PDE Midterm

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Problem (1). We see that $a=x^2, b=xy, c=y^2$. So $d=ac-b^2=x^2y^2-(xy)^2=0$, and the matrix $\begin{pmatrix} a & b \\ b & c \end{pmatrix}$ is clearly not identically zero. So the equation is parabolic on \mathbb{R}^2 .

To change into the standard form, we wish to do an invertible change of coordinate using s(x,y) and t(x,y). Under this transformation, the principal linear part becomes $L_0u=a^*u_{ss}+2b^*u_{st}+c^*u_{tt}$, where $a^*(x,y),b^*(x,y),c^*(x,y)$ are the coefficients after applying chain rules. In particular, $a^*(x,y)=as_x^2+2bs_xs_y+cs_y^2$. If we wish the u_{ss} term to vanish, we must set $a^*=0$. This means that s(x,y) satisfies the characteristic equation of L_0u . Let us find all solutions of the characteristic equation. If $\phi(x,y)$ is such solution, this means that each level set of $\phi(x,y)$ gives a characteristic curve. Locally, the level set of $\phi(x,y)$ can be expressed as y(x) (or x(y) if $\phi_y=0$) by the implicit function theorem.

$$0 = \frac{d\phi}{dx} = \phi_x + \frac{dy}{dx}\phi_y$$
$$\frac{dy}{dx} = -\frac{\phi_x}{\phi_y}$$

We can then substitute this into a^* . If x = 0, then a = 0 which forces b = 0. Then for points (0, y), the equation is already in the standard form. Likewise for y = 0. WLOG assume x > 0 and y > 0. We have

$$a\left(\frac{dy}{dx}\right)^{2} + 2b\frac{dy}{dx} + c = 0$$

$$\frac{dy}{dx} = \frac{2b \pm \sqrt{4b^{2} - 4ac}}{2a}$$

$$= \frac{b \pm \sqrt{b^{2} - ac}}{a}$$

$$= \frac{b}{a}$$

$$= \frac{\sqrt{ac}}{a}$$

$$= \sqrt{\frac{c}{a}}$$

$$= \frac{y}{x}$$

Solving this gives one family of real solutions y = Cx for each level set $\phi(x,y) = C$. Thus define $s(x,y) = \frac{y}{x}$ for $x \neq 0$. Notice that since $d^* = a^*c^* - b^{*2} = -b^{*2}$ is the determinant, which is invariant under a change of basis, $-b^{*2} = d = 0$ which forces $b^* = 0$, which also forces $c^* \neq 0$ since A is not identically zero and shall remain so under change of basis. Thus we just need to choose an t(x,y) that is independent from s(x,y) so the Jacobian is nonsingular. We see that t(x,y) = x will do, since

$$\det J = \det \begin{pmatrix} s_x & s_y \\ t_x & t_y \end{pmatrix} = \begin{pmatrix} s_x & \frac{1}{x} \\ 1 & 0 \end{pmatrix} = -\frac{1}{x} \neq 0.$$

It follows that under this change of coordinates, the equation reduces to

$$c^* u_{tt} = 0$$

$$u_{tt} = 0$$

which is in the standard form.

Problem (2). (a) This is Burger's equation so the intuition from homework is that a global solution exists iff u_x doesn't blow up in finite time. To make this precise, we solve the problem using the method of characteristic first.

The initial curve S_0 is parametrized by $t_0 = 0, x_0 = s, z_0 = u_0(s)$. It is the graph of $u_0(x)$ over the domain $\{0\} \times \mathbb{R}$. From the equation we have a = 1, b = u, c = -2u. Since

$$\det \begin{pmatrix} a & b \\ t'_0(s) & x'_0(s) \end{pmatrix} = \det \begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix} = 1 \neq 0,$$

the projection of S_0 onto (t, x)-plane C is not characteristic so u can flow out of the initial curve. Now we can solve the characteristic curves using the characteristic equation:

$$\begin{cases} \frac{dt}{d\tau} = 1, & t(0) = 0 \Rightarrow t = \tau \\ \frac{dx}{dt} = z, & x(0) = s \\ \frac{dz}{dt} = -2z, & z(0) = u_0(s) \Rightarrow u(s,t) = z(s,t) = u_0(s)e^{-2t} \end{cases}$$

For a global solution to exist, intuitively we want u to flow along characteristic curve for forever. Thus along the characteristic curve, we cannot have u_x goes to infinity in finite time, or the flow would terminate. Thus we investigate $w := u_x$. Taking partial derivative on both sides of the Burger's equation, we have

$$(u_t + uu_x)_x = -2u_x$$

$$u_{xt} + u_x^2 + u_{xx} = -2u_x$$

$$w_t + uw_x = -2w - w^2$$

$$\dot{w} = -2w - w^2$$

Moreover, since $u_0(s) \in C^1(\mathbb{R})$, we have $w(s,0) = (u_0(s)e^{-2\cdot 0})_x = u_0'(s)$. Using the method of characteristic again, we see that along the characteristic curve, w satisfies

$$\begin{cases} \frac{dw}{dt} = -2w - w^2 \\ w(s,0) = u'_0(s) \end{cases}$$

We compute

$$-\frac{dw}{w(w+2)} = dt$$

$$-\frac{1}{2} \left(\frac{1}{w} - \frac{1}{w+2}\right) dw = dt$$

$$\frac{1}{2} \ln|w+2| - \frac{1}{2} \ln|w| = t + C'$$

$$\frac{1}{2} \ln\left|\frac{w+2}{w}\right| = t + C'$$

$$1 + \frac{2}{w} = Ce^{2t}$$

$$w = \frac{2}{Ce^{2t} - 1}$$

$$w(s,0) = \frac{2}{C-1} = u'_0(s)$$

$$C = \frac{2}{u'_0(s)} + 1$$

$$w(s,t) = \frac{2}{\left(\frac{2}{u'_0(s)} + 1\right)e^{2t} - 1}$$

We see that for a fixed s (*i.e.* on the characteristic curve), w(s,t) blows up when the denominator approaches 0. This happens when

$$\left(\frac{2}{u_0'(s)} + 1\right)e^{2t} = 1$$

$$e^{2t} = \frac{u'_0(s)}{2 + u'_0(s)}$$
$$t = \frac{1}{2} \ln \left(\frac{u'_0(s)}{2 + u'_0(s)} \right)$$

This is the time of w blowing up. But if $\frac{u_0'(s)}{2+u_0'(s)} \leq 0 \Leftrightarrow -2 \leq u_0'(s) \leq 0$, or if $t < 0 \Leftrightarrow u_0'(s) \geq 0$, then t would not have a solution in the range $t \geq 0$ so w would not blow up. Therefore, the necessary condition for u_0 is that for all $x \in \mathbb{R}$, $u_0'(x) \geq -2$.

We now show that this is also the sufficient condition. Suppose the condition is true. Notice that since $u_0(x)$ has bounded C^1 -norm, u(x,0) is Lipschitz continuous in x. Along each characteristic curve, for t > 0 we have $u(x,t) = u(x,0)e^{-2t}$ which is also Lipschitz continuous in x. Now consider $\frac{dx}{dt} = u$. Picard's theorem guarantees the uniqueness of x(t) that passes through each (x,t). That is, no two characteristic curve collide in the (x,t)-plane. Since the determinant condition is always satisfied, as long as u_x doesn't blow up, we can let u(x,t) flow out via the formula $u(x,t) = u_0(x)e^{-2t}$ unobstructedly and obtain a unique global solution. We have shown that u_x doesn't blow up exactly when $u'_0(x) \ge -2$. Since $u_0(x)$ is smooth, ODE theory yields that the global solution u(x,t) is smooth as well.

We remark that this result can also be obtained via the phase line diagram of $\frac{dw}{dt} = -w(w+2)$ along a characteristic curve. We see that the critical points are 0 and 2, where the solution is constant (by uniqueness). If the initial value is at w < -2, we see that $\frac{dw}{dt} < 0$ so $w \to \infty$ as $t \to \infty$, i.e. u_x blows up. If the initial value is at -2 < w < 0 or w > 0, we see that $w \to 0$ as $t \to \infty$. Taken together, u_x doesn't blow up when the initial value $w(0) = u_0'(x) \ge -2$, which agrees with the above analysis.

(b) We see that any u(x,t) is the result of flow along some characteristic curve originated from the initial surface. Recall along any characteristic curve, $u(x,t) = u_0(x)e^{-2t}$. Since $u_0(x)$ is bounded, $|u_0(x)| \leq M$ for some M > 0. Then

$$\sup_{x \in \mathbb{R}} |u(x,t)| = \sup_{x \in \mathbb{R}} |u_0(x)e^{-2t}|$$
$$= \sup_{x \in \mathbb{R}} |u_0(x)||e^{-2t}|$$
$$< Me^{-2t}$$

Thus we see that as $t \to \infty$, $\sup_{x \in \mathbb{R}} |u(x,t)| \to 0$, as desired.

Problem (3). We guess that the solution has the form

$$u(x) = \frac{1}{|x - y|},$$

where $y \in B(0,1)$ is chosen to satisfy the initial data. We observe that

$$\frac{2}{\sqrt{7+4\sqrt{3}x_3}} = \frac{1}{\sqrt{\frac{7}{4}+\sqrt{3}x_3}}$$

whose denominator looks like a norm. So for $x \in \partial B(0,1)$, we solve

$$\frac{7}{4} + \sqrt{3}x_3 = |x - y|^2 = |x|^2 - 2x \cdot y + |y|^2 = 1 - 2x \cdot y + |y|^2$$

$$\begin{cases} 1 + y_1^2 + y_2^2 + y_3^2 = \frac{7}{4} \\ -2x \cdot y = \sqrt{3}x_3 \end{cases}$$

Thus we obtain $y_1 = y_2 = 0$ and $y_3 = -\frac{\sqrt{3}}{2}$, which indeed lies inside B(0,1). Thus,

$$u(x) = \frac{1}{\sqrt{x_1^2 + x_2^2 + (x_3 + \sqrt{3}/2)^2}}.$$

Since u(x) > 0, the bound depends on how small the denominator gets. Reverse triangle equality yields $|x - y| \ge ||x| - |y|| = ||x| - \sqrt{3}/2| \ge 1 - \sqrt{3}/2$ since $|x| \ge 1$. Thus the denominator is lower bounded by a constant, which means u(x) is upper-bounded by a constant as well. Thus it is bounded.

Problem (4). We know that $\rho^2 \cos 2\theta$ is harmonic, multiplying a harmonic function by a scalar keeps it harmonic (linearity of differentiation), and adding it by a constant keeps it harmonic. Thus,

$$u(\rho, \theta) = 1 + \frac{1}{2} \left(1 + \left(\frac{\rho}{r}\right)^2 \cos 2\theta \right)$$

is harmonic in D(0,r). When $\rho=r$, we have

$$u(r,\theta) = 1 + \frac{1}{2}(1 + \cos 2\theta) = 1 + \cos^2 \theta.$$

Thus $u(\rho, \theta)$ is the solution.

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Guideline: Please read the following carefully.

Print your name first. Remember to show all your work; including all intermediate steps and also explain in words how you are solving a problem. Partial credits are available for most problems. Correct answers without major steps will only receive minor portion of the credits of the whole problem. Work by your own, no discussion with others is permitted.

Problem 1 (30 pts) Identify the types of the following equation

$$x^2 u_{xx} + 2xy u_{xy} + y^2 u_{yy} = 0,$$

and then transfer it into standard form.

(Additional page for Problem 1)

Problem 2 (30 pts) Consider the following initial value probem

$$\begin{cases} u_t + uu_x = -2u, \ x \in \mathbf{R}, t > 0, \\ u(x, 0) = u_0(x) \in C^1(\mathbf{R}). \end{cases}$$

Here, $u_0(x)$ has bounded C^1 -norm.

a) (20 pts) Determine the sufficient and necessary conditions on the initial data $u_0(x)$ for this problem to have a unique global smooth solution.

b) (10 pts) If $u_0(x)$ satisfies the conditions found in part a), prove that the global solution u(x,t) satisfies that $||u(x,t)||_{L^{\infty}(\mathbf{R})}$ converges to zero as t goes to infinity.

Problem 3 (20 pt) Let B(0,1) be the unit ball in \mathbb{R}^3 centered at the origin. Find a bounded solution to the following Dirichlet problem outside B(0,1)

$$\begin{cases} -\Delta u(x) = 0, |x| > 1, \\ u(x) = \frac{2}{\sqrt{7 + 4\sqrt{3}x_3}}, \text{ for } |x| = 1. \end{cases}$$

Problem 4 (20 pt) Let D(0,r) be the disk on \mathbb{R}^2 centered at the origin with radius r with boundary C. Find the function $u(\rho,\theta)$ in polar coordinates so that it is harmonic on D(0,r) and $u(r,\theta)=1+cos^2(\theta)$ on C. **Hint:** $cos^2(\theta)=\frac{1}{2}(1+cos(2\theta))$.