CFD WI4011

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Estimating the global error (I)

The *global* error is directly related to the *local truncation* error:

$$L_h e = \tau$$

Try to find a *Barrier function E* such that:

$$L_h E \geq |\tau|$$

We will show that this implies $|e| \le E$.

Estimating the global error (II)

$$L_h(\pm e - E) \leq 0$$
.

We want to use the discrete maximum principle: Therefore the scheme should be of positive type or Pe < 2. The maximum occurs at Dirichlet boundary \Rightarrow

$$\pm e_j - E_j \le \pm e_1 - E_1, \quad j = 2, \cdots, J.$$
 (2.58)

Assume locally uniform grid: $h_1 = h_2 = h_{\frac{3}{2}}$:

$$L\varphi_1 = \left(\frac{u}{2} + \frac{3\varepsilon}{h_1}\right)\varphi_1 + \left(\frac{u}{2} - \frac{\varepsilon}{h_1}\right)\varphi_2 = a\varphi_1 - b\varphi_2 \le 0$$

$$0 < b < a$$

Estimating the global error (III)

This means:

$$b(\pm e_2-E_2)\leq b(\pm e_1-E_1)$$
 Using max principle $a(\pm e_1-E_1)\leq b(\pm e_2-E_2)$ Using $L_h\leq 0$

Because 0 < b < a this implies:

$$\pm e_1 - E_1 \leq 0 \Rightarrow |e_1| \leq E_1$$

And using the discrete maximum principle:

$$|e_j| \leq E_j, \quad j=1,\cdots,J.$$

Choosing/Finding a suitable E(x) is an art.



Estimating the global error (IV)

Analysis on a uniform grid:

$$E_j = M\psi(x_j), \quad \psi(x) \equiv 1 + 3x - x^2$$

$$L_h\psi(x_1) = \dots > 2\varepsilon/h$$

$$L_h\psi(x_j) = uh(3-2x_j) + 2\varepsilon h > 2\varepsilon h, \quad j=2,\dots,J-1$$

$$L_h\psi(x_J) = \dots > \varepsilon/h$$

Estimating the global error (V)

We have already shown that

$$au_1 < M_1 h$$
 $au_j < M_2 h^3, \quad j = 2, \dots, J - 1$
 $au_J < M_3 h^2$

Hence, choose:

$$M = \frac{h^2}{\varepsilon} \max\{M_1/2, M_2/2, M_3\}$$

Then
$$L_h E \geq |\tau|$$
 and $|e| < E = \mathcal{O}(h^2)$

The low-order *local truncation error* at the boundaries does not effect the accuracy(=*global error*) of the solution!

Estimating the global error (VI)

Some remarks:

- Discretisation on a nonuniform grid:
 - Similar procedure, more complicated barrier function(s)
 - ► Similar result:

$$e_j = \mathcal{O}(\Delta^2), \quad \Delta = \max(h_j)$$

- Discretisation on a vertex centered grid
 - Similar procedure
 - ► Similar result:

$$e_j = \mathcal{O}(\Delta^2), \quad \Delta = \max(h_j)$$

The fact that the local truncation error of the *vertex-centered* case is smaller than for the *cell-centered* case does not increase the global error!

Summary of the properties of the stationary convection-diffusion equation in 1D(I)

Analytical properties:

- Conservative
- ► Maximum principle
- ▶ Well-posed problem for $Pe \rightarrow \infty$ with the correct boundary conditions.
- Constant coefficient case can be solved exactly.

Summary of the properties of the stationary convection-diffusion equation in 1D(II)

Finite volume discretisation of the stationary convection-diffusion equation in 1D:

- ▶ Central scheme I:accurate solution for $Pe_h < 2$
- Upwind scheme add unwanted artificial diffusion.
- ► Solution: use local grid refinement in the region of strong gradient, of known location and thickness: Discretisation becomes *uniform* in *Pe*
- ➤ The Roughness introduced in the grid does <u>not</u> reduce the accuracy of the solution

The stationary convection-diffusion equation in 2D (I)

$$(u_lpha arphi)_{,lpha}-(arepsilon\phi_{,lpha})_{,lpha}=q({f x}),\quad ({f x})\in\Omega\equiv[0,1]^lpha$$
 with boundary conditions

$$arphi = f(\mathbf{x})$$
 on $\partial \Omega_i$ Dirichlet $arphi = f(\mathbf{x})$ on $\partial \Omega_o$ Dirichlet or $\hat{n}_{,\alpha} arphi_{,\alpha} = g(\mathbf{x})$ on $\partial \Omega_o$ Neumann

The stationary convection-diffusion equation in 2D (II)

Analytical aspects:

- Conservation form
- Maximum principle
- ▶ Well-posed problem for $Pe \rightarrow \infty$ with the correct boundary conditions.

The stationary convection-diffusion equation in 2D (III)

For efficient numerical discretisation we want to solve the same strategy of *local grid refinement*, but where are those regions of strong gradients and how thick are they going to be? These regions are called *boundary layers* and there thickness and location can be derived using *singular perturbation theory*

The stationary <u>convection</u> equation in 2D (IV)

Assume when $\varepsilon \ll 1$ we can discard diffusion effects:

$$u_{\alpha}\varphi_{,\alpha} - = \tilde{q}(x,y), \quad (x,y) \in \Omega \equiv [0,1]^{\alpha}$$

 $\tilde{q} = q(x,y) - \varphi u_{\alpha,\alpha}$

with boundary conditions

$$\varphi = f(x, y)$$
 on $\partial \Omega_i$ Dirichlet

The stationary convection equation in 2D (V)

This equation is *hyperbolic*. Parameterize a curve in Ω

$$\mathbf{r}(s) = \mathbf{\hat{e}}_{\alpha} x_{\alpha}(s)$$

 $(x_{\alpha})_{,s} = u_{\alpha}$

Then

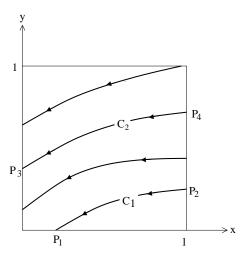
$$\varphi_{,s} = u_{\alpha}\varphi_{,\alpha} = \tilde{q}$$
 using the PDE

 $\varepsilon=0$, no source $(q=0,\ u_{\alpha,\alpha}=0))$ the solution is constant along $\mathbf{r}(s)$: characteristics

 $\varepsilon \ll 1$, no source $(q=0,\ u_{\alpha,\alpha}=0)$): define **subcharacteristics**, similar to the **characteristics** of the case $\varepsilon=0$

The stationary convection-diffusion equation in 2D (VI)

A major problem:

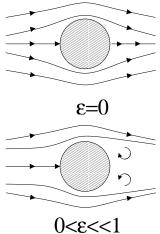


What is the value of $\varphi(P_1)$???

The stationary convection-diffusion equation in 2D (VII)

The convection-diffusion equation is a **singular perturbation problem**

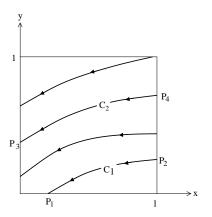
This is related to the **Paradox of d'Alembert** for high Re flow around bluff bodies:



The stationary convection-diffusion equation in 2D (IIX)

A singular perturbation problem

- ▶ The classification of the equation changes when $\varepsilon \neq 0$
- ▶ We can not neglect diffusion



 \Rightarrow the product $\epsilon \varphi_{,\alpha\alpha}$ is not negligible everywhere!! Let's go back to 1-D and find out...

The stationary convection-diffusion equation in 2D (IX)

Back to the one dimensional case:

$$u\varphi_{,x}, -\varepsilon\varphi_{,xx} = 0$$

 $\varphi(0) = a, \quad \varphi(1) = b$

We assume that in a small region near x = 1:

$$\frac{\partial^m \varphi}{\partial x^m} = \mathcal{O}(\delta^{-m}),$$

with the boundary layer thickness $\delta = \mathcal{O}(\epsilon^{\alpha})$

Introduce locally a stretched coordinate

$$\tilde{x} = (1 - x)\varepsilon^{-\alpha}$$
, α to be defined

which leads to the boundary layer equation



The stationary convection-diffusion equation in 2D (X)

$$-\varepsilon^{-\alpha}u\varphi_{,\tilde{\mathbf{x}}}-\varepsilon^{1-2\alpha}\varphi_{,\tilde{\mathbf{x}}\tilde{\mathbf{x}}}=0$$

Requirements on the solution of the boundary layer equation:

► The matching principle

$$\lim_{\tilde{\mathbf{x}} \to \infty} \varphi_{\mathsf{inner}}(\tilde{\mathbf{x}}) = \lim_{\mathbf{x} \uparrow \mathbf{1}} \varphi_{\mathsf{outer}}(\mathbf{x})$$

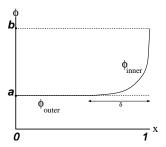
The inner solution is the solution to the boundary layer equation The outer solution is the solution to the inviscid equation ($\varepsilon=0$)

The stationary convection-diffusion equation in 2D (XI)

▶ The boundary condition on $\tilde{x} = 0$

$$\varphi_{\mathsf{inner}}(0) = b$$

Lets find out if there is an α that fulfills these conditions...



The stationary convection-diffusion equation in 2D (XII)

- $ightharpoonup \alpha < 1$: no solution
- ightharpoonup lpha = 1 : solution possible \odot
- $ightharpoonup \alpha > 1$: no solution

We refer to $\alpha = 1$ as the *distinguished* limit.

The stationary convection-diffusion equation in 2D (XIII)

Inner solution in the stretched variable for $\varepsilon \downarrow 0$, $\tilde{x} \in [0, \infty)$:

$$\varphi(\tilde{x}) = a + (b - a)e^{-u\tilde{x}} \tag{1}$$

Inner solution in the unstretched variable for $\varepsilon \downarrow 0$:

$$\varphi(x) = a + (b - a)e^{-u\frac{1-x}{\varepsilon}} \tag{2}$$

Total solution (also)

$$\varphi(x) = a + (b - a)e^{-u\frac{1-x}{\varepsilon}} \tag{3}$$

We assumed that the boundary layer occurred at x=1 why not at x=0? Lets check...