

2 Practical et theoretical challenges of

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Practical :

- In general, non-Markovian \rightarrow difficult to simulate

Theoretical

- Non-semimartingale \rightarrow can not apply Itô formula & stochastic calculus.

$$\text{N3. FBM = centered Gaussian } \mathbb{E}(W_t^H W_s^H) = \frac{1}{2}(|t|^{2H} + |s|^{2H} - |t-s|^{2H})$$

2. Problem

$$dP_t = P_t \sqrt{V_t} dB_t$$

$$V_t = V_0 + \int_0^t K_H(t-s) \theta ds + \int_0^t K_H(t-s) \xi \sqrt{V_s} dW_s \quad V_0, \theta, \xi \geq 0$$

$$K_H(t) = t^{H-\frac{1}{2}}$$

$$B = \rho W + \sqrt{1-\rho^2} W'$$

1) $\rho < 0$

2) Why $H > 0$?

If $H \leq 0$, we lose integrability of $\int K_H(t)^2 dt < \infty$ that guarantees the existence and all the theoretical results we had obtained

3) Rough Heston in practice

- Standard Heston \rightarrow characteristic function \rightarrow Fourier pricing
- Volterra Heston: combines possibility to find LF with rough models.

2.1. Reversionary Heston

$$4) \xi > 0 \quad K_{H,\xi}(t) = (t+\xi)^{H-\frac{1}{2}}, t > 0$$

$$dP_t^\xi = P_t^\xi \sqrt{V_t^\xi} dB_t$$

$$V_t^\xi = V_0 + \int_0^t K_{H,\xi}(t-s) \theta ds + \int_0^t K_{H,\xi}(t-s) \xi \sqrt{V_s^\xi} dW_s$$

$$U_t^\varepsilon = U_0 + \theta \cdot (K_{H,\varepsilon} * \mathbf{1})(t) + \xi (K_{H,\varepsilon} * \sqrt{U_0^\varepsilon} dW_t)$$

Prop $K \in C^1([0, T]) \rightarrow U_t^\varepsilon$ est une semimartingale

$$K_{H,\varepsilon}(t) = K(0) + (\mathbf{1} * K'_{H,\varepsilon})(t) \quad K_{H,\varepsilon}(0) = \varepsilon^{H-\frac{1}{2}}$$

$$U_t^\varepsilon = U_0 + \theta K_{H,\varepsilon}(0) (\mathbf{1} * dt)(t) + \theta \cdot \mathbf{1} * (K'_{H,\varepsilon} * dt)(t) + \xi \cdot K(0) \cdot (\mathbf{1} * \sqrt{U_0^\varepsilon} dW_t)(t) + \xi \cdot \mathbf{1} * (K' * \sqrt{U_0^\varepsilon} dW_t)(t)$$

$$dU_t^\varepsilon = \underline{\theta K_{H,\varepsilon}(0) dt} + \xi \left(\int_0^t K'_{H,\varepsilon}(t-s) \sqrt{U_s^\varepsilon} dW_s \right) dt + \theta \left(\int_0^t K'_{H,\varepsilon}(t-s) ds \right) dt + \xi K(0) \sqrt{U_0^\varepsilon} dW_t$$

We need $K_{H,\varepsilon}(0) < \infty$ ✓ We need $K'_{H,\varepsilon} \in L^2$ ✓

$$K'_{H,\varepsilon} = (H - \frac{1}{2})(t+\varepsilon)^{H-\frac{3}{2}} \in L^2$$

5) L_ε resolvent of the first kind. $K_{H,\varepsilon} * L_\varepsilon = 1$

$$L_\varepsilon(dt) = \frac{s_0(dt)}{K_{H,\varepsilon}(0)} + l_\varepsilon(t) dt$$

locally integrable such that: $t \mapsto (K'_{H,\varepsilon} * L_\varepsilon)(t) \in C'$

$$\text{Show that } dU_t^\varepsilon = (\varepsilon^{H-\frac{1}{2}} \theta - \left(\frac{1}{2} - H\right) \varepsilon^{-1} (U_t^\varepsilon - U_0) + \int_0^t (K'_{H,\varepsilon} * L_\varepsilon)'(t-s) (U_s^\varepsilon - U_0) ds) dt + \varepsilon^{H-\frac{1}{2}} \xi \sqrt{U_t^\varepsilon} dW_t$$

By (2), we should show that

$$\xi \int_0^t K'_{H,\varepsilon}(t-s) \sqrt{U_s^\varepsilon} dW_s + \theta \int_0^t K'_{H,\varepsilon}(t-s) ds = - \left(\frac{1}{2} - H\right) \varepsilon^{-1} (U_t^\varepsilon - U_0) + \int_0^t (K'_{H,\varepsilon} * L_\varepsilon)'(t-s) (U_s^\varepsilon - U_0) ds$$

$$K'_{H,\varepsilon} * L = K'_{H,\varepsilon} * L(0) + \mathbf{1} * (K'_{H,\varepsilon} * L)' = K_{H,\varepsilon}$$

$$K'_{H,\varepsilon} * \mathbf{1} = K'_{H,\varepsilon} * L(0) (\mathbf{1} * K_{H,\varepsilon}) + \mathbf{1} * (K'_{H,\varepsilon} * L)' * K_{H,\varepsilon}$$

$$K'_{H,\varepsilon} = (K'_{H,\varepsilon} * L)(0) * K_{H,\varepsilon} + (K'_{H,\varepsilon} * L)' * K_{H,\varepsilon}$$

ξ

$$K'_{H,\varepsilon} * \theta ds + K'_{H,\varepsilon} * \xi \sqrt{U_0^\varepsilon} dW = (K'_{H,\varepsilon} * L)(0) \underbrace{(K'_{H,\varepsilon} * \theta ds + K'_{H,\varepsilon} * \xi \sqrt{U_0^\varepsilon} dW)}_{(K'_{H,\varepsilon} * L)(0)} + (K'_{H,\varepsilon} * L)' * (U_t^\varepsilon - U_0)$$

$$(K'_{H,\varepsilon} * L)(0) = \int_0^0 K'_{H,\varepsilon}(-s) L_\varepsilon(ds) = K'_{H,\varepsilon}(0) / K_{H,\varepsilon}(0) = \frac{(H - \frac{1}{2}) \varepsilon^{H-\frac{3}{2}}}{\varepsilon^{H-\frac{1}{2}}} = (H - \frac{1}{2}) \varepsilon^{-1}$$

(6)

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$$dV_t^\varepsilon = \underbrace{\left(\varepsilon^{H-\frac{1}{2}} \mathcal{G} - \left(\frac{1}{2} - H \right) \varepsilon' (V_t^\varepsilon V_0) + \int_0^t (\kappa_{H,\varepsilon}' * b_\varepsilon)'(t-s) (V_s^\varepsilon - V_0) ds \right) dt}_{\text{Bro scaling}} + \underbrace{\varepsilon^{H-\frac{1}{2}} \varepsilon \sqrt{V_t^\varepsilon} dW_t}_{\text{Bro}}$$

Not Bro: path-dependent
Drop him!

$$\left\{ \begin{array}{l} dV_t^\varepsilon = \left(\varepsilon^{H-\frac{1}{2}} \mathcal{G} - \varepsilon' (V_t^\varepsilon - V_0) \right) dt + \varepsilon^{H-\frac{1}{2}} \varepsilon \sqrt{V_t^\varepsilon} dW_t \\ dS_t^\varepsilon = S_t^\varepsilon \sqrt{V_t^\varepsilon} dB_t \end{array} \right.$$

Bro, cool Markovian model approximating Volterra semimartingale

with shifted power-law kernel. $\varepsilon \downarrow 0$ (kernel \rightarrow power law) \rightarrow
 $\varepsilon^{-\frac{1}{2}}, \varepsilon^{H-\frac{1}{2}} \rightarrow \infty$ mean-reversion and vol-of-vol explode

2.2 Jumps at the limit

(7) $u \in \mathbb{R}$. Sketch of proof for the characteristic function

$$\mathbb{E}[e^{iu \cdot \log \frac{S_T^\varepsilon}{S_t^\varepsilon}} | \mathcal{F}_t] = \exp \left\{ \varphi_\varepsilon(T-t) + \varepsilon^{\frac{1}{2}-H} \Psi_\varepsilon(T-t) V_t^\varepsilon \right\}, \quad t \leq T \quad (*)$$

$$\mathbb{E}[e^{iu \cdot \log S_T^\varepsilon} | \mathcal{F}_t] = \exp \left\{ \varphi_\varepsilon(T-t) + \varepsilon^{\frac{1}{2}-H} \Psi_\varepsilon(T-t) V_t^\varepsilon + iu \cdot \log S_t^\varepsilon \right\} = M_t$$

Terminal condition $M_T = e^{iu \cdot \log S_T^\varepsilon} \xrightarrow{Y_t} \begin{cases} \varphi_\varepsilon(0) = 0 \\ \Psi_\varepsilon(0) = 0 \end{cases}$

If (*) holds, M_t is a martingale $M_t = e^{Y_t}$

$$\frac{dM_t}{M_t} = dY_t + \frac{1}{2} d\langle Y \rangle_t \quad \text{"dt"-part should be 0.}$$

$$dY_t = \left[-\dot{\varphi}_\varepsilon(T-t) - \varepsilon^{\frac{1}{2}-H} \dot{\Psi}_\varepsilon(T-t) \underline{V_t^\varepsilon} + \varepsilon^{\frac{1}{2}-H} \Psi_\varepsilon(T-t) \left(\varepsilon^{H-\frac{1}{2}} \mathcal{G} - \varepsilon' (V_t^\varepsilon - V_0) \right) - iu \cdot \frac{1}{2} V_t^\varepsilon + \frac{1}{2} \left[\varepsilon^{\frac{1}{2}-H} \cdot \Psi_\varepsilon(T-t)^2 \cdot \cancel{\varepsilon^{\frac{1}{2}-H}} \cdot \underline{\varepsilon \cdot V_t^\varepsilon} - u^2 \cdot \underline{V_t^\varepsilon} + iu \cancel{\varepsilon^{\frac{1}{2}-H}} \cdot \cancel{\Psi_\varepsilon(T-t)} \cdot \cancel{\varepsilon^{H-\frac{1}{2}}} \cdot \cancel{\varepsilon \cdot V_t^\varepsilon} \right] \right] dt + [\dots] dW_t$$

$$(\text{ODEs}) \quad \begin{cases} \dot{\varphi}_\varepsilon = \varphi_\varepsilon \left(\mathcal{G} + \varepsilon^{H-\frac{1}{2}} V_0 \right) \\ \varphi_\varepsilon(0) = 0 \end{cases} \quad \begin{cases} \varepsilon^{\frac{1}{2}-H} \dot{\Psi}_\varepsilon = -\Psi_\varepsilon \cdot \varepsilon^{\frac{1}{2}-H} + \frac{1}{2} \varepsilon^2 \cdot \Psi_\varepsilon^2 + \frac{1}{2} \delta \cdot iu \cdot \varepsilon \cdot \Psi_\varepsilon - \frac{1}{2} (u^2 + iu) \\ \Psi_\varepsilon(0) = 0 \end{cases}$$

We showed that (*) \rightarrow (ODEs)

For the inverse result, it is sufficient to provide the existence of the solution to (ODEs) and guarantee that M_t is a martingale.

This will imply (*).

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(8) We accept that

$$\begin{aligned}\Psi_\epsilon(t) &= (\epsilon^{-H-\frac{1}{2}}\theta + \epsilon^{\frac{1-H}{2}}V_0) \underbrace{\epsilon^{-2}((1-i\rho\epsilon^{\frac{H+1}{2}}\xi u - d)t - 2\epsilon \log(\frac{1-\rho\xi}{1+\rho\xi}))}_{\text{one can show}} e^{\frac{-\epsilon^2 t^2 d}{1-\rho^2 \xi^2}} \\ \Psi_\epsilon(1) &= \epsilon^{-H-\frac{1}{2}} \underbrace{\epsilon^{-2}(1-i\rho\epsilon^{\frac{H+1}{2}}\xi u - d)}_{\frac{1-\rho^2 \xi^2}{1-\rho\xi}} \underbrace{\frac{t-\epsilon^{-\frac{1}{2}}d}{1-\rho\xi}}_{\text{that } \Psi_\epsilon, \Psi_\epsilon' \leq 0.}\end{aligned}$$

$$g := \dots, d := \dots \quad \Re(d) > 0$$

$$\text{For } H = -\frac{1}{2}$$

$$\lim_{\epsilon \downarrow 0} \underbrace{\mathbb{E}[e^{iu \log \frac{\xi_t}{\xi_0}}]}_{= e^{\Psi_\epsilon(t) + \epsilon^{\frac{1}{2}-H} \Psi_\epsilon'(t) V_0}} = \exp\{\gamma(u) \cdot T\}$$

$$= e^{\Psi_\epsilon(T) + \epsilon^{\frac{1}{2}-H} \Psi_\epsilon'(T) V_0} \rightarrow 0$$

$$d = \sqrt{(1-i\rho\xi u)^2 + \xi^2 \left(\frac{u^2 + iu}{2}\right)}$$

$$\begin{aligned}\epsilon^{\frac{1}{2}-H} \Psi_\epsilon(T) &= \underbrace{\epsilon^{2H}}_{\epsilon} \underbrace{\xi^{-2}(1-i\rho - d)}_{\rightarrow 0} \underbrace{\frac{1-\epsilon^{-\frac{1}{2}}d}{1-\rho\xi e^{-\frac{1}{2}}d}}_{\rightarrow 0} \rightarrow 0 \quad g = \frac{1-d}{1+d} \\ \Psi_\epsilon(T) \Big|_{H=-\frac{1}{2}} &= (\theta + V_0) \underbrace{\xi^{-2}((1-i\rho\xi u - d)T - 2\epsilon \log(\frac{1-\rho\xi}{1+\rho\xi}))}_{\rightarrow 0} \rightarrow \\ &\rightarrow (\theta + V_0) \underbrace{\xi^{-2}(1-i\rho\xi u - d)T}_{\eta(u)}\end{aligned}$$

$$\text{where } \eta(u) = (\theta + V_0) \xi^{-2} \left(1 - i\rho\xi u - \sqrt{(1-i\rho\xi u)^2 + \xi^2 \left(\frac{u^2 + iu}{2}\right)}\right)$$

We can identify $\alpha, \beta, S, \zeta, \gamma$, but we do not want...

(9) (X_t) a mysterious process if $X_0=0$ and it has independent stationary increments s.t. $\mathbb{E}[e^{iuX_1}] = e^{\eta(u)}$

We have shown that $\mathbb{E}[e^{iu \log \frac{\xi_t}{\xi_0}}] \xrightarrow{\epsilon \downarrow 0} e^{\eta(u)}$ for $H = -\frac{1}{2} \rightarrow$

\rightarrow by continuity theorem $\log \xi_t - \log \xi_0 \xrightarrow{d} X_1$

$$\mathbb{E}[e^{iu \log \frac{\xi_T}{\xi_0}}] \rightarrow e^{\eta(u)T}$$

It's enough to show that $\mathbb{E}[e^{iuX_T}] = e^{\eta(u) \cdot T}$

For $T \in \mathbb{N}$ $X_T = (X_T - X_{T-1}) + (X_{T-1} - X_{T-2}) + \dots + (X_1 - X_0)$ $\xrightarrow{\text{indep \& stationary}} \mathbb{E}[e^{iuX_T}] = e^{\eta(u)T}$

Similar argument for $T = \frac{T}{m}$: $X_T - X_0 = (X_T - X_{T-1}) + (X_{T-1} - X_{T-2}) + \dots + (X_1 - X_0)$

So, it's true for $T \in \mathbb{Q}$. By continuity of X and dominated convergence it's true $\forall T \in \mathbb{R}$.

- (10) ATM skew is exploding (power law?) and there are jumps in
the limiting mysterious process. So, it's a Lévy process with jumps
(Normal inverse gaussian-inverse Gaussian proc.)
- (11) Main messages: after an approximation we get a model that
can describe both the rough behavior and approximate the
Volterra process for HGO, but also can converge to the jump
processes for HGO.