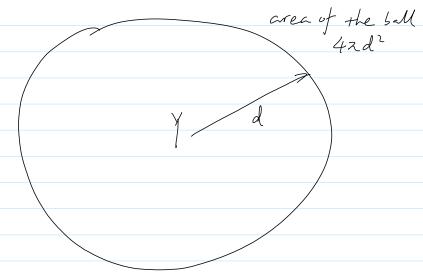
_ec4-mwf					
Monday, January 14, 2008 9:51 PM					
- Lec4supp					
- Fig 2.4-2.5 of schwartz.					

Free-space model-15min

Saturday, January 12, 2008 10:53 AM - No obstades



First, assume that power emitted from the transmitting antenna isotropically. (omni-directional)

PT = transhissin power

From the conservation of power, the receiver power density at distance d is

Spr (d) = 1212

If the receiving antenna has an effective area (or operture) AR, the received power 13

MR<1 is the efficiency parameter of the receiving attenna, due to

- transmission line attenuation

- filter loss - Jantenna loss, etc.

Effective area of isotropic antenna is $A_R = \frac{\lambda^2}{42}$, $P_R(d) = l_T(\frac{\lambda}{4\lambda d})$ - Note that this decreases as IV (on fg) $\lambda = c/f$ c=3*10^8 m/s Gain due to directional Antennas Directional antennas can provide a gain factor over omnidirectional antennas energy transed to directional antenna gain in one direction t egnal-strength curve among the cancel received signeds of each antenna element \$ 0.7. phase difference $=\frac{r-s.in\theta}{\lambda}$ directional antenna built from

Both transmitter of the receiver can have such type of gains.

At the tomsmitter side,

G7 = gain factor of the transmitting

Got is proportional to the effective radiating area AT (the antenna size in wavelengths) of the toansmitting antenna.

GT = 42NT AT

\(\text{vavelength} \)

NT = effeciency factor for transmitting antenna.

(Similar to NR)

Isotropic Antenna

42. AT = 1

The received poner is then

Pr(d)= PT GTARTR

The effective area AR obeys a similar volationship.

If me define the attenna gain at the receiver end

Then

$$Pr(d) = P_T G_T G_R \left(\frac{\lambda}{4\pi d}\right)^2$$

$$Pr(d) = P_T G_T G_R \left(\frac{\lambda}{4\pi d}\right)^2$$
- Friis free-space equation.

Notes:

The Triis equation is good only for the far field (or Frankoefer I region)

Far-field distance $df = \frac{20}{\lambda^2} - in \# of wavelengths$ where D is the longest physical linear dimension of the antenna.

- $\langle c \rangle = \frac{c}{f}$
- (3) Pr(d) can be written as $Pr(d) = Pr(do) \left(\frac{do}{d}\right)^{L}$

In dB (decibel), 10 log10 X the power decreases by 20 dB as the distance is increased by 10 dB.

2-ray model - 10min

Saturday, January 12, 2008 11:21 AM

In the free-space model, only I path between sender and receiver. This is often not the case for terrestrial communication, where the ground becomes a natural reflector.

In the 2-ray model, the transmitted

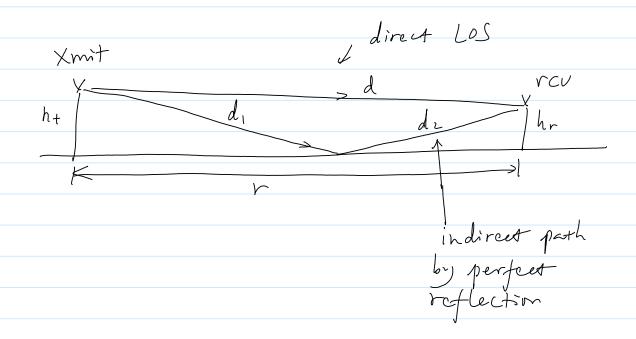
BM wave reaches the receiver through

2 paths

- the direct line of sight pash

- indirectly by perfect reflection

from a flat ground Surface



We will show that in the 2-ray model $Pr(d) \propto \frac{1}{d4}$

The transmitted signal is Eye just Direct-pash received signal

ET e j'wc (+- d)

Indirect-pash received signal

To juc (t-di+di)

And e juc (t-di+di)

perfect reflection leads to

phase shift of 7.

Total received signel $\frac{E_7}{d} e^{j\omega_c(t-\frac{d}{c})} \left[1 - \frac{d}{d_1 + d_2} e^{-j\omega_c(\frac{d_1 + d_2 - d}{c})} \right]$

Total received power $\frac{\overline{E_7}}{d^2} \left| 1 - \frac{d}{d_1 + d_2} e^{-j\omega_c} \left(\frac{d_1 + d_2 - d}{c} \right) \right|^2$

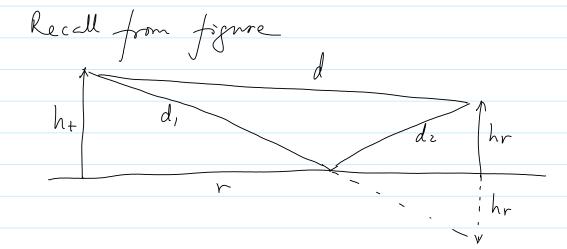
equivalent to free-space due to two rays PTGTGR (Aad)

The cancelation effect of 2-ray - 15min

Saturday, January 12, 2008 11

We will show that when
$$h_{+}$$
, $h_{r} \ll d$,
$$\left| \left(\frac{d}{d_{1}+d_{2}} \right) e^{-j\omega_{c}} \left(\frac{d_{1}+d_{2}-d}{c} \right) \right|^{2}$$

$$\left| \frac{d}{d_{1}+d_{2}} \right| \approx \left(\frac{4z}{\lambda} \frac{h_{+}h_{r}}{\lambda} \right)$$



$$d = \int r^2 + (h_t - h_r)^2$$

$$d_1 + d_2 = \int r^2 + (h_t + h_r)^2$$

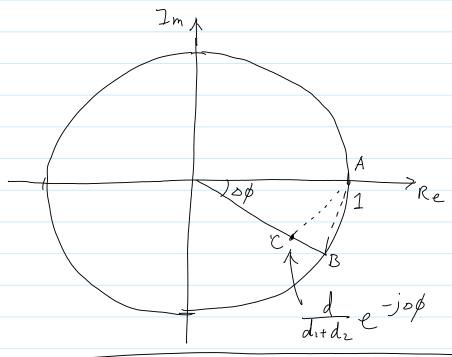
$$\Rightarrow d_1 + d_2 = \int d^2 + 4h + hr$$

$$= d \cdot \int 1 + \frac{4h + hr}{d^2}$$

$$2f \quad A \quad 4h + hr << d^2,$$

$$+hen$$

$$d_1 + d_2 \approx d + \frac{2h + hr}{d}$$



We know $|BC| = 1 - \frac{d}{d + dx}$ $= 1 - \frac{d}{d} + \frac{2h+hr}{d} + \frac{2h+hr}{d}$ $|AB| = 2s - \frac{3b}{2}$ $|AB| = \frac{4z}{d} + \frac{h+hr}{d} < < 1$ then

1AB = 0\$ = 42h+hr When $d \gg \lambda$, $|AB| \gg |BC|$ $\left| 1 - \frac{d}{d_1 + d_2} e^{-j\delta\phi} \right| = |Ac|$ $P_{R} = P_{T} G_{T} G_{R} \left(\frac{\lambda}{4zd}\right)^{2} \left(\frac{4zh_{t}h_{r}}{\lambda d}\right)^{2}$ = PT GT GR (h+hr)2

Key print: It is the phase difference
that determines the combined
effect of 2 rays

- We will see this again in
multi path fading

(35

The cancelation effect of 2-ray - supp

Saturday, January 12, 2008

11:31 AM

We will show that when h_{+} , $h_{r} << d$, $\left| \left(\frac{d}{d_{1} + d_{2}} \right) e^{-j\nu_{c}} \left(\frac{d_{1} + d_{2} - d}{c} \right) \right|^{2}$ $\left| \frac{d}{d_{1} + d_{2}} \right| = \frac{4z h_{+} h_{r}}{\lambda d}$

Recall from figure

h+

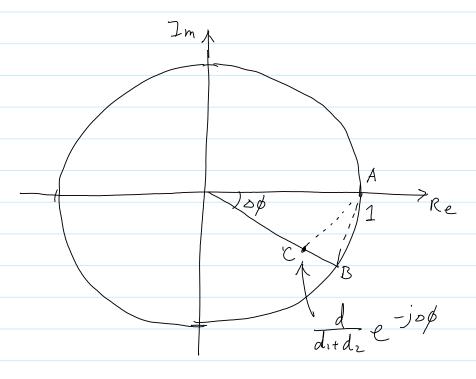
d,

hr

ihr

$$\begin{cases}
d = \int r^2 + (h_t - h_r)^2 \\
d_1 + d_2 = \int r^2 + (h_t + h_r)^2
\end{cases}$$

Next, let $\omega = \omega = \frac{d_1 + d_2 - d_1}{c}$



When $d \gg \lambda$, $|AB| \gg |BC|$ Hence, $\begin{vmatrix} 1 - \frac{d}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC| \\ \frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC| \\ \frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi} | = |AC|$ $\frac{1}{d_1 + d_2} e^{-j \diamond \varphi}$

Key print: It is the phase difference
that determines the combined
effect of 2 rays

- We will see this again in
andsi path fading

Validity of Assumptions - 5min

Sunday, January 13, 2008 11:57 AM

Assumption (A)

4 hthr << d2

Example: $h_t = tom$, $h_r = 2m$ $\Rightarrow d >> 20m$

Assumption (B)

42 hthr <<1

Example: f = 800 MHz, $\lambda = \frac{3}{8} \text{ m}$ $\Rightarrow \lambda > 3.3 \text{ km}$

f=30GHz, \ \ = 0.0/m

d)) /26 km

Assuption B is wondly more strongent than assumption A.

More problematic for high prepuencies. Instead of seeing smooth I yan behavion, more likely to see small-scale multi-path fading.

The actual power received, measured over relatively long distances of many wavelengths, is found to vary randomly about the area - mean - power P_G-GR & (d).

Shadow tadiy is in general caused by variations in signal power due to signal attenuated by terrain obstructions, such as hills, brildigs, or even leaves

The measured signal power thus may differ substantially at different location, even though they may be at the same distance from the xmitter

A good approximation for shadow fading is to assume that power measured in dB follows a Eansian distribution centered about its average value.

 $10^{8/10}$ — ly-normal 6-10 dB $\times N(0, \sigma^2)$

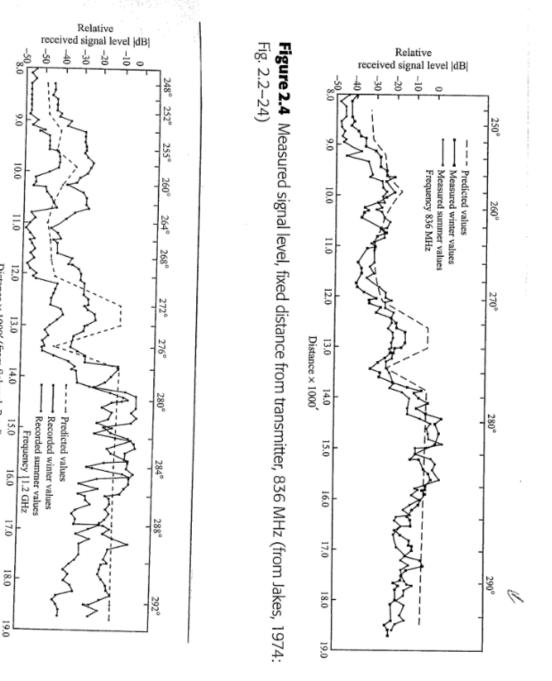
where

$$(3) = \frac{1}{\sqrt{220}} = \frac{x^2}{20^2}$$

Rational: A form of central limit theorem There are often many obstacles between the xmitter and receiver. The existence of an obstacle (e.g. a building) with width d, attenuates the signal by $- \propto d_1$ where dis some constant, d, =0 if there does not exist such an obstacle With n potental obstacles, the attenuation is Important: Shadow fading occurs over distances of tens/hundreds of wavelengths (>>) -> large-scale fading see Tig 2.4. Tig 2.5 in Schwartz

(55)

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lec4-mwf-new Page 18

Fig. 2.2-25)

Figure 2.5 Measured signal level, fixed distance from transmitter, 11.2 GHz (from Jakes, 1974:

Distance × 1000' (from Schanck Road)

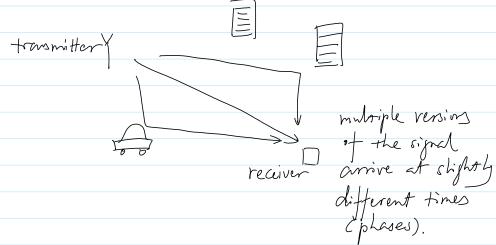
Multi-path fading - 10min

Sunday, January 13, 2008 2:25 PM

Multi-path fading (small-scale fading) is used to describe the rapid function of the amphitude of a radio vignal over a short period of time or over a short travel distance (on the order of wavelengths).

- Variations of path-loss or large-scale tading should be negligible over such what distances

Multipath fading is caused by interference between 2 or more versions of the transmitted signal, which arrive at the receiver at shightly different times.



These multiple paths combine at the receiver antenna to give a resultant signal that varies wildly in amplitude & phase

- Similar to what happens in the

Multipesh fedig is modelled by the 2them in PR = 0210 X/108(d) PTGTGR

When # of multipath signals is large (>16), and none of them dominate, & follows Raleigh distribution with pdf

 $f_{\alpha}(\alpha) = \frac{\alpha}{\sigma_{\alpha}^{2}} e^{-\left(\frac{\alpha^{2}}{2\sigma_{\alpha}^{2}}\right)}$

When only one component dominates (e.g. when there is a strong direct path), & follows Ricean distorbution

 $f_{\mathcal{A}}(a) = \frac{a}{\sigma_{\mathcal{R}}^{2}} e^{-\frac{a^{2} + A^{2}}{2\sigma_{\mathcal{R}}^{2}}} I_{o}(\frac{aA}{\sigma_{\mathcal{R}}^{2}})$

Where A: amplitude of the dominant signal

To (oc): Bessel function of the first kind 2 0-order

 $J_{o}(x) = \frac{1}{22} \int_{0}^{22} x \cos \theta d\theta$

Note: 2f A → 0, Ricean → Raleigh.

- More recent models assume 2 dominant protes. (see suggested readings), which may match better nichtle pet of toot-tading compounts

