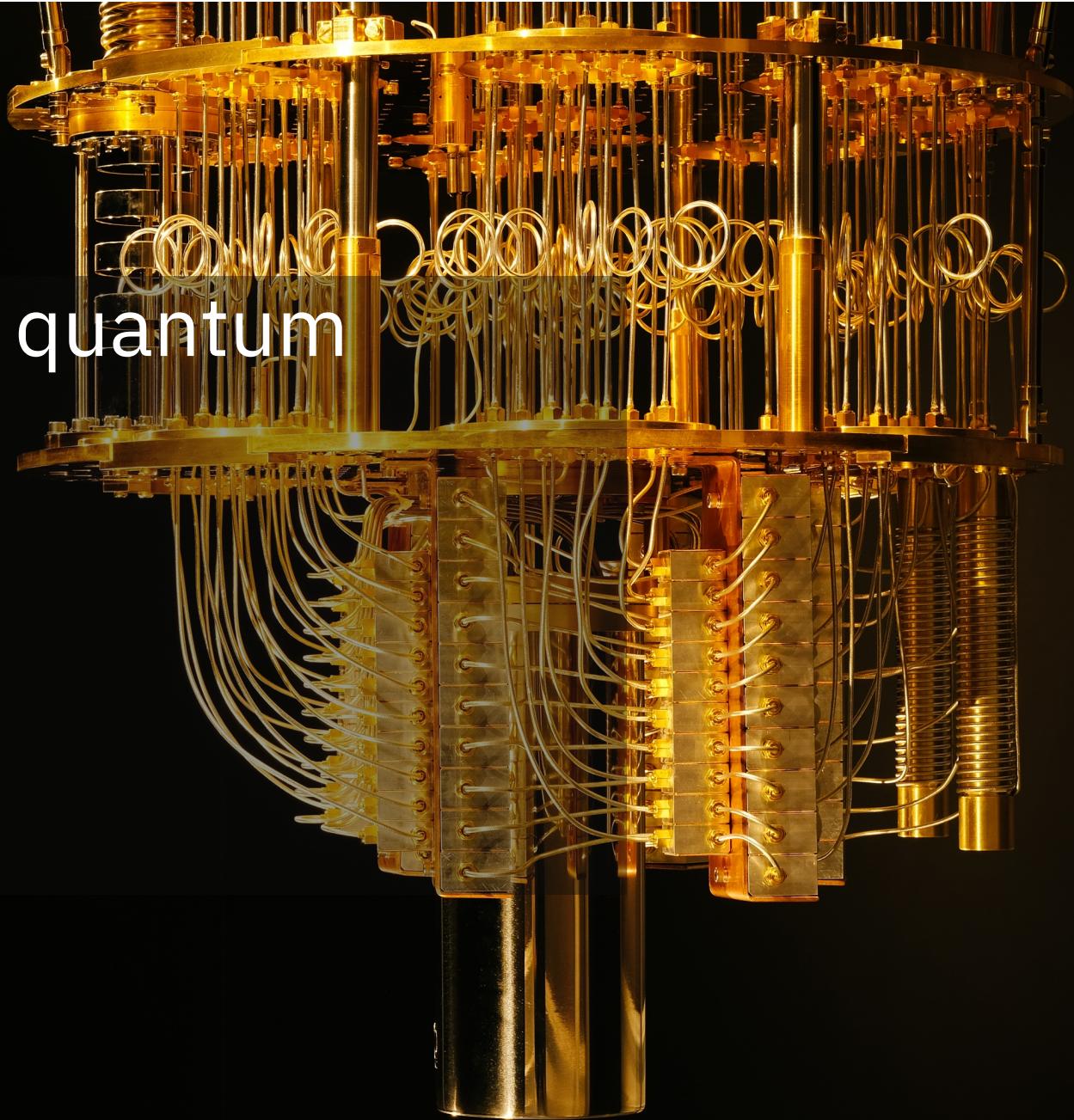


Simulating chemistry on a quantum computer II

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IBM Quantum

QISKIT Global Summer School
July 29, 2020



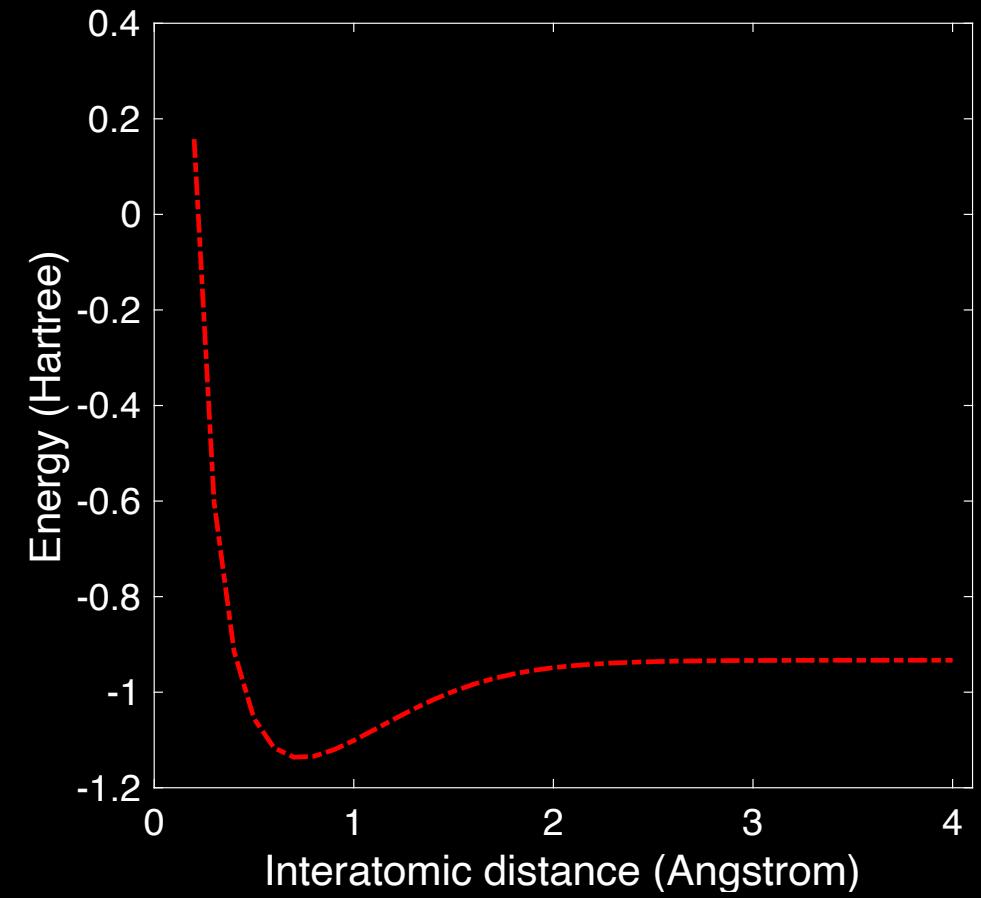
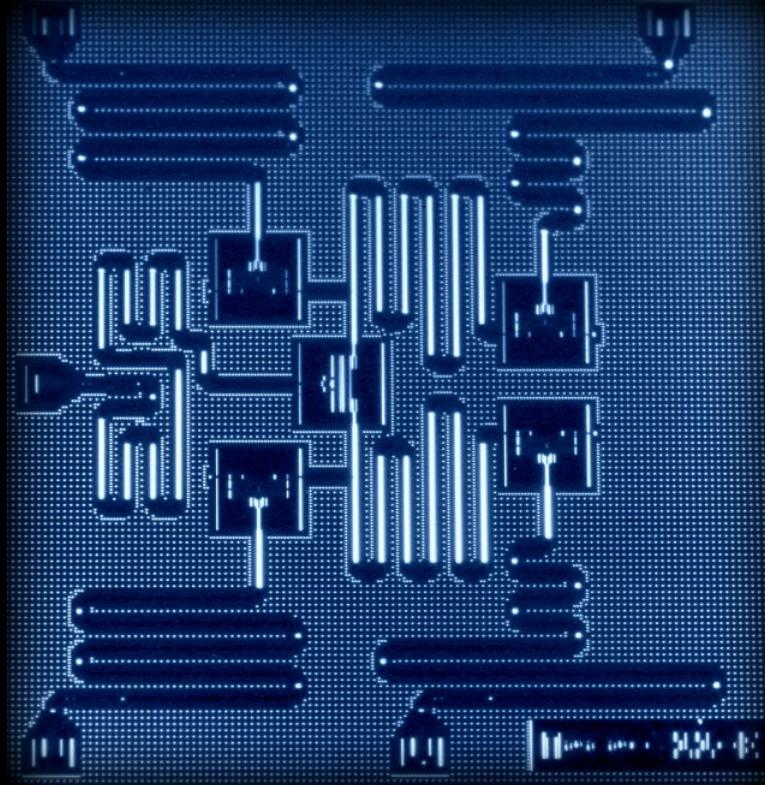
IBM Quantum

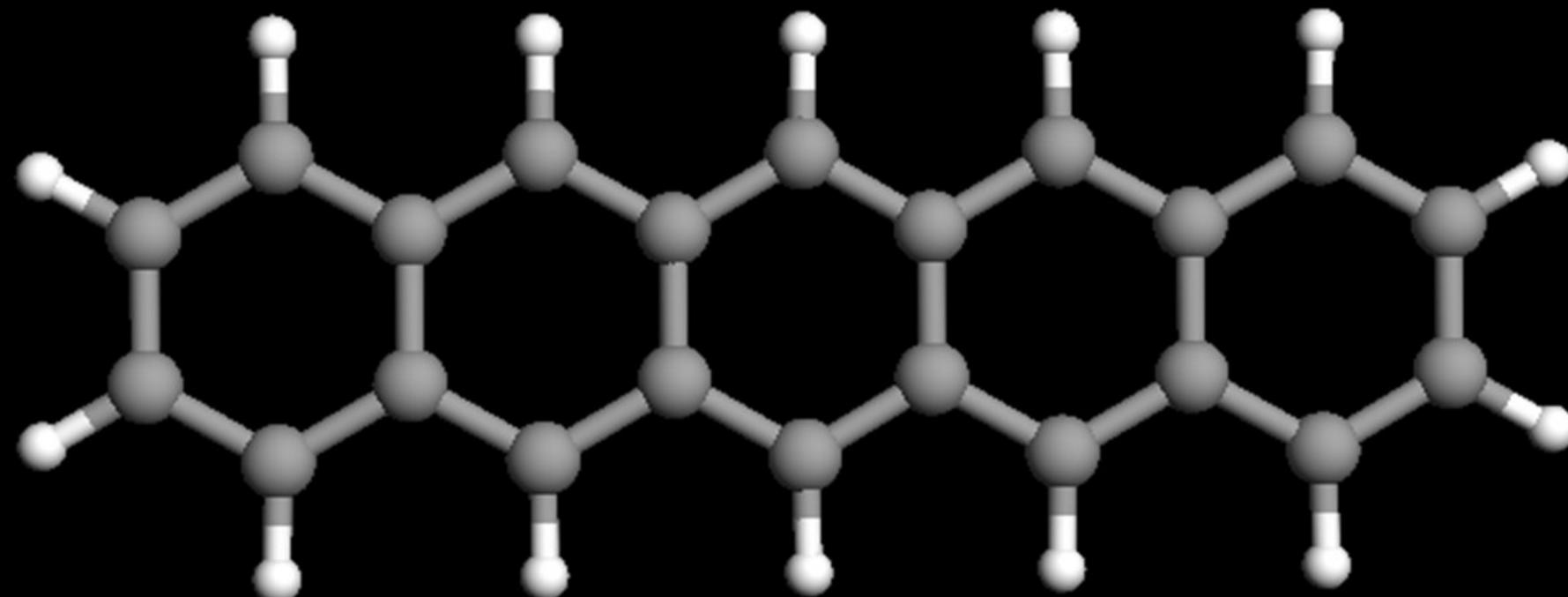


IBM Research Frontiers
Institute

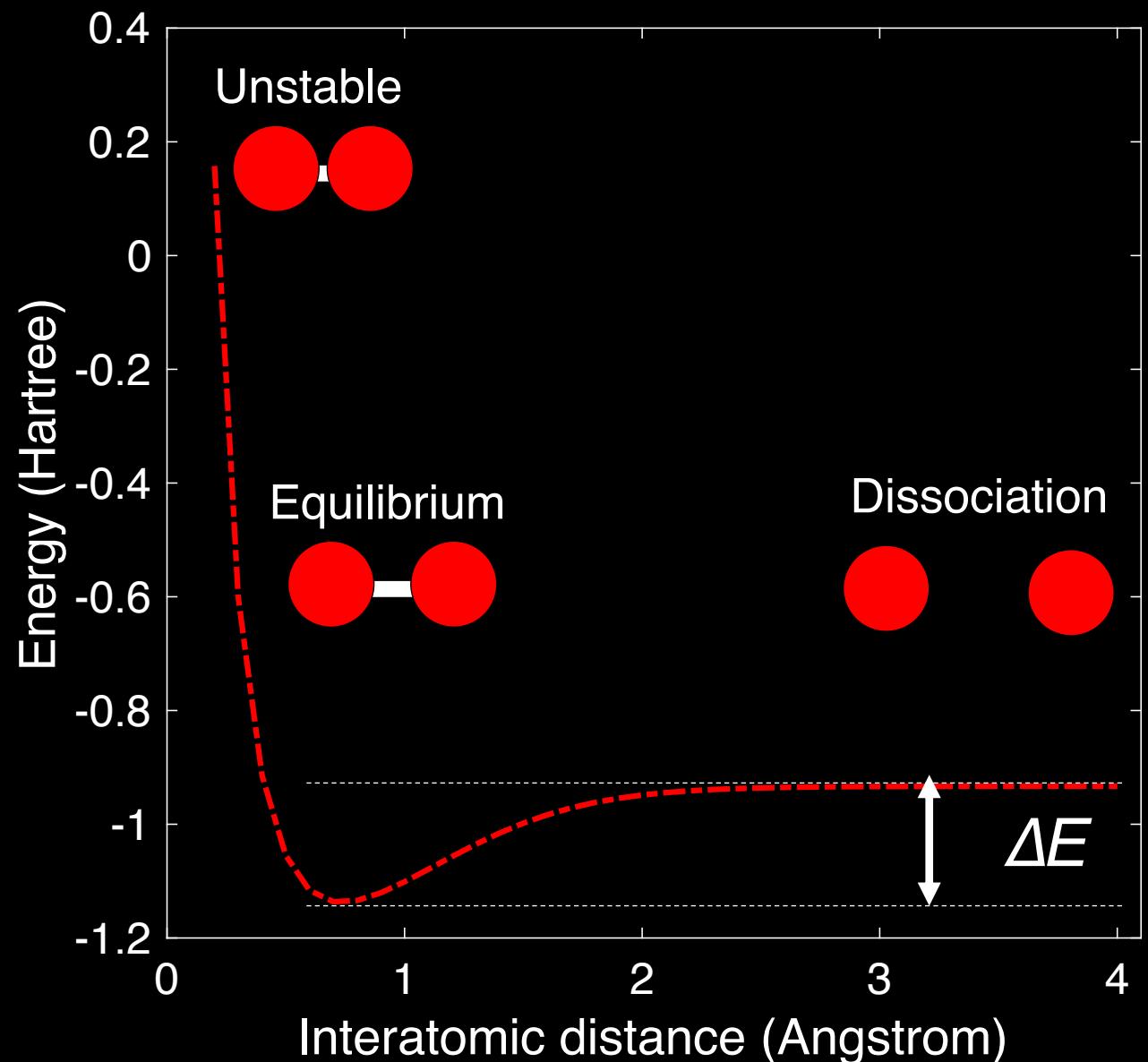
IARPA

Motivations: Quantum hardware meets quantum algorithms





Potential energy curve of a molecule



Rate $\propto \text{Exp}(-\Delta E/kT)$

The Electronic Structure Problem

Interacting fermionic problems: A core challenge in modern computational physics and HPC

$$H_e = - \sum_{i=1}^N \frac{1}{2} \nabla_i^2 - \sum_{i=1}^N \sum_{A=1}^M \frac{Z_A}{r_{iA}} + \sum_{j>i} \frac{1}{r_{ij}}$$

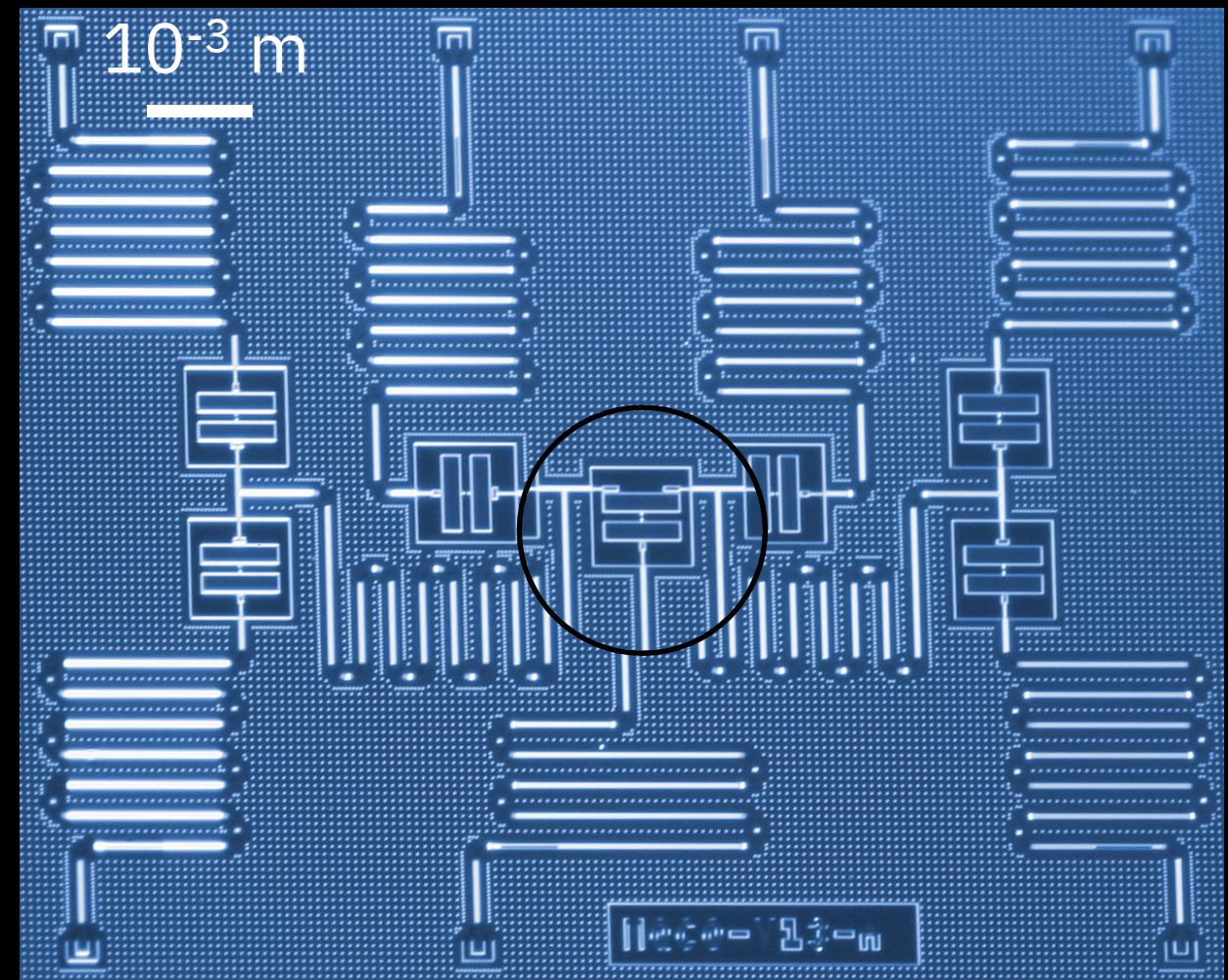
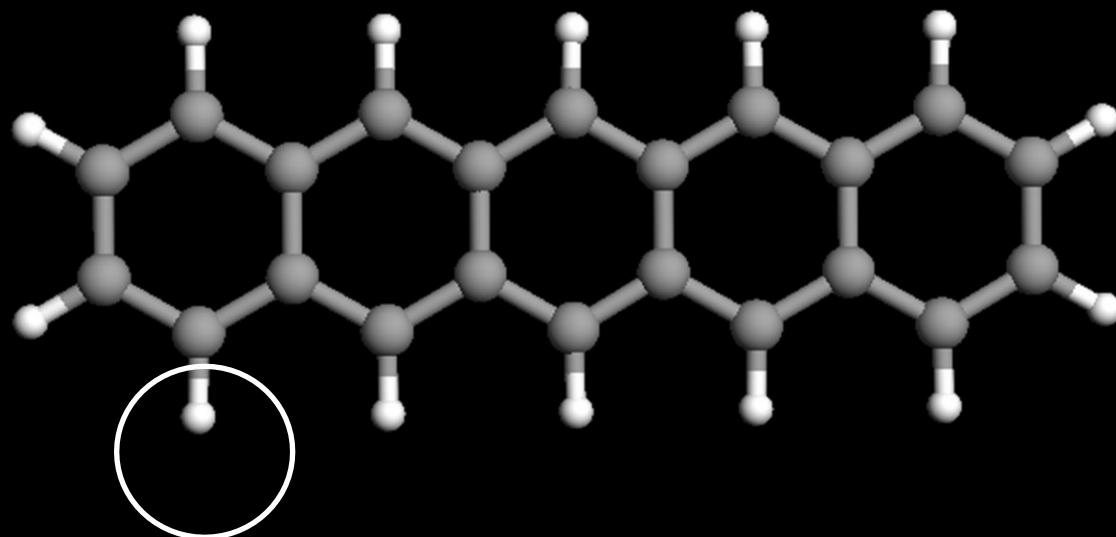
$$H|\psi_G\rangle = E_G|\psi_G\rangle$$

**Can we build a programmable,
well-controlled quantum system to
simulate the properties of other
natural quantum systems, like
molecules?**

Natural atoms

Artificial Atoms

10^{-9} m



Classical

0



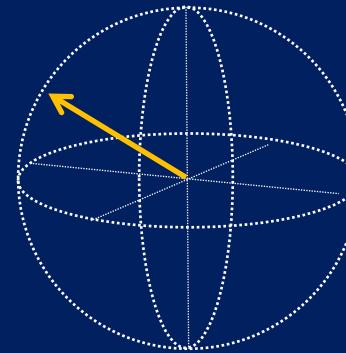
1

10110...1001

N bit number: 2^N possible states
Always in 1 of those states

Quantum

0



1

Quantum Bits (“Qubits”) can be in a superposition of 0 and 1

10110...1001
11010...1010
~~00100...1101~~
~~10110...1111~~

Can be in a superposition of all 2^N states

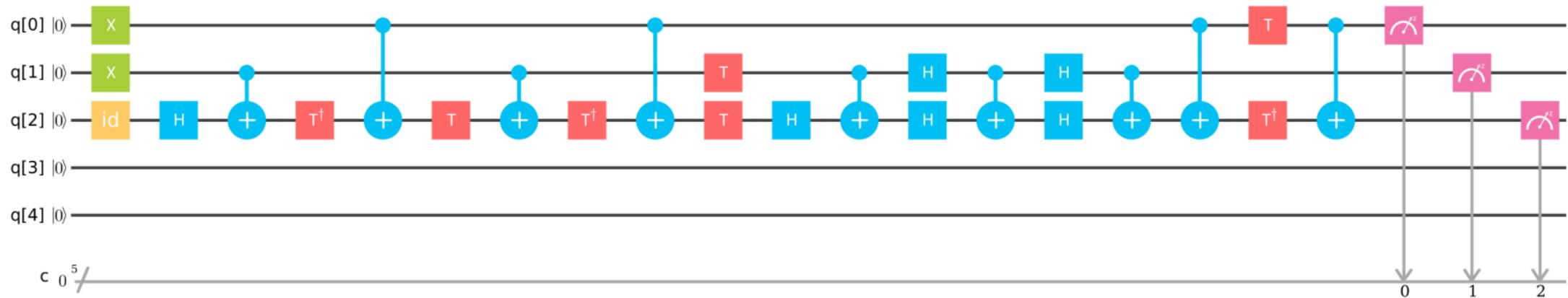
Quantum Circuits

Decompose unitary

$$U \in \mathbf{U}(2^N)$$

$$|\psi\rangle = U|0\dots 0\rangle$$

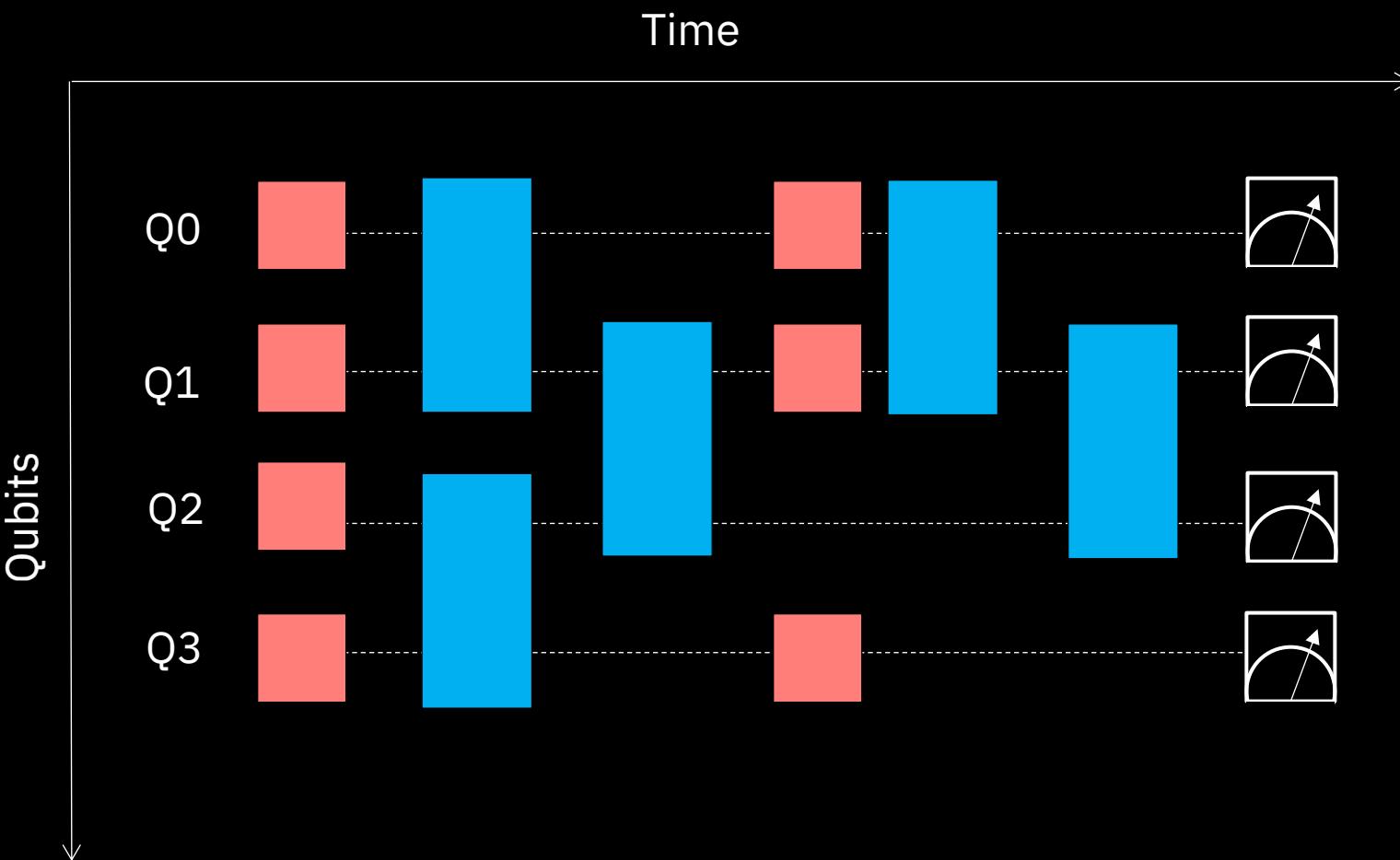
$$U = \text{CNOT}_{0,2} T_0 T_2^\dagger \dots X_0 X_1$$



Universal gate set: CNOT (two-qubit gate) and single qubit rotations

Short depth circuits

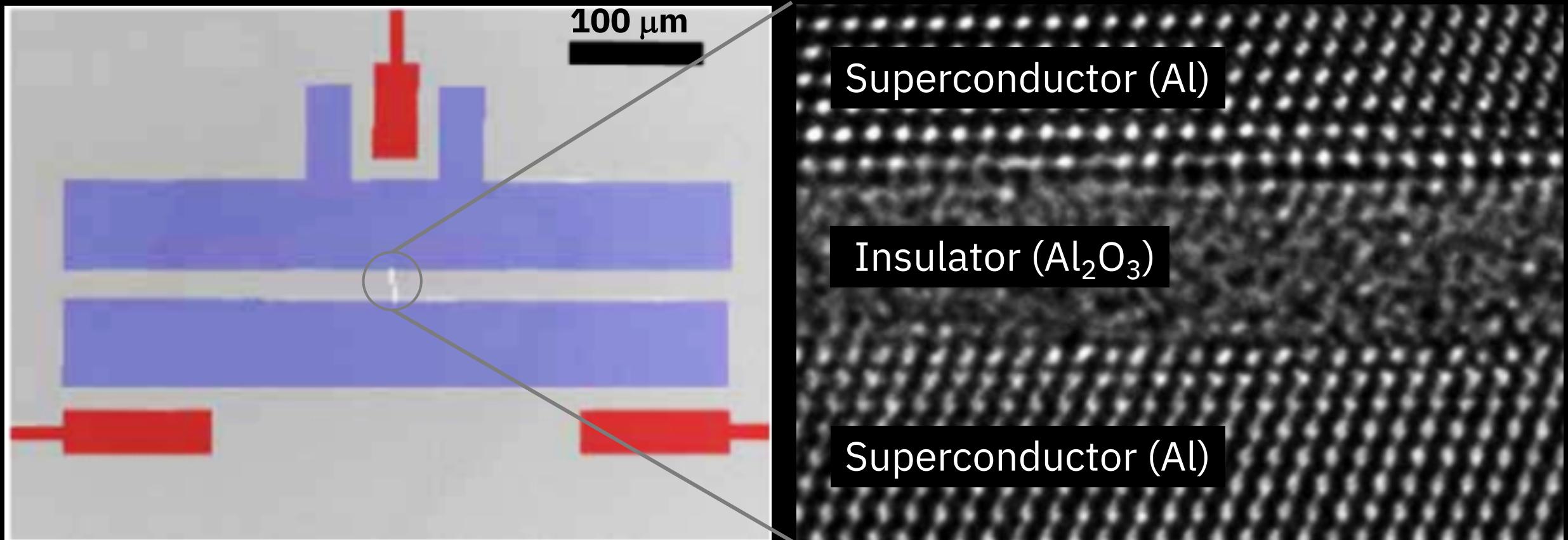
Approximate quantum computing with noisy quantum hardware



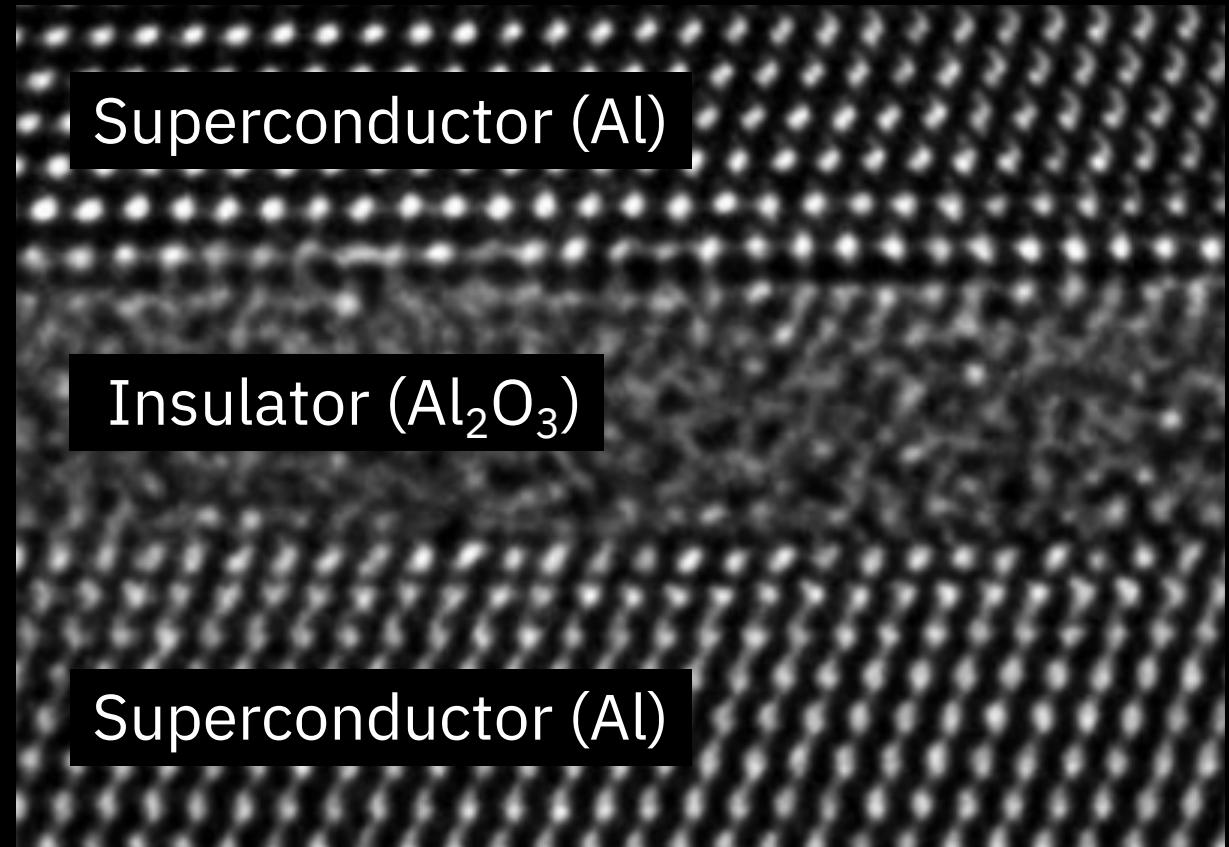
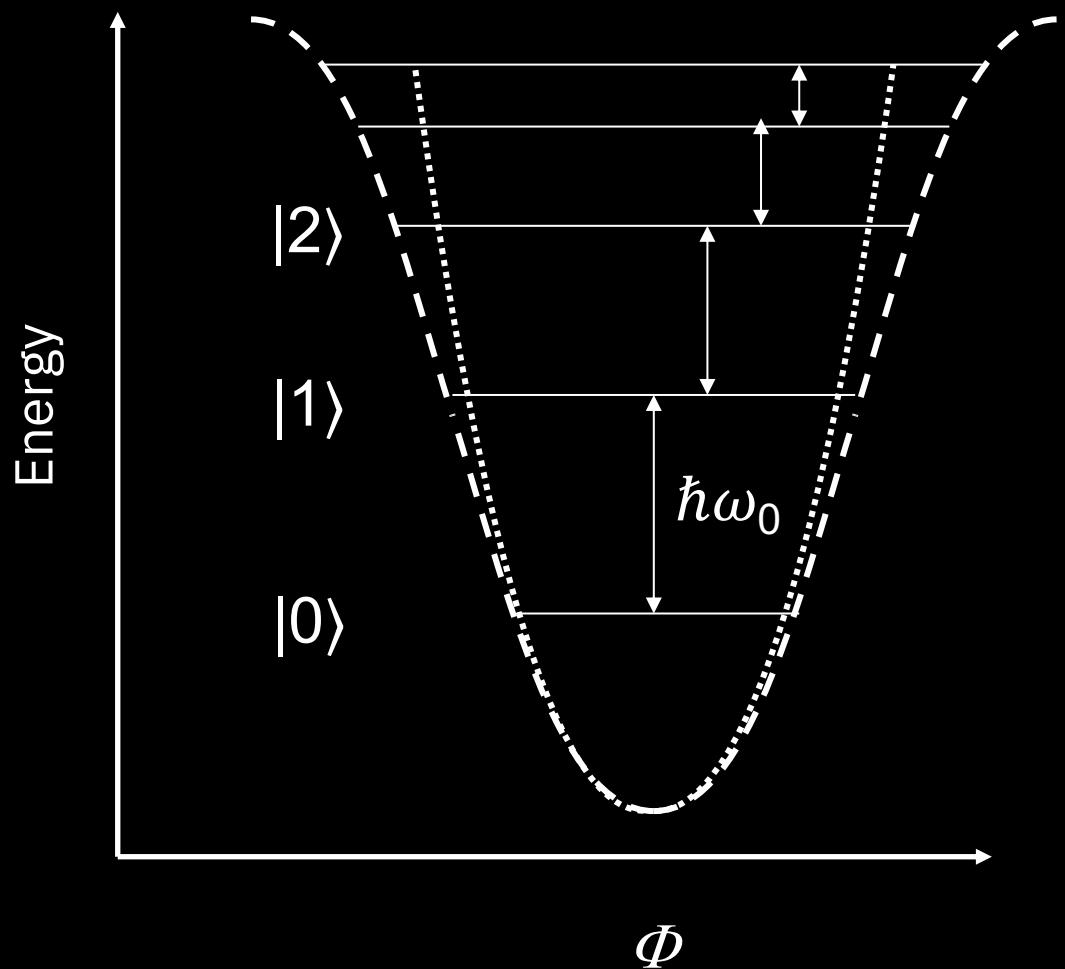
- Near-term quantum hardware is noisy
- Fault tolerant architectures not immediately accessible
- Novel algorithms focused on using short depth circuits that make best use of coherence window

Quantum Hardware recap ..

Artificial atom: Transmon



Artificial atom: Transmon



Single qubit control: Rabi Oscillations

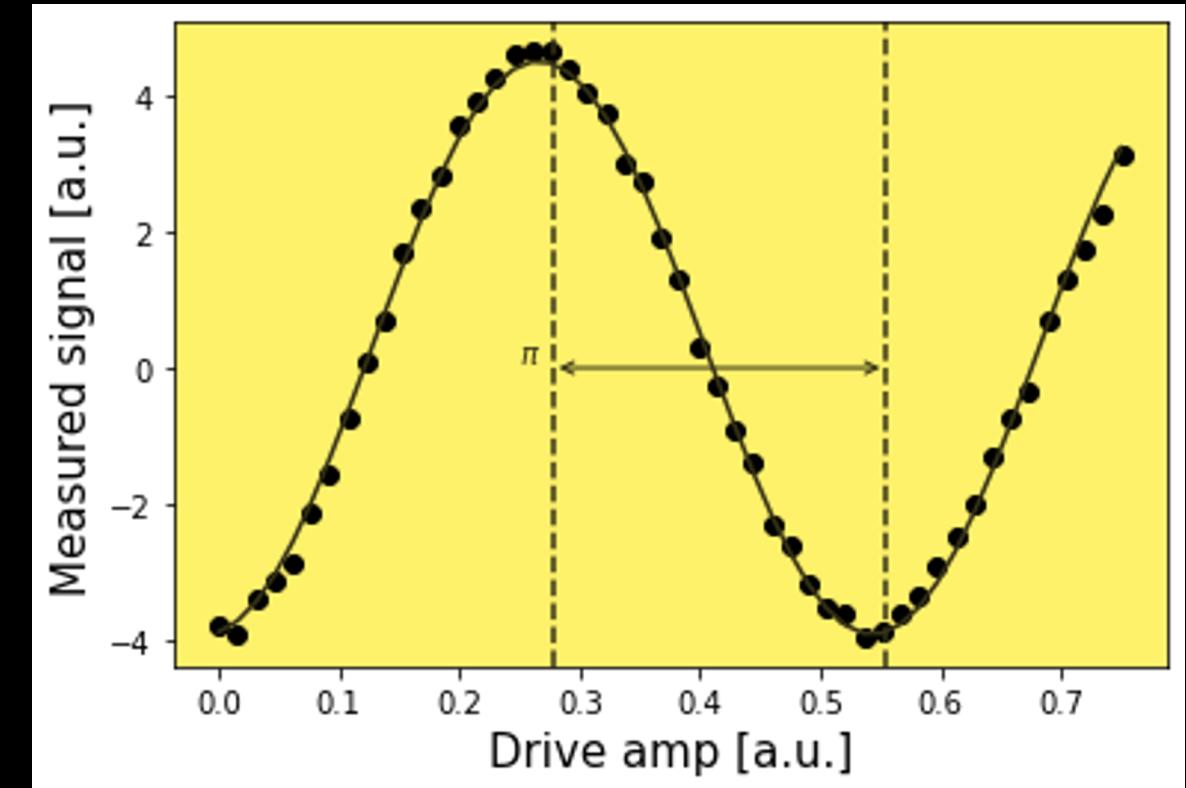
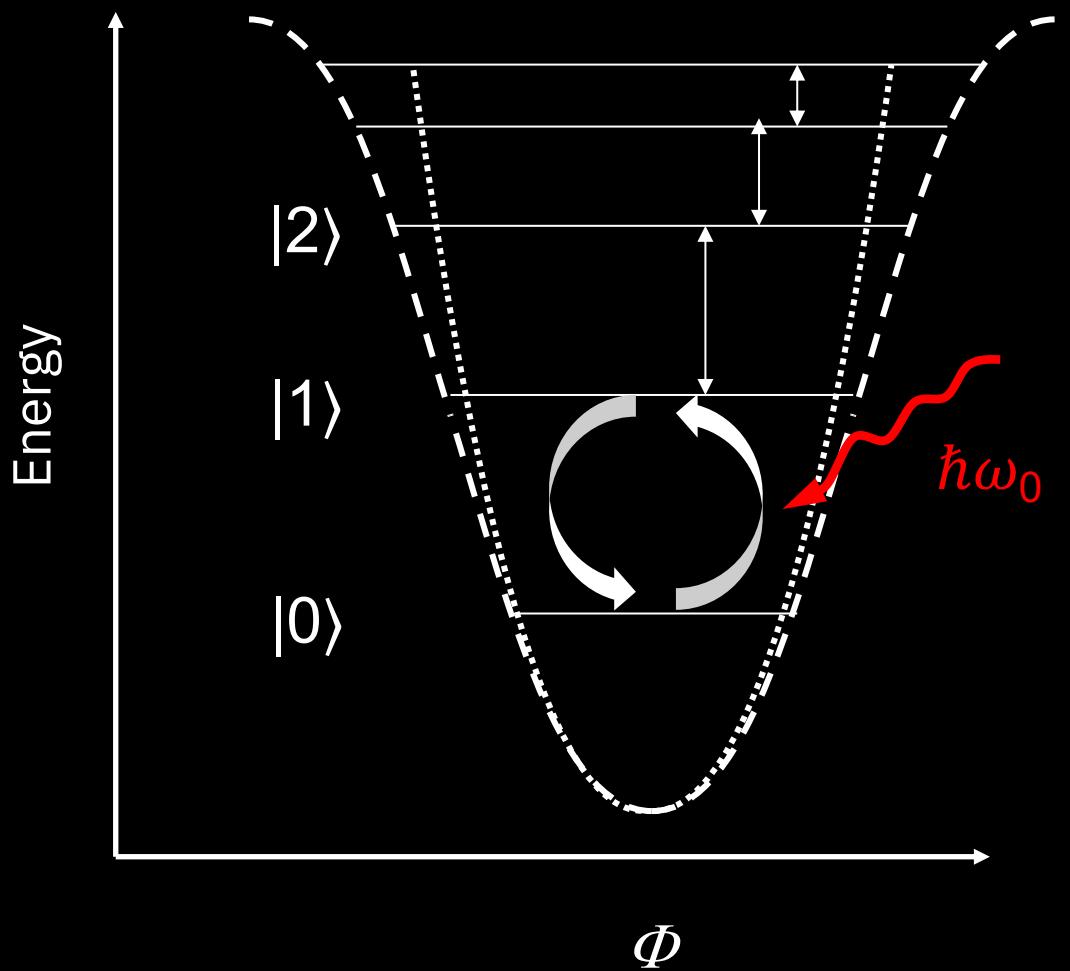
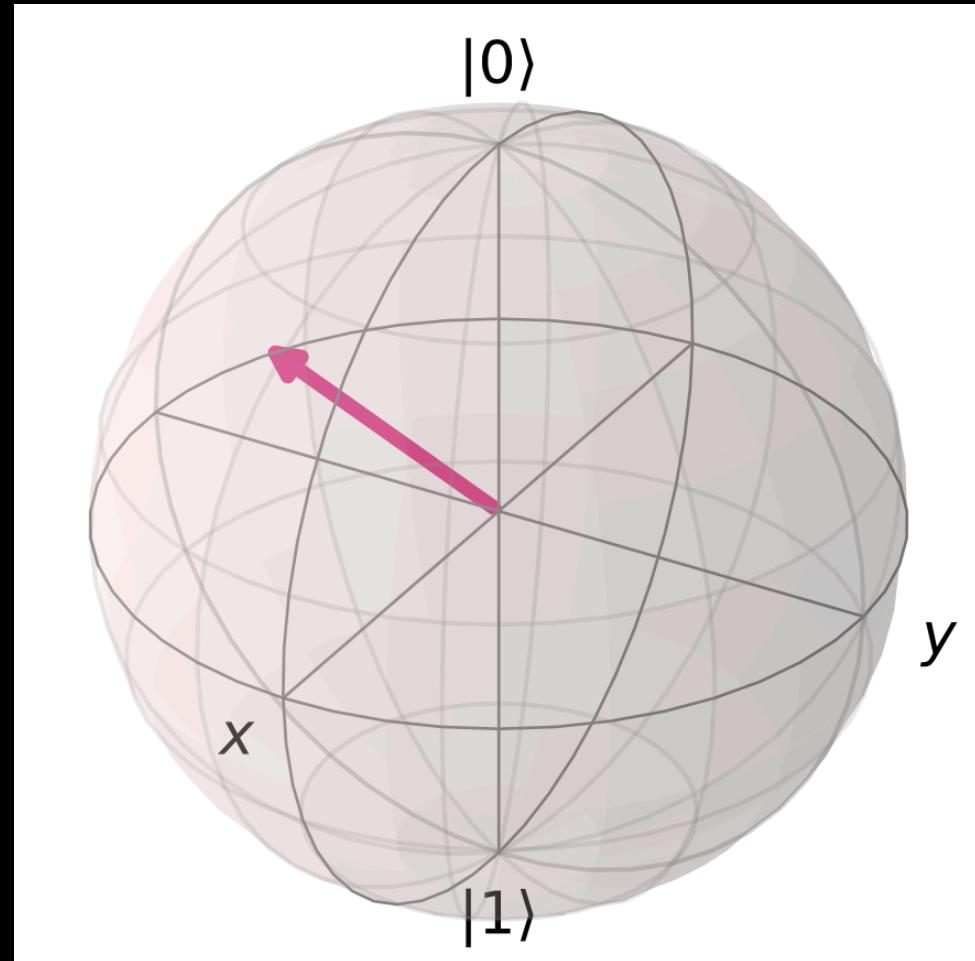
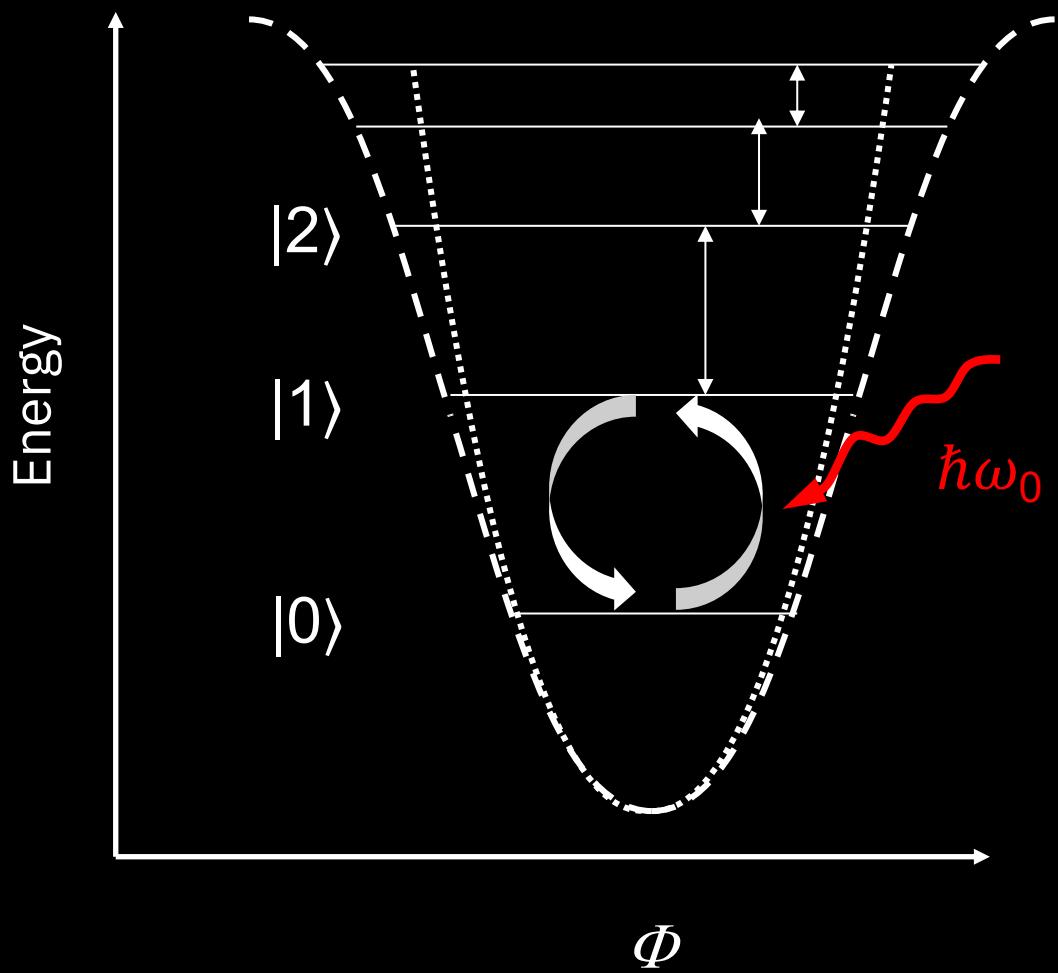


Image: QISKit textbook

Single qubit control: Rabi Oscillations

Image: QISKit textbook



- Amplitude/duration of pulse controls angle of rotation.
- Phase of pulse controls axis of rotation

Two-qubit entangling gates

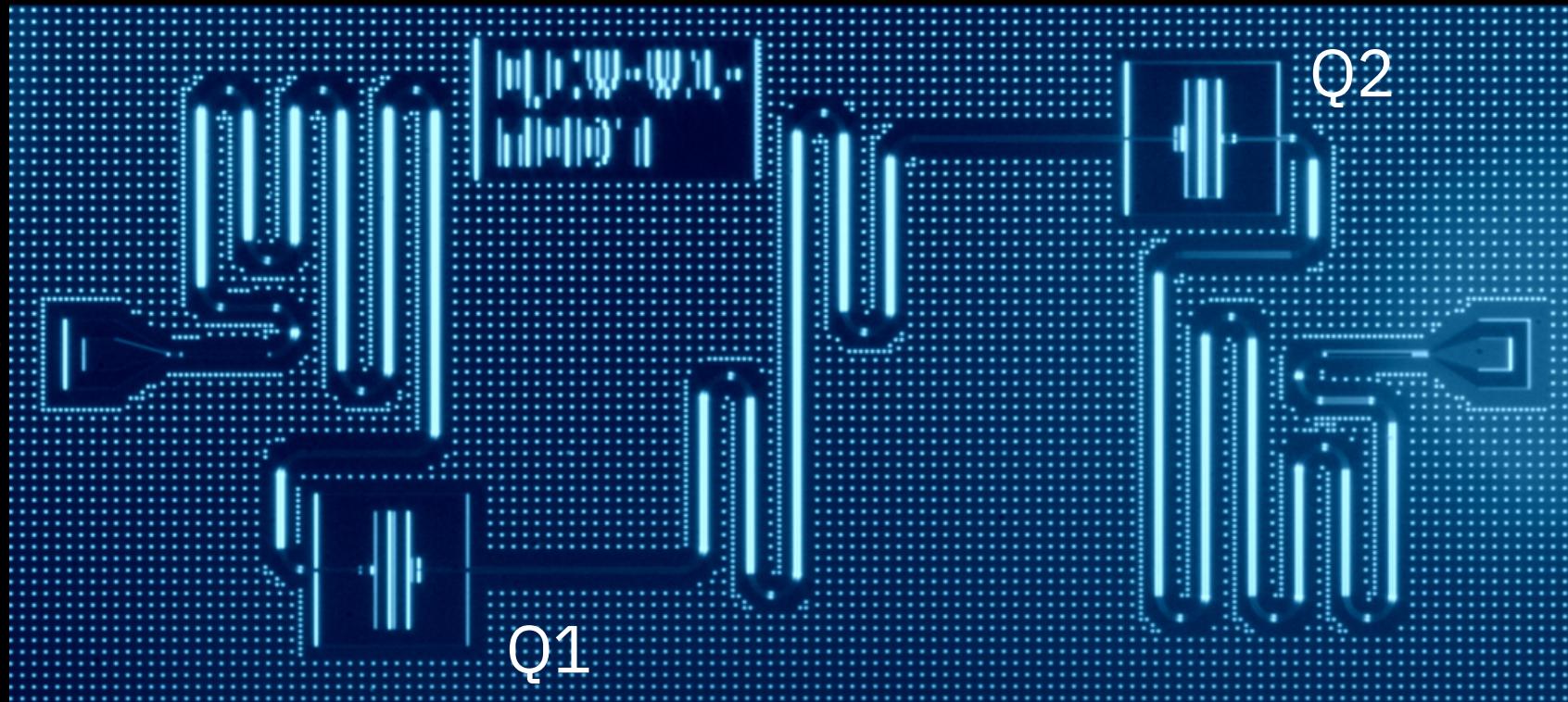
Apply operations on two qubit product states to create an entangled state, i.e. one that cannot be factored into its individual qubit components.

Entanglement generated by *conditional* rotations i.e rotations of a qubit (target) dependent on the state of the other (control).

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)_A |0\rangle_B$$

$$U_{\text{CNOT}} |\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle_A |0\rangle_B + |1\rangle_A |1\rangle_B) \neq (\dots)_A (\dots)_B$$

Two-qubit entangling gates: Cross resonance

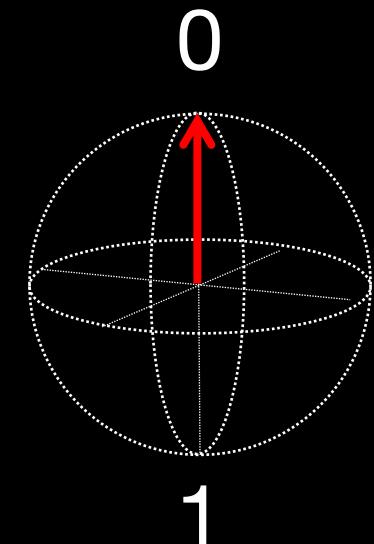
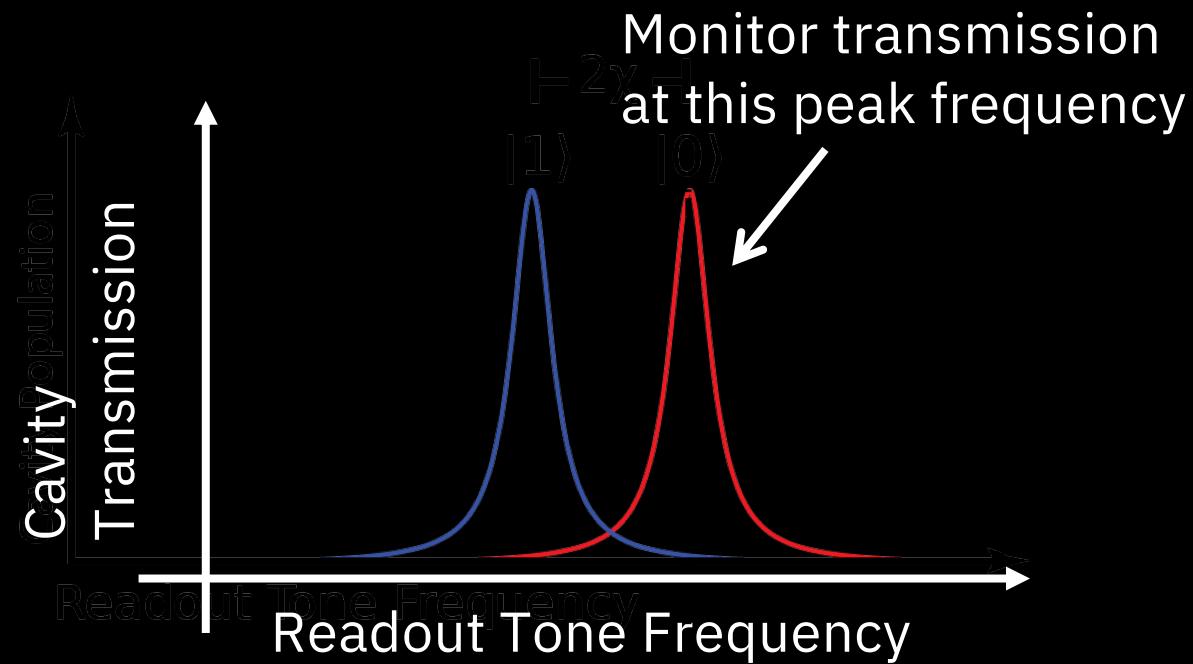


Can generate a CNOT with only microwave pulses

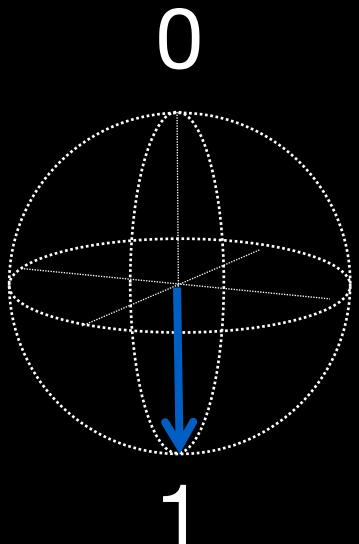
Dispersive readout

$$H = \frac{\hbar\omega'_{01}}{2}\hat{\sigma}_z + \hbar(\omega'_r + \chi\hat{\sigma}_z)\hat{a}^\dagger\hat{a}$$

Qubit State dependent resonator frequency

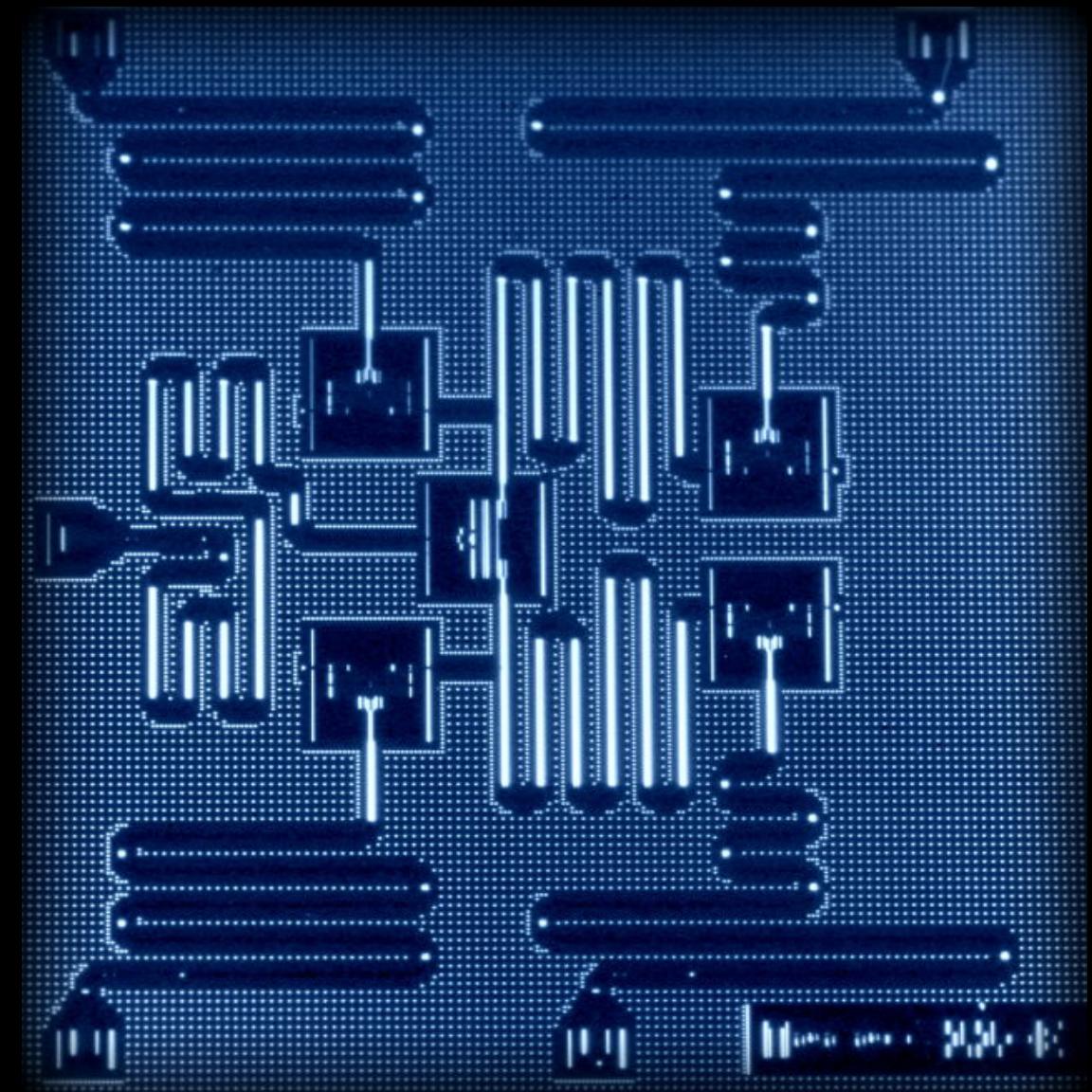
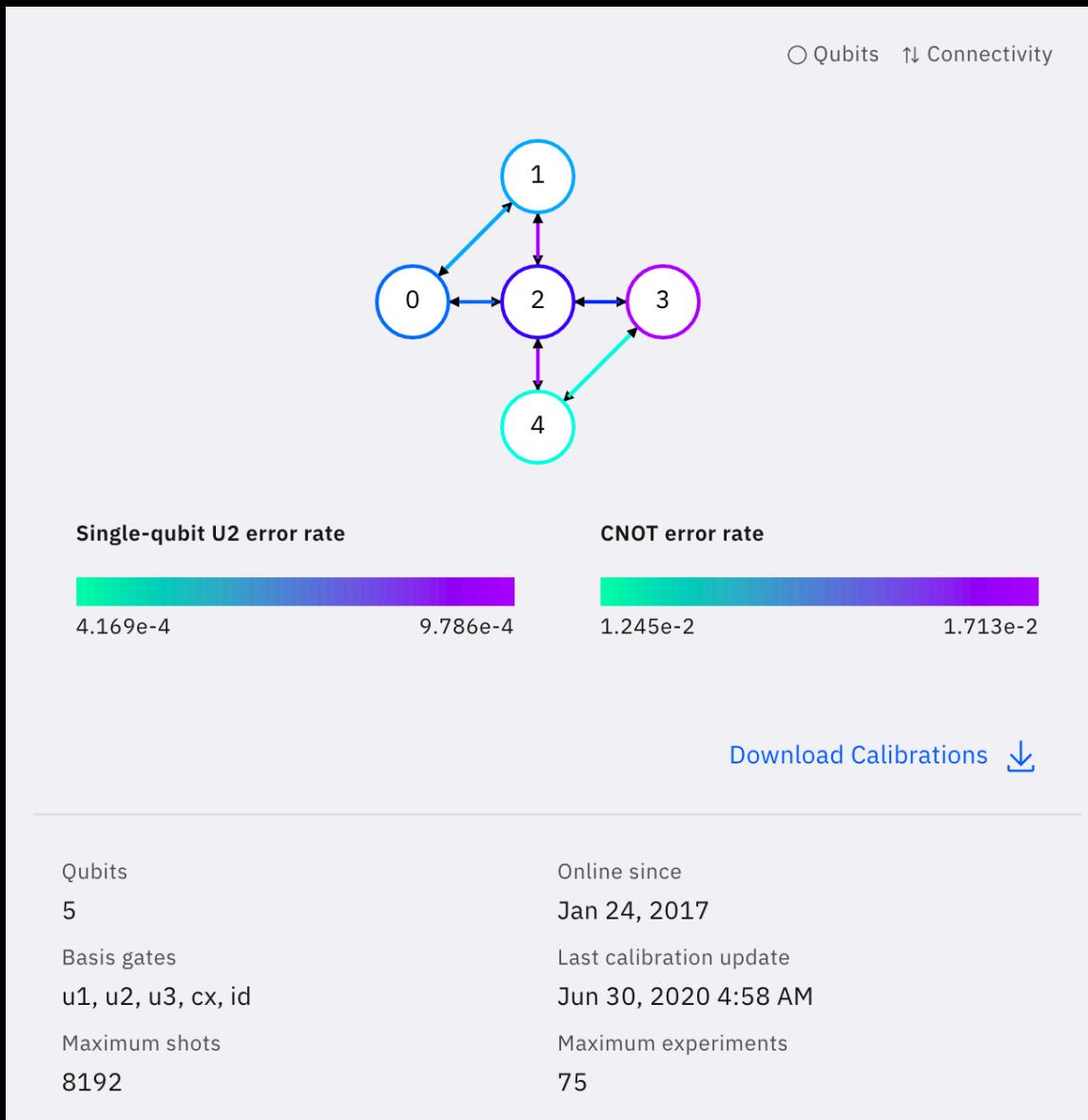


High transmission



Low transmission

Superconducting quantum processor: ibmq_5_yorktown



Device characterization

Qubit	T1 (μs)	T2 (μs)	Frequency (GHz)	Readout error	Single-qubit U2 error rate	CNOT error rate
Q0	64.2	74.54	5.286	2.45E-02	6.49E-04	cx0_1: 1.380e-2, cx0_2: 1.430e-2
Q1	77.35	83.17	5.238	1.75E-02	5.72E-04	cx1_0: 1.380e-2, cx1_2: 1.713e-2
Q2	50.9	41.91	5.031	2.75E-02	8.32E-04	cx2_0: 1.430e-2, cx2_1: 1.713e-2, cx2_3: 1.470e-2, cx2_4: 1.708e-2
Q3	60.04	47.72	5.296	3.35E-02	9.79E-04	cx3_2: 1.470e-2, cx3_4: 1.245e-2
Q4	57.19	64.79	5.084	1.45E-02	4.17E-04	cx4_2: 1.708e-2, cx4_3: 1.245e-2

Q: Which 2 qubits would you choose for a H₂ simulation?

Device characterization

Qubit	T1 (μs)	T2 (μs)	Frequency (GHz)	Readout error	Single-qubit U2 error rate	CNOT error rate
Q0	64.2	74.54	5.286	2.45E-02	6.49E-04	cx0_1: 1.380e-2, cx0_2: 1.430e-2

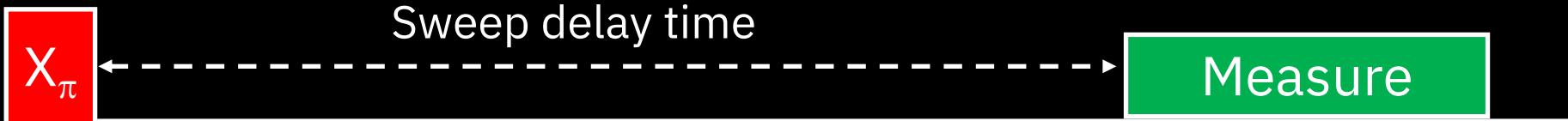
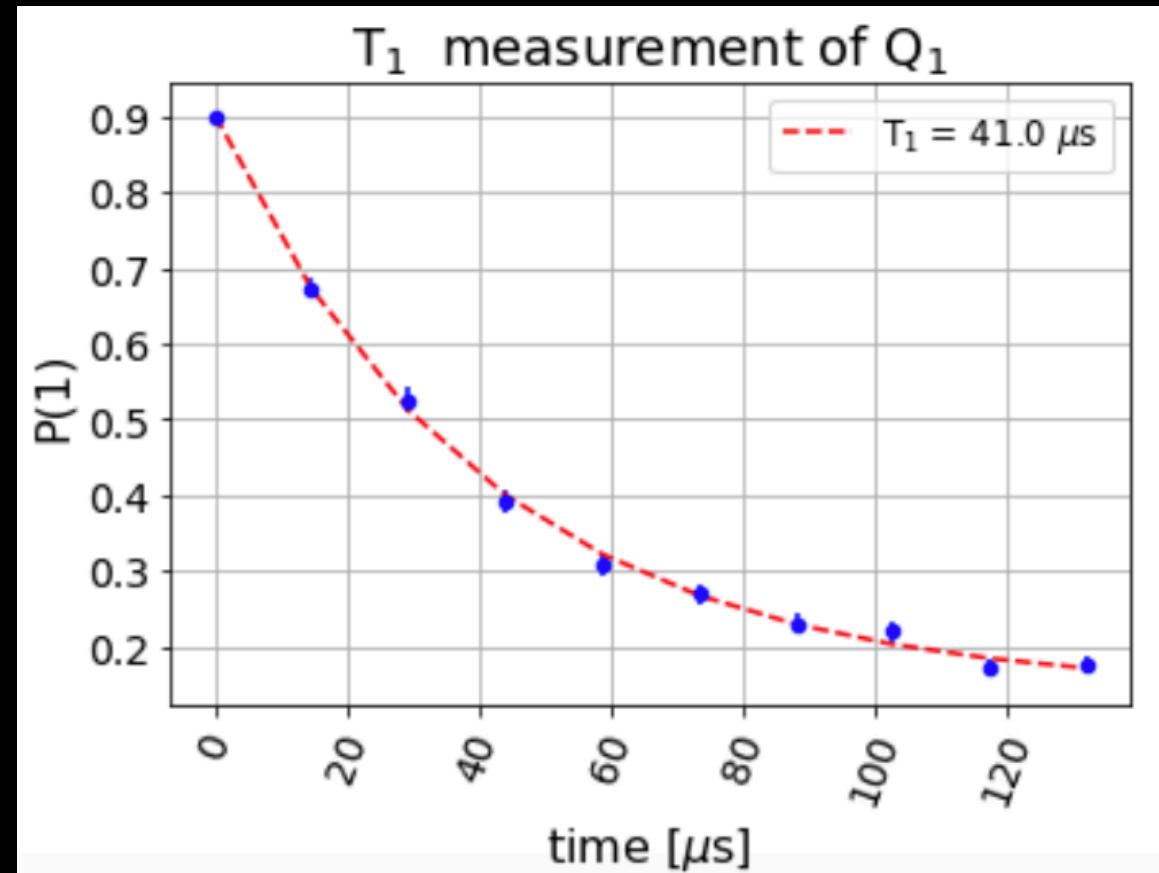
Hardware errors

Errors

- Incoherent
- Leakage
- Coherent
- Measurement

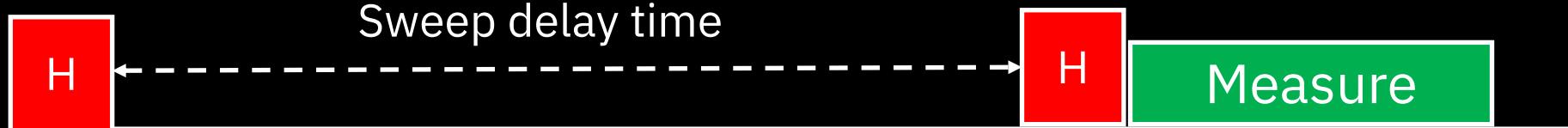
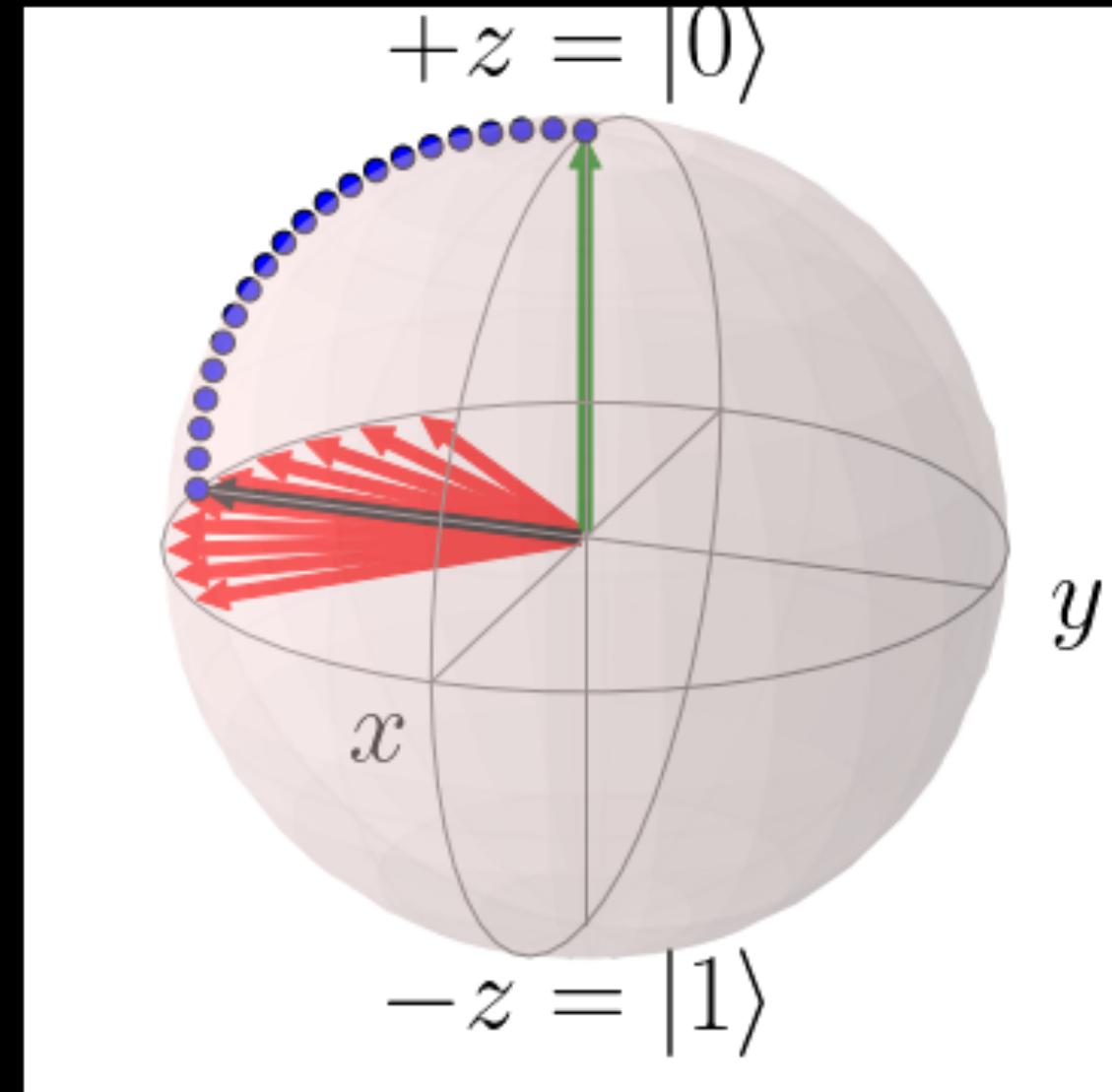
Incoherent errors: Energy relaxation

- Initialize qubit in $|1\rangle$ state and probe its decay to $|0\rangle$
- How quickly does the qubit lose its energy?
- $P(|1\rangle) = \exp\{-t/T_1\}$



Incoherent errors: Dephasing

- Initialize qubit in superposition state
- How quickly does the qubit lose its phase coherence?
- Characteristic time: T_2



Coherence budget

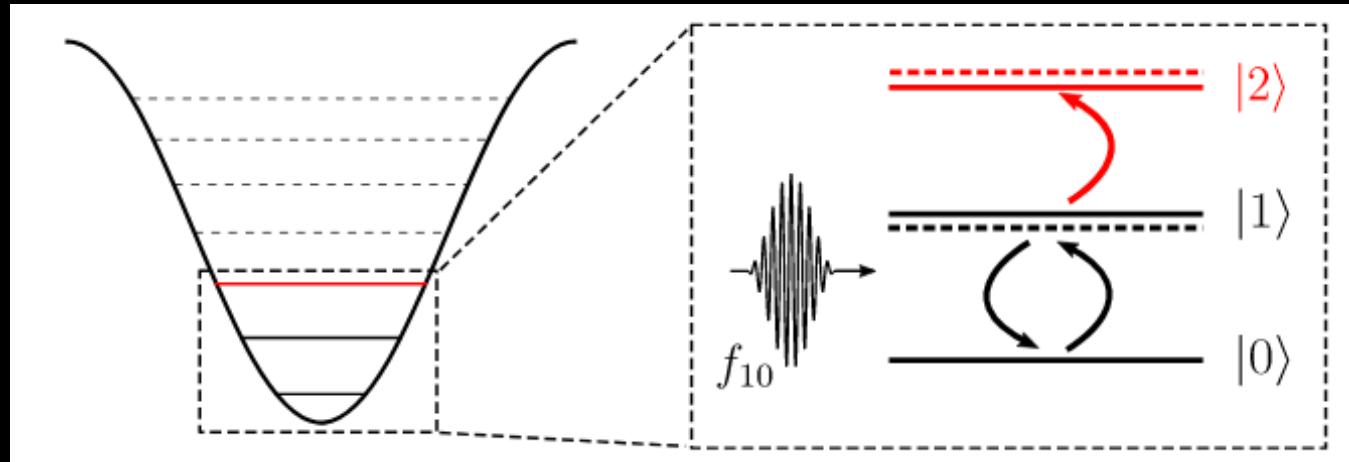
Single qubit gate time $\sim 0.01 \times 10^{-6}$ s

Two qubit gate time $\sim 0.1 \times 10^{-6}$ s

$T_1, T_2 \sim 100 \times 10^{-6}$ s

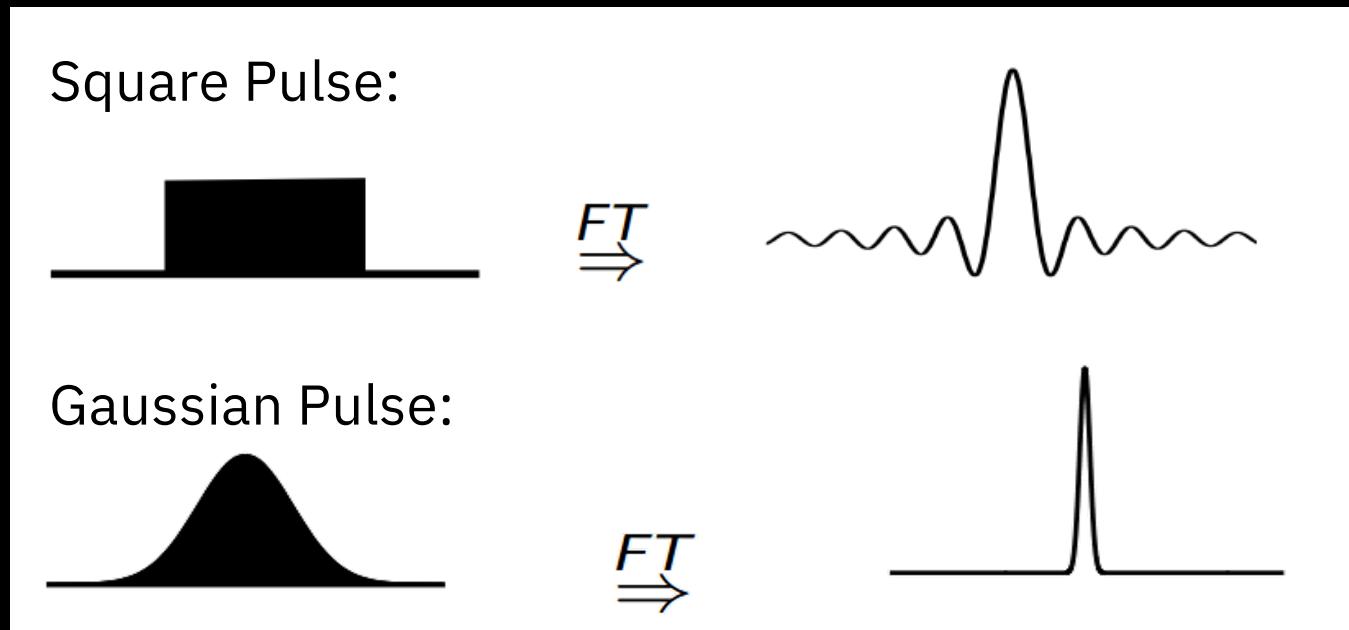
Leakage errors

- Transmon qubits are weakly anharmonic ($\omega_{01} - \omega_{12} \sim 0.330$ GHz). Not true two-level systems.

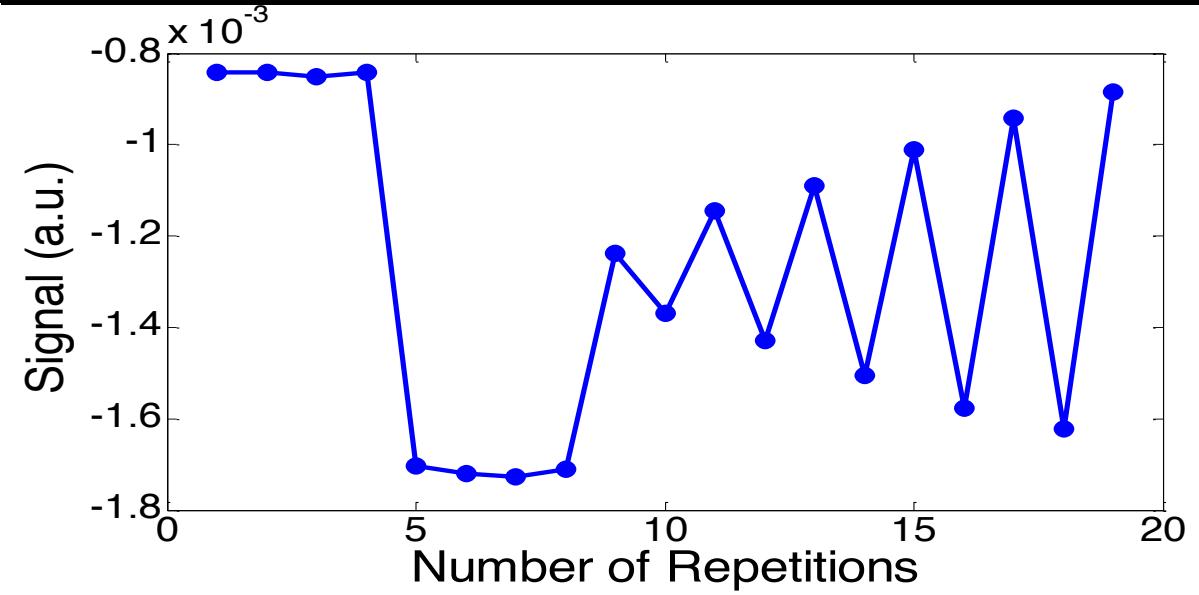
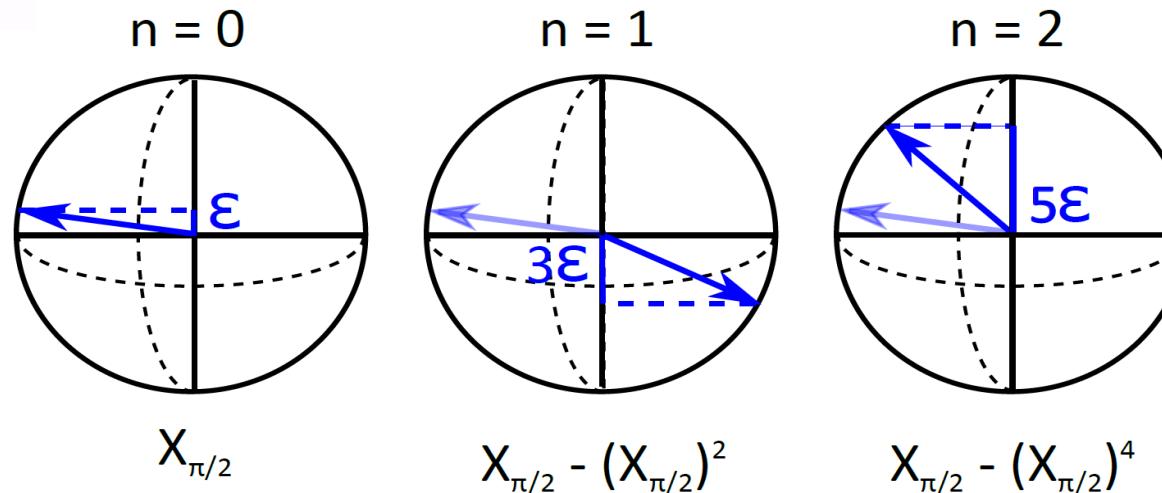


PRL 116,020501 (2016)

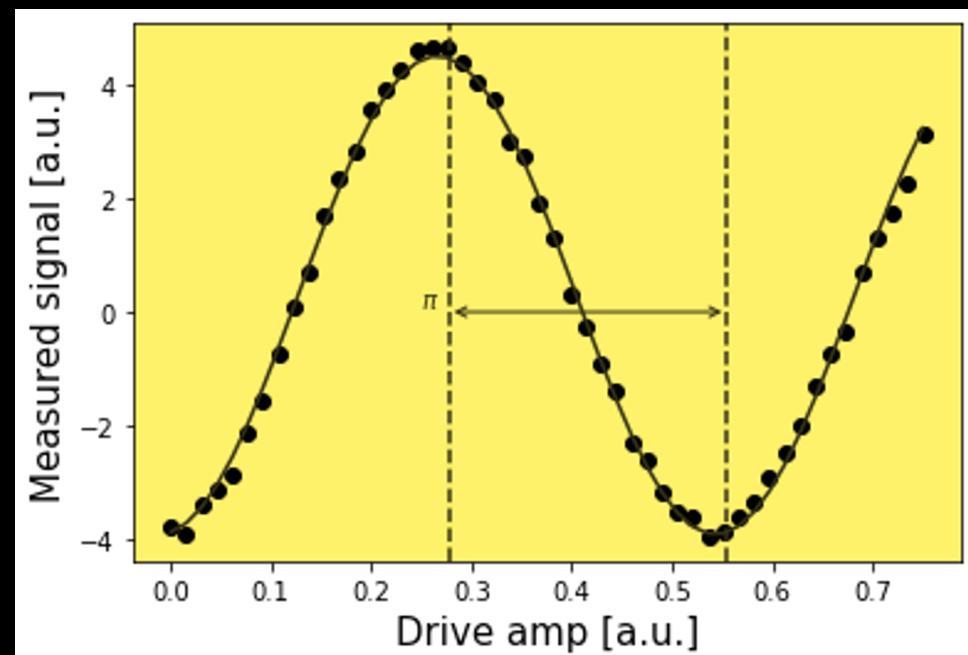
- Can place limits on gate speed
- **Leakage out of computational subspace**
- Pulse shaping is important



Coherent errors



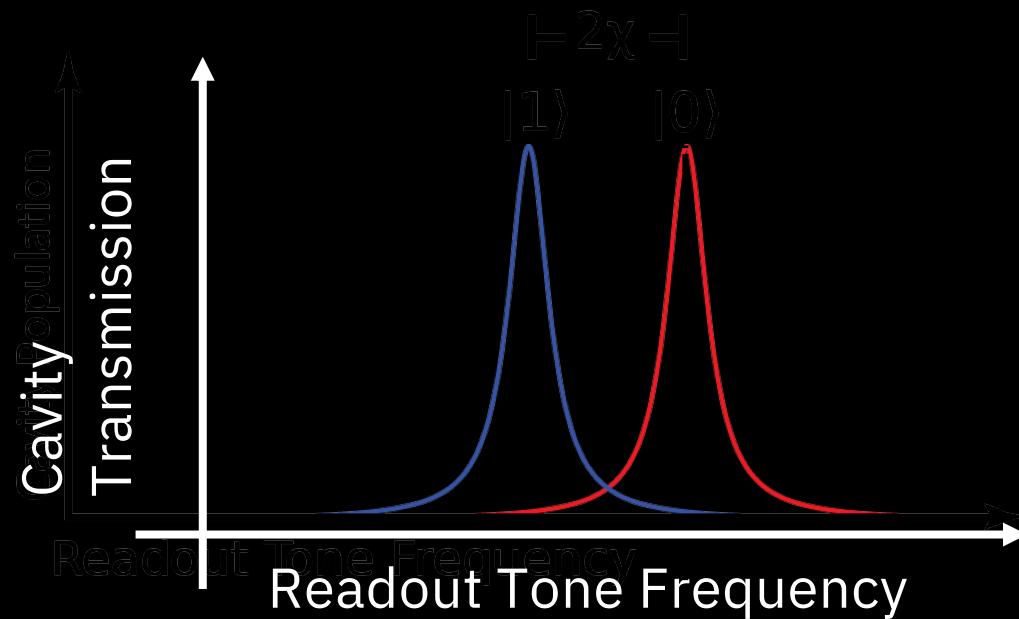
Example: Under/Over rotations



Measurement errors

$$H = \frac{\hbar\omega'_{01}}{2}\hat{\sigma}_z + \hbar(\omega'_r + \chi\hat{\sigma}_z)\hat{a}^\dagger\hat{a}$$

Qubit State dependent resonator frequency



Monitor transmission
at this peak frequency

Example:

Prepare ‘0’

Measured {‘0’ : 99% ‘1’: 1%}

Prepare ‘1’

Measured {‘0’ : 5% ‘1’: 95%}

Measurement fidelity: 97%

Back to chemistry

The Electronic Structure Problem

Interacting fermionic problems: A core challenge in modern computational physics and HPC

$$H_e = - \sum_{i=1}^N \frac{1}{2} \nabla_i^2 - \sum_{i=1}^N \sum_{A=1}^M \frac{Z_A}{r_{iA}} + \sum_{j>i} \frac{1}{r_{ij}}$$

$$H|\psi_G\rangle = E_G|\psi_G\rangle$$

The problem and its qubit representation

$$H_e = - \sum_{i,j} \frac{1}{2} \langle i | \nabla_i^2 | j \rangle a_i^\dagger a_j + \sum_{i,j} \langle i | \frac{Z_A}{r_{iA}} | j \rangle a_i^\dagger a_j + \sum_{i,j,k,m} \langle i, j | \frac{1}{r_{ij}} | k, m \rangle a_i^\dagger a_j^\dagger a_k a_m$$

$$\{a_i^\dagger, a_j\} = \delta_{ij}$$

$$a_j \rightarrow \left(\prod_{i=1}^{j-1} \sigma_i^z \right) \sigma_j^+$$

In general, problem of M electrons in N spin orbitals ($M < N$) mapped onto N qubit- Hamiltonian with $O(N^4)$ Pauli strings

$$a_i^\dagger \rightarrow \left(\prod_{i=1}^{j-1} \sigma_i^z \right) \sigma_j^-$$

Jordan Wigner, Brayvi-Kitaev, Parity mapping ...

$$H = \sum_{\alpha} h_{\alpha} \sigma(\alpha)$$

$$\sigma_A \in \pm \{I, \sigma^x, \sigma^y, \sigma^z\}^{\otimes M}$$

Variational Quantum Eigensolver (VQE)

Variational principle: the energy of any trial wave-function is greater than or equal to the exact ground state energy

$$\frac{\langle \Psi(\vec{\theta}) | H | \Psi(\vec{\theta}) \rangle}{\langle \Psi(\vec{\theta}) | \Psi(\vec{\theta}) \rangle} \geq E_G$$

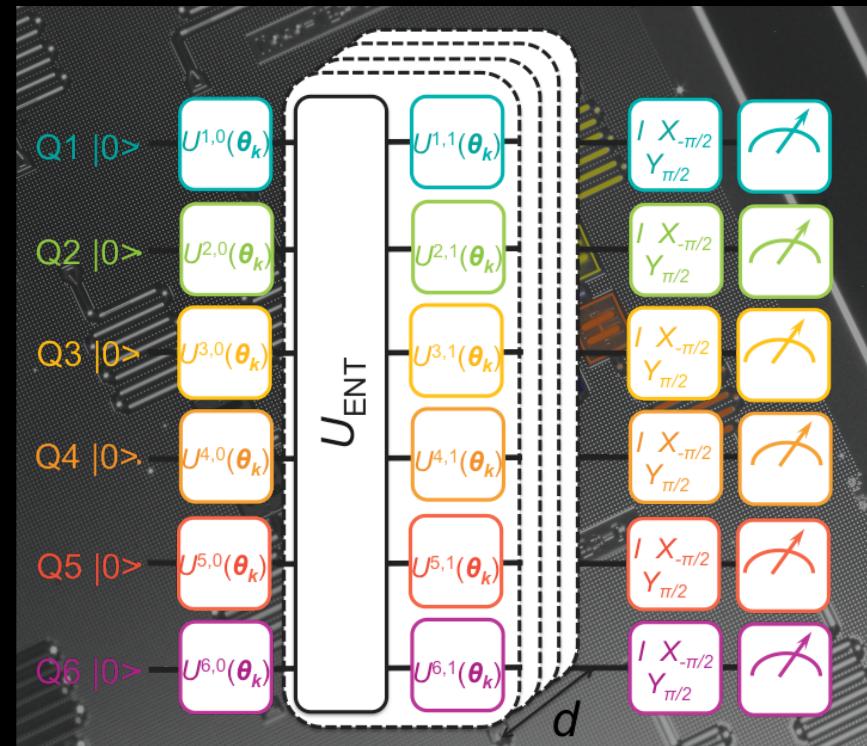
Variational principle

Variational principle

Variational Quantum Eigensolver (VQE)

$$H = \sum_{\alpha} h_{\alpha} \sigma(\alpha)$$

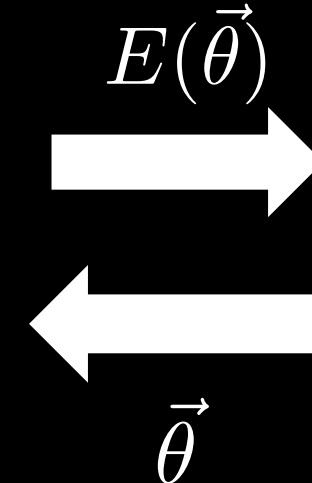
Map problem onto Paulis



Prepare guess state $|\psi_G(\vec{\theta})\rangle$

Measure its energy

$$E(\vec{\theta}) = \sum_{\alpha} h_{\alpha} \langle \psi_G(\vec{\theta}) | \sigma_{\alpha} | \psi_G(\vec{\theta}) \rangle$$



Use classical optimizer

Mapping fermions to qubits

$$H = h_1 \cancel{IXIY} + h_2 XZZX + h_3 XYZ$$

Pauli string



Number of non-identity single qubit Pauli operators : weight of the Pauli string

Larger weight Pauli strings are increasingly sensitive to measurement error

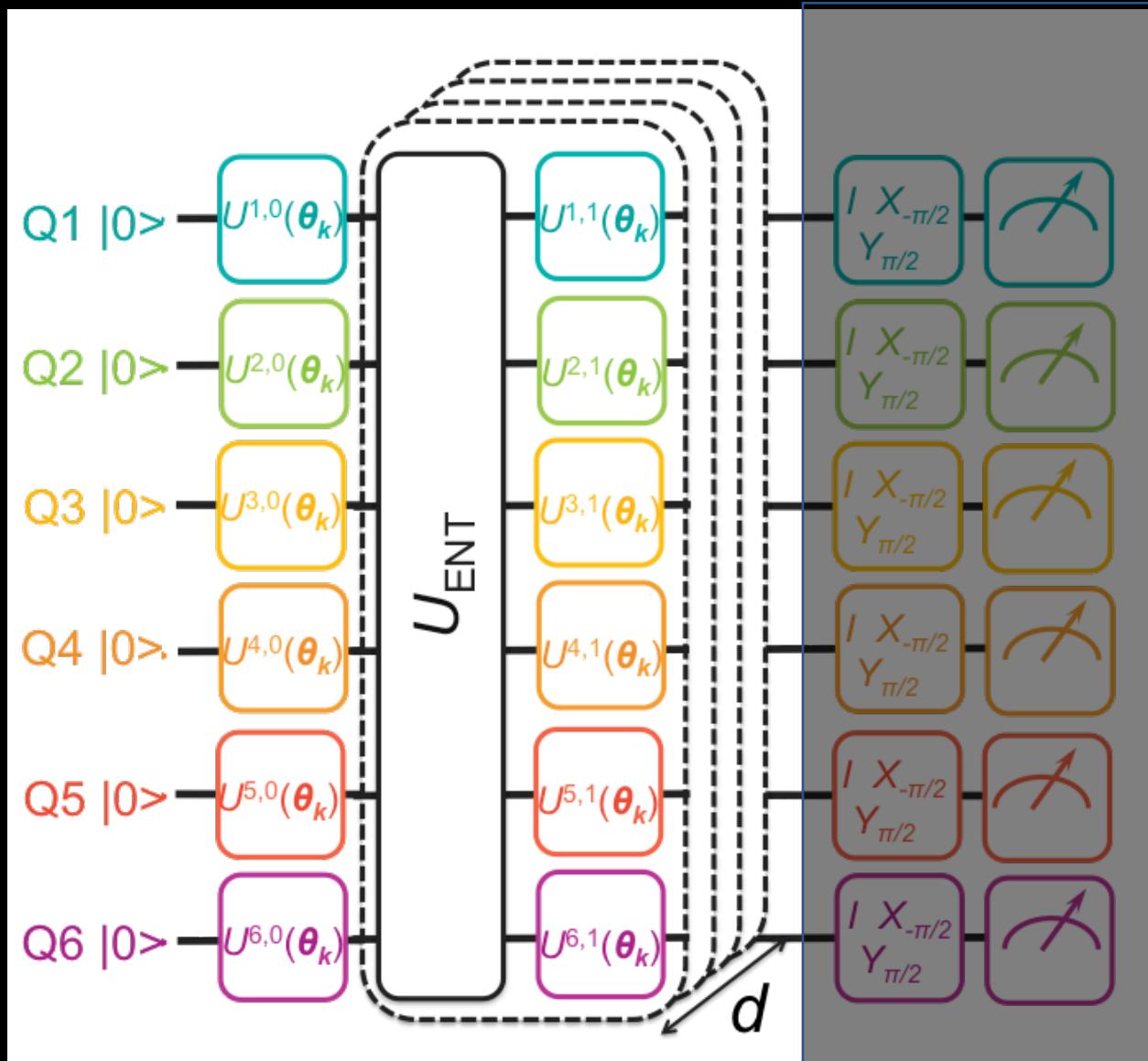
Choice of fermion-qubit mapping important: Each fermionic operator maps to $O(N)$ qubits for JW, $O(\log N)$ for BK

Tapering of qubits

$$H = h_1 IXIY + h_2 XZZX + h_3 XYZ$$

Tapering of qubits

Trial State Preparation

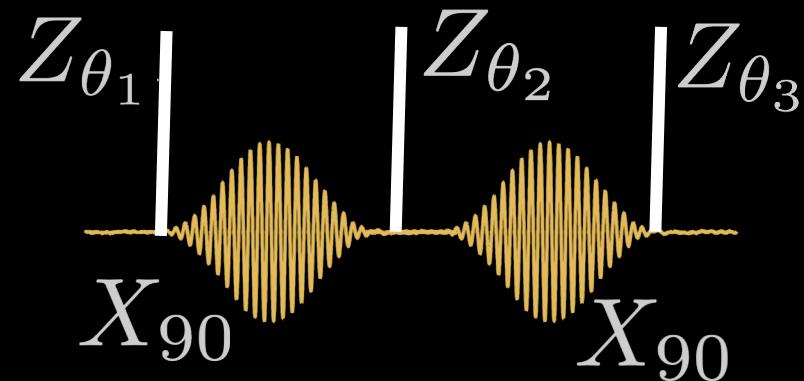


- Unitary Coupled Cluster ansatz : Offers physical intuition, focuses ground state search in "good" space
- However, translates to very expensive gate count
- Alternative approach: Keeping hardware limitations in mind, do what hardware can do best!

Trial State Preparation

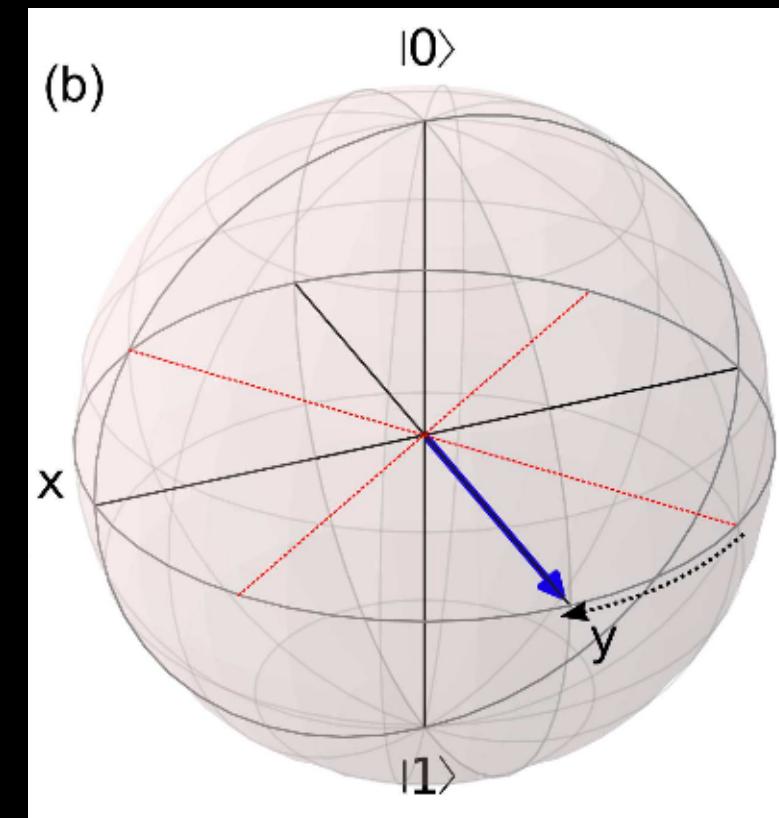
Trial state preparation: Arbitrary single qubit rotation

$$U(\vec{\theta}) = Z_{\theta_1} X_{90} Z_{\theta_1} X_{90} Z_{\theta_1}$$

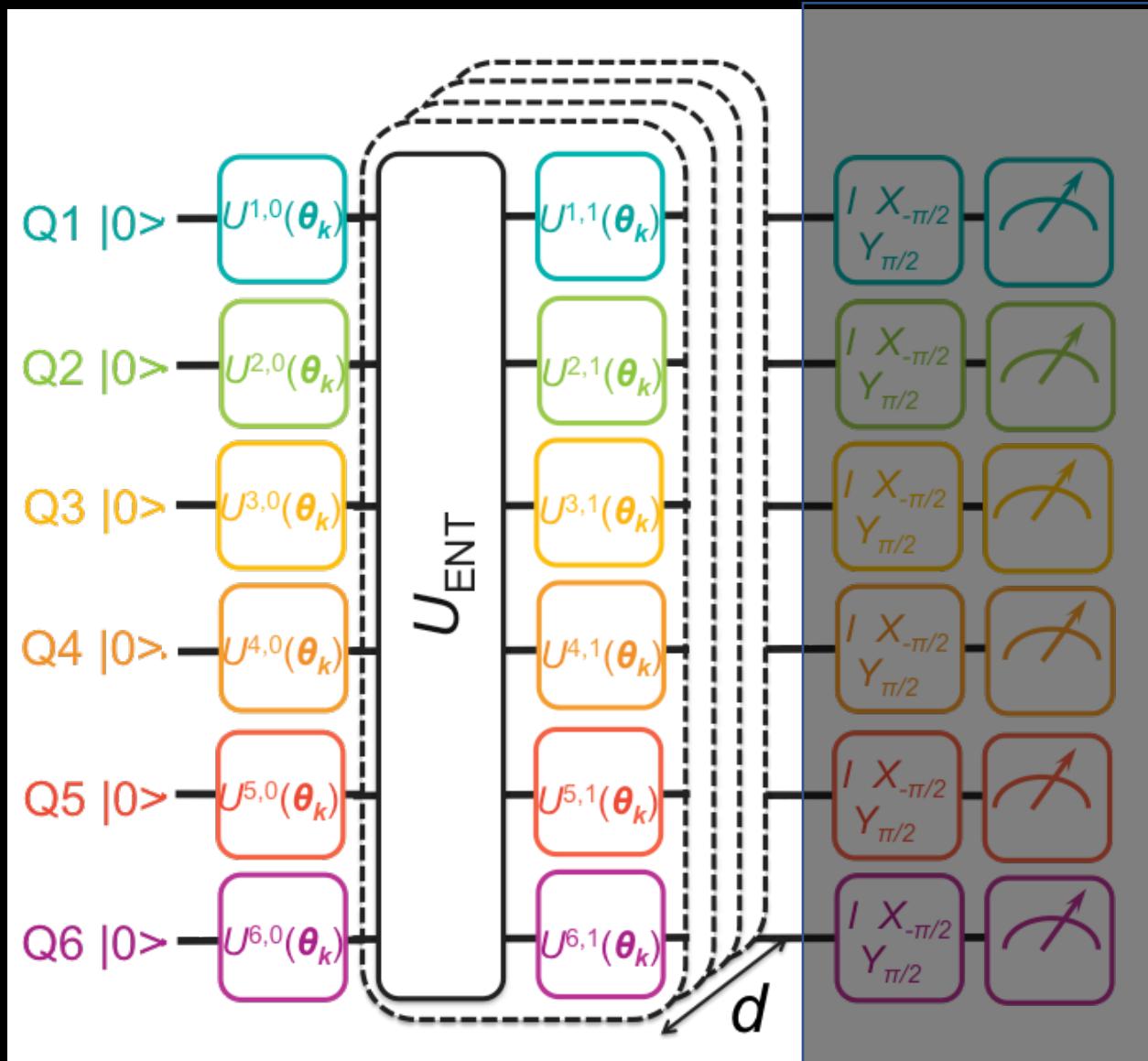


- Variational parameters:
- Pulse envelope
 - Pulse length
 - Amplitude
 - Phase

Software implemented Z-gates
(Perfect, zero time)

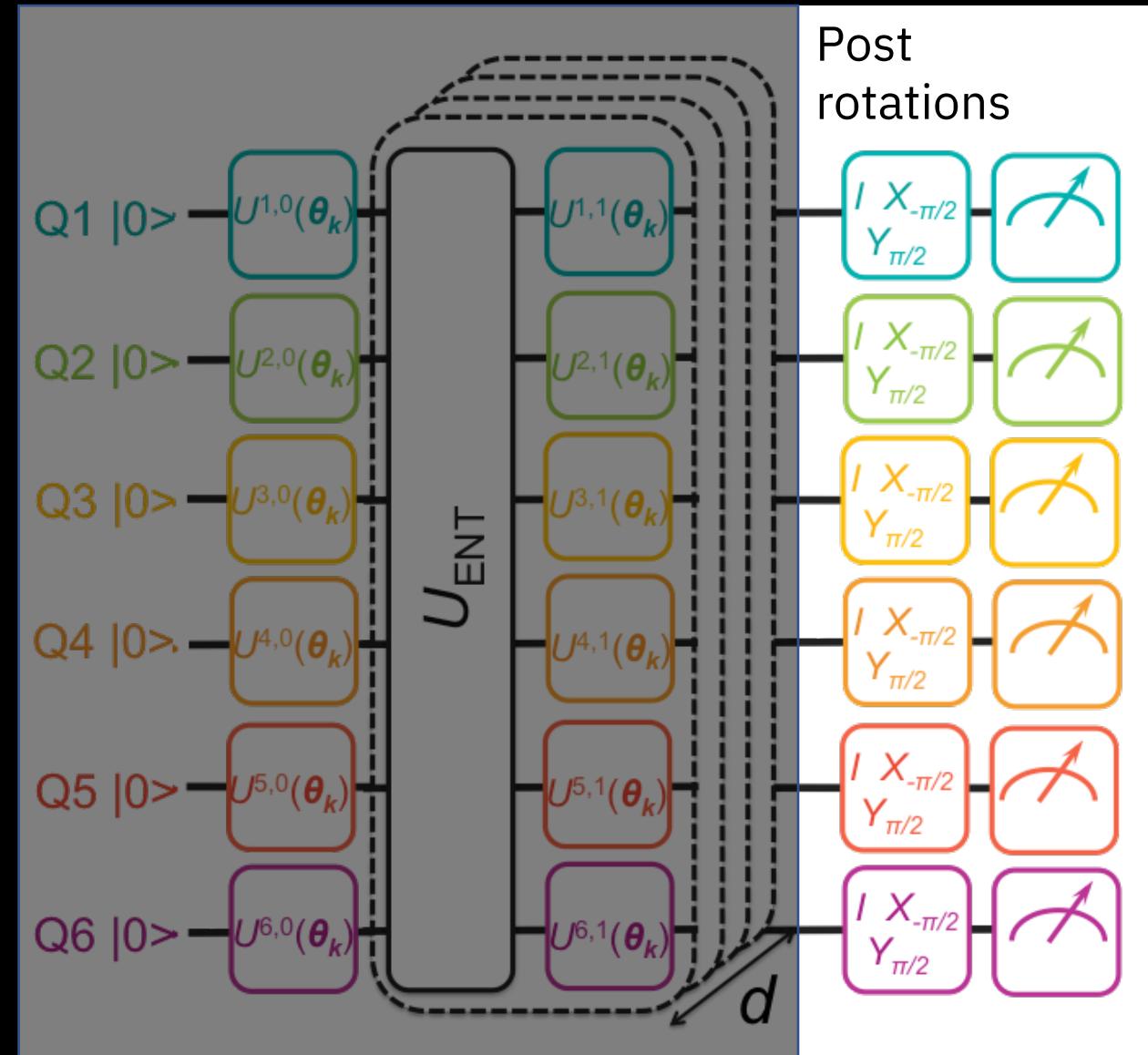


Trial State Preparation



- Sequence of interleaved arbitrary single qubit rotations and naturally available entangling gates
- Depth set by available quantum coherence (incoherent errors)
- $N(3d+2)$ variational parameters
- This structure offers some robustness to coherent errors

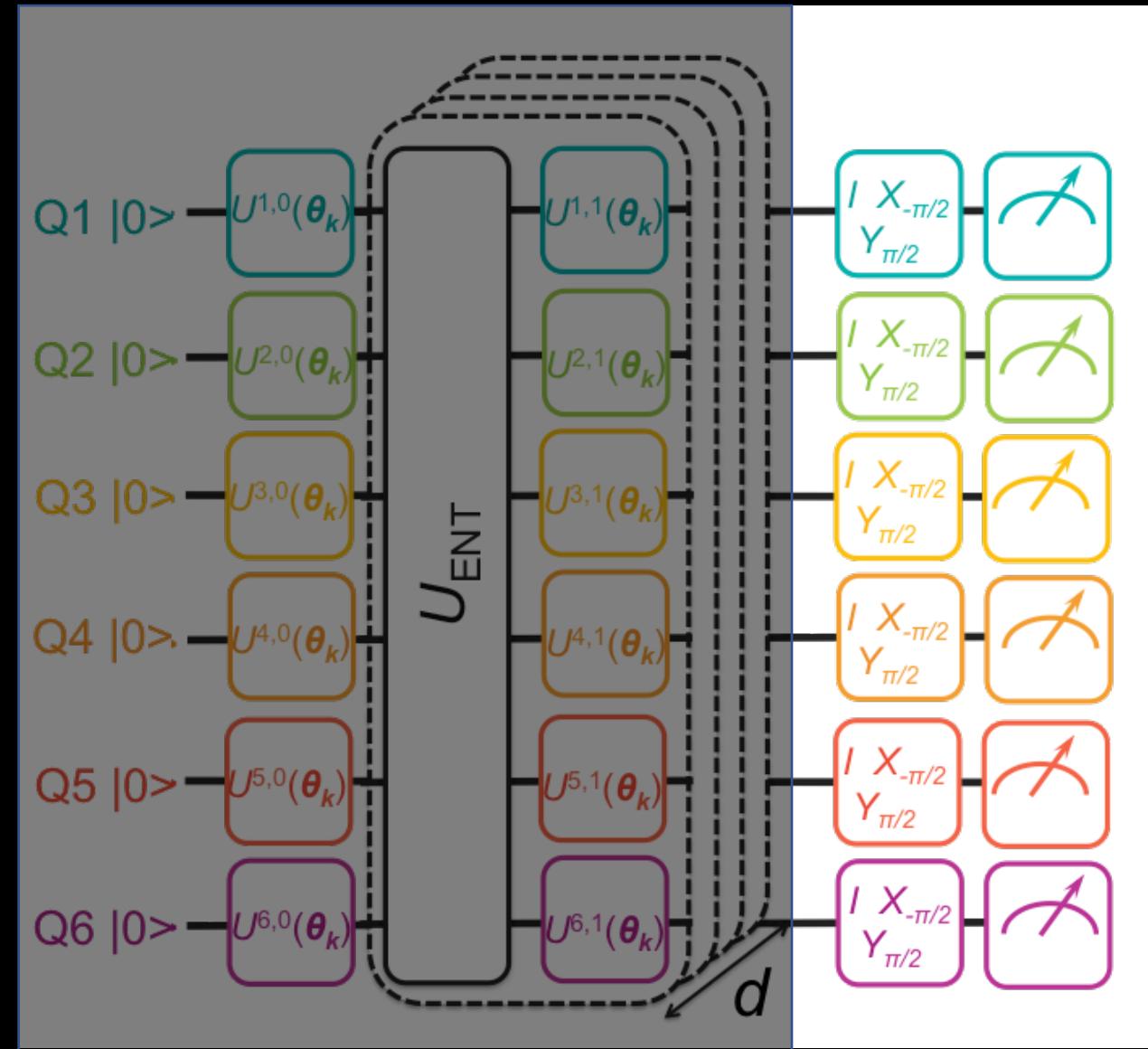
Energy measurement : Measuring expectation values



Energy measurement : Measuring expectation values

Energy measurement : Measuring expectation values

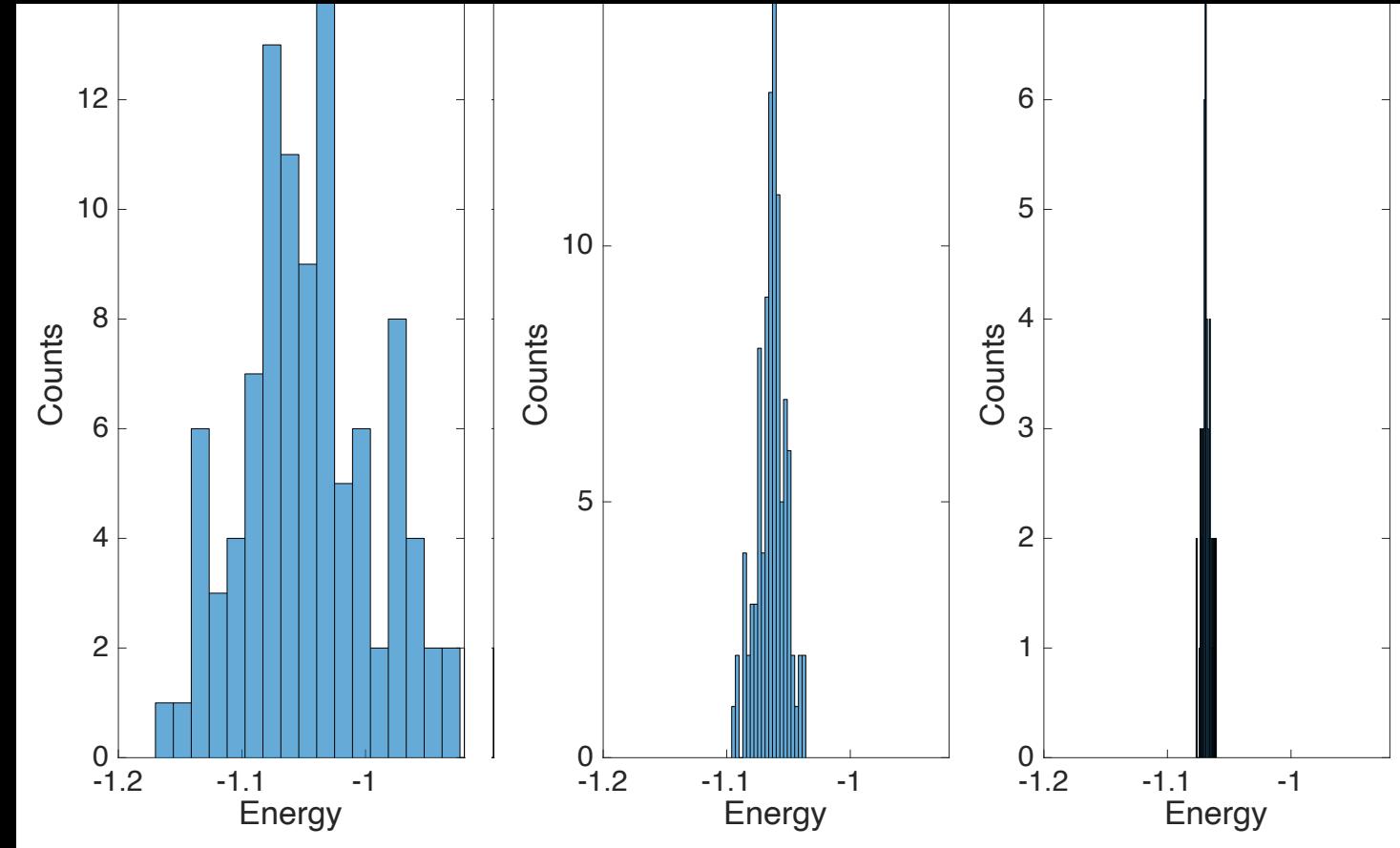
Energy measurement: Measurement Basis



Energy measurement: Measurement Basis

Energy measurement: Pauli Grouping

Energy measurement: Shot noise



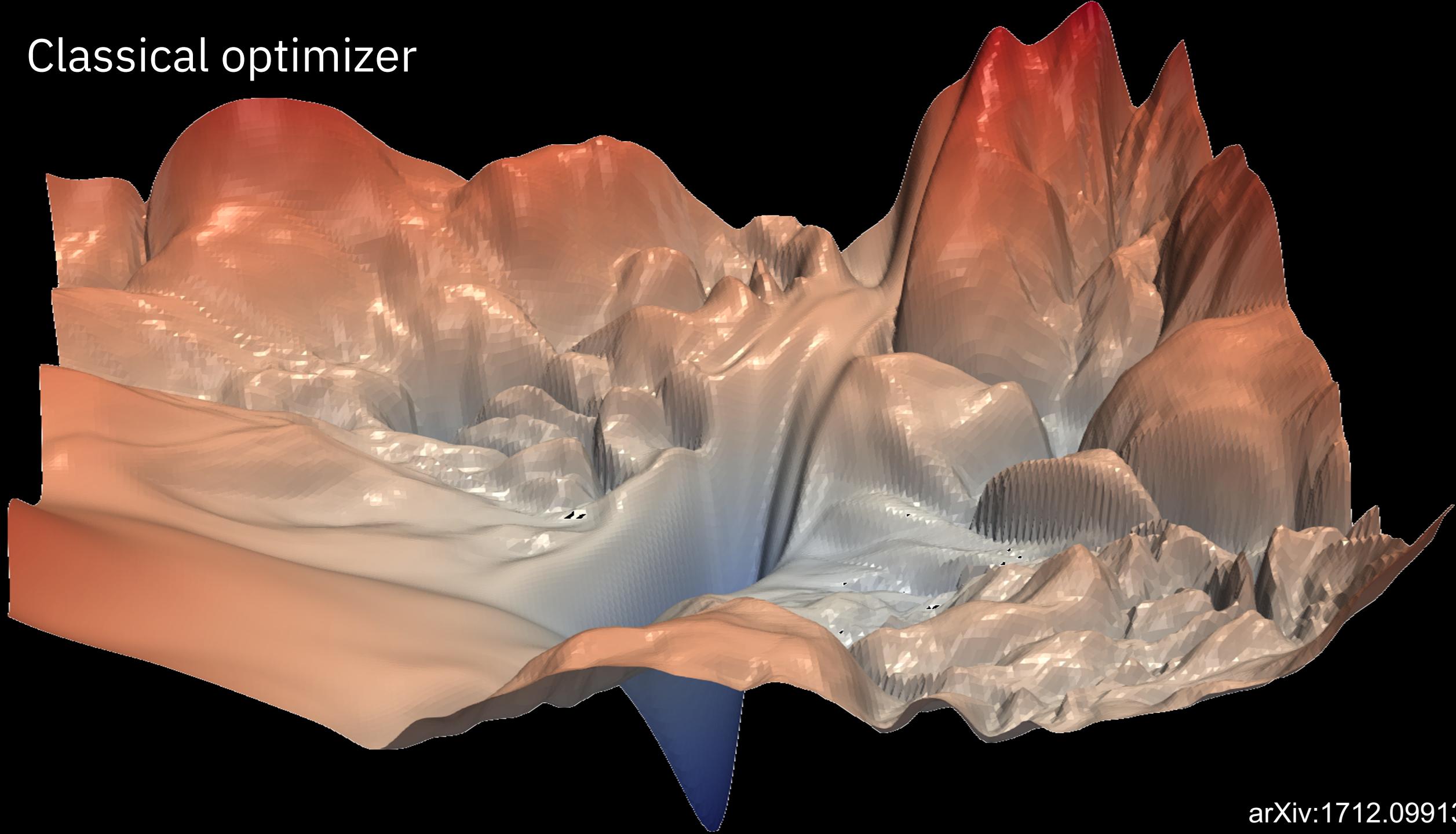
Number of shots N

Even if ground state preparation is *perfect*, estimating $\langle O \rangle$ will still have error

$$\epsilon_{\langle O \rangle} = \sqrt{\frac{\langle O^2 \rangle - \langle O \rangle^2}{S}}$$

Can affect accuracy of classical optimizer

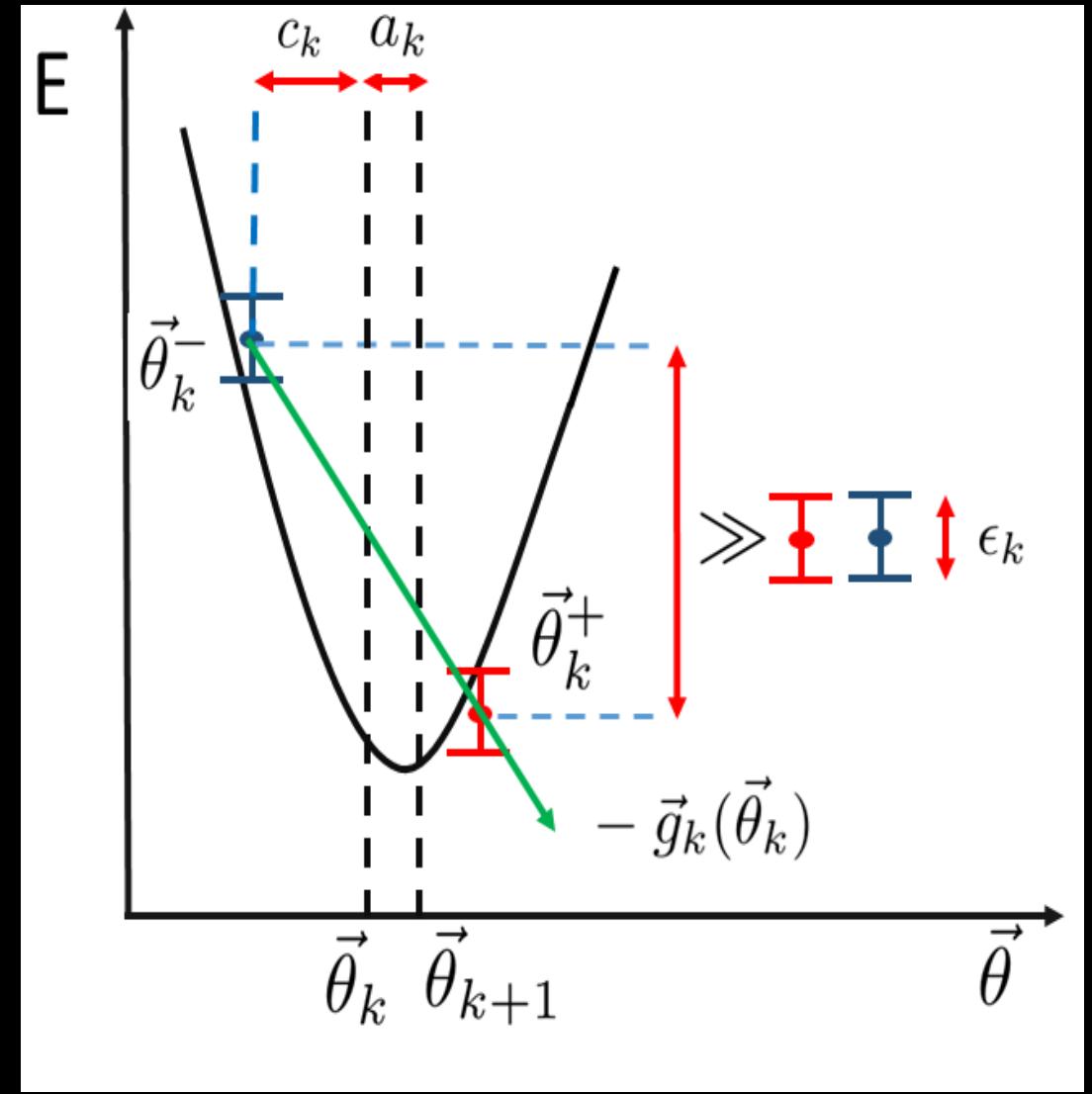
Classical optimizer



Classical optimizer: Simultaneous Perturbation Stochastic Approximation (SPSA)

Need to reduce of calls to Quantum computer

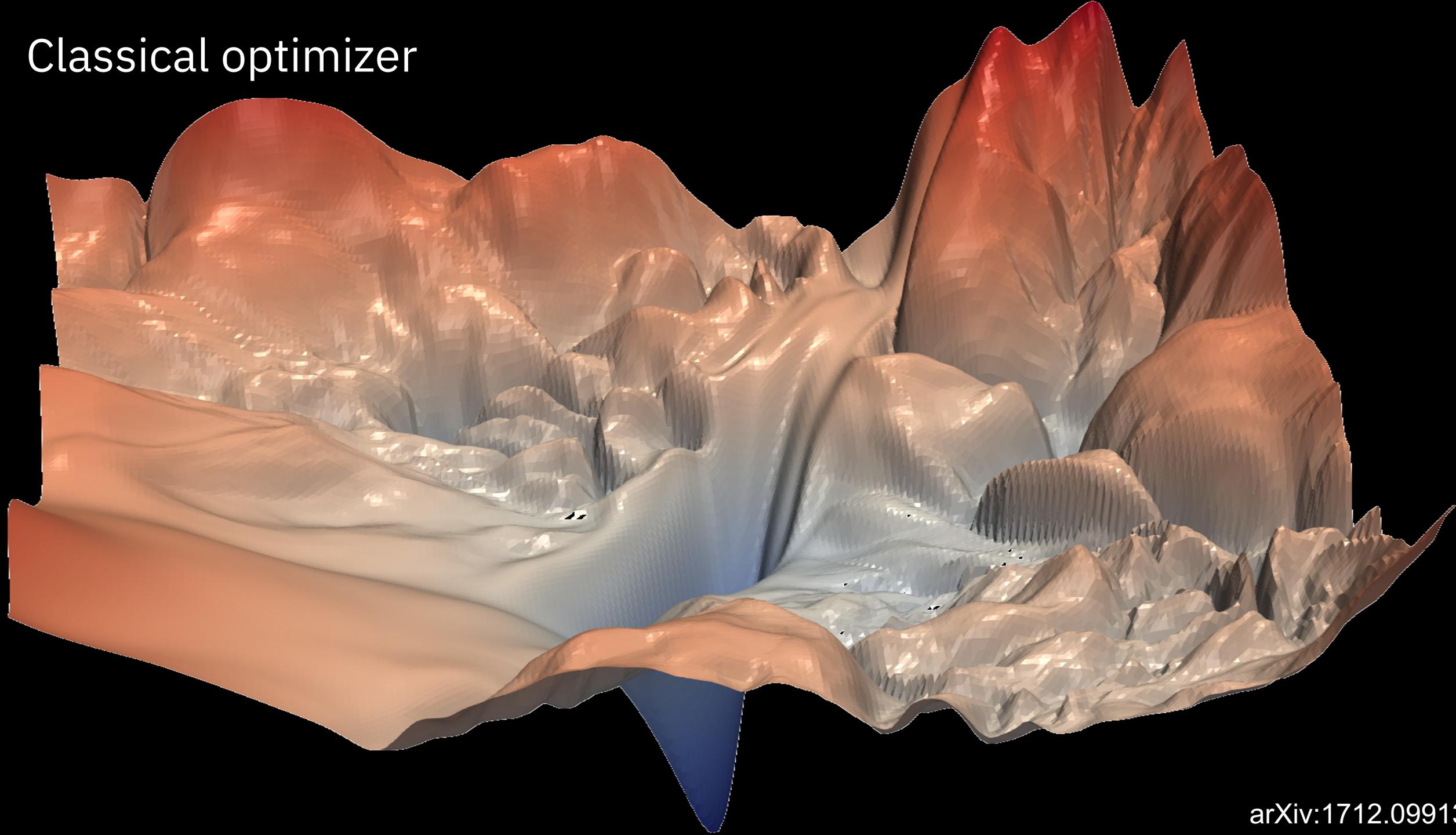
Gradient approx.: Regardless of dimension of optimization problem, utilizes only **two** measurements per iteration



[1] J. C. Spall, Multivariate stochastic approximation using a simultaneous perturbation gradient approximation, IEEE Transactions on Automatic Control 37, 332 (1992)

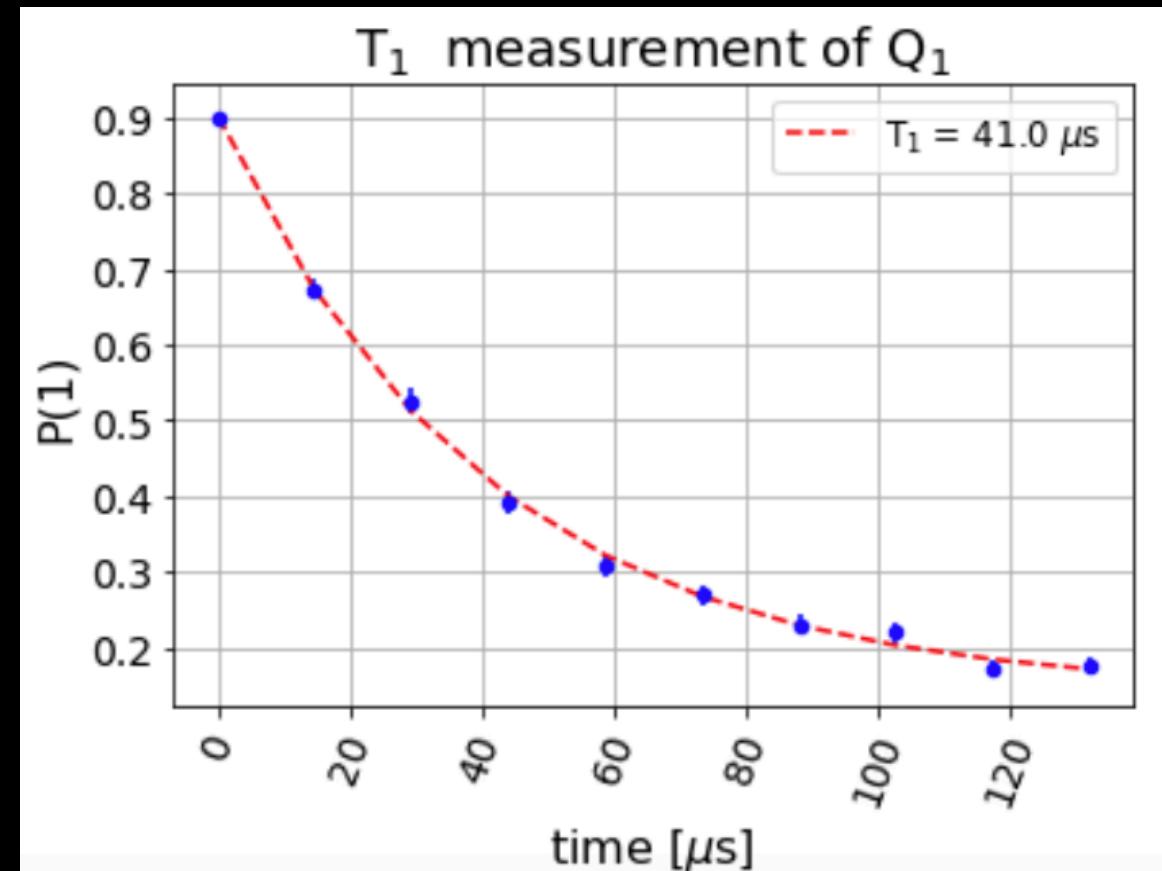
A. Kandala, A. Mezzacapo, et al,
Nature 549, 242-246 (2017)

Classical optimizer



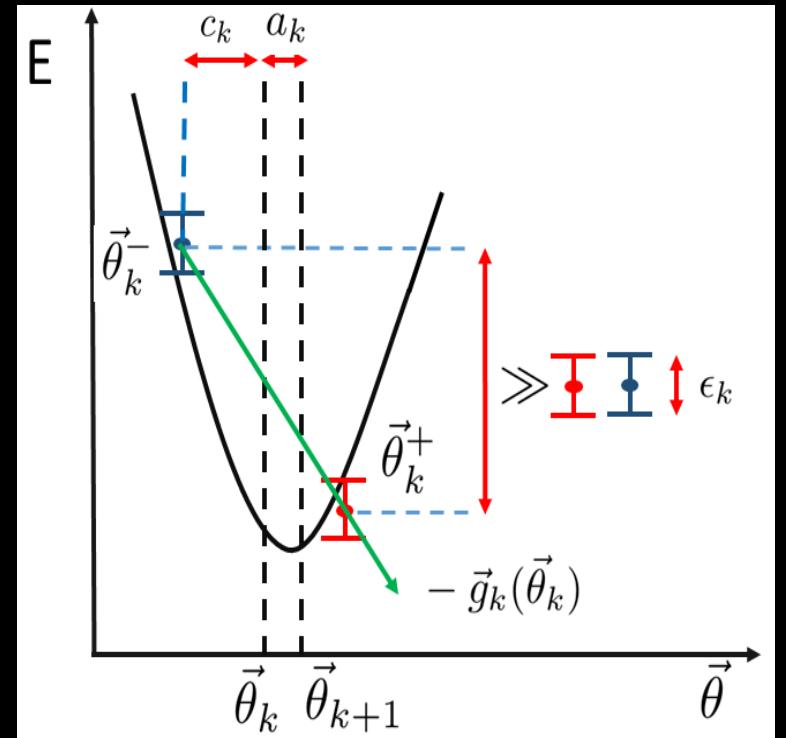
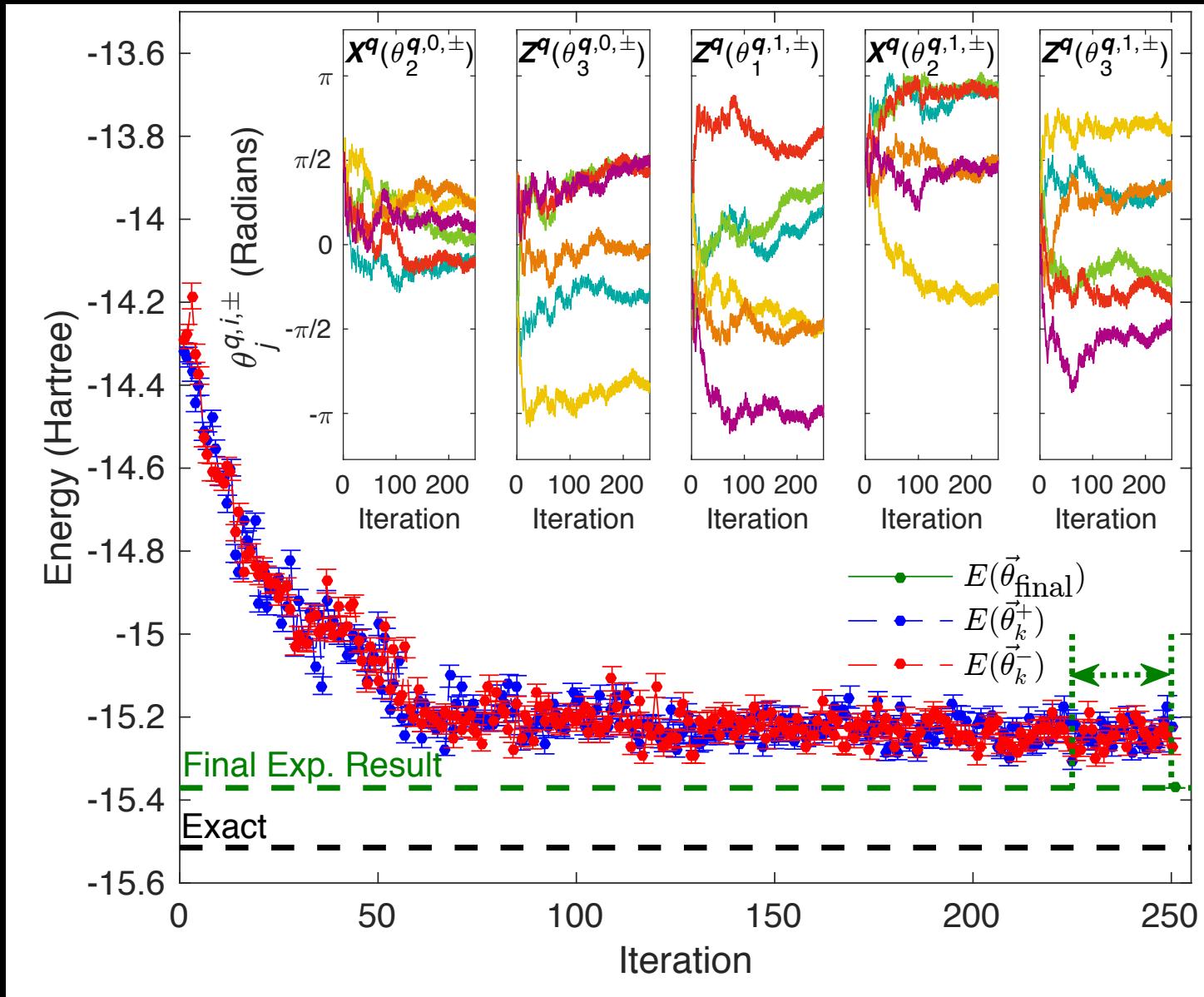
Run time overhead

- Number of energy estimations per iteration
- Number of Pauli strings in the Hamiltonian
- Number of shots
- System initialization



Putting it all together ..

Experimental optimization of a 6 Qubit Hamiltonian



Application to quantum chemistry : H₂

H_A: 1s¹ H_B: 1s¹

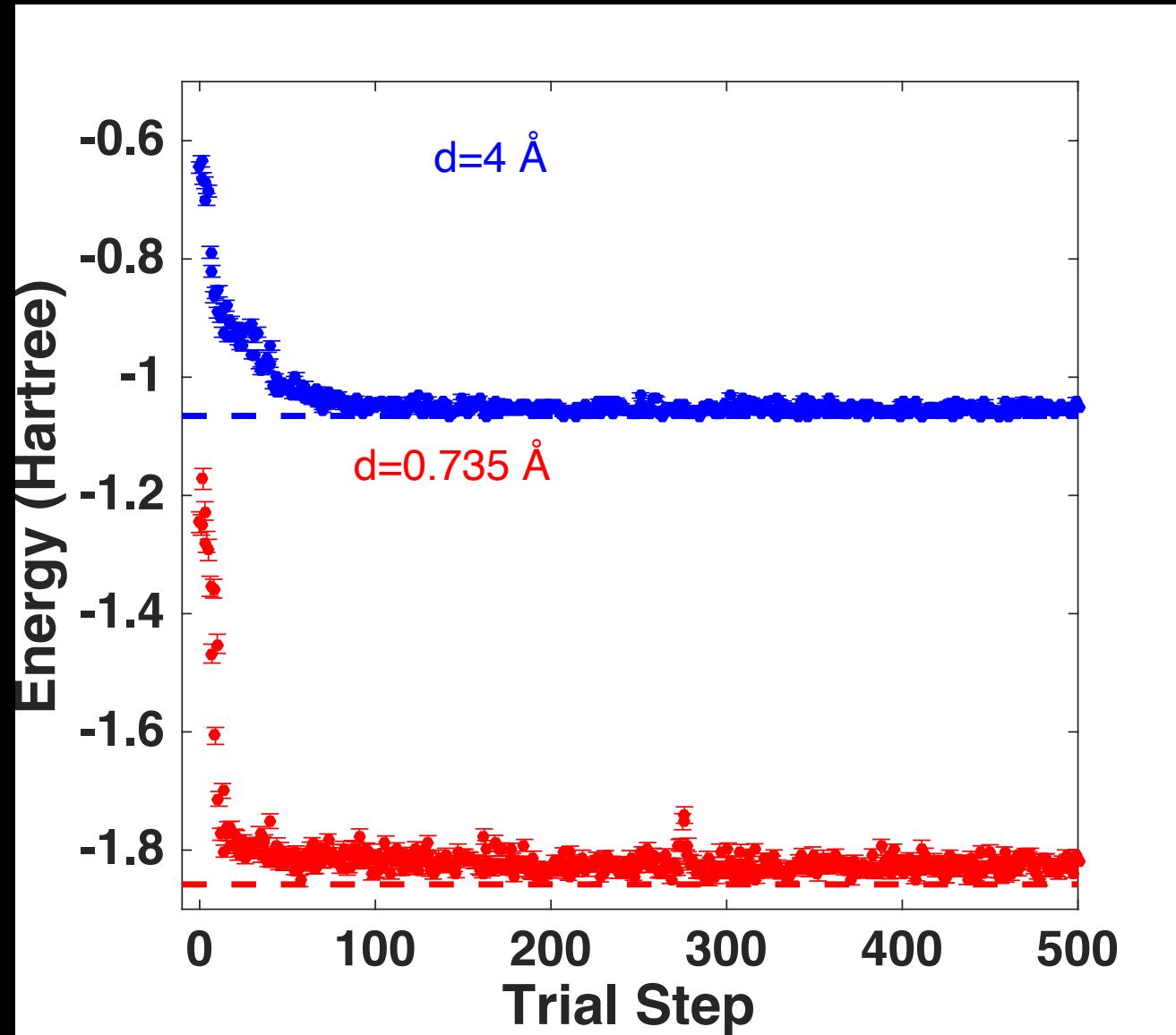
4 spin orbitals mapped to 2 qubits

Equilibrium d = 0.735 Å

$H = (-1.05237)II + (0.39735)ZI + (0.39735)IZ +$
 $(0.11279)ZZ + (0.18093)XX$

Dissociation d = 4 Å

$H = (-0.70461)II + (0.00012)ZI + (0.00012)IZ +$
 $(1.6673e-10)ZZ + (0.33438)XX$



Application to quantum chemistry : H₂

H_A: 1s¹ H_B: 1s¹

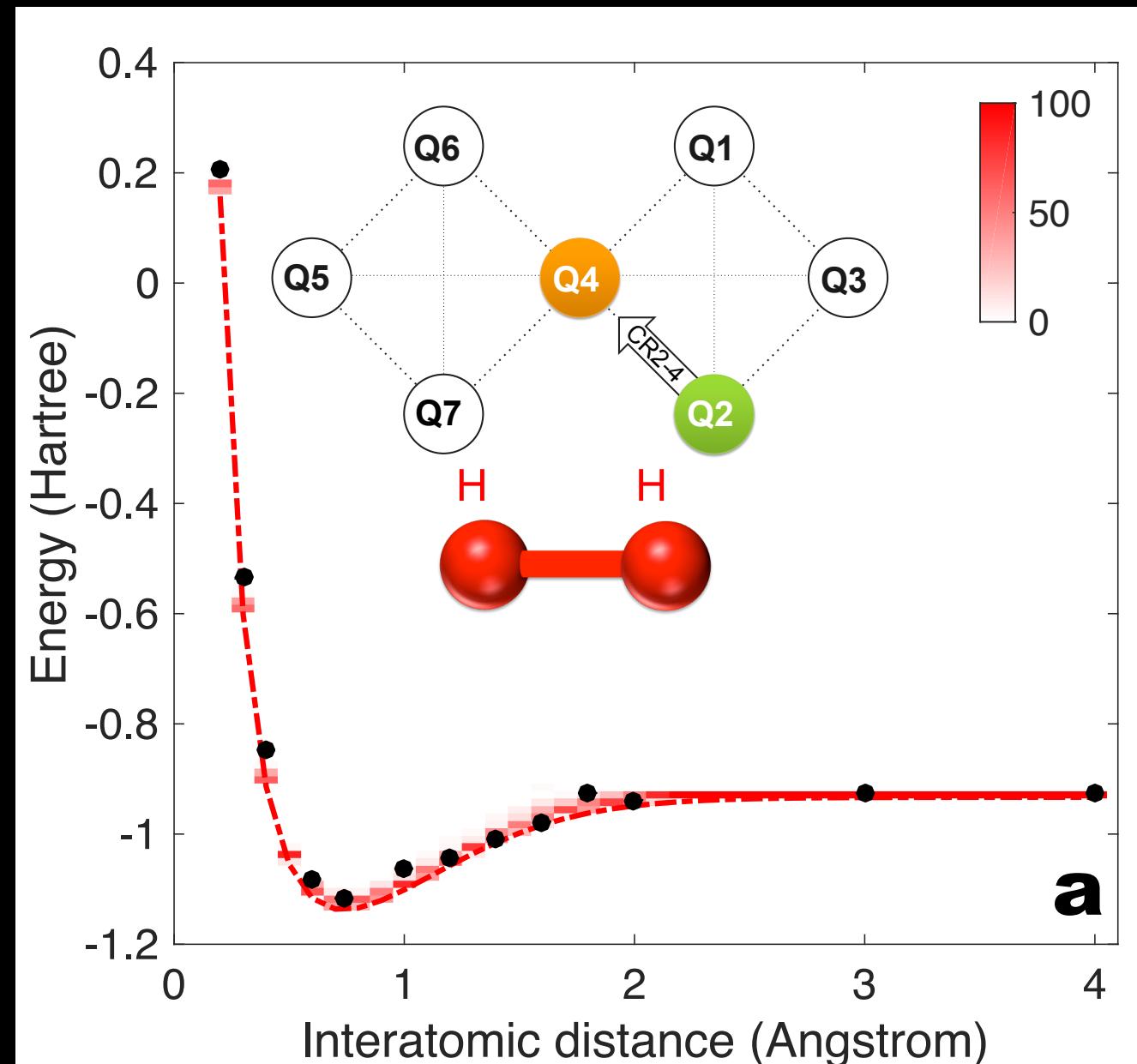
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Dissociation d = 4 Å

$H = (-0.70461)II + (0.00012)ZI + (0.00012)IZ +$
 $(1.6673e-10)ZZ + (0.33438)XX$



Application to quantum chemistry : Going beyond period 1 ...

II
-1.05237
ZI
0.39735
IZ
0.39735
ZZ
0.11279

XX
0.18093

H₂: 5 Pauli terms, 2 sets

ZZZZ
0.080333850925909334
IIII
-0.20665748323683217
ZIII
-0.096022179507889355
ZZII
-0.20612801807987652
IZII
0.36474560607027956
IIIZ
0.096022179507889077
IIIZZ
-0.20612801807987641
IIIZ
-0.36474560607027956
ZIZI
-0.14543788191435422
ZIZZ
0.056040328985617657
ZIIZ
0.11081066572357191
ZZZI
-0.056040328985617643
ZZZI
0.063672720404946023
IZZI
0.11081066572357193
IZZZ
-0.063672720404946037
IZIZ
-0.095215796658694896

XZXX
0.0081947444613071196
XZIX
0.0012708408559466877

ZZXZ
-0.0026669416291266101
XZII
-0.012585096432460004
XIII
0.012585096432460002
IIXZ
0.012585096432460008
IIXI
0.012585096432460006
XZXI
-0.0026669416291266101
XIXZ
0.0026669416291266101
XIXI
0.0026669416291266101
XZIZ
0.0072647375176923996
XIIZ
-0.0072647375176923996
IZXZ
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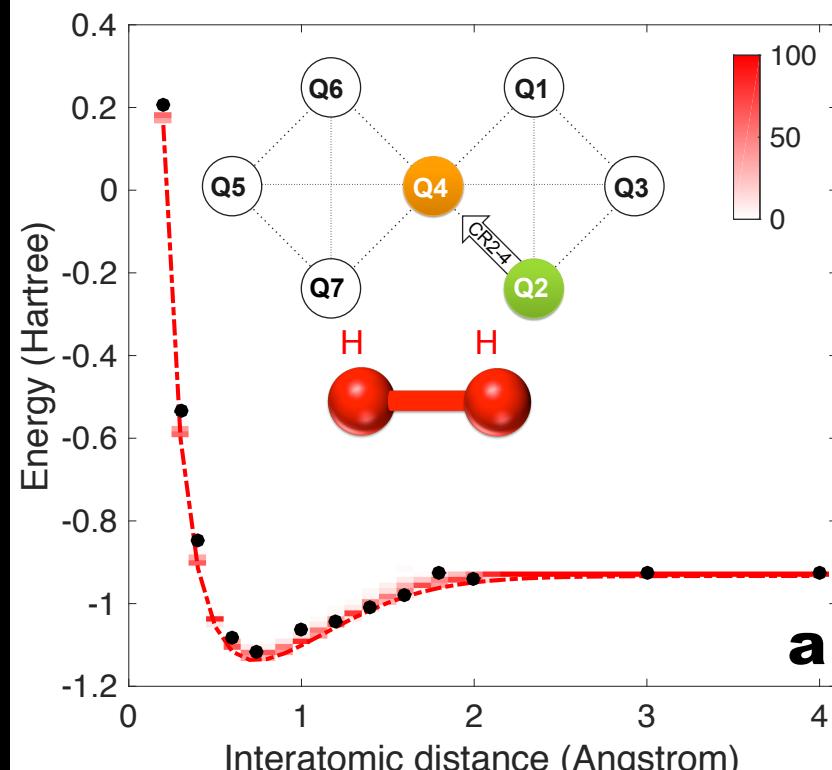
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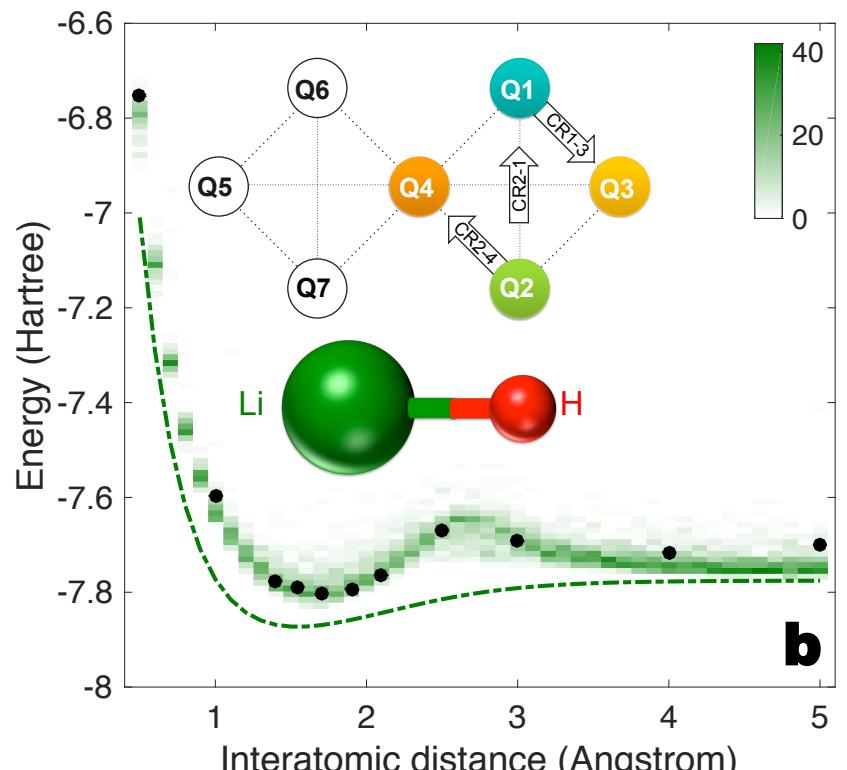
LiH: 100 Pauli terms, 25 sets

VQE for quantum chemistry



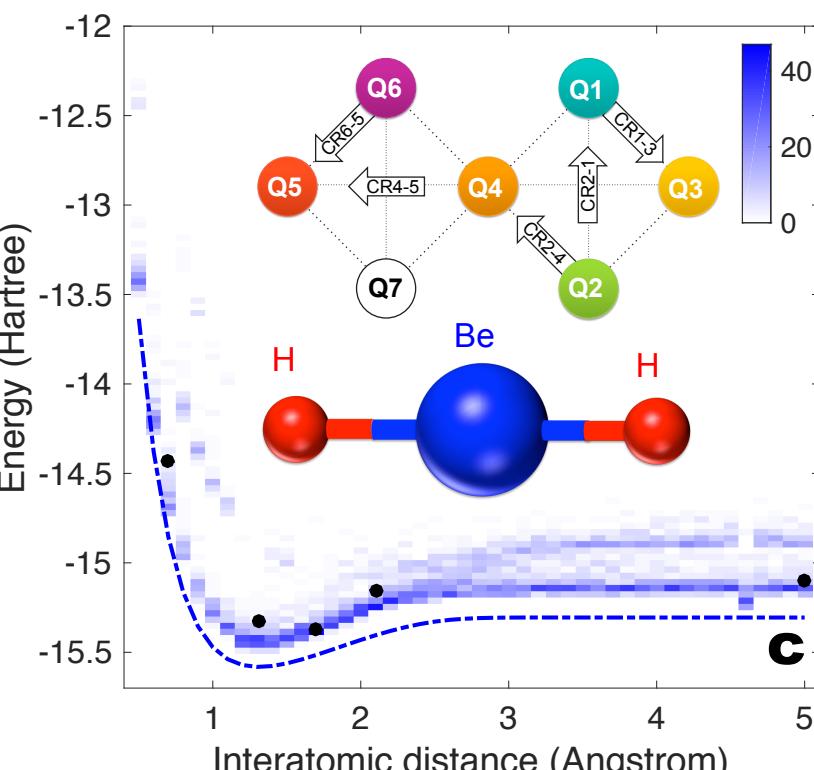
H_2 : 2 qubits
5 pauli terms, 2 sets

- Decoherence
- Sampling error
- Limited iterations



LiH : 4 qubits
100 pauli terms, 25 sets

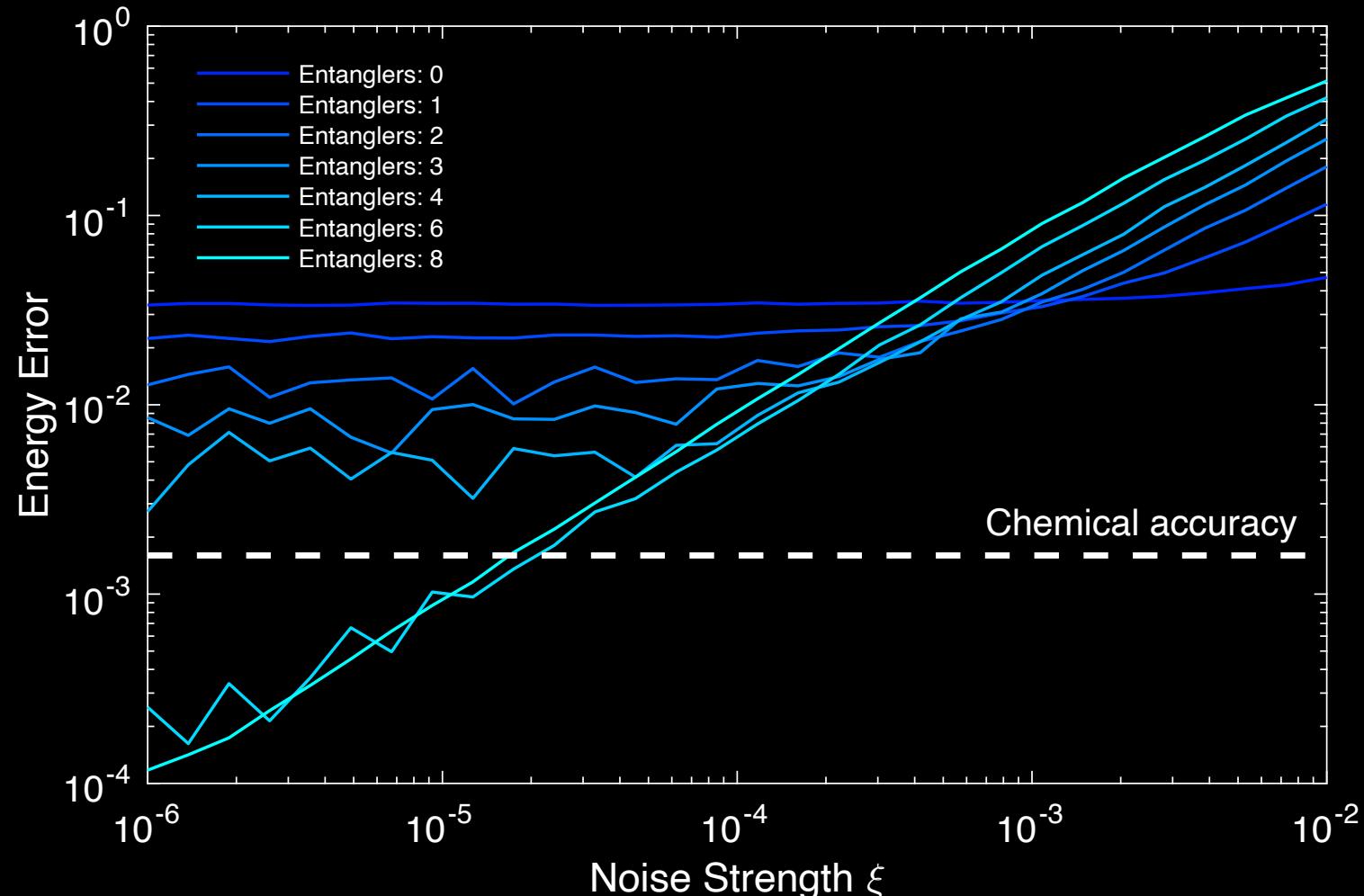
- Accuracy of the classical optimizer
- Insufficient depth



BeH_2 : 6 qubits
165 pauli terms, 44 sets

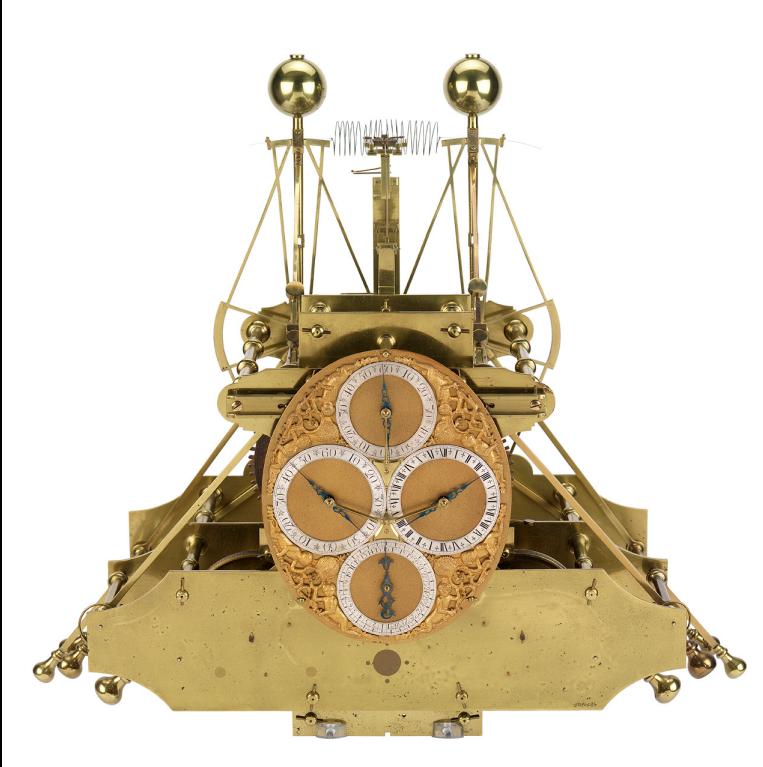
A. Kandala, A. Mezzacapo, et al,
Nature **549**, 242-246 (2017)

Performance trade-off: Decoherence v/s Circuit depth



The longitude problem

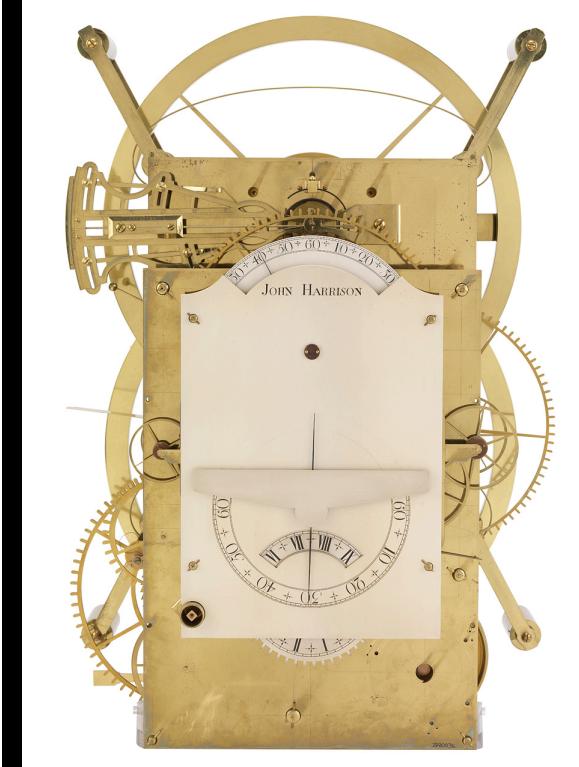
IBM Quantum



H1
1730-1735



H2
1737-1739



H3
1740-1759
Invention of
Bimetallic strip



H4
1755-1759

Error mitigation

Zero-noise extrapolation

Expectation value of observable of interest:

$$E_K(\lambda) = E^* + \sum_{k=1}^n a_k \lambda^k + R_{n+1}(\lambda, \mathcal{L}, T)$$

Assume experimentalist has exquisite control over incoherent noise (T_1, T_2)

$$E_K(\lambda) = E^* + a_K \lambda + \mathcal{O}(\lambda^2)$$

$$E_K(c\lambda) = E^* + a_K c \lambda + \mathcal{O}(c^2 \lambda^2)$$

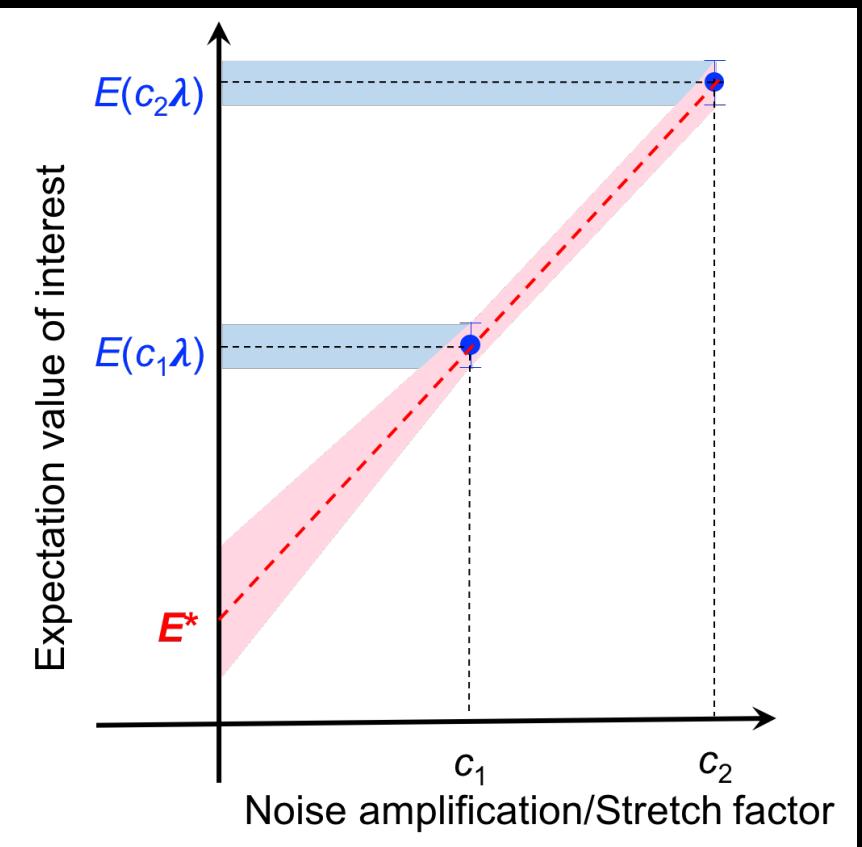
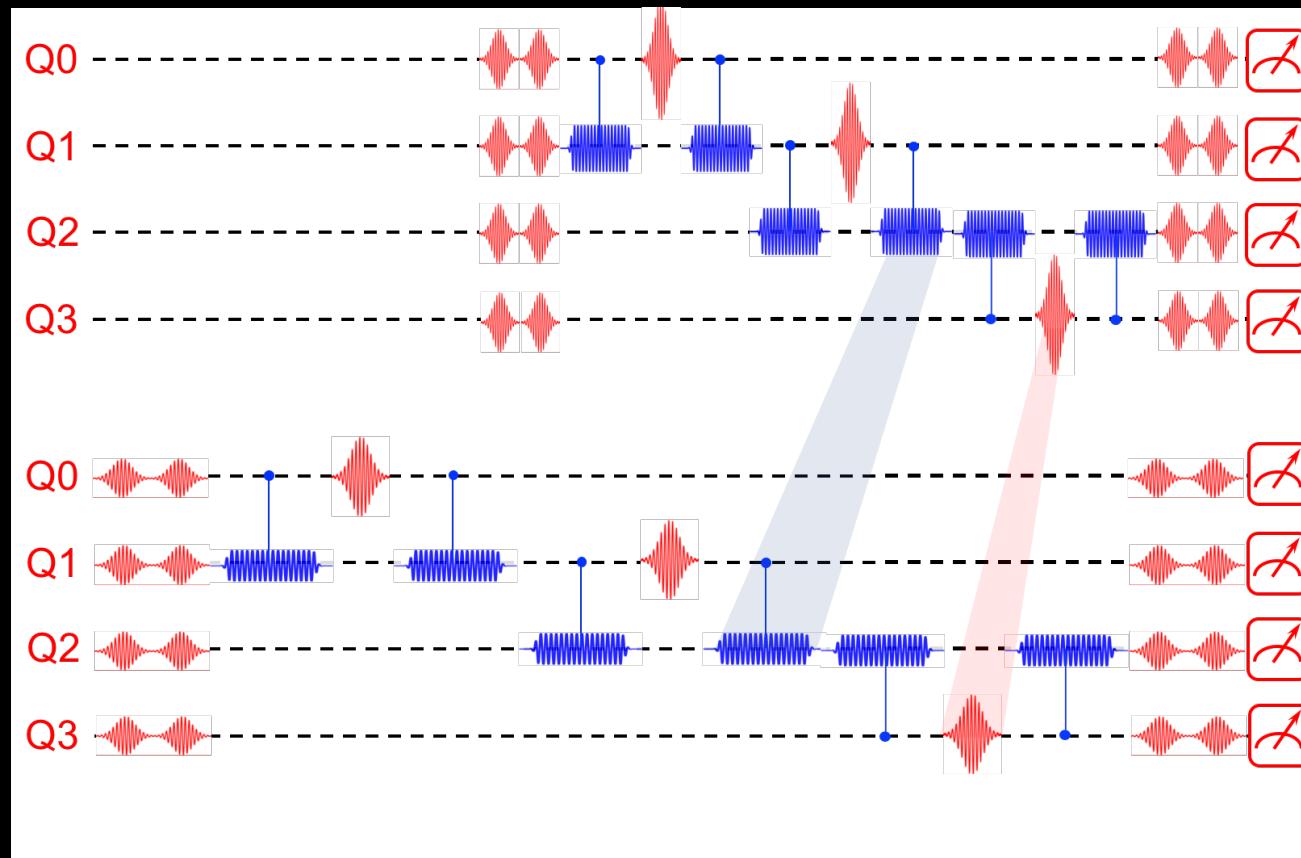
$$\hat{E}_K^2(\lambda) = \frac{cE_K(\lambda) - E_K(c\lambda)}{c - 1} = E^* + \mathcal{O}(\lambda^2)$$

With N measurements, can reduce error in estimate to $\mathcal{O}(\lambda^N)$

How does one measure $E(c_i \lambda)$??

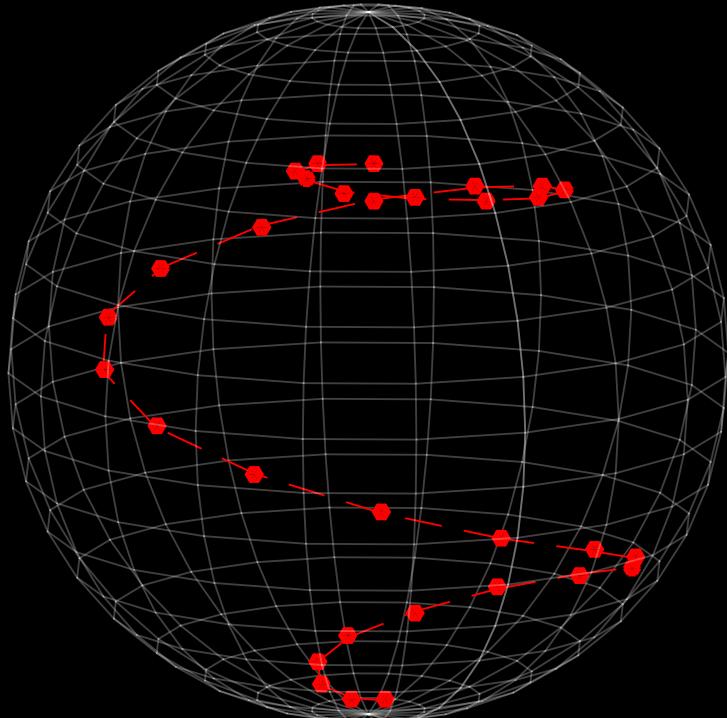
When two wrong's make a right

Amplifying noise strength equivalent to rescaling dynamics under the assumption of time invariant noise.

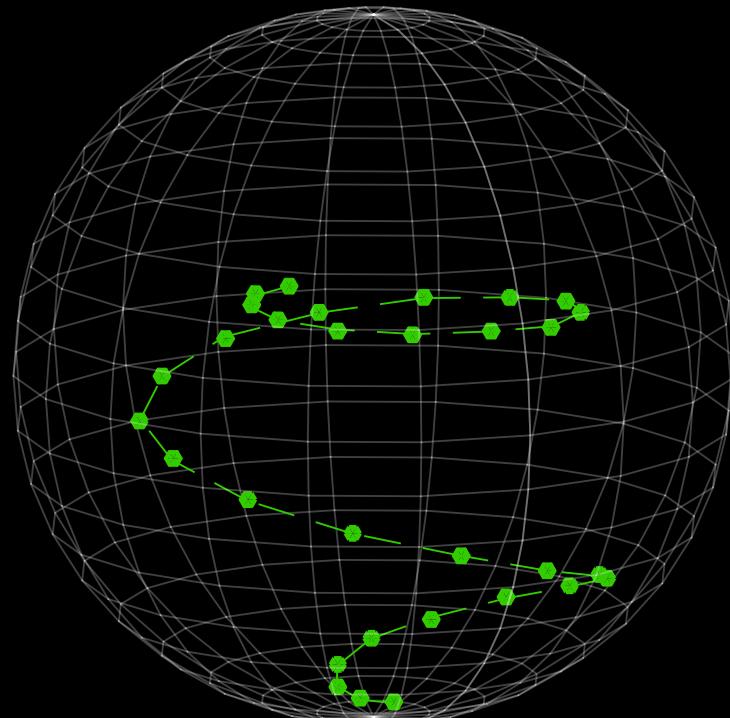


Error Mitigation: 1Q trajectory

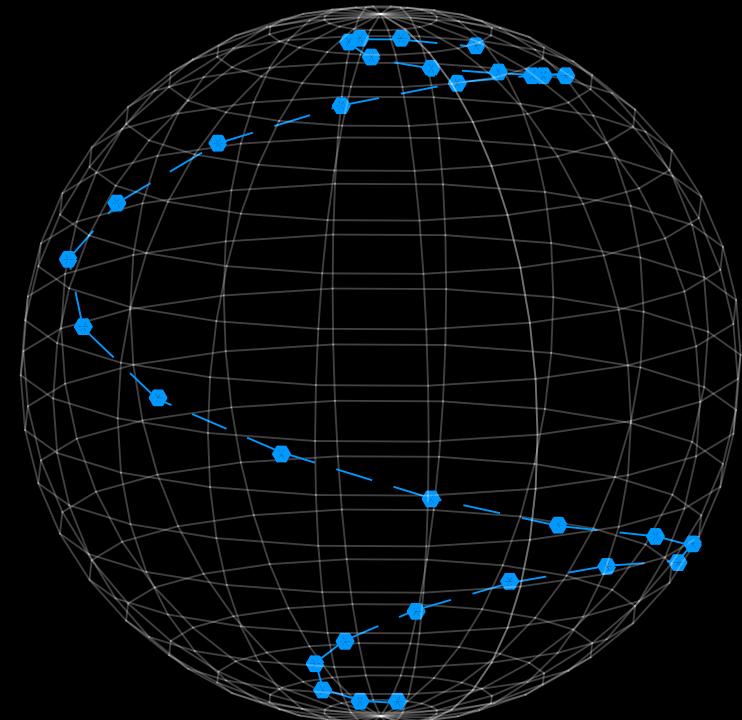
$|1\rangle$



$|1\rangle$



$|1\rangle$

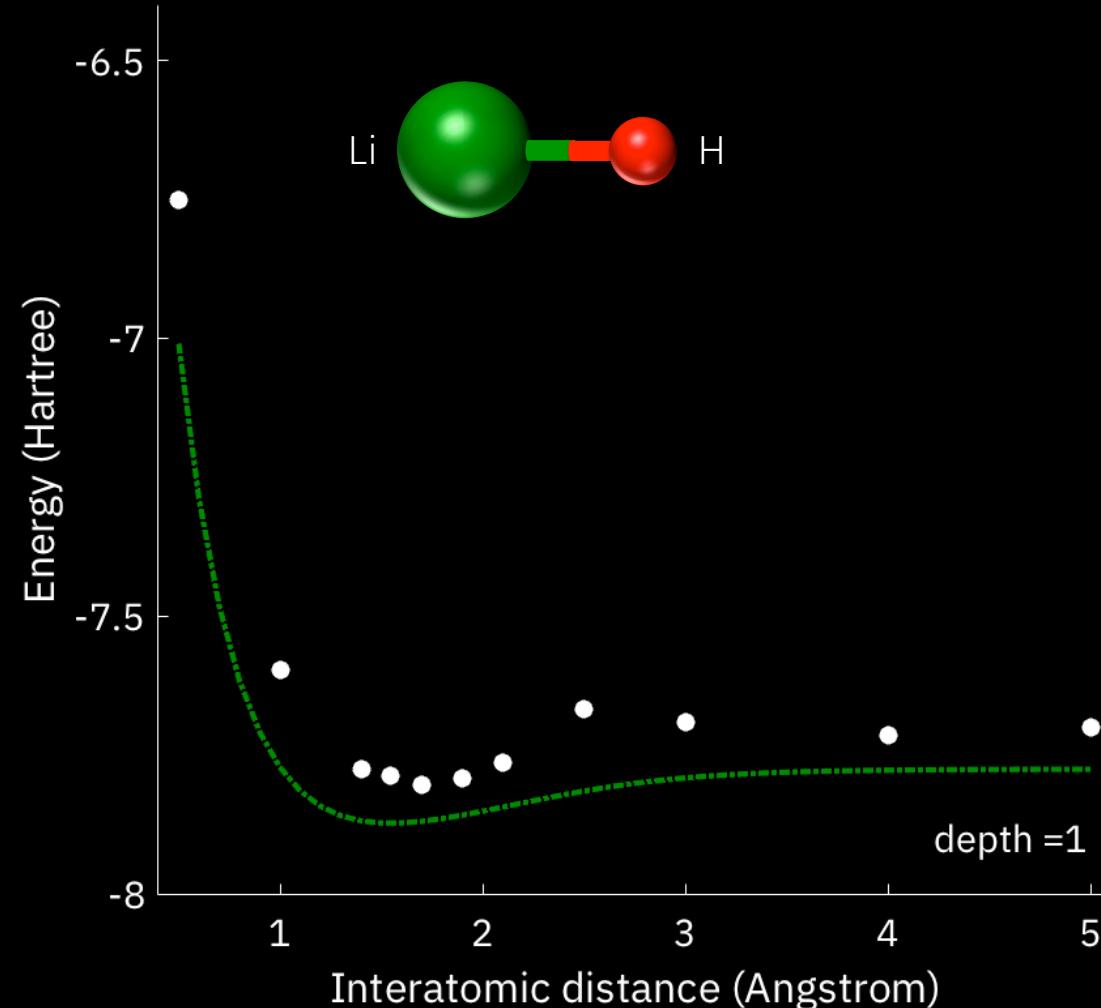


$|0\rangle$

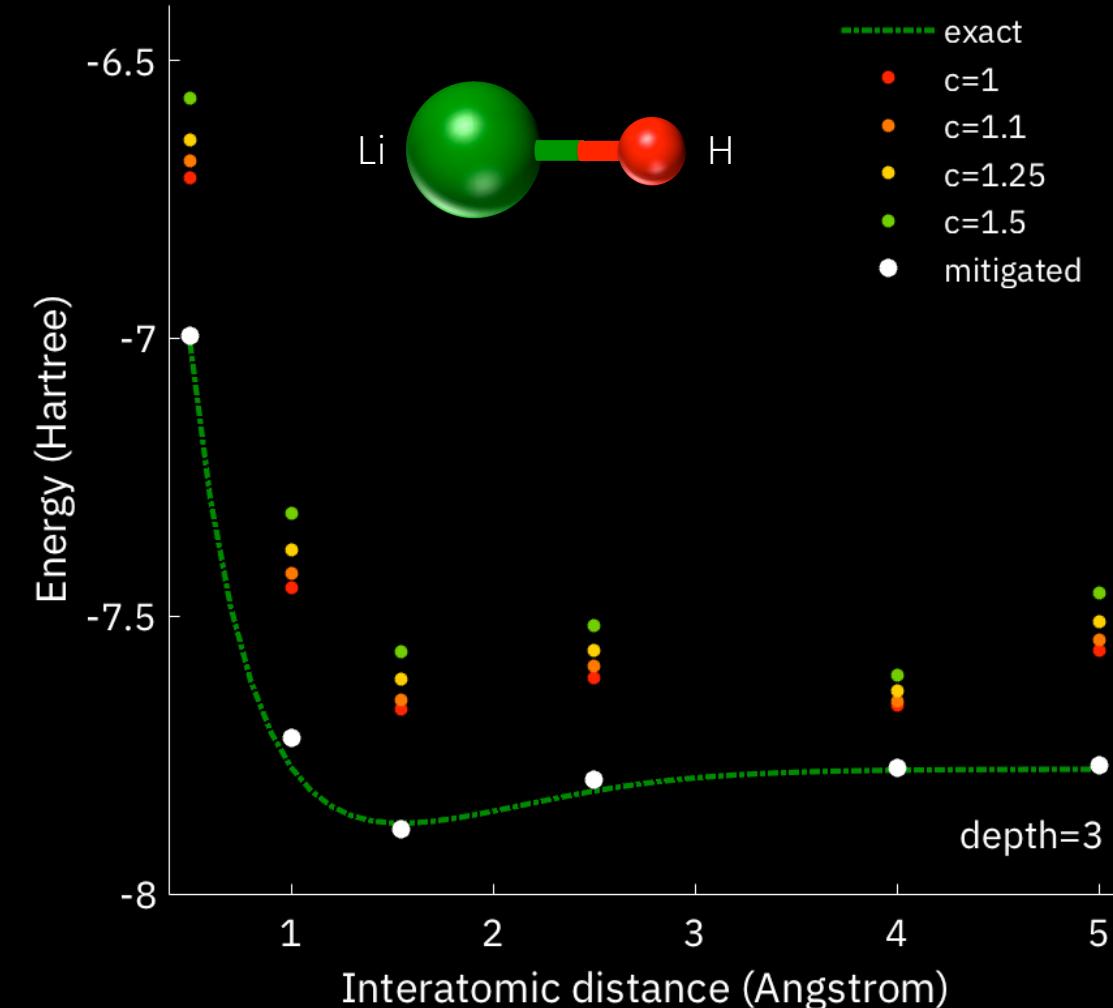
$|0\rangle$

$|0\rangle$

Error Mitigation in molecular simulation: 4Q LiH

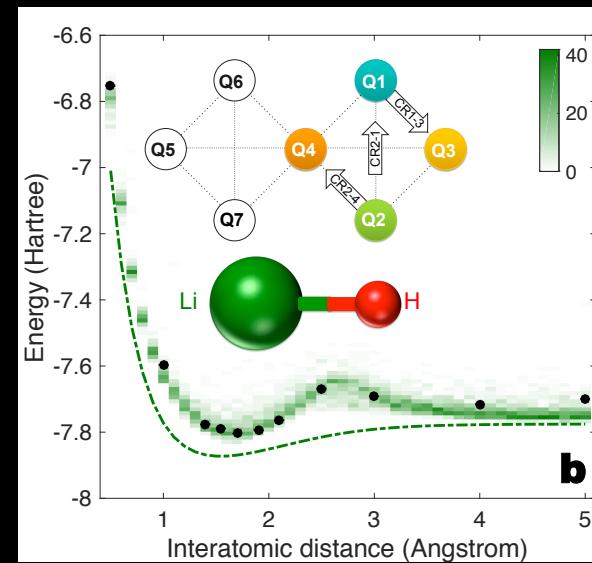
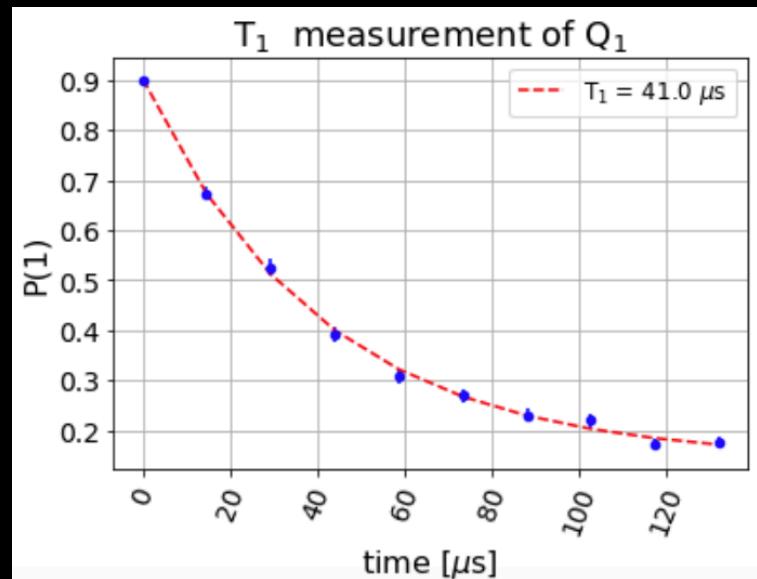
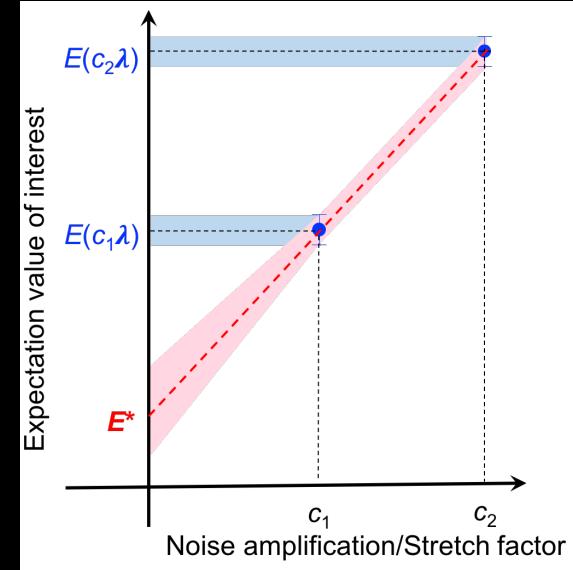
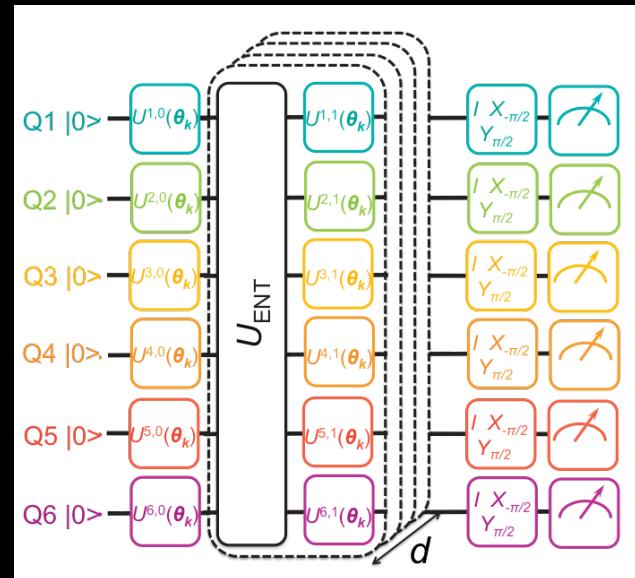
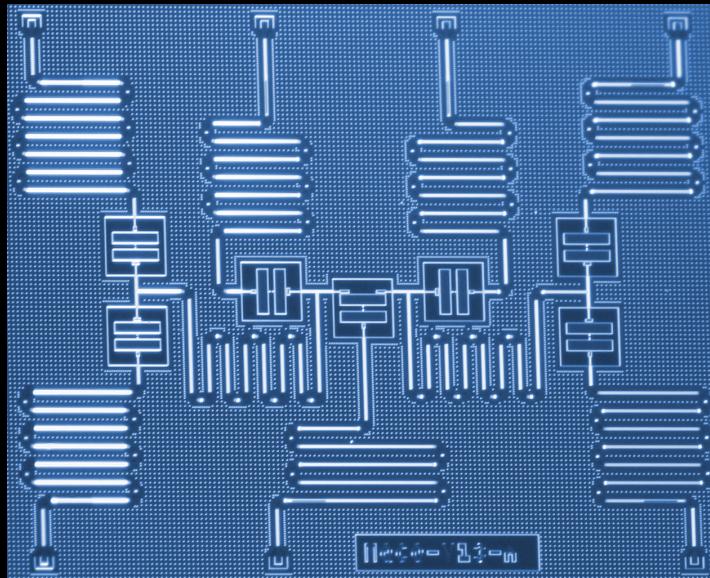


Kandala et al
Nature **549**, 242–246 (2017)



Kandala et al
Nature **567**, 491–495 (2019)

What we've covered today



Techniques go beyond just chemistry ...

A photograph of a modern, multi-story building at night. The building features a long section with floor-to-ceiling glass windows, through which the interior rooms are brightly lit. Above this glass section is a layer of stone walls. The sky is dark blue, indicating it's nighttime. In the foreground, there's a grassy area and a set of steps leading up to the building.

THANK YOU