

# VIBRATION OF POINT MASSES

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## Aims

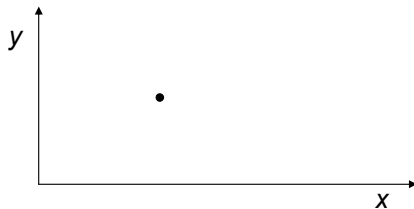
- To explore the motion of a point mass with forces exerting on it.
- To construct and formulate, from first principles the mathematical formulas that govern one- and two-mass oscillators.
- To analyse and interpret the oscillatory solutions of these systems.
- To derive useful concepts for analysis and representation of these systems.

## Content

- Point Masses
- Forces and Acceleration
- One-mass Oscillator
- Resonance
- Damped Oscillator
- Driven Oscillator
- Time- and Frequency Response
- Two-Mass Systems
- Normal Modes
- Multi-Mass & Continuous Systems

## Point Mass

A geometric (0-dimensional) point that may be assigned a finite mass ( $m$ ). Since a point has zero volume, the density of a point mass having a finite mass is infinite, so point masses do not exist in reality!

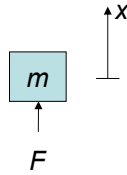


However, it is often a useful simplification in real problems. We represent the point mass here as follows:



## Force and Acceleration

Force	$F$
Mass	$m$
Position	$x$
Acceleration	$\ddot{x} = \frac{\partial^2 x}{\partial t^2}$

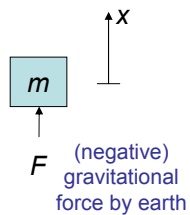


### Newton's Second Law:

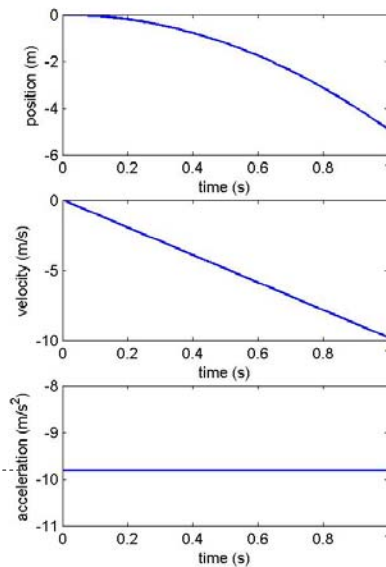
$$F = m \ddot{x}$$

total net force equals mass times acceleration

## A Falling Mass



gravitational  
acceleration  
 $g = -9.81 \text{ m/s}^2$



# One-Mass Oscillator

## Spring-mass system

Hooke's law :  $F_r = -k x$

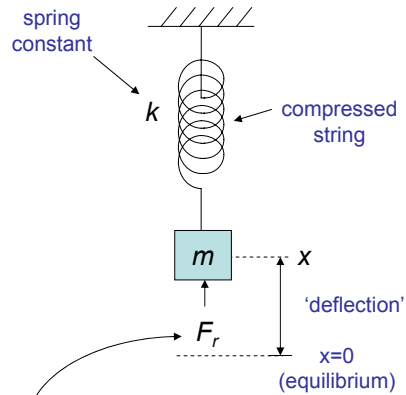
Newton's 2nd law :  $F_r = m \ddot{x}$

equation of motion:

$$\Rightarrow m \ddot{x} + k x = 0$$

Ordinary differential equation (ODE)

restoring force exerted by the spring on the mass



# One-Mass Oscillator: Solutions

Look for solution to the equation of motion (differential equation)

Try solution of the form:

$$x = e^{st}$$

thus:

$$\dot{x} = s e^{st} = s x$$

$$\ddot{x} = s^2 e^{st} = s^2 x$$

Substitute into:

$$m \ddot{x} + k x = 0$$

gives

$$m s^2 x + k x = 0$$

$$\Leftrightarrow (m s^2 + k) x = 0$$

for all  $x$ , thus

$$m s^2 + k = 0$$

$$s^2 + \left(\frac{k}{m}\right) = 0$$

$$s = \pm \sqrt{-\frac{K}{M}} = \pm \sqrt{j^2 \frac{K}{M}} = \pm j \sqrt{\frac{K}{M}}$$

$(j = \sqrt{-1})$

## One-Mass Oscillator: Solutions (cont.)

Dimension of sqrt(k/m)?

$$K \rightarrow \frac{\text{Force}}{\text{displacement}} \rightarrow \frac{N}{m}$$

$$M \rightarrow \frac{\text{Force}}{\text{Acceleration}} \rightarrow \frac{N}{(m/\text{sec}^2)} = \frac{N}{m} \text{sec}^2$$

$$\sqrt{\frac{K}{M}} \rightarrow \sqrt{\frac{Nm}{Nm \text{sec}^2}} = \frac{1}{\text{sec}} \quad \Rightarrow \quad \text{dimension of frequency!}$$

Turns out it is the (angular) **natural** frequency of the system:

$$\omega_0 = 2\pi f_0 = \sqrt{\frac{K}{M}} \quad (\text{rad/sec})$$

resonance frequency in Hz

## One-Mass Oscillator: Solutions (cont.)

General solution of the differential equation?

By trying, found the solutions:

$$x = e^{j\omega_0 t} \quad \text{and} \quad x = e^{-j\omega_0 t}$$

Turns out that general (complex-valued) solution is:

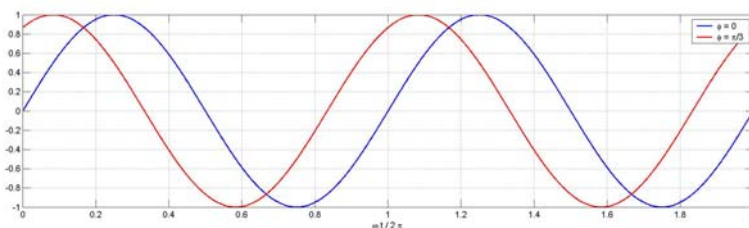
$$x = C_1 e^{+j\omega_0 t} + C_2 e^{-j\omega_0 t}$$

General **real-valued** solution can be written:

$$x = A \sin(\omega_0 t + \Phi)$$

amplitude

phase



## Damped Oscillator

### Adding a damping term

Damping force:

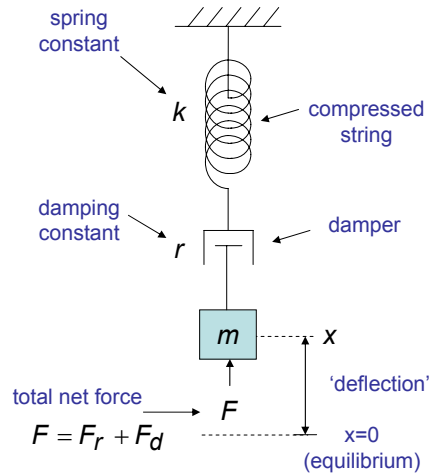
$$F_d = -r \dot{x}$$

Apply Newton's 2<sup>nd</sup> Law:

$$m \ddot{x} = F = F_r + F_d$$

Equation of motion:

$$m \ddot{x} + r \dot{x} + k x = 0$$



## Damped Oscillator: Solutions

Solution to the equation of motion:

$$Ms^2x + Rsx + Kx = 0$$

$$\left( Ms^2 + Rs + K \right) x = 0$$

for all  $x$ , thus

$$Ms^2 + Rs + K = 0$$

$$s^2 + \left( \frac{R}{M} \right) s + \left( \frac{K}{M} \right) = 0$$

$$\omega_0 = \sqrt{\frac{K}{M}} \quad \text{"natural frequency"}$$

$$\alpha = \frac{R}{2M} \quad \text{"decay rate"}$$

$$\omega_r = \sqrt{\omega_0^2 - \alpha^2} \quad \text{"resonance frequency"}$$

$$s = \frac{-\left( \frac{R}{M} \right) \pm \sqrt{\left( \frac{R}{M} \right)^2 - 4 \left( \frac{K}{M} \right)}}{2} = -\alpha \pm j\omega_r$$

## Damped Oscillator: Solutions (cont.)

$$s = \frac{-\left(\frac{R}{M}\right) \pm \sqrt{\left(\frac{R}{M}\right)^2 - 4\left(\frac{K}{M}\right)}}{2} = -\alpha \pm j\omega_r$$

As long as damping is sufficiently small:

$$\left(\frac{R}{M}\right)^2 < 4\left(\frac{K}{M}\right)$$

If the damping (R) is so high that the square root operation is performed on a real number, s becomes real, which means no oscillation! The system is then called **overdamped**.

## Damped Oscillator: Solutions (cont.)

General solution of the differential equation

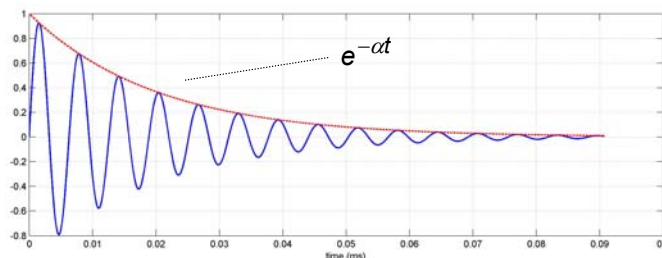
$$x = C_1 e^{-\alpha t} e^{+j\omega_r t} + C_2 e^{-\alpha t} e^{-j\omega_r t}$$

$$= e^{-\alpha t} \underbrace{\left( C_1 e^{+j\omega_r t} + C_2 e^{-j\omega_r t} \right)}_{\text{un-damped solution with } \omega_r}$$

↙  
exponential decay

Real-valued solution:

$$x = A e^{-\alpha t} \sin(\omega_r t + \Phi)$$



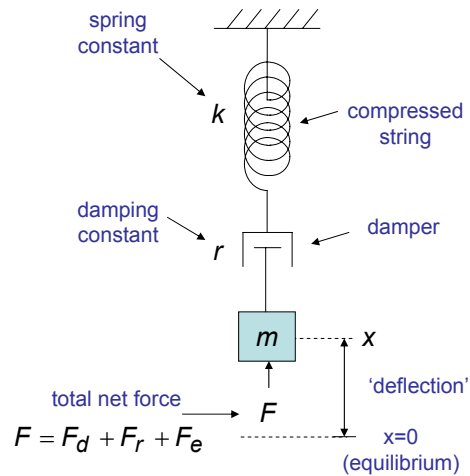
## Driven Oscillator

Add external force:

$$m \ddot{x} = F_d + F_r + F_e$$

Equation of motion:

$$m \ddot{x} + r \dot{x} + k x = F_e$$



## Frequency Response

Sinusoidal driving force:

$$m \ddot{x} + r \dot{x} + k x = A e^{j\omega t}$$

Try

$$x = B e^{j\omega t}$$

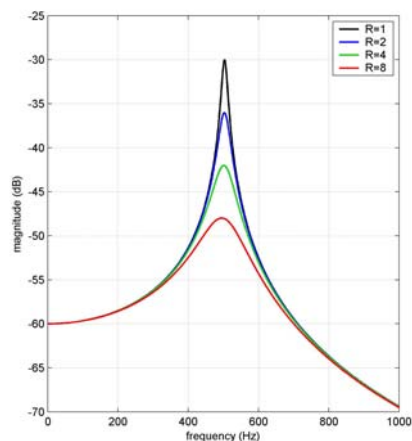
$$\left( -m \omega^2 + j\omega r + k \right) B e^{j\omega t} = A e^{j\omega t}$$

Frequency response:

$$H(\omega) = \frac{x}{F} = \frac{B e^{j\omega t}}{A e^{j\omega t}}$$

$$= \frac{1}{k + j\omega r - m \omega^2}$$

Frequency response plot (magnitude)





## Impulse Response (Time-Domain)

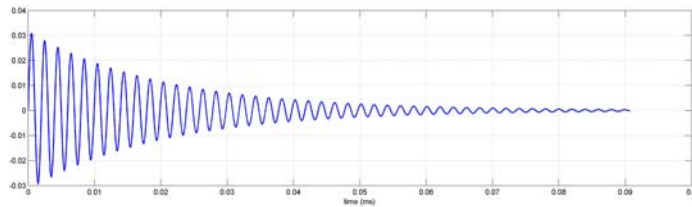
Frequency response = Fourier transform (FT) of impulse response function

$$H(\omega) = \frac{X(\omega)}{F(\omega)} = \frac{1}{k + j\omega r - m\omega^2}$$

$$h(t) = FT^{-1}\{H(\omega)\} \text{ (Inverse Fourier Transform)}$$

$$h(t) = \left( \frac{e^{-\alpha t}}{m\omega_r} \right) \sin(\omega_r t)$$

Impulse response = output response (x) to a single input impulse at  $t=0$ .



## Two-Mass Systems

Total net force on first mass:

$$F_1 = -k x_1 + k(x_2 - x_1)$$

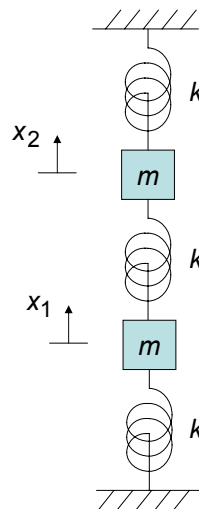
Total net force on second mass:

$$F_2 = -k x_2 + k(x_1 - x_2)$$

Thus the **equations of motion** are:

$$m \ddot{x}_1 + k x_1 + k(x_1 - x_2) = 0$$

$$m \ddot{x}_2 + k x_2 + k(x_2 - x_1) = 0$$



## Two-Mass Systems (cont.)

Assume that both masses oscillate with the same frequency  $\omega$  :

$$\begin{aligned} x_1 &= A_1 e^{st} \\ x_2 &= A_2 e^{st} \end{aligned} \quad (s = j\omega)$$

Substituting into equations of motion, we get, after some algebra:

$$\begin{aligned} (s^2 + 2\omega_0^2)A_1 - A_2\omega_0^2 &= 0 \\ (s^2 + 2\omega_0^2)A_2 - A_1\omega_0^2 &= 0 \end{aligned} \quad \left( \omega_0^2 = k/m \right)$$

By eliminating  $A_1$  and  $A_2$ , we find:

$$s^4 + 4\omega_0^2 s^2 + 3\omega_0^4 = 0$$

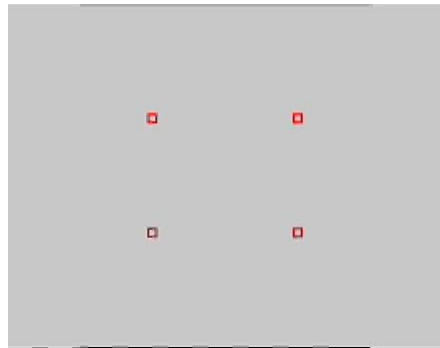
## Two-Mass Systems (cont.)

This can be written:

$$(\omega^2)^2 - 4\omega_0^2(\omega^2) + 3\omega_0^4 = 0$$

and has the solutions:

$$\omega^2 = 2\omega_0^2 \pm \omega_0^2$$



There are thus two possible **modes of vibration**:

$$\begin{aligned} \omega = \omega_0 : (\omega_0^2 - 2\omega_0^2)A_1 + \omega_0^2 A_2 &= 0 \Rightarrow A_1 = A_2 \quad \rightarrow \text{masses moving in: same direction} \\ \omega = \sqrt{3}\omega_0 : (3\omega_0^2 - 2\omega_0^2)A_1 + \omega_0^2 A_2 &= 0 \Rightarrow A_1 = -A_2 \quad \rightarrow \text{opposite direction} \end{aligned}$$

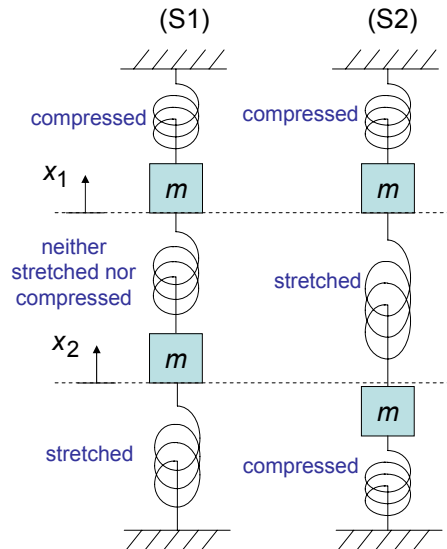
## Two-Mass Systems (cont.)

General real-valued solution is a superposition of the two **normal modes of vibration**:

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} A & B \\ A & -B \end{bmatrix} \cdot \begin{bmatrix} \sin(\omega_0 t) \\ \sin(\sqrt{3}\omega_0 t) \end{bmatrix}$$

Where the **modal weights**  $A$  and  $B$  depend on initial conditions (in case of free vibration), or on excitation position (in case of forced vibration)

**Example:** in initial scenario (S1) we have  $B = 0$ , and in initial scenario (S2) we have  $A = 0$ . In all other scenarios (except  $A=B=0$ ), the motion of the each mass is a combination of two modal vibrations.



## More Complex Systems

Two-mass systems can be more complex; different spring constants, different masses, transversal as well as longitudinal motion.

Multi-mass systems possible  $\Rightarrow N$  masses means  $N$  modes of vibration. Solution strategy is the same as for two-mass systems.

Continuous systems can be described as a system with infinite number of mass elements. They have an infinite number of modes.

Modes can be tested experimentally: modal analysis.

Modes of complex geometrical systems can also be determined via numerical simulation: finite element or finite difference method.

## Summary

- Learned how basic forces act on a point mass.
- Understood how an equation of motion for a point mass (with spring and/or damper) is constructed from Newton's 2<sup>nd</sup> law.
- Realised how an essential physical meaning of a solution-parameter of a physical system (one-mass oscillator) can be derived from dimensional analysis.
- Saw that a one-mass-oscillator exhibits 'free vibration' at its resonance frequency, which equals  $\sqrt{K/M}$ .
- Noticed and understood that damping forces impose an exponential decay onto the solution.
- Learned the formulas for the frequency response and impulse response.
- Learned about the formulation and solutions of a two-mass system.
- Learned that  $N$ -mass systems have  $N$  resonances.