

Boundary States for Strings in Non-constant EM Background

Consider the action:

$$S = -\frac{1}{4\pi\alpha'} \int d\tau d\sigma \left[-\partial_\tau X^\mu \partial_\tau X_\mu + \partial_\sigma X^\mu \partial_\sigma X_\mu \right]$$

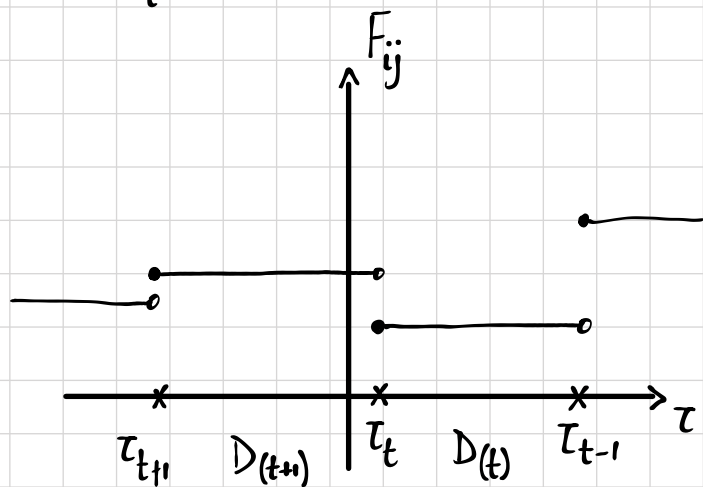
and add EM potential to OPEN STRING ENDPOINTS (suppose space-filling D9 or D25):

$$S = -\frac{1}{4\pi\alpha'} \int d\tau d\sigma \left(-\partial_\tau X^\mu \partial_\tau X_\mu + \partial_\sigma X^\mu \partial_\sigma X_\mu \right) + \int d\tau A_\mu \partial_\tau X^\mu \Big|_{\sigma=\pi} - \int d\tau A_\mu \partial_\tau X^\mu \Big|_{\sigma=0}$$

Now consider N D-branes, each with a different value of the potential A_μ . Let $F_{\mu\nu}$ be its field strength such that:

$$\partial_\tau F_{\mu\nu}(\tau) = 0 \quad \text{if} \quad \tau_t < \tau < \tau_{t+1} \quad (\text{i.e.: when the string is on the brane } D_{(t)})$$

and $\lim_{\tau \rightarrow \tau_t^+} F_{\mu\nu}(\tau) \neq \lim_{\tau \rightarrow \tau_t^-} F_{\mu\nu}(\tau)$ (i.e.: $F_{\mu\nu}$ has discontinuities):



NB: t defined mod N

$$F_{\mu\nu}(\tau) = \sum_{t=1}^N F_{\mu\nu(t)}(\tau) \theta(\tau - \tau_t) \theta(\tau_{t+1} - \tau)$$

$$\Rightarrow A_\mu(\tau, \sigma) = F_{\mu\nu(t)}(\tau) X^\nu(\tau, \sigma) \quad \text{for} \quad \tau_t < \tau < \tau_{t+1}.$$

Define the potential on the brane $D_{(t)}$:

$$A_{\mu(t)} = F_{\mu\nu(t)} X^\nu_{(t)}$$

where

$$X^\mu_{(t)} = X^\mu \theta(\tau - \tau_t) \theta(\tau_{t+1} - \tau).$$

Now we take $\tau = i\tau_E$:

$$\begin{aligned} \delta S &= \frac{1}{2\pi\alpha'} \int d\tau_E d\sigma \left(\partial_{\tau_E} X^\mu \cdot \partial_{\tau_E} \delta X_\mu + \partial_\sigma X^\mu \cdot \partial_\sigma \delta X_\mu \right) - \sum_{t=1}^{\infty} \int_{\tau_t}^{\tau_{t+1}} d\tau_E \left(F_{\mu\nu}(t) \delta X_{(t)}^\nu \partial_{\tau_E} X_{(t)}^\mu \Big|_{\sigma=0} + F_{\mu\nu}(t) X_{(t)}^\nu \partial_{\tau_E} \delta X_{(t)}^\mu \Big|_{\sigma=\pi} - \dots \Big|_{\sigma=0} \right) = \\ &= \frac{1}{2\pi\alpha'} \int d\tau d\sigma \left(\partial_\tau \partial_\tau X^\mu + \partial_\sigma \partial_\sigma X^\mu \right) \delta X_\mu + \frac{1}{2\pi\alpha'} \int d\tau \partial_\sigma X^\mu \cdot \delta X_\mu \Big|_{\sigma=0, \pi} + \sum_{t=1}^{\infty} \int d\tau \left(F_{\mu\nu}(t) \partial_\tau X_{(t)}^\mu \Big|_{\sigma=0} - F_{\mu\nu}(t) \partial_\tau X_{(t)}^\mu \Big|_{\sigma=\pi} \right) \delta X_{(t)}^\nu = \end{aligned}$$

$$\Rightarrow \text{EOM: } \square X^\mu = 0$$

$$\text{BC: } \partial_\sigma X_\mu(t) - 2\pi\alpha' F_{\mu\nu}(t) \partial_\tau X_{(t)}^\nu \Big|_{\sigma=0} = 0.$$

Now look at closed strings: $X^\mu(\tau, \sigma) = X_L^\mu(\tau + i\sigma) + X_R^\mu(\tau - i\sigma)$:

$$X_L^\mu(\tau + i\sigma) = \frac{1}{2} \tilde{X}_0^\mu + \ell^2 p^\mu (\tau + i\sigma) + i\ell \sum_{n \neq 0} \frac{\alpha_n^\mu}{n} e^{-in(\tau + i\sigma)}$$

$$X_R^\mu(\tau - i\sigma) = \frac{1}{2} X_0^\mu + \tilde{\ell}^2 \tilde{p}^\mu (\tau - i\sigma) + i\tilde{\ell} \sum_{n \neq 0} \frac{\tilde{\alpha}_n^\mu}{n} e^{-in(\tau - i\sigma)}$$

Then we impose $X^\mu(\tau, \sigma + 2\pi) = X^\mu(\tau, \sigma)$:

$$\ell^2 p^\mu \cdot 2i\pi - \tilde{\ell}^2 \tilde{p}^\mu \cdot 2i\pi = 0 \quad \text{and} \quad n \in \mathbb{Z}$$

Then: $\ell^2 = \tilde{\ell}^2$, $p^\mu = \tilde{p}^\mu$ and $n \in \mathbb{Z}$:

$$X_{(c)}^\mu(\tau, \sigma) = \frac{1}{2} (X_0^\mu + \tilde{X}_0^\mu) + 2\ell^2 p^\mu \tau + i\ell \sum_{n \neq 0} \frac{1}{n} (\alpha_n^\mu e^{-in(\tau + i\sigma)} + \tilde{\alpha}_n^\mu e^{-in(\tau - i\sigma)}).$$

Then we look at the centre-of-mass momentum:

$$\frac{1}{2\pi} \int_0^{2\pi} d\sigma \mathcal{P}^\mu = \frac{1}{2\pi} \left(\frac{1}{2\alpha'} \right) \int_0^{2\pi} d\sigma \partial_\tau X^\mu = \frac{1}{2\alpha'} \cdot 2\ell^2 p^\mu = \pi_{CM}^\mu \leftrightarrow \ell^2 = \alpha', \quad p^\mu = \pi_{CM}^\mu$$

$$\Rightarrow X_{(c)}^\mu(\tau, \sigma) = \frac{1}{2} (X_0^\mu + \tilde{X}_0^\mu) + 2\alpha' \pi_{CM}^\mu \tau + i\sqrt{2\alpha'} \sum_{n \neq 0} \frac{1}{n} (\alpha_n^\mu e^{-in(\tau + i\sigma)} + \tilde{\alpha}_n^\mu e^{-in(\tau - i\sigma)}).$$

We try now to impose the Neumann conditions in the closed string channel:

$$\partial_\tau X_\mu(t) - 2\pi\alpha' F_{\mu\nu}(t) \partial_\sigma X^\nu(t) \big|_{\tau=0} |B_{\text{eff}}\rangle = 0.$$

where:

$$\partial_\tau X_\mu(t) \big|_{\tau=0} = 2\alpha' \pi_{\mu\eta} + i\sqrt{2\alpha'} \sum_{n \neq 0} (\alpha_n^\nu e^{-in\sigma} + \tilde{\alpha}_n^\nu e^{in\sigma}) \eta_{\mu\nu} \quad (\eta_{\mu\nu} = \text{diag}(-, +, +, \dots))$$

$$\partial_\sigma X^\nu(t) \big|_{\tau=0} = i\sqrt{2\alpha'} \sum_{n \neq 0} (-i\alpha_n^\nu e^{-in\sigma} + i\tilde{\alpha}_n^\nu e^{in\sigma})$$

$$\hookrightarrow 2\alpha' \pi_{\mu\eta} + i\sqrt{2\alpha'} \sum_{n \neq 0} \left[(\alpha_n^\nu \eta_{\mu\nu} - i2\pi\alpha' F_{\mu\nu}(t) \alpha_n^\nu) e^{-in\sigma} + (\tilde{\alpha}_n^\nu \eta_{\mu\nu} + i2\pi\alpha' F_{\mu\nu}(t) \tilde{\alpha}_n^\nu) e^{in\sigma} \right] |B_{\text{eff}}\rangle = 0$$

$$\Rightarrow 2\alpha' \pi_{\mu\eta} + i\sqrt{2\alpha'} \sum_{n \neq 0} \left(\alpha_n^\nu \underbrace{(\eta_{\mu\nu} - 2\pi i \alpha' F_{\mu\nu}(t))}_{F_{\mu\nu}(t)} e^{-in\sigma} + \tilde{\alpha}_n^\nu \underbrace{(\eta_{\mu\nu} + 2\pi i \alpha' F_{\mu\nu}(t))}_{F_{\mu\nu}(t)} e^{in\sigma} \right) |B_{\text{eff}}\rangle = 0$$

$$\Rightarrow \pi_{\mu\eta} |B_{\text{eff}}\rangle = 0 \quad (\text{state is localized on the brane})$$

$$\sum_{n \neq 0} (\alpha_n^\nu (\eta_{\mu\nu} - F_{\mu\nu}(t)) + \tilde{\alpha}_n^\nu (\eta_{\mu\nu} + F_{\mu\nu}(t))) |B_{\text{eff}}\rangle = 0$$

As in the usual case

$$\sum_{n \neq 0} (j_n + \tilde{j}_{-n}) |B_0\rangle = 0 \rightarrow |B_0\rangle \propto \exp\left(-\sum_{n=1}^{\infty} \frac{1}{n} j_n^\mu \tilde{j}_{-n}^\nu \eta_{\mu\nu}\right),$$

then, in this case:

$$|B_{\text{eff}}\rangle = \mathcal{N} \delta^9(y - x_0) \exp\left(-\sum_{n=1}^{\infty} \frac{1}{n} \alpha_n^\mu (\eta_{\mu\nu} - F_{\mu\nu}(t))^\dagger (\eta_{\rho\sigma} + F_{\rho\sigma}(t)) \alpha_{-n}^\rho \eta_{\mu\rho}\right).$$