\mathbf{f}_B \cdot \frac{d\mathbf{p}_B}{dq}} \\ & + \mathbf{F}_g \cdot \frac{d\mathbf{r}_G}{dq} - m\ddot{\mathbf{r}}_G \cdot \frac{d\mathbf{r}_G}{dq} \delta q = 0 \end{aligned}$$

1. $\lambda_N^A \hat{\mathbf{n}}_A \cdot \frac{d\mathbf{p}_A}{dq} = 0$ (Normal force is perpendicular to virtual displacement: does no work)
2. $\lambda_f^A \hat{\mathbf{t}}_A \cdot \frac{d\mathbf{p}_A}{dq} = \lambda_f^A \cdot \left\| \frac{d\mathbf{p}_A}{dq} \right\|$, because $\hat{\mathbf{t}}_A \parallel \frac{d\mathbf{p}_A}{dq}$
3. $\left(\lambda_{\parallel}^A \cdot \frac{d\mathbf{p}_A}{dq} + \lambda_{\parallel}^B \cdot \frac{d\mathbf{p}_B}{dq} \right) \cdot \hat{\mathbf{e}}_{\beta}$ (Combined contribution of spine-direction forces at A and B)
4. $\left(\lambda_{\perp}^A \cdot \frac{d\mathbf{p}_A}{dq} + \lambda_{\perp}^B \cdot \frac{d\mathbf{p}_B}{dq} \right) \cdot \hat{\mathbf{e}}_{\beta}^{\perp}$ (Combined contribution of perpendicular forces along the direction orthogonal to the spine)

5.

$$\hat{\mathbf{t}}_B = \frac{1}{\sqrt{1 + (h'(x_B))^2}} \begin{bmatrix} 1 \\ h'(x_B) \end{bmatrix}, \quad \hat{\mathbf{n}}_B = \frac{1}{\sqrt{1 + (h'(x_B))^2}} \begin{bmatrix} -h'(x_B) \\ 1 \end{bmatrix}.$$

2.4 Zig-Zag Terrain

The zig-zag terrain is defined by a periodic piecewise function $y = h(x^{\mathcal{I}})$, with period $T = 2r + d$ and height $h = r \tan \theta$. For each point $x^{\mathcal{I}} \in \mathbb{R}$, define $x_k = x^{\mathcal{I}} \bmod T$. Then:

$$h(x^{\mathcal{I}}) = \begin{cases} \tan \theta \cdot x_k, & x_k \in [0, r] \\ \tan \theta \cdot (2r - x_k), & x_k \in [r, 2r] \\ 0, & x_k \in [2r, 2r + d] \end{cases}$$

Here, in order to make the terrain differentiable at any point, we can use fourier series to approximate the terrain like:

$$h(x^{\mathcal{I}}) = \frac{8r \tan \theta}{\pi^2} \sum_{n=1,3,5,\dots}^N \frac{1}{n^2} \cos \left(\frac{2\pi n x}{2r + d} \right)$$