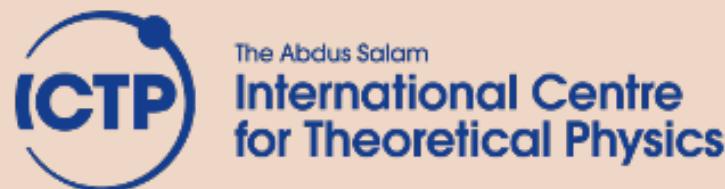




MAX School on Advanced Materials and Molecular Modelling
with QUANTUM ESPRESSO

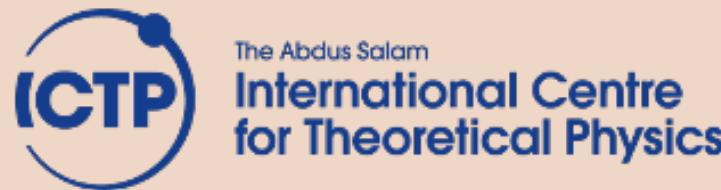
QUANTUM ESPRESSO: overview and basic
functionalities. The self-consistent cycle. PBC:
supercells and k-point sampling

Ralph Gebauer

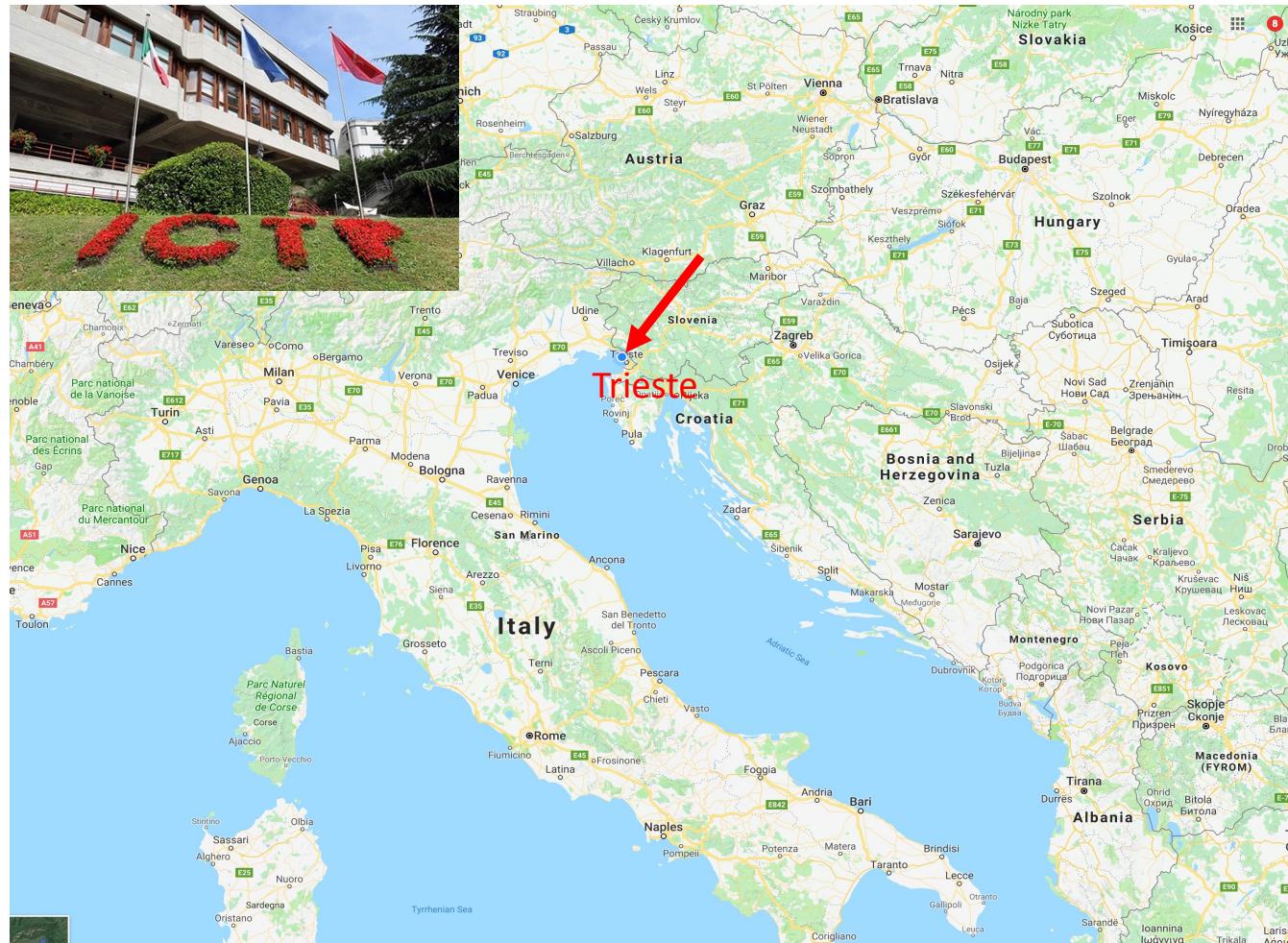


Trieste, May 17th, 2021

This school is a large-scale collaboration:

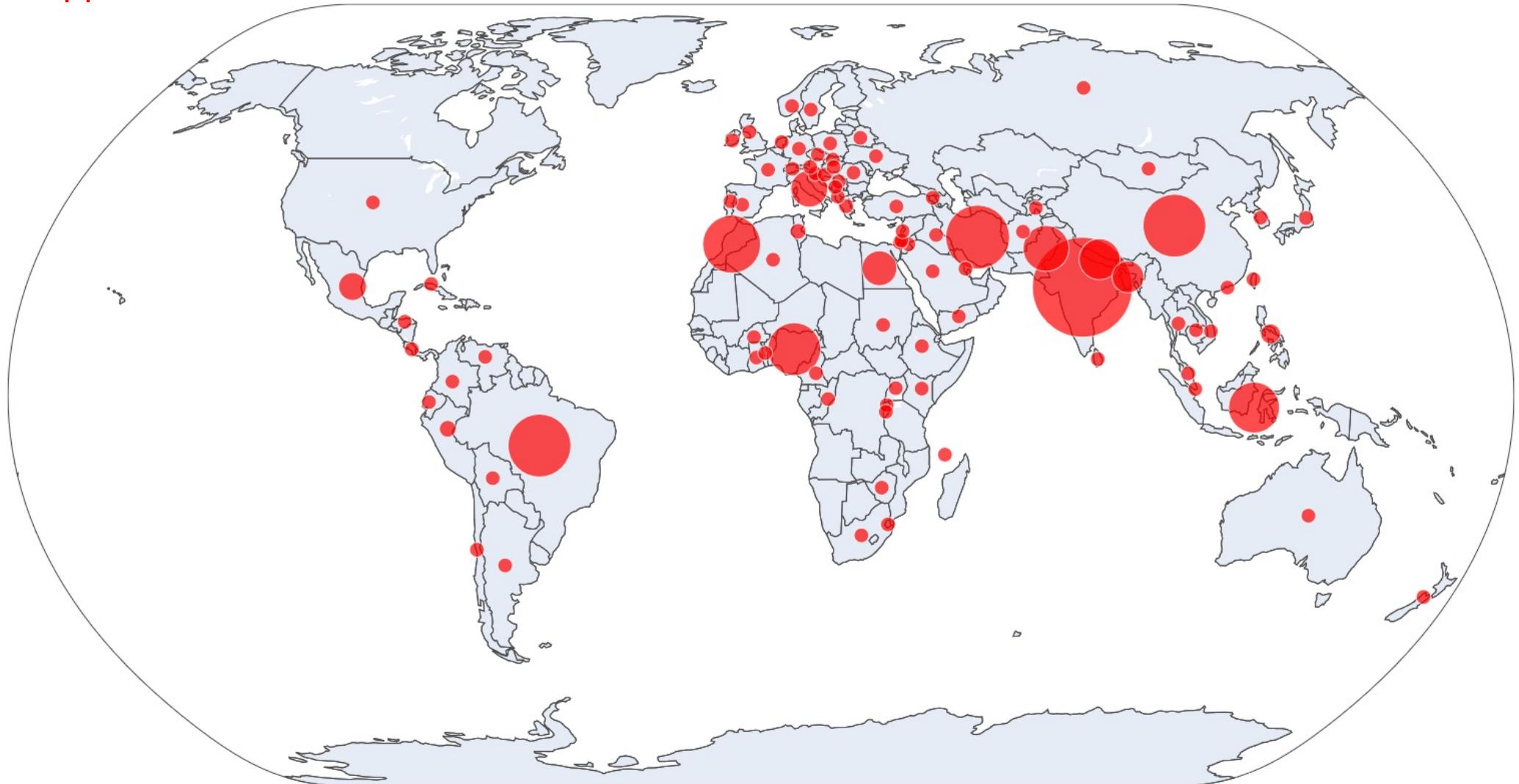






Founded in 1964, the Abdus Salam International Centre for Theoretical Physics (ICTP) operates under the aegis of two United Nations Agencies: UNESCO and IAEA and is regularized by a seat agreement with the Government of Italy.

1292 applicants!



Special thanks to the secretaries: Monica Ancuta, Viktoriya Lvova and Adriana Pinto

The star of this school:
Quantum-ESPRESSO

Quantum ESPRESSO: Quantum opEn-Source Package for Research in Electronic Structure, Simulation, and Optimization, is a distribution (an integrated suite) of software for first-principle simulations, i.e., atomistic calculations based on electronic structure, using density-functional theory, a plane-wave basis set, pseudopotentials.

QE exists since 2002, resulting from the merge of pre-existing packages; some core components have been under development for ~ 30 years

The main goals of Quantum-ESPRESSO are:

- innovation in theoretical methods and numerical algorithms
- efficiency on modern computer architectures





The star of this school:

Quantum-ESPRESSO License:

QE is distributed under the GNU (Gnu's Not Unix) General Public License (GPL v.2), probably the most common free-software license.

- The source code is available.
- You can do whatever you want with the sources, but if you distribute any derived work, you have to distribute under the GPL the sources of the derived work.

Advantages:

- Everybody – including commercial entities – can contribute.
- Nobody can “steal” the code and give nothing back to the community.

The most successful example is probably the Linux Kernel.



The star of this school:
Quantum-ESPRESSO

You can get QE help, answers and general info from various sources:

the documenting papers:

J. Phys.: Condens. Matter 21, 395502 (2009) and

J. Phys.: Condens. Matter 29, 465901 (2017)

the web site: www.quantum-espresso.org

the documentation in the Doc subdirectories

the developers' portal: github.com/QEF/q-e/releases

the mailing lists:

users@lists.quantum-espresso.org

developers@lists.quantum-espresso.org

(use first the rich ARCHIVE of these lists)



Quantum-ESPRESSO as a distribution

QE is not a single, executable code. It is composed of several packages:

PWscf: self-consistent electronic structure, (variable-cell) structural optimization, molecular dynamics

CP: Car-Parrinello molecular dynamics, also with variable cell

They share a common installation method, input format, pseudopotential format, data output format, large parts of the basic code.

Further codes:

PHphon: linear-response calculations (phonons, dielectric properties)

PostProc: graphical and postprocessing utilities (density of states, STM, etc.)

PWneb: Nudged Elastic Band (NEB) for reaction pathways and barriers

atomic: pseudopotential generation code

PWGui: a Graphical User Interface for production of input files

PWcond: ballistic conductance

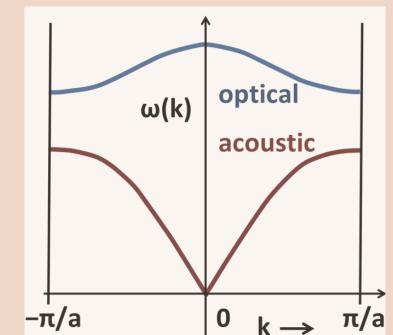
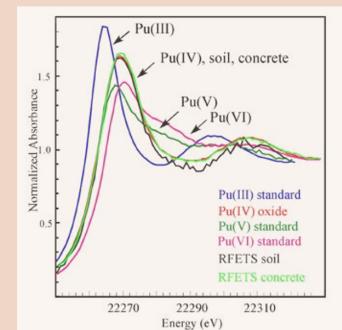
Xspectra: Calculation of X-ray near-edge adsorption spectra (XANES)

GWL: GW band structure with ultralocalized Wannier functions

TD-DFPT: Time-Dependent Density-Functional Perturbation Theory

EPW: Electron-phonon coefficients and related properties

HP: Hubbard parameters from linear response



Density Functional Theory

The ground-state energy of a system of N electrons is a functional of the electronic density $n(\mathbf{r})$:

$$E^{DFT} [n] \quad n(\mathbf{r}) \geq 0 \quad \int d\mathbf{r} n(\mathbf{r}) = N$$

The exact functional is unknown. Within the Kohn-Sham (KS) formalism, we write:

$$E^{DFT} [n] = T_s [\{\psi_i\}] + E_{ext}[n] + E_{Har}[n] + E_{xc}[n] + E_{Ions}$$

Where the KS orbitals have been introduced to approximate the kinetic energy:

$$T_s [\{\psi_i\}] = -\frac{\hbar^2}{2m} \sum_{i=1}^N \int d\mathbf{r} \psi_i^*(\mathbf{r}) \nabla^2 \psi_i(\mathbf{r})$$

$$\langle \psi_i | \psi_j \rangle = \delta_{i,j} \quad i, j = 1 \cdots N$$

$$n(\mathbf{r}) = \sum_{i=1}^N \psi_i^*(\mathbf{r}) \psi_i(\mathbf{r})$$

$$E_{ext} [n] = \int d\mathbf{r} n(\mathbf{r}) V_{ext}(\mathbf{r})$$

$$E_{Har} [n] = \frac{e^2}{2} \int d\mathbf{r} d\mathbf{r}' \frac{n(\mathbf{r}) n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

$$E_{Ions} = \sum_{IJ} \frac{Z_I Z_J e^2}{|\mathbf{R}_I - \mathbf{R}_J|}$$

The main difficulty lies in approximating the exchange-correlation functional $E_{xc}[n]$

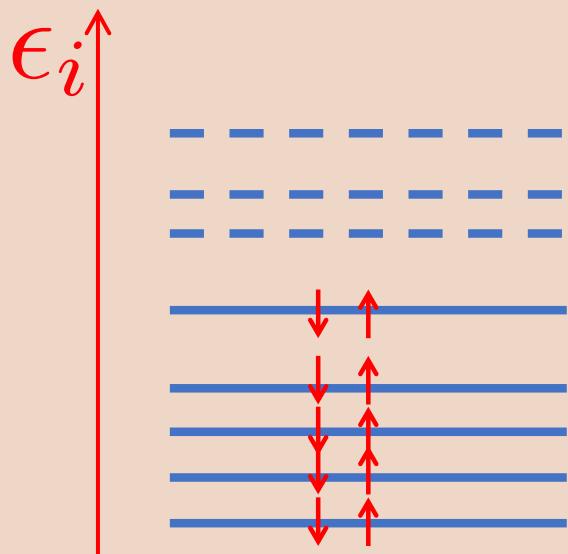
Density Functional Theory

Minimization of $E^{\text{DFT}}[n]$ with respect to n (actually, with respect to the KS orbitals) leads to the KS equations:

$$H^{KS} \psi_i(\mathbf{r}) = \left(-\frac{\hbar^2}{2m} \nabla^2 + V_{ext}(\mathbf{r}) + V_{Har}(\mathbf{r}) + V_{xc}(\mathbf{r}) \right) \psi_i(\mathbf{r}) = \epsilon_i \psi_i(\mathbf{r})$$

$$V_{xc}(\mathbf{r}) = \frac{\delta E_{xc}}{\delta n(\mathbf{r})}$$

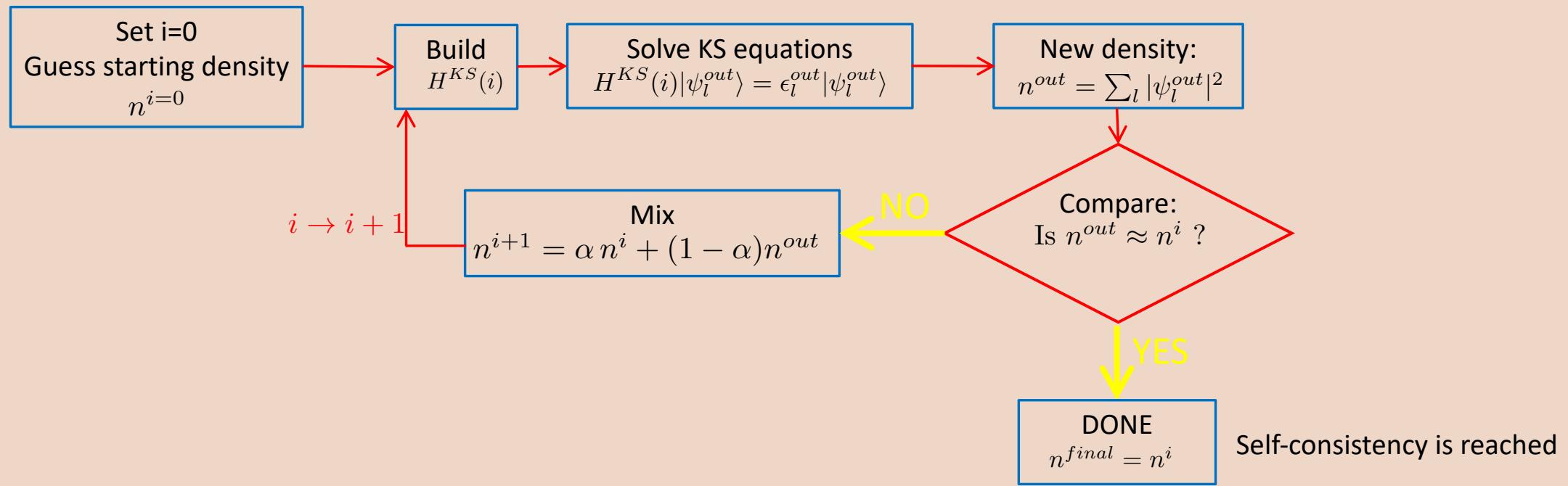
$$V_{Har}(\mathbf{r}) = \frac{\delta E_{Har}}{\delta n(\mathbf{r})} = e^2 \int d\mathbf{r}' \frac{n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$



Problem: H^{KS} depends on the charge density

Density Functional Theory

The self-consistent cycle



The self-consistent cycle

Initial potential from superposition of free atoms

starting charge 29.42246, renormalised to 30.00000

negative rho (up, down): 7.551E-05 0.000E+00

Starting wfcs are 30 randomized atomic wfcs

total cpu time spent up to now is 1.5 secs

Self-consistent Calculation

iteration # 1 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 1.00E-02, avg # of iterations = 3.0

negative rho (up, down): 2.325E-05 0.000E+00

total cpu time spent up to now is 2.4 secs

total energy = -75.13217408 Ry

estimated scf accuracy < 1.92237710 Ry

iteration # 2 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 6.41E-03, avg # of iterations = 2.0

negative rho (up, down): 1.858E-05 0.000E+00

total cpu time spent up to now is 3.1 secs

total energy = -75.22539882 Ry

estimated scf accuracy < 0.08211827 Ry

iteration # 3 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 2.74E-04, avg # of iterations = 6.0

negative rho (up, down): 5.570E-05 0.000E+00

iteration # 10 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 6.88E-09, avg # of iterations = 7.0

negative rho (up, down): 3.572E-04 0.000E+00

total cpu time spent up to now is 9.8 secs

total energy = -75.28039726 Ry

estimated scf accuracy < 0.00000156 Ry

iteration # 11 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 5.20E-09, avg # of iterations = 3.0

negative rho (up, down): 3.554E-04 0.000E+00

total cpu time spent up to now is 10.5 secs

End of self-consistent calculation

k = 0.0000 0.0000 0.0000 (4318 PWs) bands (ev):

21.1685 -18.3817 -18.3806 -14.7643 -14.7544 -12.8025 -11.1007 -10.8244

10.1315 -10.1240 -8.9981 -8.1363 -8.1213 -6.2696 -6.2696 -1.1442

highest occupied, lowest unoccupied level (ev): -6.2696 -1.1442

total energy = -75.28039766 Ry

estimated scf accuracy < 0.00000008 Ry

The total energy is the sum of the following terms:

one-electron contribution = -518.40332837 Ry

hartree contribution = 262.27253441 Ry

xc contribution = -25.13397296 Ry

ewald contribution = 205.98436926 Ry

convergence has been achieved in 11 iterations

Representing charge densities, KS orbitals in a computer:

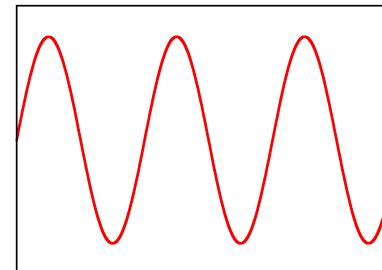
Basis sets

$$f(\mathbf{r}) = \sum_{\alpha=1}^M c_{\alpha} b_{\alpha}(\mathbf{r})$$

Function to be represented Expansion coefficient Basis function Size of basis set

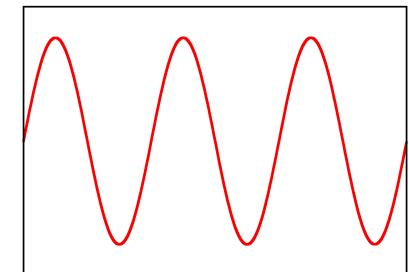
In Quantum-ESPRESSO, plane waves are used as basis set:

$$b_{\alpha}(\mathbf{r}) = \frac{1}{\sqrt{V}} \exp(i \mathbf{G}_{\alpha} \cdot \mathbf{r})$$



Advantages of plane waves as basis set

$$b_\alpha(\mathbf{r}) = \frac{1}{\sqrt{V}} \exp(i \mathbf{G}_\alpha \cdot \mathbf{r})$$



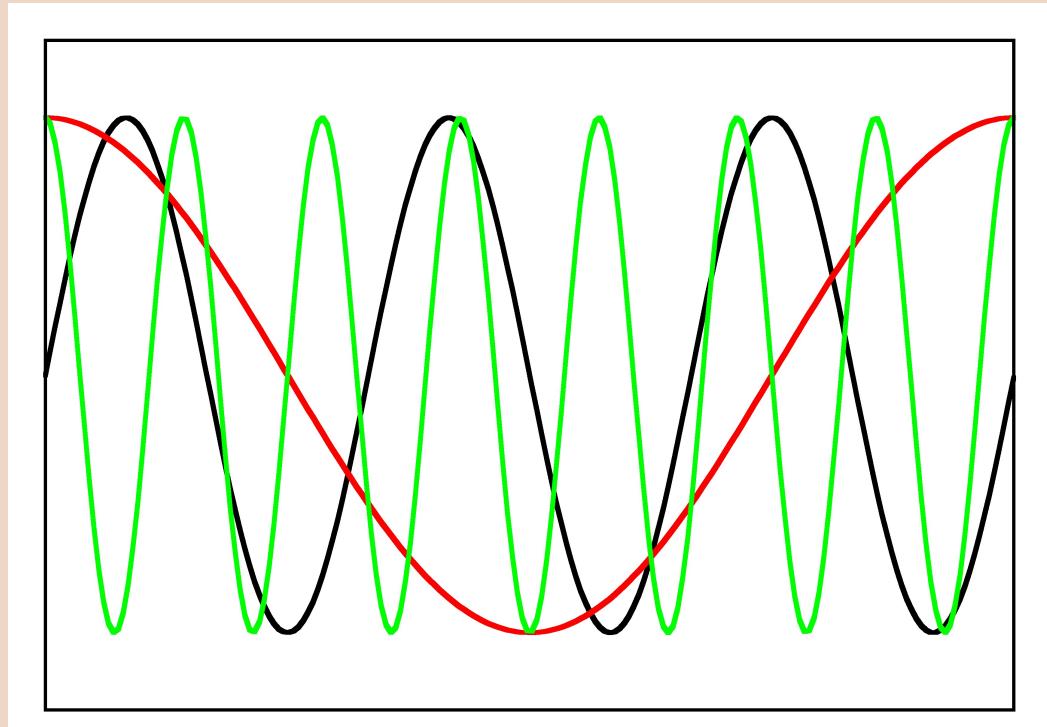
Using plane-waves as a basis set has many advantages:

- Simple analytical form: derivatives and integrals are easy to perform
- PWs are orthonormal
- Unbiased: no assumptions where charges/electrons are localized
- Independent of atomic positions: no “Pulay forces”
- Easy to control convergence of basis set size (see following slides ...)
- Straightforward use of FFTs: easy use of R- and G-space dualities

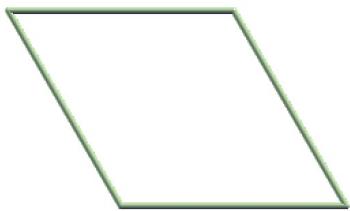
However: PW basis set are typically MUCH bigger than other sets of
(localized) basis functions

Periodicity

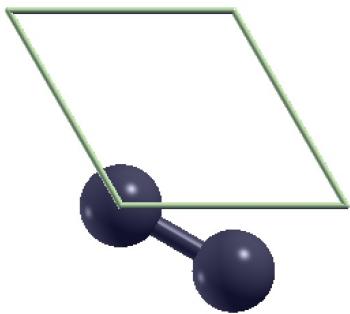
Due to the periodic nature of PWs, the use us this basis set is closely linked with the periodicity of the physical system.



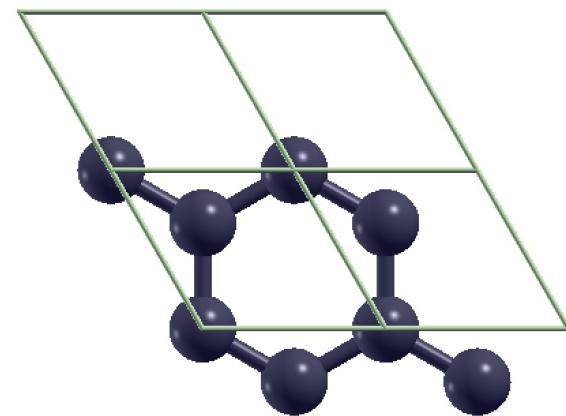
Periodicity



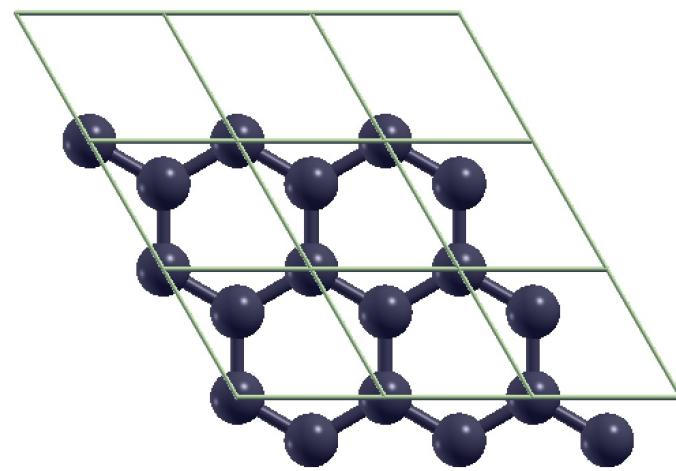
Periodicity



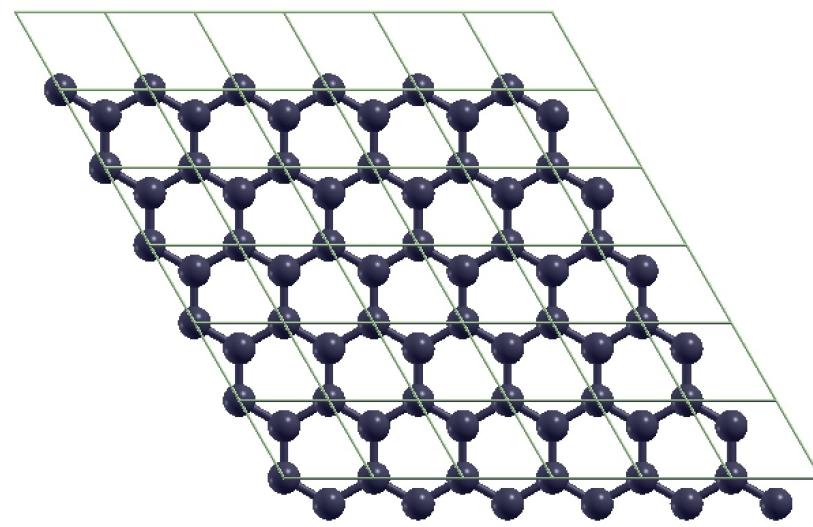
Periodicity



Periodicity



Periodicity



Periodicity

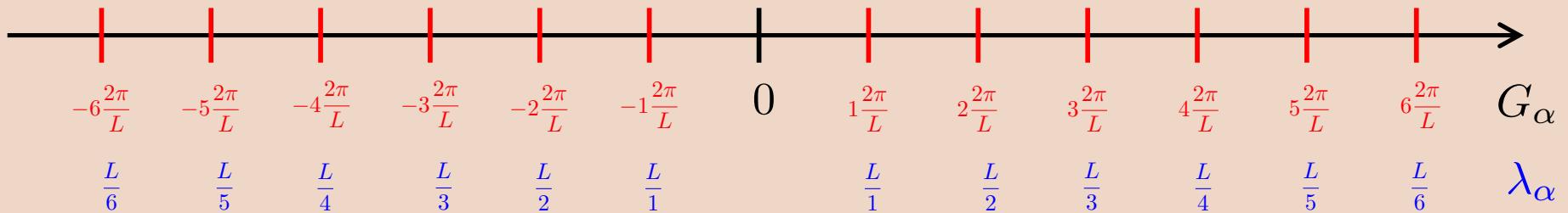
If a function is periodic in real space, its Fourier transform is non-zero only for discrete wavevectors

$$\text{In 1-D: } \exp(iG_\alpha x) \rightarrow G_\alpha = n \times \frac{2\pi}{L} \quad n \in \mathbb{Z}$$

$$\text{In 3-D: } \exp(i\mathbf{G}_\alpha \cdot \mathbf{r}) \rightarrow \mathbf{G}_\alpha = m \times \mathbf{B}_1 + n \times \mathbf{B}_2 + p \times \mathbf{B}_3 \quad m, n, p \in \mathbb{Z}$$

(where $\mathbf{B}_{1,2,3}$ are reciprocal lattice vectors)

1-D illustration:



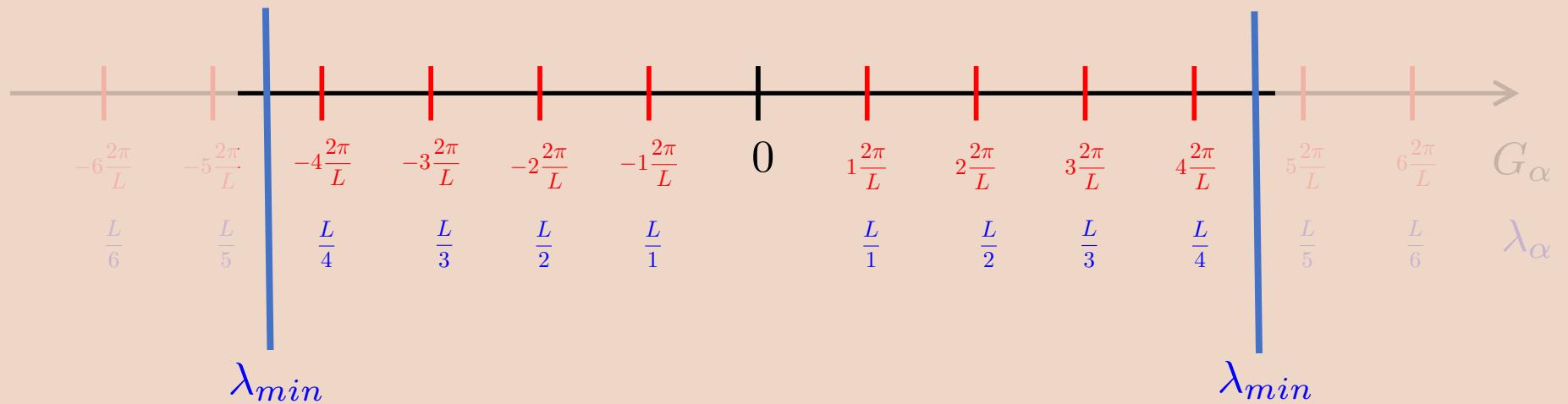
Wavelength λ_α is inversely proportional to $|\mathbf{G}_\alpha|$

Periodicity

The number of PWs compatible with a given periodicity is **infinite**.

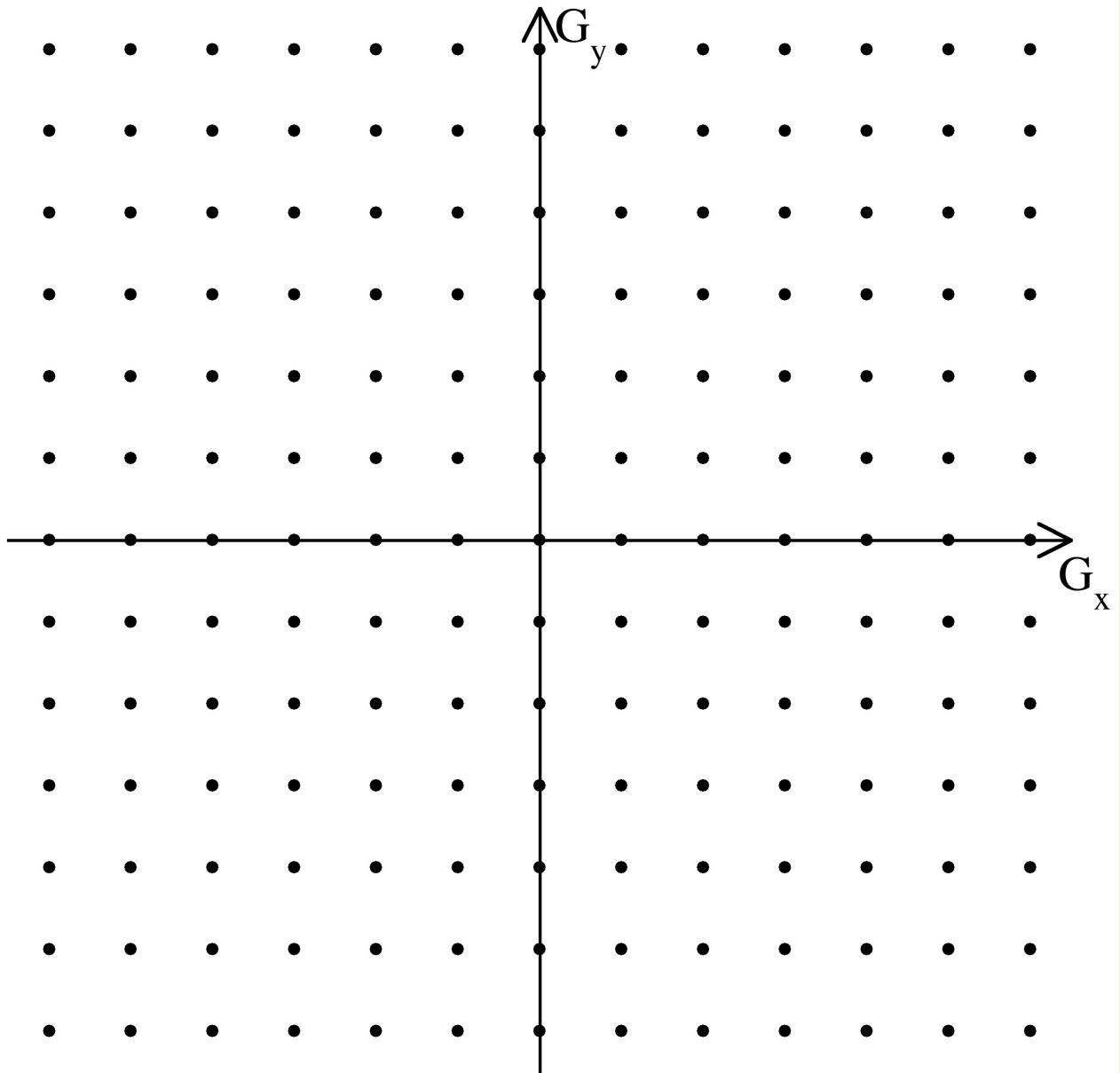
In order to obtain a **finite** number of basis functions, one fixes a “smallest” feature size: λ_{min}

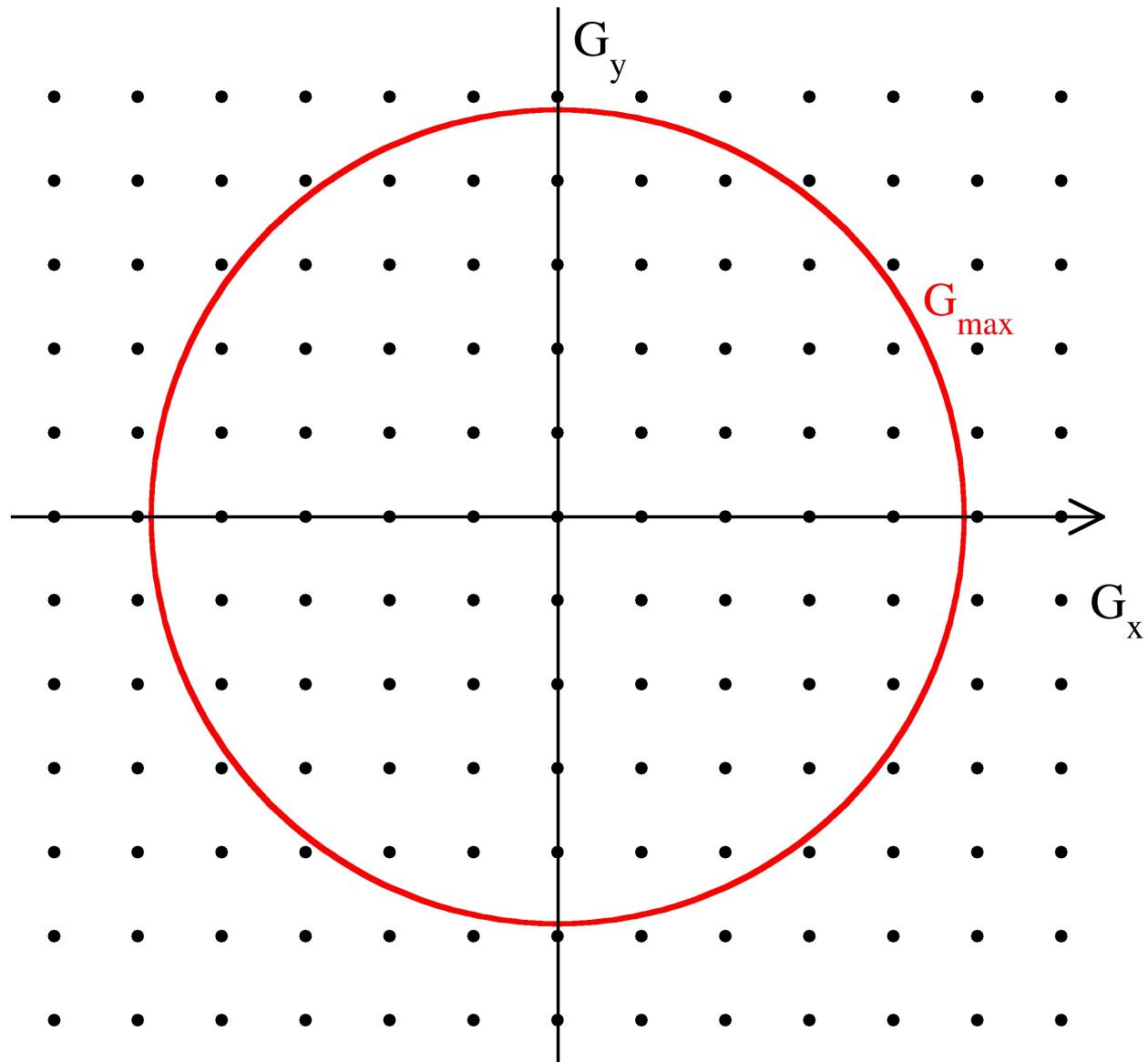
1-D illustration:

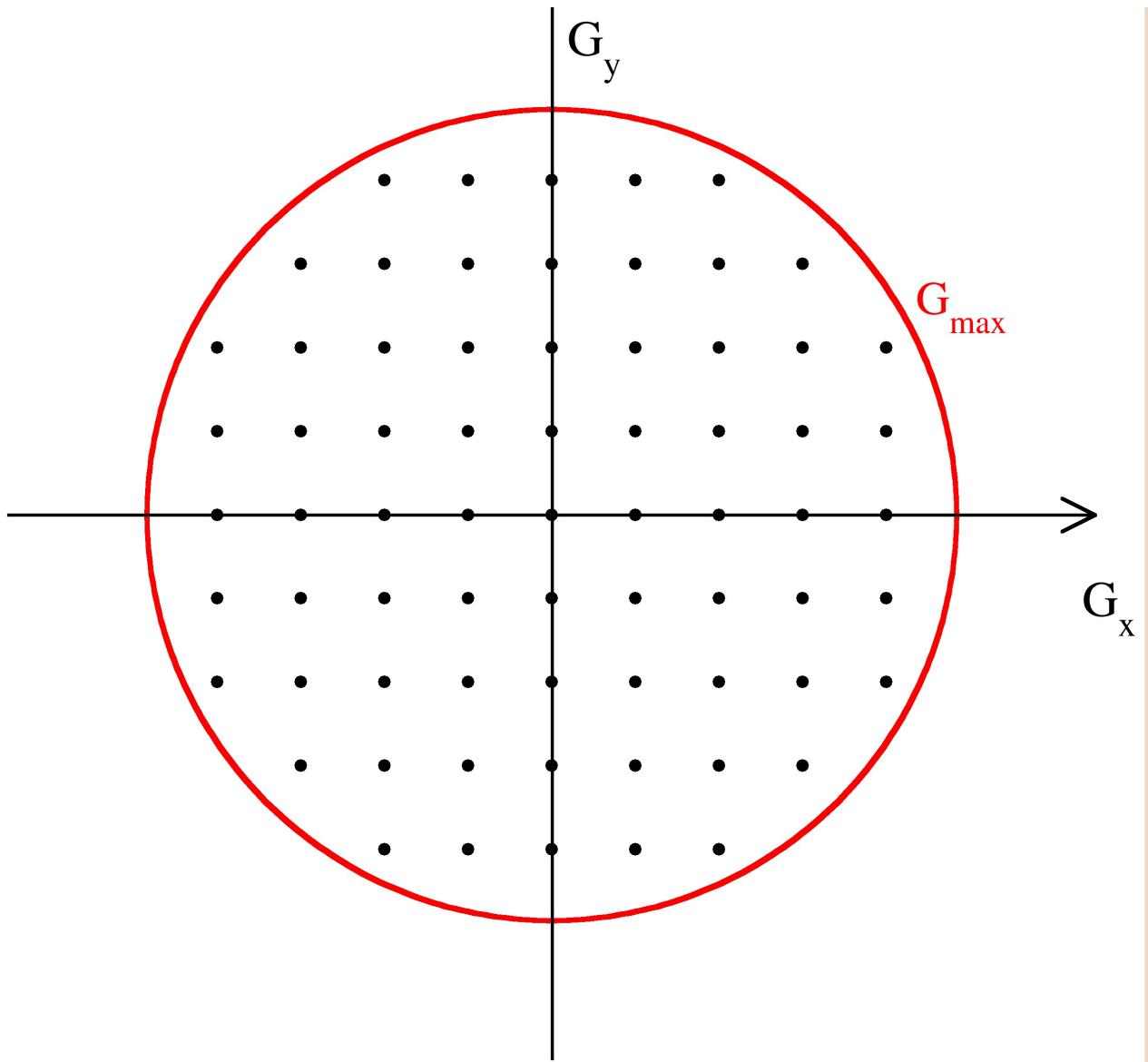


In practice, one describes a **maximal norm** for the wavevector: $|\mathbf{G}|_{max} = \frac{2\pi}{\lambda_{min}}$

By setting a **cutoff kinetic energy** for the plane-waves: $E_{cut} = \frac{\hbar^2}{2m} |\mathbf{G}|_{max}^2$







From KS orbitals to the charge density

Imagine that orbitals are represented in PWs, with a cutoff G_{\max}

$$\psi_l(\mathbf{r}) = \sum_{|\mathbf{G}| < G_{\max}} c(\mathbf{G}) \exp(i\mathbf{G} \cdot \mathbf{r})$$

The charge density if given by the squared modulus of the orbitals:

$$n(\mathbf{r}) = \sum_l \psi_l^*(\mathbf{r}) \psi_l(\mathbf{r})$$
$$\tilde{n}(\mathbf{G}) = \sum_l \sum_{|\mathbf{G}'| < G_{\max}} \tilde{\psi}_l^*(\mathbf{G}') \tilde{\psi}_l(\mathbf{G} - \mathbf{G}')$$

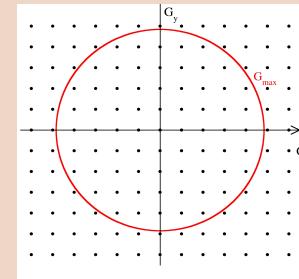
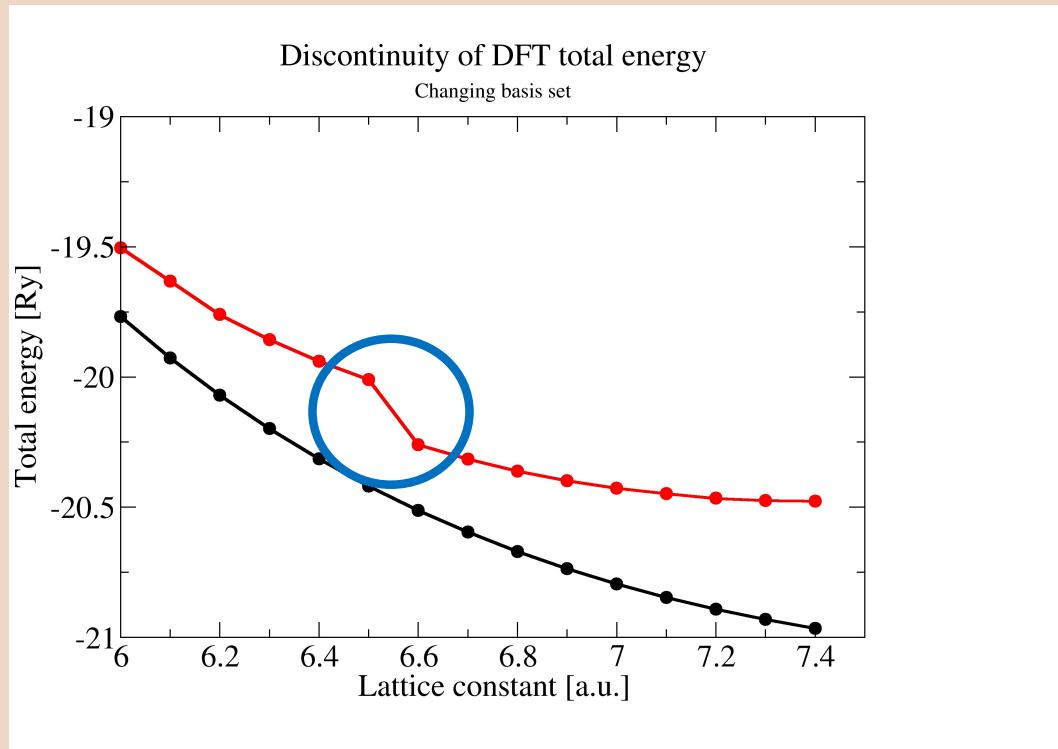
If the orbitals are represented with a cutoff of G_{\max} , then the charge density has non-zero Fourier components up to $2 * G_{\max}$!!! This means: $E_{\text{cut}}(\rho) = 4 * E_{\text{cut}}(\text{wfc})$

From KS orbitals to the charge density

```
&control
  calculation='scf'
  prefix='diamond'
  pseudo_dir='./'
/
&system
  ibrav = 2
  celldm(1)=7.4
  nat=2
  ntyp=1
  ecutwfc=60.
/
&electrons
/
```

```
bravais-lattice index      =          2
lattice parameter (alat)   =    7.4000 a.u.
unit-cell volume           = 101.3060 (a.u.)^3
number of atoms/cell       =          2
number of atomic types     =          1
number of electrons         =      8.00
number of Kohn-Sham states =          4
kinetic-energy cutoff       =    60.0000 Ry
charge density cutoff      = 240.0000 Ry
scf convergence threshold  = 1.0E-06
mixing beta                 =    0.7000
number of iterations used  =          8 plain      mixing
Exchange-correlation= SLA PW PBX PBC
                           ( 1 4 3 4 0 0 0)
```

Beware: changing the lattice constant means changing the PW basis set



Larger cell size
→ grid of G-vectors more dense
→ at fixed G_{max} , "new" plane waves
enter the calculation discontinuously
→ sudden decrease of energy when basis set
is enlarged

Curing the discontinuity: Check for input variables `ecfixed`, `qcutz`, `q2sigma`

So have seen how periodic functions are expressed in PWs

BUT: orbitals in periodic systems are NOT lattice periodic!

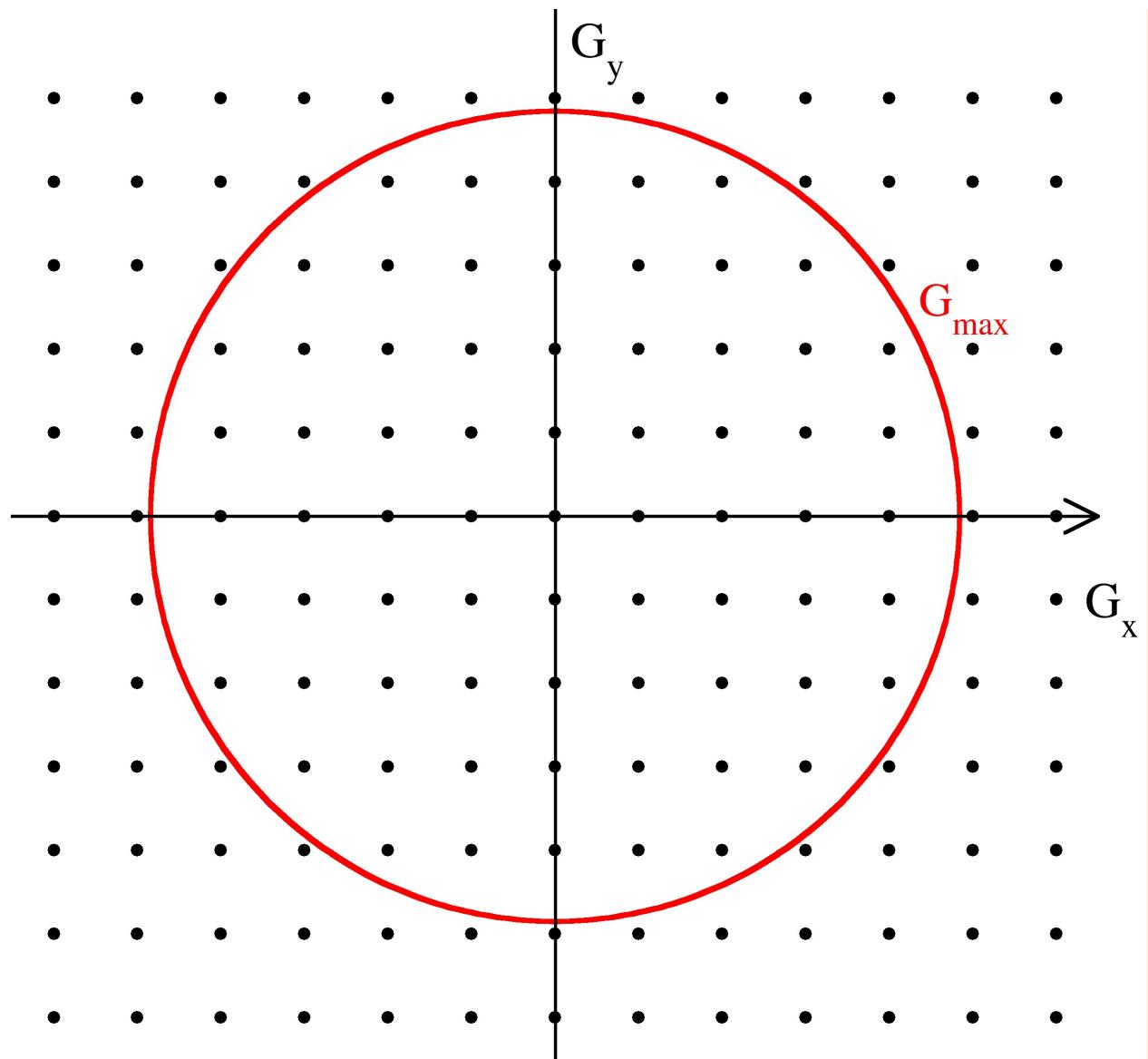
Bloch's theorem: $\psi_{\mathbf{k},l}(\mathbf{r}) = \exp(i\mathbf{k} \cdot \mathbf{r}) \frac{u_{\mathbf{k},l}(\mathbf{r})}{\text{periodic}}$

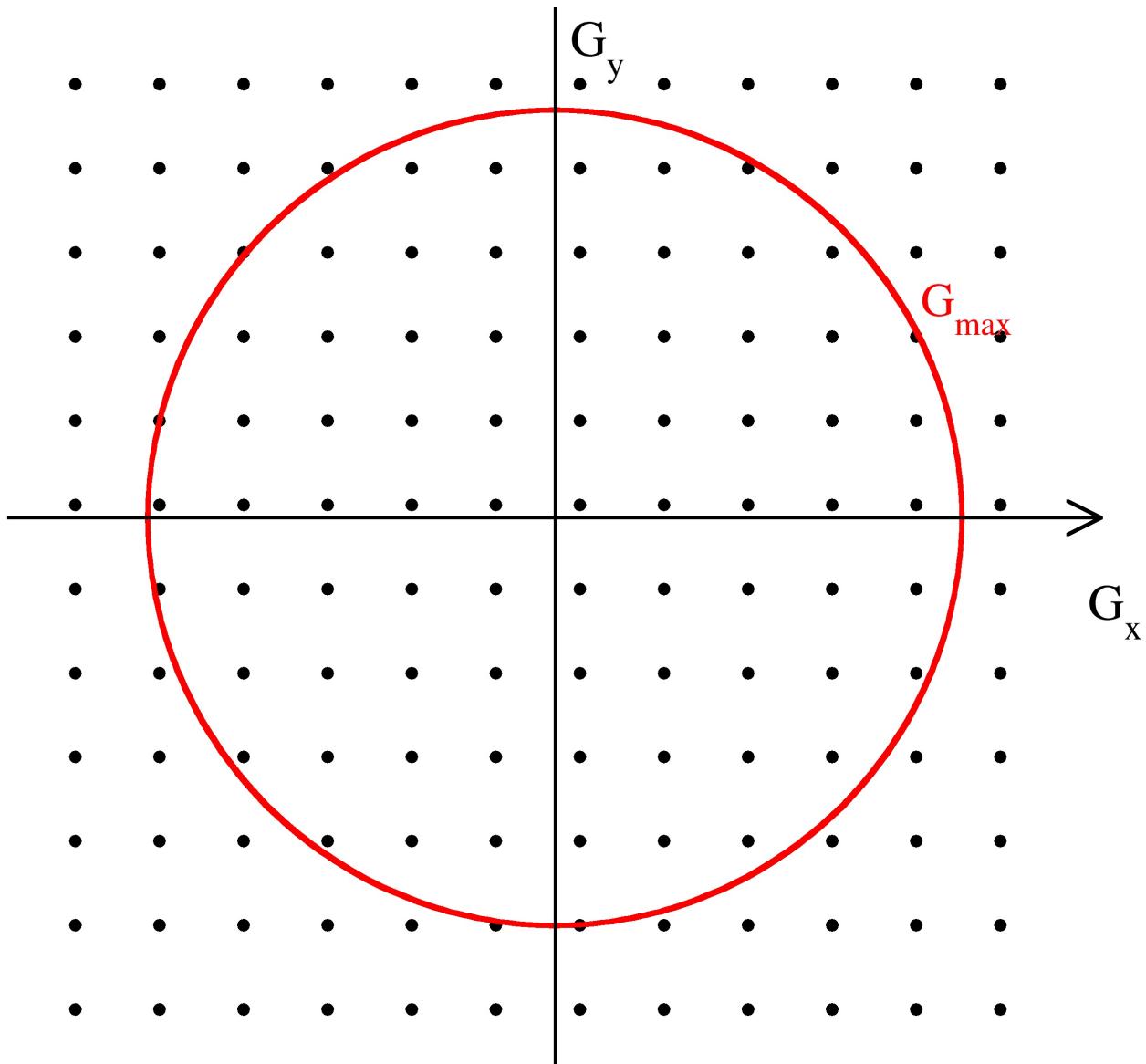
$$u_{\mathbf{k},l}(\mathbf{r}) = \sum_{|\mathbf{G}| < G_{max}} c_{\mathbf{k},l} \exp(i\mathbf{G} \cdot \mathbf{r})$$

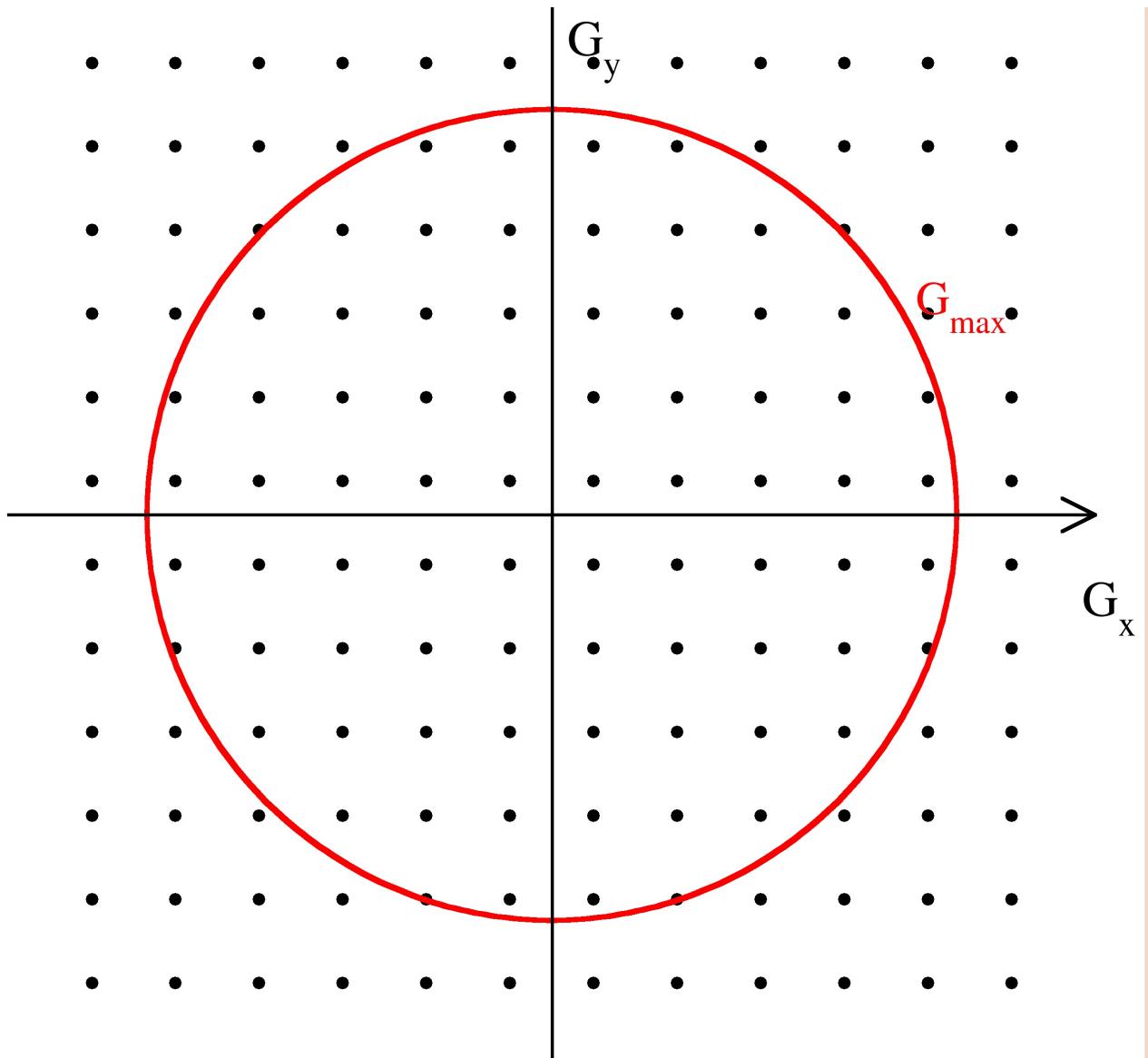
grid of G-vectors like before

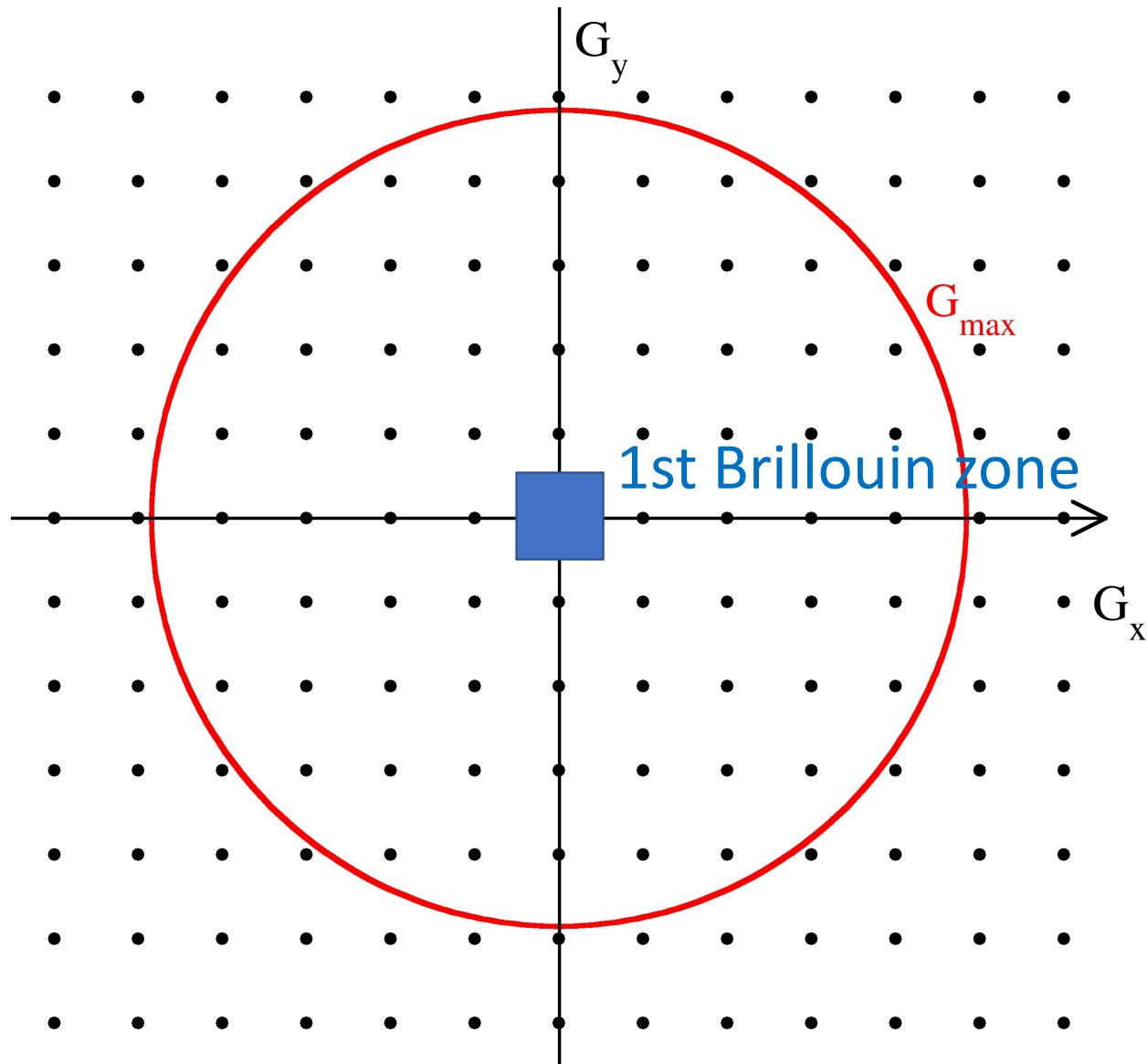
$$\psi_{\mathbf{k},l}(\mathbf{r}) = \sum_{|\mathbf{G}| < G_{max}} c_{\mathbf{k},l} \exp(i(\mathbf{k} + \mathbf{G}) \cdot \mathbf{r})$$

shifted grid of $(\mathbf{k} + \mathbf{G})$ -vectors

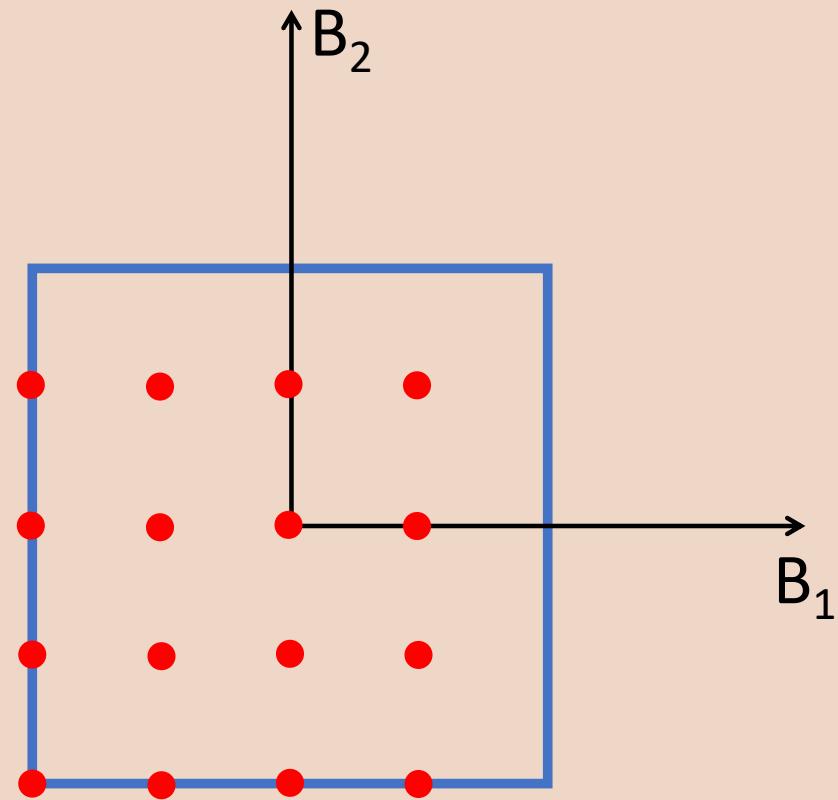






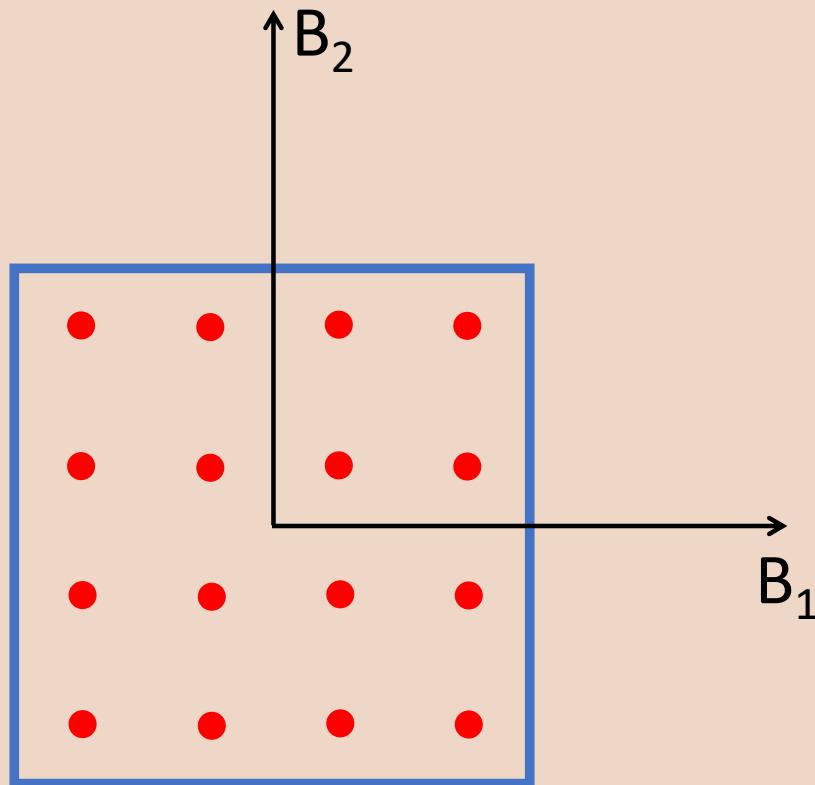


Sampling the Brillouin zone



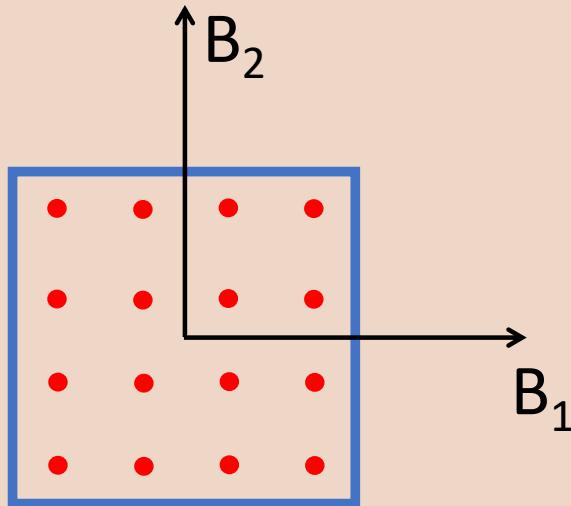
BZ sampled with a 4x4 k-point grid

Sampling the Brillouin zone



BZ sampled with a shifted 4x4 k-point grid
Shifted k-point grids usually converge faster (symmetry!)

Sampling the Brillouin zone

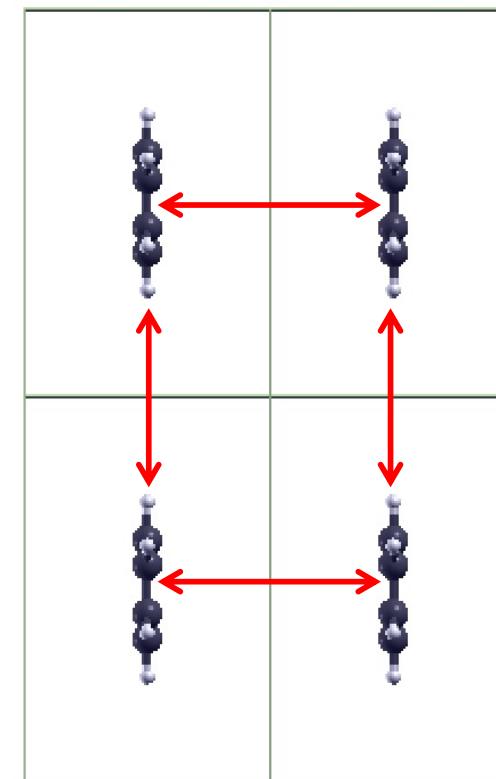
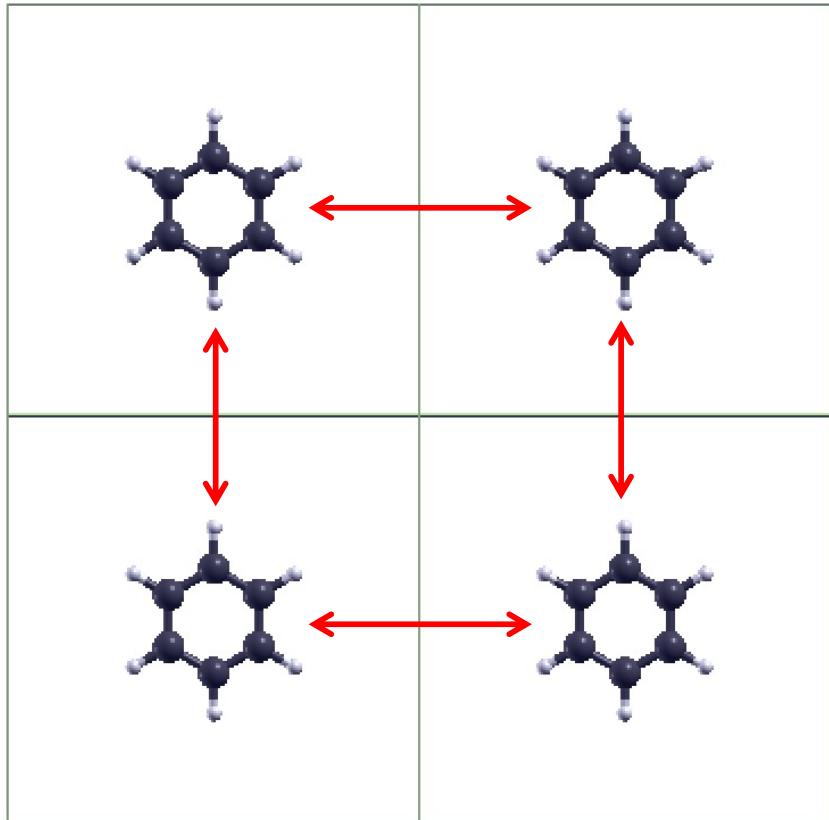


```
K_POINTS automatic  
9 9 1 1 1 0
```

```
number of k points= 25  
cart. coord. in units 2pi/alat  
k( 1) = ( 0.0555556 0.0962250 0.0000000), wk = 0.0493827  
k( 2) = ( 0.0555556 0.2245251 0.0000000), wk = 0.0987654  
k( 3) = ( 0.0555556 0.3528252 0.0000000), wk = 0.0987654  
k( 4) = ( 0.0555556 0.4811252 0.0000000), wk = 0.0987654  
k( 5) = ( 0.0555556 -0.5452753 0.0000000), wk = 0.0987654  
k( 6) = ( 0.0555556 -0.4169752 0.0000000), wk = 0.0987654  
k( 7) = ( 0.0555556 -0.2886751 0.0000000), wk = 0.0987654  
k( 8) = ( 0.0555556 -0.1603751 0.0000000), wk = 0.0987654  
k( 9) = ( 0.0555556 -0.0320750 0.0000000), wk = 0.0493827  
k( 10) = ( 0.1666667 0.2886751 0.0000000), wk = 0.0493827  
k( 11) = ( 0.1666667 0.4169752 0.0000000), wk = 0.0987654  
k( 12) = ( 0.1666667 0.5452753 0.0000000), wk = 0.0987654  
k( 13) = ( 0.1666667 -0.4811252 0.0000000), wk = 0.0987654  
k( 14) = ( 0.1666667 -0.3528252 0.0000000), wk = 0.0987654  
k( 15) = ( 0.1666667 -0.2245251 0.0000000), wk = 0.0987654  
k( 16) = ( 0.1666667 -0.0962250 0.0000000), wk = 0.0493827  
k( 17) = ( 0.2777778 0.4811252 0.0000000), wk = 0.0493827  
k( 18) = ( 0.2777778 0.6094253 0.0000000), wk = 0.0987654  
k( 19) = ( 0.2777778 -0.4169752 0.0000000), wk = 0.0987654  
k( 20) = ( 0.2777778 -0.2886751 0.0000000), wk = 0.0987654  
k( 21) = ( 0.2777778 -0.1603751 0.0000000), wk = 0.0493827  
k( 22) = ( 0.3888889 0.6735753 0.0000000), wk = 0.0493827  
k( 23) = ( 0.3888889 -0.3528252 0.0000000), wk = 0.0987654  
k( 24) = ( 0.3888889 -0.2245251 0.0000000), wk = 0.0493827  
k( 25) = ( -0.5000000 -0.8660254 0.0000000), wk = 0.0246914  
Dense grid: 3243 G-vectors FFT dimensions: ( 15, 15, 40)
```

What if a system is NOT 3D-periodic ?

Use a **supercell**: introduce enough space between periodic replicas that the interaction is negligible



What if a system is NOT 3D-periodic ?

Use a **supercell**: introduce enough space between periodic replicas that the interaction is negligible

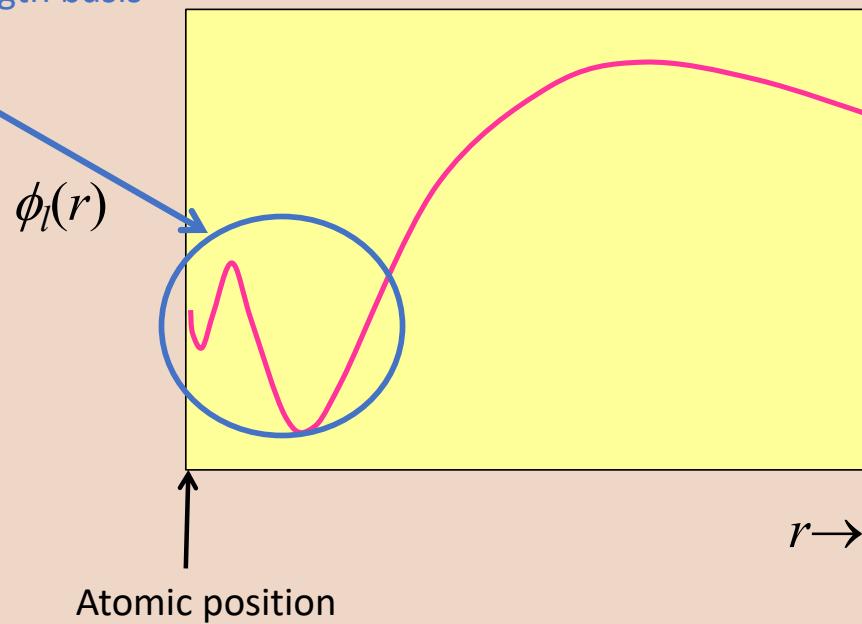
Can be problematic if the object is charged or has a dipole!

Solution:

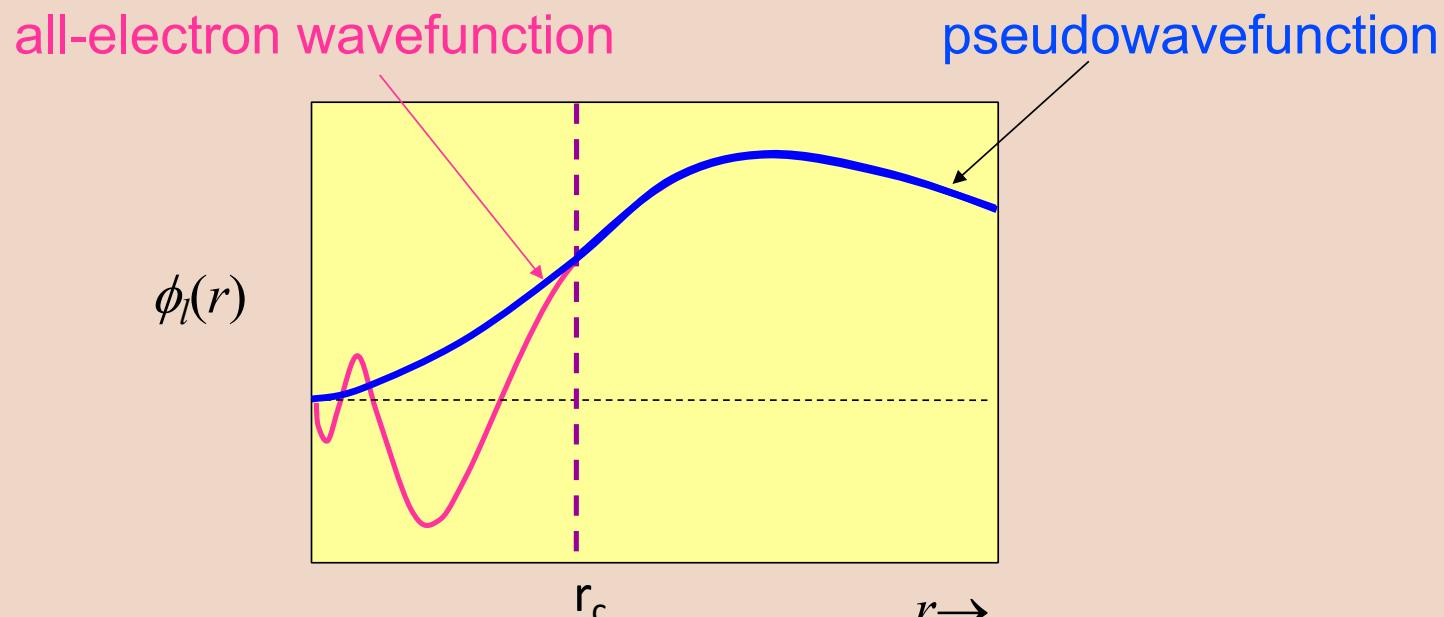
- input variable **assume_isolated**: useful in 0-D (molecules, clusters, etc)
- introduce a dipole layer in the vacuum between surfaces (**dipfield**)

The need for pseudopotentials

rapid oscillations would need
a very small wavelength basis
functions



Pseudopotentials

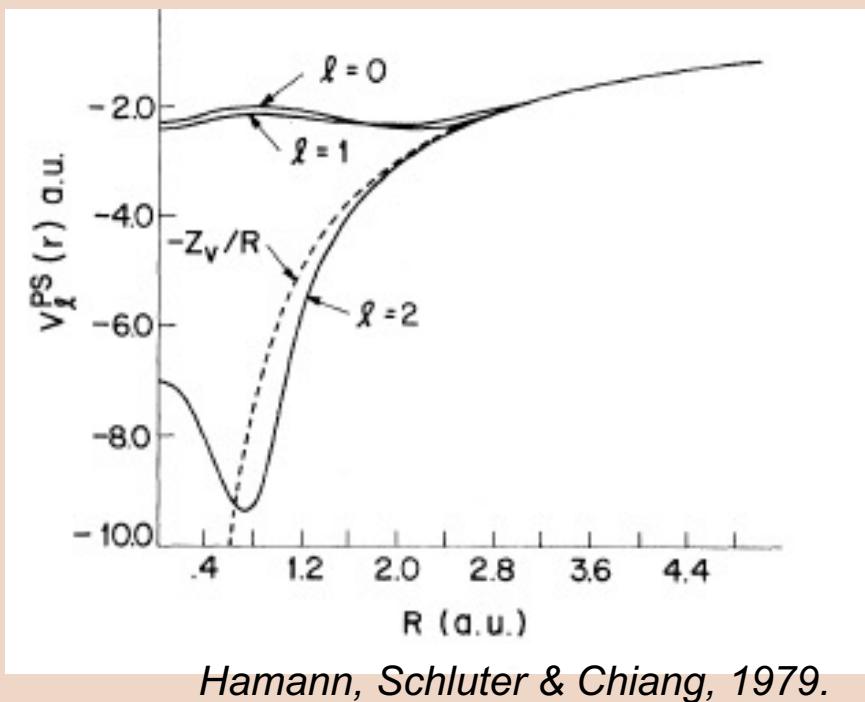


Necessary condition for the pseudopotential to be transferable: **NORM-CONSERVATION**

$$\int_0^{r_c} \phi^{*AE}(r) \phi^{AE}(r) dr = \int_0^{r_c} \phi^{*PS}(r) \phi^{PS}(r) dr$$

What does a pseudopotential look like?

Example for Mo:



- Weaker than full Coulomb potential
- No singularity at $r=0$
- Different pseudopotential for each l (example of semilocal pseudopotential)
- Will be V_{ion} (replacing nuclear potential)

Pseudopotentials in quantum-ESPRESSO:

PPs in QE are provided by the user in the form of files, which contain the PP on a radial grid:
in the PW input file, the PP filenames are specified for every element in the system

```
ATOMIC_SPECIES
Mo 95.96 Mo.pbe-spn-rrkjus_psl.1.0.0.UPF
W 183.8 W.pbe-spn-rrkjus_psl.1.0.0.UPF
S 32.07 S.pbe-n-rrkjus_psl.1.0.0.UPF
O 16.00 O.pbe-n-rrkjus_psl.1.0.0.UPF
H 1.00 H.pbe-rrkjus_psl.1.0.0.UPF
```

PPs are DFT-functional specific. The code will by default apply the functional to your system with
which the PP has been calculated

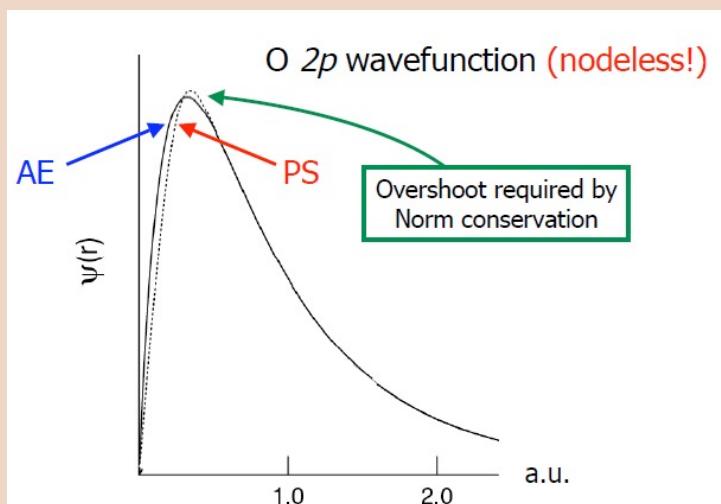
It is generally a bad idea to mix pseudopotentials generated with different functionals!

Beyond norm-conservation:

Cut-offs still higher than we would like, especially for

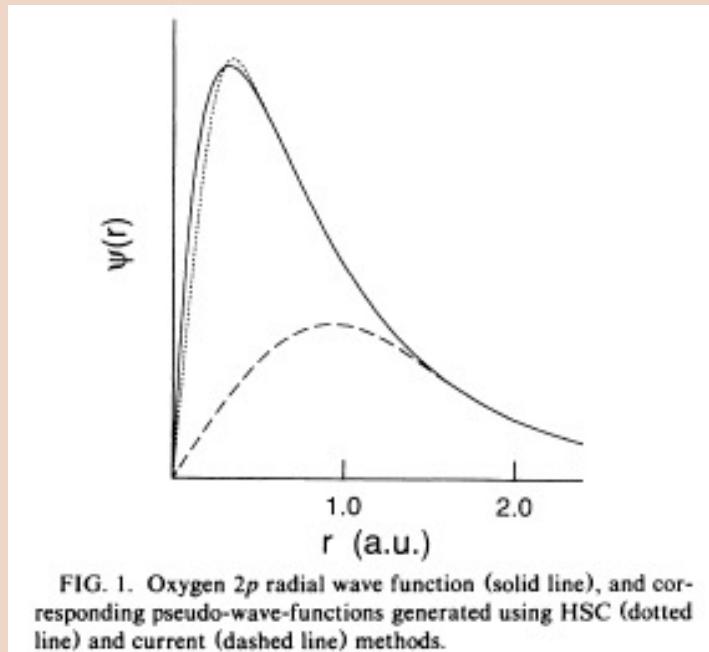
- > first row elements (1s, 2p nodeless)
- > transition metals (3d nodeless)
- > rareearths (4f nodeless)

This is because of the constraint of norm conservation...



Beyond norm-conservation:

Vanderbilt's ultrasoft pseudopotentials:



Remove the constraint of norm-conservation

The price to be paid: a more complicated expression for the electronic charge density:

$$n(\mathbf{r}) = \sum_{i=1}^N \psi_i^*(\mathbf{r})\psi_i(\mathbf{r}) + n^{aug}(\mathbf{r})$$

Augmentation charges, localized close to the nucleus
account for “missing norm”
expressed using projection functions

NB: Ultrasoft (US) and PAW pseudopotentials are (from the user’s perspective) very similar

Beyond norm-conservation:

Vanderbilt's ultrasoft pseudopotentials:

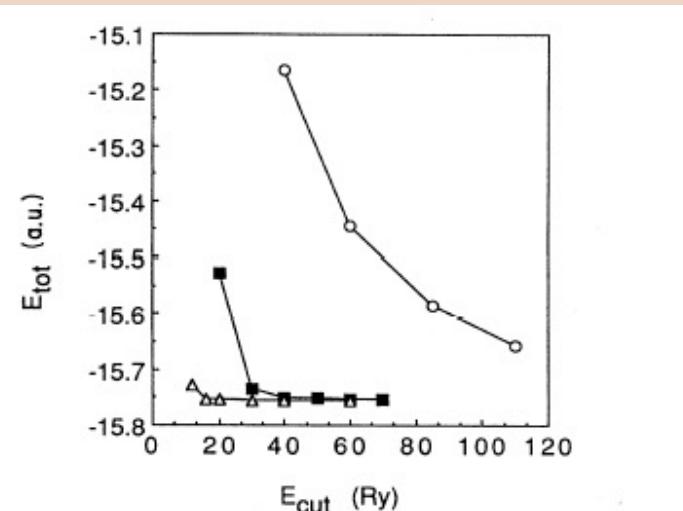
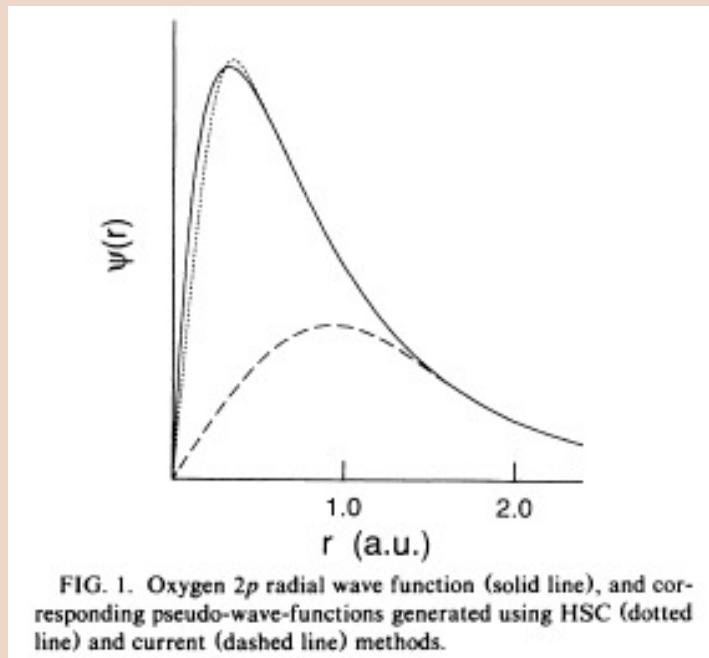


FIG. 1. Total energy of ground-state oxygen atom vs plane-wave cutoff for Bachelet-Hamann-Schlüter pseudopotential (open circles) and for Vanderbilt pseudopotential with $r_c = 1.2$ a.u. (solid squares) and $r_c = 1.8$ a.u. (open triangles).

As a user: BEWARE!

Using US or PAW pseudopotentials, you need a cutoff energy E_{cut} for the cgarge density which is
MORE THAN $4*E_{\text{cut}}$ for the orbitals

ATOMIC_SPECIES

Mo 95.96	Mo.pbe-spn-rrkjus_psl.1.0.0.UPF
W 183.8	W.pbe-spn-rrkjus_psl.1.0.0.UPF
S 32.07	S.pbe-n-rrkjus_psl.1.0.0.UPF
O 16.00	O.pbe-n-rrkjus_psl.1.0.0.UPF
H 1.00	H.pbe-rrkjus_psl.1.0.0.UPF

Employing US or PAW pseudopotentials and NOT specifying
ecutrho is the most common beginner's mistake in QE!

```
&system
  ibrav=0,
  nat=78,
  ntyp=5,
  ecutwfc = 40. ,
  ecutrho = 400.
  nspin=2,
  starting_magnetization(1) = 0.
  starting_magnetization(2) = 0.1
  starting_magnetization(3) = 0.
  starting_magnetization(4) = 0.2
  starting_magnetization(5) = 0.2
  vdw_corr='grimme-d3'
  occupations='smearing'
  smearing='cold'
  degauss=0.001
/
```

HAPPY QE-COMPUTING!