# Supplement to: Global sonde datasets do not support a mesoscale transition in the turbulent energy cascade

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### S1 A generalized anisotropic structure function

The two-dimensional structure function represented by Eqn. 9 in the main text is not the unique choice if we simply require Eqns. 7 to be recovered when  $\Delta x = 0$  or  $\Delta z = 0$ . To obtain Eqn. 9, we added the requirement that isotropic turbulence is recovered when  $H_h = H_v$  and  $\varphi_v = \varphi_h$ . Here we drop this assumption by adding an additional parameter  $\eta$  such that

$$\langle \Delta v(\Delta x, \Delta z)^2 \rangle = \left( \varphi_h^{\eta/H_h} \Delta x^{2\eta} + \varphi_v^{\eta/H_h} \Delta z^{H_v \eta/H_h} \right)^{H_h/\eta}. \tag{S1}$$

When  $\eta = 1$ , Eqn. 9 in the main text is recovered. Otherwise, Eqns. 7 are still obtained when  $\Delta x = 0$  or  $\Delta z = 0$ , but isotropic turbulence cannot be recovered. Visually, the parameter  $\eta$  makes the isolines of  $\Delta v$  more or less "boxy," as shown in Fig. S1.

Fitting Eqn. S1, rather than Eqn. 9, to the empirical structure function (Fig. S2) resulted in best-fit values of  $\phi_h = 0.015 \pm 0.007$ ,  $\phi_v = 0.006 \pm 0.001$ ,  $H_h = 0.34 \pm 0.02$ ,  $H_v = 0.65 \pm 0.01$ , and  $\eta = 0.81 \pm 0.07$ , which are close to the values reported in the main text for  $\eta = 1$ . The contour lines of the fit in Fig. S2 also provide a slightly better match to the empirical contour lines as compared to Fig. 7.

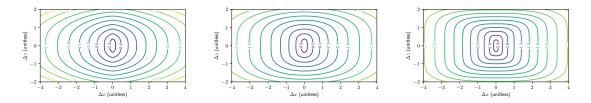


Figure S1: Countour plot of dimensionless form of Eqn. S1, as in Fig. 2 but for  $\eta = 0.75$ ,  $\eta = 1$ , and  $\eta = 1.5$ . The value  $\eta = 1$  was used in the main text.

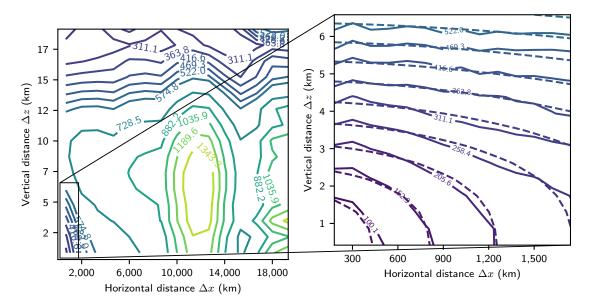


Figure S2: As in Fig. 7, but using Eqn. S1 for the fit.

#### S2 First and third order structure functions

In the main text we prefer second-order structure functions of velocity increments, which represent kinetic energy  $\Delta v^2$ , because they are most commonly examined in prior studies of turbulence. However, structure functions may be defined for arbitrary orders q as

$$\langle \Delta v^q \rangle \sim \Delta r^{\zeta(q)}$$
. (S2)

For a nonintermittent turbulent cascade, the structure function exponent  $\zeta(q)$  is simply equal to qH. For the more realistic intermittent case, in general  $\zeta(q) < qH$  for q>1 (Lovejoy and Schertzer, 2013). If the intermittency is weak, we can reasonably infer H for structure functions with q>1 using the equation  $H\approx \zeta(q)/q$  with the expectation that this method may slightly underestimate H. The main text used this method for q=2. In Figs. S3-S6 we reproduce the Figs. 4-8 for first order structure functions, where it is expected that  $\zeta=H$  for a multiplicative cascade (Lovejoy and Schertzer, 2013), and for q=3 for comparison with other prior studies. Table S1 lists the calculated exponents derived from structure functions calculated for q=1, q=2, and q=3.

Table S1: Structure function exponents for orders q = 1, 2, 3 (Eqn. S2) and estimated values for the Hurst exponent.

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Dataset	Direction	Order $(q)$	Structure Function Exponent $\zeta(q)$	Inferred $H = \zeta(q)/q$
IGRA	Vertical	1	$0.62 \pm 0.02$	$0.62 \pm 0.02$
		2	$1.24 \pm 0.04$	$0.62 \pm 0.02$
		3	$1.80 \pm 0.06$	$0.60 \pm 0.02$
ACTIVATE	Vertical	1	$0.70 \pm 0.01$	$0.70 \pm 0.01$
		2	$1.43 \pm 0.03$	$0.71 \pm 0.01$
		3	$2.12 \pm 0.04$	$0.71 \pm 0.01$
Hurricanes	Vertical	1	$0.523 \pm 0.007$	$0.523 \pm 0.007$
		2	$1.03 \pm 0.02$	$0.513 \pm 0.008$
		3	$1.51 \pm 0.03$	$0.504 \pm 0.009$
IGRA	Horizontal	1	$0.49 \pm 0.02$	$0.49 \pm 0.02$
		2	$1.00 \pm 0.04$	$0.50 \pm 0.02$
		3	$1.44 \pm 0.04$	$0.48 \pm 0.01$

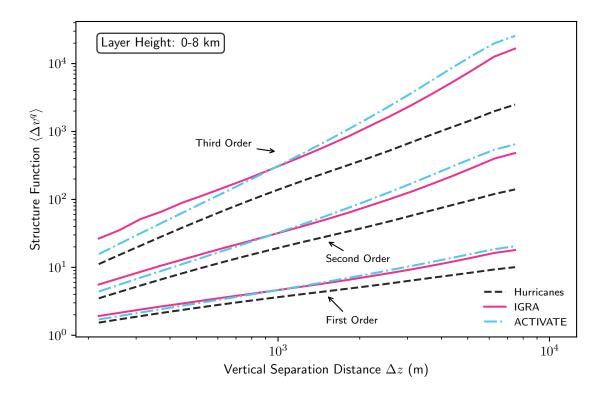


Figure S3: As in Fig. 4, but for first-, second-, and third-order structure functions  $\langle \Delta v^2 \rangle$  (lower),  $\langle \Delta v^2 \rangle$  (middle) and  $\langle \Delta v^3 \rangle$  (upper).

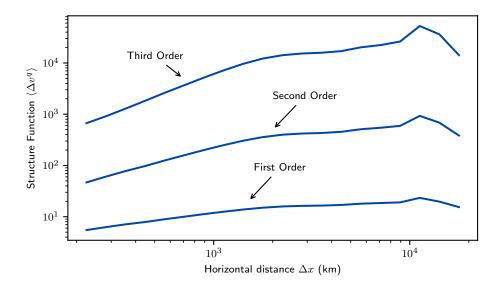


Figure S4: As in Fig. 6, but for first-, second-, and third-order structure functions  $\langle \Delta v^2 \rangle$  (lower),  $\langle \Delta v^2 \rangle$  (middle) and  $\langle \Delta v^3 \rangle$  (upper).

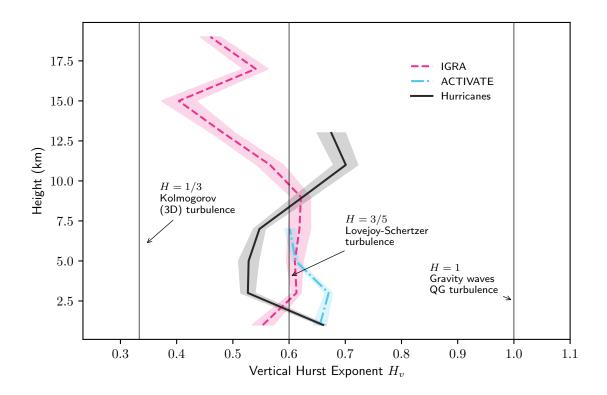


Figure S5: As in Fig. 5, but for first-order structure functions  $\langle \Delta v \rangle$ .

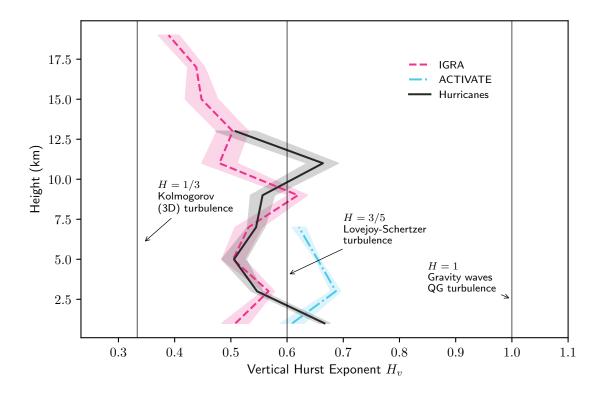


Figure S6: As in Fig. 5, but for third-order structure functions  $\langle \Delta v^3 \rangle$ .

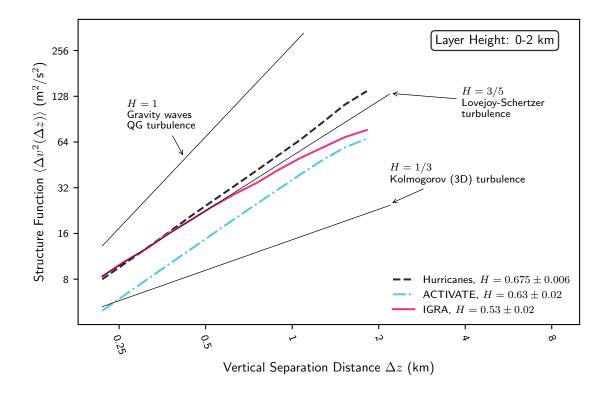


Figure S7: As in Fig. 4, but for altitudes between 0 and 2 km.

## S3 Structure functions for individual 2 km-thick layers

Figures S7 to S16 display the structure functions for each altitude layer that were used to calculate  $H_v$  as a function of height in Fig. 5.

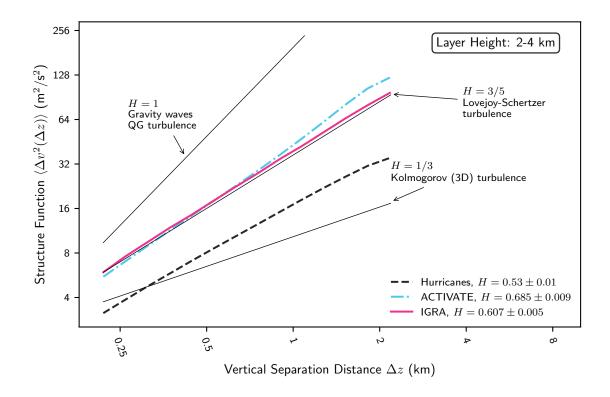


Figure S8: As in Fig. 4, but for altitudes between 2 and 4 km.

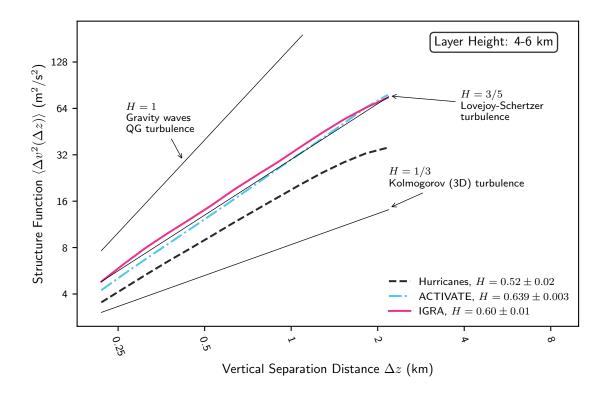


Figure S9: As in Fig. 4, but for altitudes between 4 and 6 km.

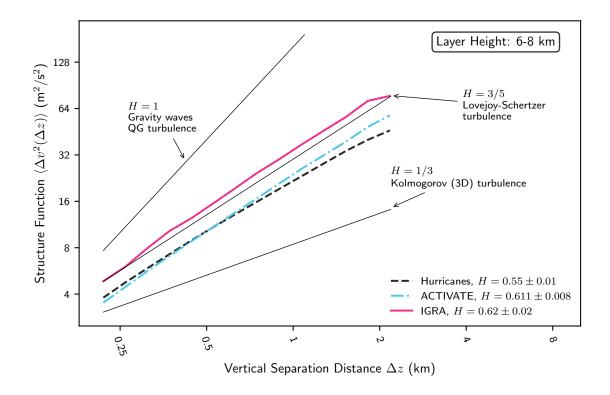


Figure S10: As in Fig. 4, but for altitudes between 6 and 8 km.

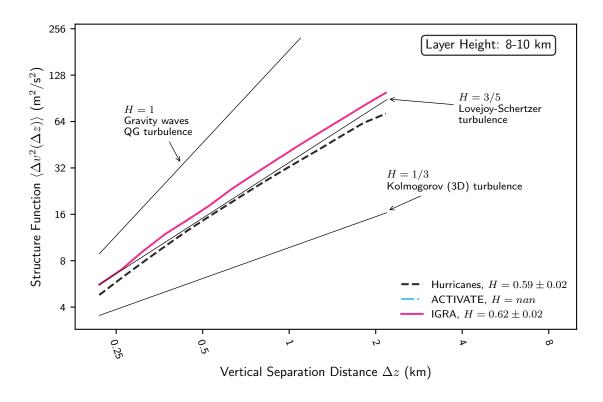


Figure S11: As in Fig. 4, but for altitudes between 8 and 10 km.

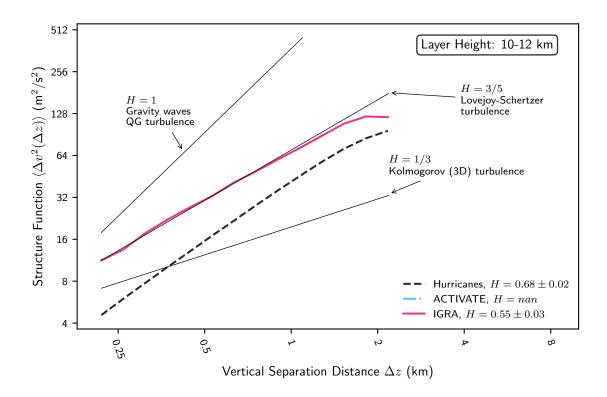


Figure S12: As in Fig. 4, but for altitudes between 10 and 12 km.

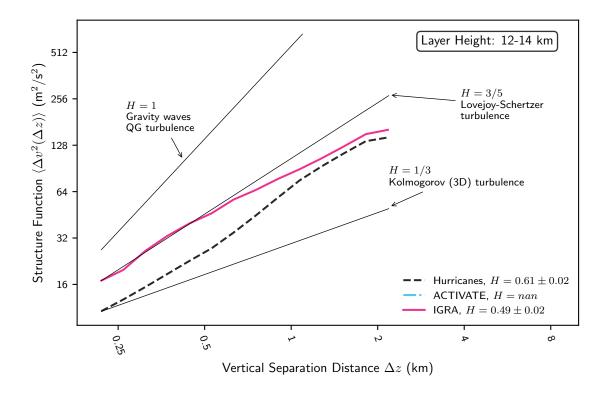


Figure S13: As in Fig. 4, but for altitudes between 12 and 14 km.

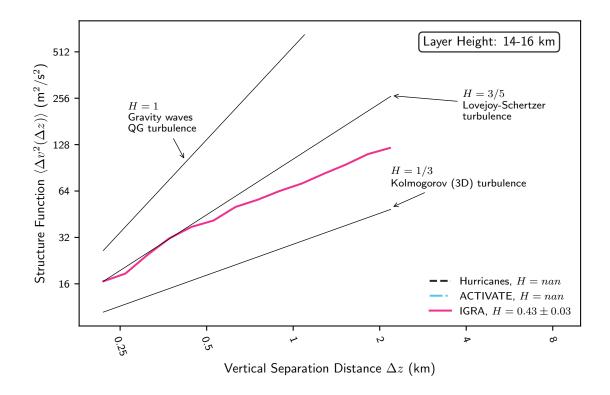


Figure S14: As in Fig. 4, but for altitudes between 14 and 16 km.

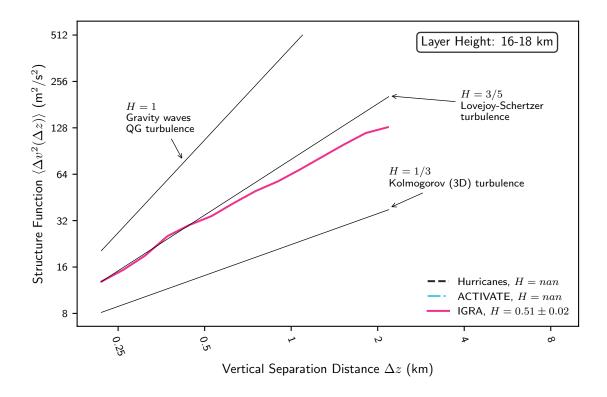


Figure S15: As in Fig. 4, but for altitudes between 16 and 18 km.

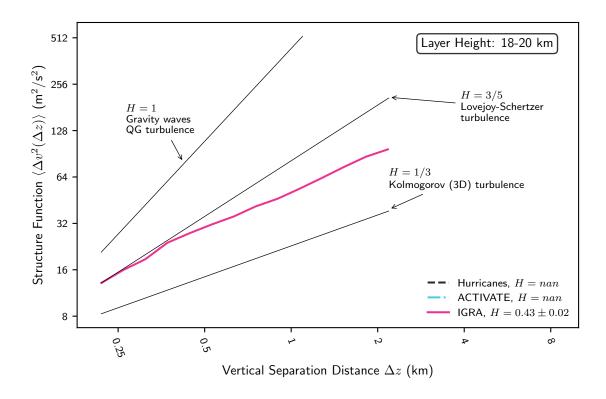


Figure S16: As in Fig. 4, but for altitudes between 18 and 20 km.

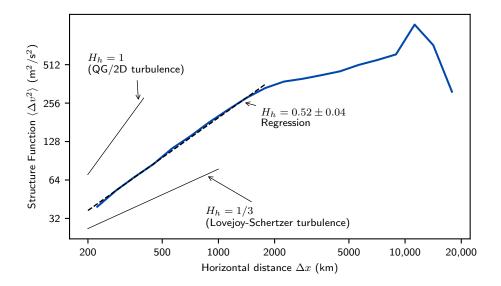


Figure S17: As in Fig. 6, but for observation pairs separated by at most 5 minutes and 5 m in altitude.

#### S4 Additional horizontal structure functions

In this section we report four additional horizontal structure functions calculated from IGRA data. Figure S17 shows a horizontal structure function as in Fig. 6 but for observation pairs separated by at most 5 minutes and 5 m in altitude as described in Section 3.1.

Figure S18 shows a structure function calculated for observation pairs between -20° and 20° in latitude, while Fig. S19 shows a structure function calculated for observation pairs between 45° and 90° in latitude.

Figures S20 and S21 show horizontal structure functions for temperature variance  $T^2$  and pressure variance  $p^2$ , respectively, rather than kinetic energy.

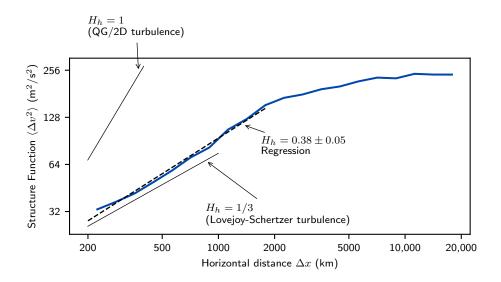


Figure S18: As in Fig. 6, but for observation pairs between -20° and 20° in latitude.

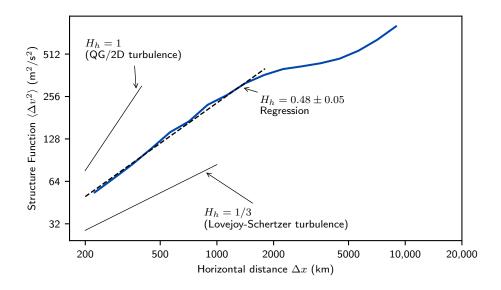


Figure S19: As in Fig. 6, but for observation pairs between 45° and 90° in latitude.

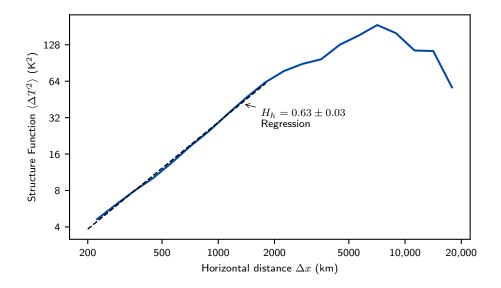


Figure S20: As in Fig. 6, but for temperature T rather than kinetic energy, where  $\Delta T^2 \sim \Delta x^{2H_h}$ .

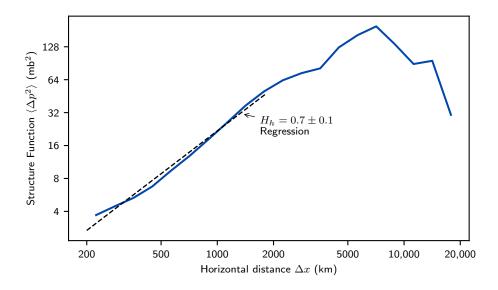


Figure S21: As in Fig. 6, but for pressure p rather than kinetic energy, where  $\Delta p^2 \sim \Delta x^{2H_h}$ .

## References

Lovejoy, S. and Schertzer, D.: The weather and climate: emergent laws and multifractal cascades, Cambridge University Press, 2013.