Solutions

$C:\Omega^0_t \stackrel{\partial_t}{\longrightarrow} \dot{C}:\Omega^0_t$

Not limits, but lifts

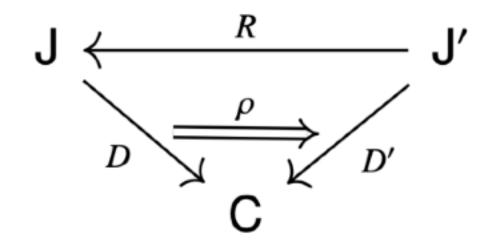
The heat equation:
$$\frac{\partial C}{\partial t} = k \nabla C$$

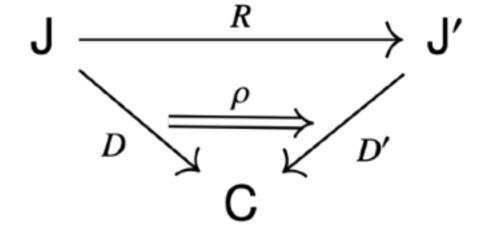
- Instead, solutions correspond to lifts of the diagram D through the codomain functor cod: El_S(C) → C
 - when working with (P)DEs, we usually take $S = \mathbb{R}$ (the constant sheaf of \mathbb{R} on M), since morphisms $\mathbb{R} \to \mathcal{F}$ of sheaves (of real vector spaces on M) are in bijection with global sections of \mathcal{F}
- Note that we can recover limits from lifts, since a lift through cod is exactly a cone over D with apex S

Morphisms of diagrams

Two constructions

• Given a category \mathscr{C} , the **backwards category of diagrams in** \mathscr{C} has objects $(\mathscr{J}, D: \mathscr{J} \to \mathscr{C})$ and morphisms $(R: \mathscr{J}' \to \mathscr{J}, \rho: D \circ R \Rightarrow D')$; the **forwards** category has morphisms $(R: \mathscr{J} \to \mathscr{J}', \rho: D \Rightarrow D' \circ R)$





Forwards

Backwards

- If ho is the identity, then we recover (the opposite of) the slice category Cat/ $\mathscr C$
- Morphisms in the backwards category send solutions to solutions
 - thus these tend to go in the direction of *increasing generality*, e.g. from static Maxwell–Faraday *with* potentials to static Maxwell–Faraday *without* potentials