



Finite Strains and Hyper-elasticity

Brief Introduction

July 10th 2019

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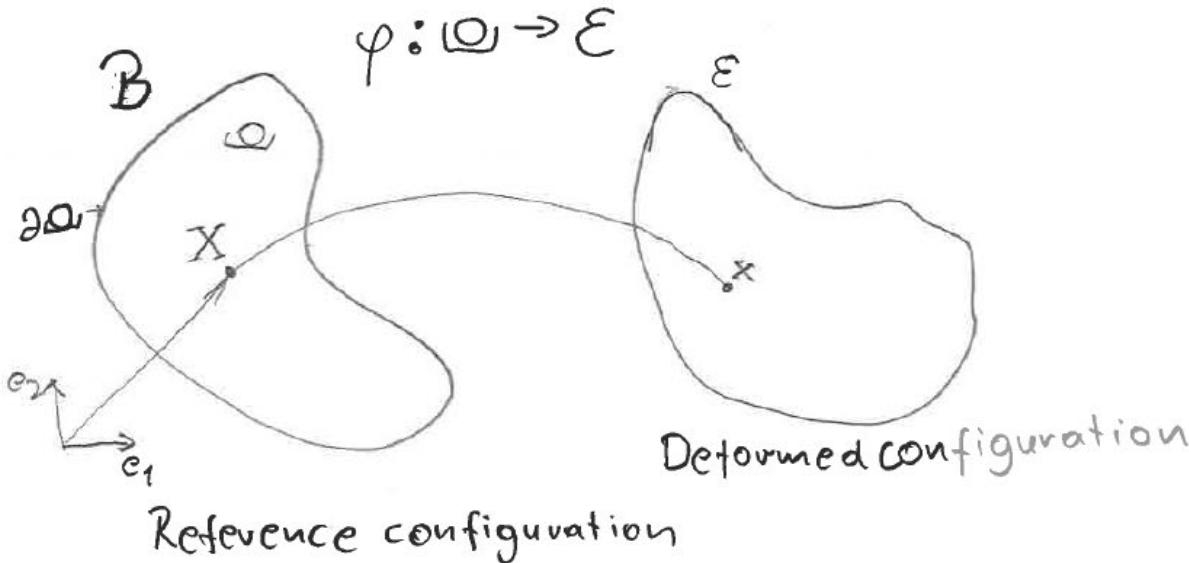
Literature

1. Non-Linear Elastic Deformations, R.W. Ogden
2. Computational Methods for Plasticity, de Souza Neto, D. Perić, DRJ Owen
3. Introduction to the Mechanics of a Continuous Medium, L.E. Malvern

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Material and Spatial Configurations

Deformation Gradient



$\varphi \rightarrow$ is a smooth one to one function that maps each material point X into a point $x = \varphi(X)$

Displacement of X is defined: $u(\bar{X}) = \varphi(\bar{X}) - X$

Special cases:

$$\varphi(X) = X + u \quad \text{Rigid translation}$$

$$\varphi(X) = q + R(X-q) \quad \text{Rigid rotation about } q; \\ R \text{ is proper orthogonal tensor}$$

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Deformation Gradient: What Is It?

Deformation gradient is the second order tensor F defined by:

$$F(X) = \nabla_X \varphi(X) = \frac{\partial X}{\partial \tilde{X}}$$

if you choose
to remember
just one equation
from this class,
remember this one.
BB

Gradient Definition:
gradient of scalar ϕ :

$$\frac{d\vec{\phi}}{ds} = \nabla \vec{\phi} \cdot \vec{u}$$

where $\vec{u} = \frac{d\vec{p}}{ds}$
(unit vector)

- Note:
1. In the definition above, we calculate gradient of a vector field
 2. Gradient of a vector field is a tensor

From the definition:

$$dx = F d\tilde{x}$$

deformation gradient is the linear operator that relates $d\tilde{x}$ (vector) into dx (vector). This is a definition of second order tensor!

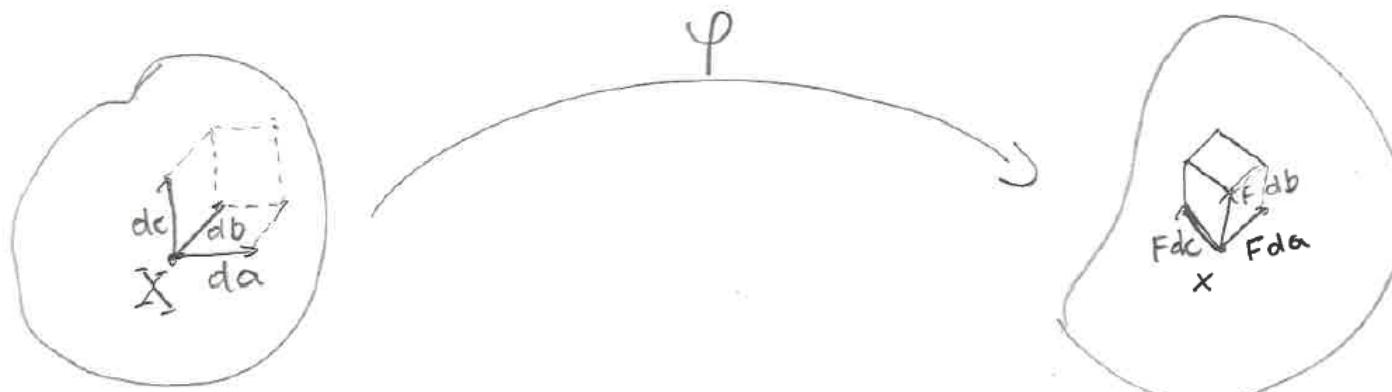
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Volume Change

Let us start with:

$$dx = F dx$$

← this equation defines how
“fiber” element is deformed
(remember it can include
rotations and translations)



$$dv_0 = (\vec{da} \times \vec{db}) \cdot \vec{dc}$$

Special cases:

$\gamma = 1$ - isochoric deformation

$F = \alpha I$ - volumetric def.

$$\gamma = \det F = \frac{dv}{dv_0} \quad (\text{see next page})$$

$$\int dv_0 = dv$$

$$dv = (F \vec{da} \times F \vec{db}) \cdot F \vec{dc}$$

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Volume Change

Exercise 1

Show that for any tensor T and set $\vec{u}, \vec{v}, \vec{w}$ of
Linearly independent vectors:

$$\det T = \frac{(T\vec{u} \times T\vec{v}) \cdot T\vec{w}}{(\vec{u} \times \vec{v}) \cdot \vec{w}} \quad (1)$$

Write (1) in index notation.

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Polar Decomposition

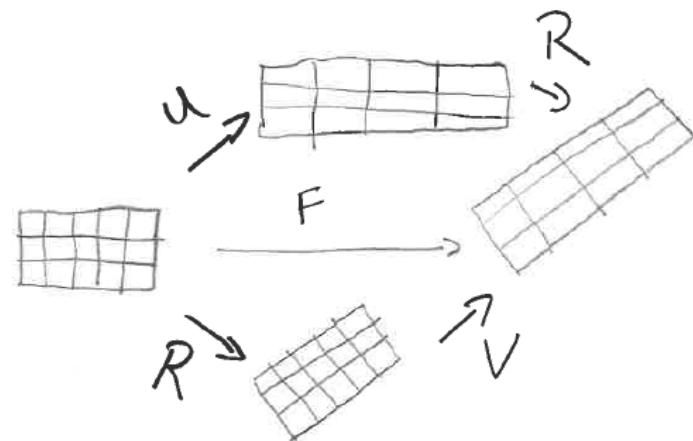
Polar Decomposition of Def. Gradient F :

$$F = R U = V R$$

R - Local rotation tensor

U - right stretch tensor

V - Left stretch tensor



$$U = \sqrt{C} \quad V = \sqrt{B}$$

$$C = F^T F \quad B = F F^T$$

Right Cauchy-Green
strain tensor

Left Cauchy-Green
strain tensor

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Polar Decomposition: Principal Stretches

Exercise 2

Assume the following deformation:

$$\varphi: \begin{cases} x_1 = X_1 \lambda_1 \cos\alpha - X_2 \lambda_2 \sin\alpha \\ x_2 = X_1 \lambda_1 \sin\alpha + X_2 \lambda_2 \cos\alpha \end{cases}$$

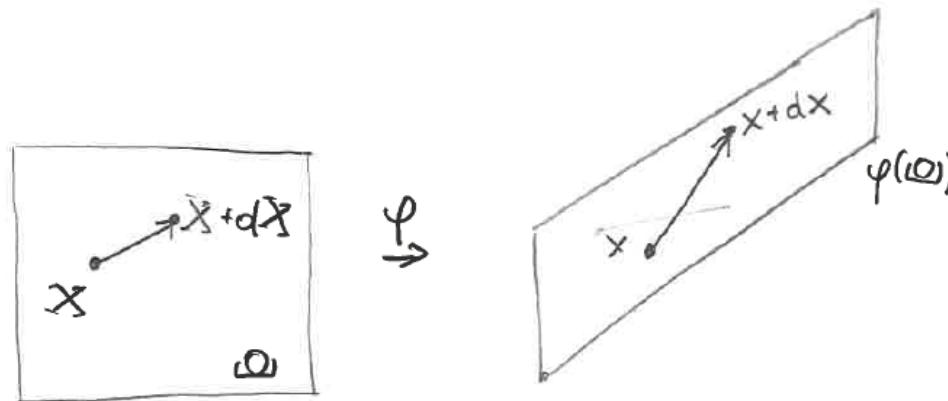
Calculate \bar{F}, R, U, V .

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Strain Measures

$$\|\boldsymbol{x}\|^2 = \mathbf{F} d\mathbf{X} \cdot \mathbf{F} d\mathbf{X}$$

Deformed Length of the material "fibre"



$$\begin{aligned}\|\boldsymbol{x}\|^2 &= C d\mathbf{X} d\mathbf{X} \\ &= (\mathbf{I} + 2 \mathbf{E}^{(2)}) d\mathbf{X} d\mathbf{X}\end{aligned}$$

$$E^{(2)} = \frac{1}{2} (C - \mathbf{I}) \leftarrow \text{Green Lagrange strain}$$

In general, Lagrangian strain tensors are defined:

$$E^{(m)} = \begin{cases} \frac{1}{m} (\mathbf{U}^m - \mathbf{I}) & m \neq 0 \\ \ln(\mathbf{U}) & m = 0 \end{cases}$$

Similar method can be used to define Eulerian strain tensors with respect to \mathbf{V} .

$m = 1 \leftarrow \text{Biot strain}$

$m = 0 \leftarrow \text{Logarithmic strain}$

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Stress Measures

$$\text{Cauchy stress : } \vec{t}(\vec{x}, \vec{n}) = \bar{\sigma}(\vec{x}) \vec{n} \quad \leftarrow \begin{array}{l} \text{Cauchy theorem} \\ (\text{momentum balance}) \end{array}$$

\vec{t} - surface force
 \vec{n} - area element normal
 $\bar{\sigma}$ - second order tensor (Cauchy stress)

Note that \vec{t} is the force exerted across a material surface per unit deformed area.

\vec{f} → force that acts across any surface whose normal is \vec{n} in the deformed configuration per unit reference area

$$\vec{t} = \frac{da}{da_0} \vec{t}$$

$$\bar{P} = \bar{\gamma} G F^{-T} \leftarrow \text{Piola-Kirchhoff stress (nominal stress)}$$

$\bar{\gamma} = \det F$

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Green Elasticity (Hyperelasticity)

In hyperelasticity, there exists a strain-energy function defined on the space of deformation gradients:

$$P = \frac{\partial W}{\partial F} \quad \leftarrow \text{nominal stress}$$

$$\bar{\sigma} = J^{-1} \frac{\partial W}{\partial F} F^T \quad \leftarrow \text{Cauchy stress}$$

$W(F)$ ← strain energy function (stored energy function)

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Strain Energy Function Invariant Formulation

Isotropic tensor-valued function of a tensor in 3D can be represented if the function is isotropic:

$$Y(X) = \alpha_0 I + \alpha_1 X + \alpha_2 X^2$$

if where α_i are scalar valued functions of the principal invariants of X .

Using an assumption that the material is isotropic $F = FQ$ it can be shown:

$$W(F) = \tilde{W}(B)$$

$$\tilde{W} = \tilde{W}(I_0(B), I_1(B), I_3(B))$$

representation of
energy function
by principal invariants
 I_1, I_2, I_3

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Strain Energy Function Principal Stretches Formulation

Similarly, stress energy density function can be represented using principal stretches :

$$\underline{W(F) = \hat{W}(\lambda_1, \lambda_2, \lambda_3)}$$

λ_i - principal stretches

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Strain Energy Function

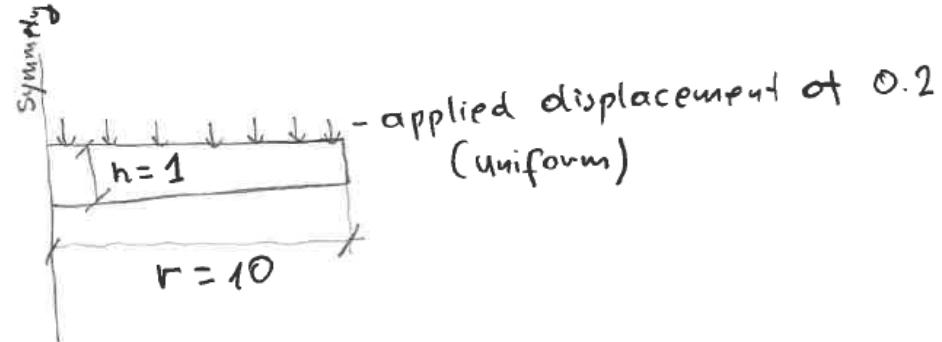
Exercise 3

Calculate Cauchy stress if W is represented using invariant and principal stretches.

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Homework 1

Use Abaqus FE package, to model unidirectional stress compression of a thin disk ($r=10$, $h=1$). Apply uniform displacement to the top surface $\Delta h = 0.2$.



Use hyperelastic Mooney-Rivlin material with:

- $C10 = 10$, $C01 = 1$, assume incompressibility
- $C10 = 10$, $C01 = 1$, $D = 0.01$

Calculate principal stretches in a disk after deformation and lateral displacement.

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