

# Numerical solution of the stationary convection-diffusion equation in 1D (XIV)

## The numerical scheme

$$\begin{aligned}F_{j+1/2} &= \beta_j^0 \varphi_j + \beta_{j+1}^1 \varphi_{j+1}, \quad j = 1, \dots, J-1, \\F_{1/2} &= \beta_1^1 \varphi_1 + \gamma_0, \quad F_{J+1/2} = \beta_J^0 \varphi_J + \gamma_1.\end{aligned}$$

$$L_h \varphi_j = F_{j+1/2} - F_{j-1/2} \Rightarrow (2.31) :$$

$$L_h \varphi_j = \alpha_j^{-1} \varphi_{j-1} + \alpha_j^0 \varphi_j + \alpha_j^1 \varphi_{j+1} = \tilde{q}_j, \quad j = 1, \dots, J,$$

with  $\alpha_1^{-1} = \alpha_J^1 = 0$ .

This is called the numerical or finite volume scheme.

# Numerical solution of the stationary convection-diffusion equation in 1D (XV)

General form of a linear scheme:

$$L_h \varphi_j = \sum_{k \in K} \alpha_j^k \varphi_{j+k} = \tilde{q}_j ,$$

with  $K$  some index set; in our case

$$K = \{-1, 0, 1\}$$

**Stencil**  $[L_h]$  of the operator  $L_h$ :

$$[L_h]_j = [ \alpha_j^{-1} \quad \alpha_j^0 \quad \alpha_j^1 ] .$$

# Numerical solution of the stationary convection-diffusion equation in 1D (XVI)

$$Ay = b, \quad y = \begin{bmatrix} \varphi_1 \\ \vdots \\ \varphi_J \end{bmatrix}, \quad b = \begin{bmatrix} \tilde{q}_1 \\ \vdots \\ \tilde{q}_J \end{bmatrix},$$

$A$  is tridiagonal matrix:

$$A = \begin{bmatrix} \alpha_1^0 & \alpha_1^1 & 0 & \cdots & 0 \\ \alpha_2^{-1} & \alpha_2^0 & \alpha_2^1 & & \vdots \\ 0 & \ddots & \ddots & \ddots & 0 \\ \vdots & & \alpha_{J-1}^{-1} & \alpha_{J-1}^0 & \alpha_{J-1}^1 \\ 0 & \cdots & 0 & \alpha_J^{-1} & \alpha_J^0 \end{bmatrix}.$$

# Numerical solution of the stationary convection-diffusion equation in 1D (XVII)

## Two important questions

- How well does the numerical solution approximate the exact solution?
- How accurately and efficiently can we solve the linear algebraic system?

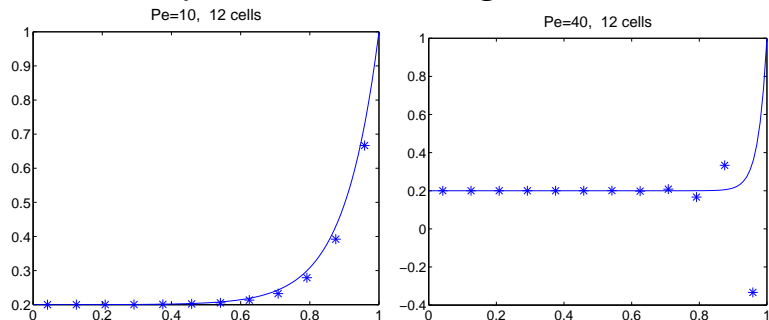
# Numerical solution of the stationary convection-diffusion equation in 1D (XII)

Lets zoom in on the behaviour of the algorithm for different values of the Péclet number and see where the behaviour changes...

- ▶ Two-sided Dirichlet boundary condition
- ▶ Uniform mesh
- ▶ Convection: Central scheme I (averaging)
- ▶ Diffusion: Central discretisation.

# Numerical solution of the stationary convection-diffusion equation in 1D (XIX)

## Numerical experiments on uniform grid



Exact solution (—) and numerical solution (\*).

Why the *wiggles* for  $Pe = 40$ ?

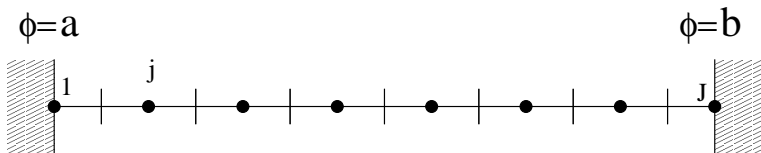
# Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (I)

A sidestep to the vertex-centered case:

For a conservative scheme  $[L_h]_j = [\alpha_j^{-1} \quad \alpha_j^0 \quad \alpha_j^1]$  on a uniform grid the dependence on  $j$  disappears and  $\alpha^0 = -(\alpha^{-1} + \alpha^{+1})$

$$\alpha^{-1}\varphi_{j-1} - (\alpha^{-1} + \alpha^{+1})\varphi_j + \alpha^{+1}\varphi_{j+1} = 0, \quad j = 1..J-1$$

$$\varphi_0 = a, \quad \varphi_J = b$$



In the uniform grid case the only difference between cell/vertex centered is in treatment of the boundary conditions..

## Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (II)

The analytical solution of this difference equation with the Dirichlet boundary conditions is given by:

$$\varphi_j = a + (b - a) \frac{1 - z^j}{1 - z^J}, \quad z = \frac{\alpha^{-1}}{\alpha + 1}, \quad j = 0, 1, \dots, J$$

Question: Depending on the coefficients  $a, b, \alpha^{-1}, \alpha^{+1}$ , when will the solution have oscillations?



# Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (III)

Central scheme I:

$$\begin{aligned}\alpha^{j-1} &= -\left(\frac{u}{2} + \frac{\varepsilon}{h}\right) = -\frac{\varepsilon}{h} \left(\frac{1}{2} \frac{uh}{\varepsilon} + 1\right) = \\ &\quad -\frac{\varepsilon}{h} \left(\frac{1}{2} Pe_h + 1\right) = -\frac{\varepsilon}{h} \left(1 + \frac{1}{2} Pe_h\right) \\ \alpha^{j+1} &= \left(\frac{u}{2} - \frac{\varepsilon}{h}\right) = -\frac{\varepsilon}{h} \left(-\frac{1}{2} \frac{uh}{\varepsilon} + 1\right) = \\ &\quad -\frac{\varepsilon}{h} \left(-\frac{1}{2} Pe_h + 1\right) = -\frac{\varepsilon}{h} \left(1 - \frac{1}{2} Pe_h\right)\end{aligned}$$

Divide by the common factor  $-\frac{\varepsilon}{h}$ ,  $Pe_h$  is the *mesh*-Péclet number.

$$\varphi_j = a + (b - a) \frac{1 - \left(\frac{1 + \frac{Pe}{2}}{1 - \frac{Pe}{2}}\right)^j}{1 - \left(\frac{1 + \frac{Pe}{2}}{1 - \frac{Pe}{2}}\right)^N}$$

## Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (IV)

The condition for an oscillatory solution that:

$$\frac{\alpha^{-1}}{\alpha^{+1}} < 0$$

Corresponds to

$$\frac{1 + \frac{1}{2}Pe_h}{1 - \frac{1}{2}Pe_h} < 0 \quad \Rightarrow \quad 1 - \frac{1}{2}Pe_h < 0 \quad \Rightarrow \quad Pe_h > 2$$

Lets check that...

## Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (V)

For large values of  $Pe_h$  we can even predict how the solution will look. Substitute  $\epsilon = 2/Pe_h$ .

$$\varphi_j = a + (b - a) \frac{1 - \left(\frac{1+1/\epsilon}{1-1/\epsilon}\right)^j}{1 - \left(\frac{1+1/\epsilon}{1-1/\epsilon}\right)^J}$$

$Pe_h$  is *large* and  $\epsilon$  very small: make an expansion in  $\epsilon$ .

## Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (VI)

$$\varphi_j = a + (b - a) \frac{1 - (-1)^j(1 + 2j\epsilon)}{1 - (-1)^J(1 + 2J\epsilon)} + \mathcal{O}(\epsilon^2)$$

Case 1:  $J$  odd:

$$\varphi_j \approx a + (b - a) \frac{1 - (-1)^j - 2(-1)^j j\epsilon}{2(1 + J\epsilon)}$$

$$j \text{ even } \phi_j \approx a - (b - a)j\epsilon; \quad j \text{ odd } \varphi_j \approx b - (b - a)(J - j)\epsilon$$

## Analytical solution of the *discretised* stationary convection-diffusion equation in 1D (VII)

$$\varphi_j = a + (b - a) \frac{1 - (-1)^j(1 + 2j\epsilon)}{1 - (-1)^J(1 + 2J\epsilon)} + \mathcal{O}(\epsilon^\infty)$$

Case 2:  $J$  even:

$$\varphi_j \approx a + (b - a) \frac{1 - (-1)^j(1 + 2j\epsilon)}{-2J\epsilon}$$

$$j \text{ even: } \phi_j \approx a - (b - a) \frac{j}{J}; \quad j \text{ odd: } \varphi_j \approx a - (b - a) \left( \frac{-1}{J\epsilon} - \frac{j}{J} \right)$$

# The discrete maximum principle (I)

We can also explain the behavior for  $P_h > 2$  from the properties of the operator: no wiggles means no local extrema: the discrete counterpart of the maximum principle.

$$\begin{aligned} c_n, x_n \in \mathbb{R} \quad n = 1, 2, \dots, N : \\ \left\{ \sum_{n=1}^N c_n = 0, \quad c_n < 0 \mid n > 1, \quad \sum_{n=1}^N c_n x_n \leq 0 \right\} \Rightarrow \\ x_n = x_1, \quad n = 1, \dots, N \quad \vee \quad x_1 < \max\{x_n : n > 1\} \end{aligned}$$

$x_j$  is the average of the neighbouring  $x_j$ 's...

# The discrete maximum principle (II)

Scheme in stencil notation:

$$L_h \varphi_j = \sum_{k \in K} \alpha_j^k \varphi_{j+k} = \tilde{q}_j, \quad j = 1, \dots, J.$$

The operator  $L_h$  is *of positive type* if

$$\sum_{k \in K} \alpha_j^k = 0, \quad j = 2, \dots, J-1$$

and

$$\alpha_j^k < 0, \quad k \neq 0, \quad j = 2, \dots, J-1$$

# The discrete maximum principle (III)

*Discrete maximum principle.*

If  $L_h$  is of positive type and

$$L_h \varphi_j \leq 0, \quad j = 2, \dots, J-1,$$

then  $\varphi_j \leq \max\{\varphi_1, \varphi_J\}$ .

Extrema can only occur on the boundary of the domain, just like in the continuous case!



## The discrete maximum principle (IV)

When is our scheme (Central Scheme I) of positive type?

$$[L_h] = \begin{bmatrix} -\frac{1}{2}u - \frac{\varepsilon}{h} & 2\frac{\varepsilon}{h} & \frac{1}{2}u - \frac{\varepsilon}{h} \end{bmatrix}$$

Of pos. type iff  $\alpha^{+1} < 0$ :

$$\frac{1}{2}u - \frac{\varepsilon}{h} < 0 \Rightarrow \frac{uh}{\varepsilon} < 2 \Rightarrow Pe_h < 2$$

So to avoid wiggles in the solution we should reduce the meshwidth such that  $h < 2/Pe$ . This is totally impossible in the limit  $Pe \Rightarrow \infty$   
Alternatives? How'bout this upwind scheme?

# The discrete maximum principle (V)

When is the upwind scheme of positive type?

$$u > 0 : \quad [L_h] = \begin{bmatrix} -u - \frac{\varepsilon}{h} & u + 2\frac{\varepsilon}{h} & -\frac{\varepsilon}{h} \end{bmatrix} .$$

$$u < 0 : \quad [L_h] = \begin{bmatrix} -\frac{\varepsilon}{h} & -u + 2\frac{\varepsilon}{h} & u - \frac{\varepsilon}{h} \end{bmatrix} .$$

The *Upwind Scheme* is of positive type for all  $Pe_h \odot$ ,  
but heavily overpredicts the thickness of the boundary layer in the  
solution  $\odot$

The *Upwind Scheme* corresponds to *Central Scheme I*, with  $\varepsilon$   
replaced by  $\varepsilon' = \varepsilon + \varepsilon_a$ ,  $\varepsilon_a = uh/2$ , which dominates the solution  
for  $Pe \gg 1$ .

# Local grid refinement (I)

A second alternative is to use *local grid refinement*

Reduce the grid size **only** in the business part of the domain:

In the **boundary layer** at the outflow boundary.

$$\text{Exact sol.} \quad \Rightarrow \quad \delta = \mathcal{O}(\varepsilon) = \mathcal{O}(\text{Pe}^{-1}).$$

Of course we locally need  $h < \delta$  to capture the boundary layer

Take  $\delta \equiv 6\varepsilon$ .

Put 6 equal cells in the boundary layer region  $\Omega_1 = [1 - \delta, 1]$   
and 6 equal cells in the rest:  $\Omega_2 = [0, 1 - \delta]$ , with  $\Omega = \Omega_1 \cup \Omega_2$ .

## Local grid refinement (II)

Lets compare the performance of the different schemes:

- ▶ Central Scheme I

$$\varphi_{j+1/2} = \frac{\varphi_{j+1} + \varphi_j}{2}$$

- ▶ Central Scheme II

$$\varphi_{j+1/2} = \frac{h_j \varphi_{j+1} + h_{j+1} \varphi_j}{h_j + h_{j+1}}$$

- ▶ Upwind Scheme

$$\varphi_{j+1/2} = \varphi_j, \quad (u > 0);$$

The coefficients of the stencil are now dependent on  $j$ ...

Central Scheme II will be more accurate, right?

## Local grid refinement (III)

Summary of results:

| Scheme            | Pe=10 | Pe=40 | Pe=400 | Pe=4000 |
|-------------------|-------|-------|--------|---------|
| Central Scheme I  | .0607 | .0852 | .0852  | .0852   |
| Upwind            | .0785 | .0882 | .0882  | .0882   |
| Central Scheme II | .0607 | .0852 | .0856  | .3657   |

Maximum error norm; 12 cells.

# Local grid refinement (IV)

## Central Scheme I

$$[L_h] = \begin{bmatrix} - & \frac{1}{2} - \frac{\varepsilon}{H} \\ & \frac{\varepsilon}{H} + \frac{2\varepsilon}{h+H} \\ & \frac{1}{2} - \frac{2\varepsilon}{h+H} \end{bmatrix}.$$

## Central Scheme II

$$[L_h] = \begin{bmatrix} - & \frac{1}{2} - \frac{\varepsilon}{H} \\ - & \frac{1}{2} + \frac{h}{h+H} + \frac{\varepsilon}{H} + \frac{2\varepsilon}{h+H} \\ & \frac{H}{h+H} - \frac{2\varepsilon}{h+H} \end{bmatrix}.$$

# Global and local truncation error (I)

**Definition** The *global truncation error* is defined as

$$e_j \equiv \varphi(x_j) - \varphi_j, \quad j = 1, \dots, J,$$

with  $\varphi(x)$  the exact solution.

Only in very few cases we can determine the global error directly, so we use another quantity to estimate the global error.

## Global and local truncation error (II)

**Definition** The *local truncation error* of the discrete operator  $L_h$  is defined as

$$\tau_j \equiv L_h(\varphi(x_j) - \varphi_j) \equiv L_h e_j, \quad j = 1, \dots, J.$$

$\varphi(x)$  is an *exact* solution of the continuous problem  $L\varphi(x) = q(x)$

$\varphi_j$  is an *exact* solution of the discretized problem  $L_h\varphi_j = q_j$

But not the otherway around...



## Global and local truncation error (III)

Let  $e = (e_1, \dots, e_J)^T$ ,  $\tau = (\tau_1, \dots, \tau_J)^T$ .

$$e = L_h^{-1} \tau.$$

$$\|e\| \leq \|L_h^{-1}\| \|\tau\|.$$

*This means that a scheme with smaller  $\|\tau\|$  will have necessarily a smaller  $\|e\|$ , because  $\|L_h^{-1}\|$  may be larger!!!*

This explains why it is possible that Central Scheme II (linear interpolation for convection) may be worse than averaging ☹

## Global and local truncation error (IV)

Estimate of local truncation error for the following scheme  $L_h \varphi_j$

$$L_h \varphi_j \equiv - \left( \frac{u}{2} + \frac{\varepsilon}{h_{j-1/2}} \right) \varphi_{j-1} + \varepsilon \left( \frac{1}{h_{j-1/2}} + \frac{1}{h_{j+1/2}} \right) \varphi_j \\ + \left( \frac{u}{2} - \frac{\varepsilon}{h_{j+1/2}} \right) \varphi_{j+1} = h_j q_j, \quad j = 2, \dots, J-1,$$

Although the exact sol. is unknown, nevertheless  $\tau$  can be estimated.

## Global and local truncation error (V)

Use **Taylor's formula** to approximate  $\varphi_{j\pm 1}$  to express  $L_h\varphi(x_j)$  in  $\varphi$  and  $\frac{d^n f(\xi)}{dx^n}(x_j)$

$$f(x) = f(x_0) + \sum_{k=0}^{n-1} \frac{1}{k!} \frac{d^k f(x_0)}{dx^k} (x - x_0)^k + \frac{1}{n!} \frac{d^n f(\xi)}{dx^n} (x - x_0)^n$$

## Global and local truncation error (VI)

$$\begin{aligned}\varphi(x_{j\pm 1}) = & \varphi(x_j) \pm h_{j\pm 1/2}\varphi^{(1)}(x_j) + \frac{1}{2}h_{j\pm 1/2}^2\varphi^{(2)}(x_j) \\ & \pm \frac{1}{6}h_{j\pm 1/2}^3\varphi^{(3)}(x_j) + \frac{1}{24}h_{j\pm 1/2}^4\varphi^{(4)}(x_j) + \mathcal{O}(h_{j\pm 1/2}^5) .\end{aligned}$$

Here  $\mathcal{O}$  is *Landau's order symbol*:

A function  $f(h) = \mathcal{O}(h^p)$  if there exist a constant  $M$  independent of  $h$  and a constant  $h_0 > 0$  such that

$$\frac{|f(h)|}{h^p} < M, \quad \forall h \in (0, h_0) .$$

## Global and local truncation error (VII)

Combining Taylor with the continuous operator:

$$L\varphi(x) = u\varphi_{,x} - \varepsilon\varphi_{,xx} = q(x)s \quad (1)$$

We can show that up to  $\mathcal{O}(\Delta^4)$  the exact solution of  $L$  satisfies  $\tilde{L}_h$

$$L_h\varphi(x_j) = \tilde{L}_h\varphi(x_j) + \mathcal{O}(\Delta^4) ,$$

$\varphi(x)$  is the exact sol.

$$\begin{aligned} \tilde{L}_h\varphi(x_j) &= \frac{1}{2}q_j(h_{j-1/2} + h_{j+1/2}) \\ &+ \left( \frac{1}{4}u\varphi^{(2)} - \frac{1}{6}\varepsilon\varphi^{(3)} \right) (h_{j+1/2}^2 - h_{j-1/2}^2) \\ &+ \left( \frac{1}{12}u\varphi^{(3)} - \frac{1}{24}\varepsilon\varphi^{(4)} \right) (h_{j+1/2}^3 + h_{j-1/2}^3) + \mathcal{O}(\Delta^4) . \end{aligned}$$

## Global and local truncation error (II)

The local truncation error is given as

$$\tau_j = L_h e_j = L_h[\varphi(x_j) - \varphi_j] = \tilde{L}\varphi(x_j) - h_j q_j + \mathcal{O}(\Delta^4) .$$

$$\begin{aligned} \tau_j &= \frac{1}{2} q_j (h_{j-1/2} - 2h_j + h_{j+1/2}) \\ &+ \left( \frac{1}{4} u \varphi^{(2)} - \frac{1}{6} \varepsilon \varphi^{(3)} \right) (h_{j+1/2}^2 - h_{j-1/2}^2) \\ &+ \left( \frac{1}{12} u \varphi^{(3)} - \frac{1}{24} \varepsilon \varphi^{(4)} \right) (h_{j+1/2}^3 + h_{j-1/2}^3) + \mathcal{O}(\Delta^4) . \end{aligned}$$

## Global and local truncation error (IX)

*Smooth* grid:

$$\begin{aligned} |h_{j+1/2} - h_{j-1/2}| &= \mathcal{O}(\Delta^2) \quad \text{and} \\ |h_{j-1/2} - 2h_j + h_{j+1/2}| &= \mathcal{O}(\Delta^3). \end{aligned}$$

Any grid which does not fulfill the conditions for a *Smooth* grid is defined to be a *Rough* grid.

$$\begin{aligned} \tau_j &= \frac{1}{2} q_j (h_{j-1/2} - 2h_j + h_{j+1/2}) \\ &+ \left( \frac{1}{4} u \varphi^{(2)} - \frac{1}{6} \varepsilon \varphi^{(3)} \right) (h_{j+1/2}^2 - h_{j-1/2}^2) \\ &+ \left( \frac{1}{12} u \varphi^{(3)} - \frac{1}{24} \varepsilon \varphi^{(4)} \right) (h_{j+1/2}^3 + h_{j-1/2}^3) + \mathcal{O}(\Delta^4). \end{aligned}$$

What will be the order of the local truncation error on a *Smooth* grid?

## Global and local truncation error (X)

Smooth grid:  $\tau_j = \mathcal{O}(\Delta^3)$ .

Rough grid:  $\tau_j = \mathcal{O}(\Delta)$ .

This is cause for common misunderstanding that smooth grids are required for good accuracy. *But this is not necessary in general.* Locally refined grid used in the preceding numerical experiments is rough, but nevertheless the accuracy was satisfactory. Again: local truncation error is not a good indicator of the accuracy of the solution.



# Global and local truncation error (XI)

Estimate of local truncation error at the boundaries (assume locally uniform grid)

$$\tau_1 = \frac{h}{4}\varepsilon\varphi^{(2)} + \mathcal{O}(h^2)$$

Interior, uniform grid:  $\tau_j = \mathcal{O}(h^3)$ .

*However, it is not necessary to improve the local accuracy near a Dirichlet boundary!*

Neumann at  $x = 1$ :

$$\tau_J = \frac{h^2}{24}\varepsilon\varphi^{(3)} + \mathcal{O}(h^3).$$

This also does not require improvement!

This follows from estimate of *global* truncation error.

# Estimating the global error (I)

The *global* error is directly related to the *local truncation* error:

$$L_h e = \tau$$

Try to find a *Barrier function*  $E$  such that:

$$L_h E \geq |\tau|$$

We will show that this implies  $|e| \leq E$ .

## Estimating the global error (II)

$$L_h(\pm e - E) \leq 0 .$$

We want to use the discrete maximum principle:

Therefore the scheme should be of positive type or  $Pe < 2$ .

The maximum occurs at Dirichlet boundary  $\Rightarrow$

$$\pm e_j - E_j \leq \pm e_1 - E_1, \quad j = 2, \dots, J. \quad (2.58)$$

Assume locally uniform grid:  $h_1 = h_2 = h_{\frac{3}{2}}$ :

$$L\varphi_1 = \left( \frac{u}{2} + \frac{3\varepsilon}{h_1} \right) \varphi_1 + \left( \frac{u}{2} - \frac{\varepsilon}{h_1} \right) \varphi_2 = a\varphi_1 - b\varphi_2 \leq 0$$

$$0 < b < a$$

## Estimating the global error (III)

This means:

$$b(\pm e_2 - E_2) \leq b(\pm e_1 - E_1) \quad \text{Using max principle}$$

$$a(\pm e_1 - E_1) \leq b(\pm e_2 - E_2) \quad \text{Using } L_h \leq 0$$

Because  $0 < b < a$  this implies:

$$\pm e_1 - E_1 \leq 0 \Rightarrow |e_1| \leq E_1$$

And using the discrete maximum principle:

$$|e_j| \leq E_j, \quad j = 1, \dots, J.$$

Choosing/Finding a suitable  $E(x)$  is an art.