# Theoretical modeling of the collective tunneling of a Wigner necklace

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#### October 25, 2022

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1 Path Integral Formalism (introduction)

## 2 Wigner Crystals

#### 3 Instanton Formalism

- 3.1 Free Particle
- 3.2 Harmonic Oscillator
- 3.3 One particle tunneling
- 3.3.1 One Particle Hamiltonian
- 3.3.2 Action Integral Calculation
- 3.3.3 Tunneling calculation in quartic potentials

- 4 Experimental Considerations
- 4.1 Modeling the Potential and Fitting
- 4.2 Polarization data
- 4.3 Scaling

## 5 Many Body Tunneling Calculations - Instanton

- 5.1 Quartic Hamiltonian model
- 5.1.1 Rescaling and Dimensionless Hamiltonian
- 5.2 Milnikov method
- 5.3 Classical Equilibrium Positions
- 5.4 Harmonic Oscillator Eigenfrequencies adn Eigenmodes
- 5.5 Trajectory Calculation
- 5.6 Arc Lengt Paramterization
- 5.7 Egzact Diagonalization (ED)

- 6 Polarization
- 6.1 Classical Polarization Calculation
- 6.2 Polarization using ED

#### $\mathbf{A}$ Mathematics

In this appendix I will write down some useful mathematical expressions, proofs or identities that are used in the discussion above

#### A.1Gaussian Integrals

#### Besel Problem

Fourier Way Let's consider first an arbitrary function f(x). The fourier series of this function is:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left\{ a_n \cos\left(\frac{n\pi l}{L}\right) + b_n \sin\left(\frac{n\pi l}{L}\right) \right\}$$
 (A.1)

where  $l \in [l_i, l_f]$  and  $L = l_f - l_i$ let  $f(x) = x^2$  and  $l_i = -l_f = -\pi$  so that  $L = 2\pi$ 

Now the coefficients can be calculated:

$$a_0 = \frac{2}{L} \int_{-\pi}^{\pi} f(x) dx$$
 (A.2)

$$= \frac{2}{\pi} \int_0^{\pi} dx \, x^2 \tag{A.3}$$

$$=\frac{2}{3}\pi^2\tag{A.4}$$

$$a_n = \frac{2}{L} \int_{-\pi}^{\pi} dx \, f(x) \cos\left(\frac{n\pi 2x}{L}\right) \tag{A.5}$$

$$= \frac{2}{\pi} \left( \frac{x^2 \sin(nx)}{n} \right) \Big|_0^{\pi} - \frac{2}{\pi} \int_0^{\pi} dx \, \frac{2x \sin(nx)}{n}$$
 (A.6)

$$=$$
 another partial integration later (A.7)

$$=\frac{4(-1)^n}{n^2} \tag{A.8}$$

 $b_n = 0$  because  $x^2$  is even. So the function now looks like this:

$$x^{2} = \frac{2\pi^{2}}{3 \cdot 2} + \sum_{n=1}^{\infty} \frac{4(-1)^{n}}{n^{2}} \cos\left(\frac{2n\pi x}{L}\right)$$
 (A.9)

$$= \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4(-1)^n}{n^2} (-1)^n \tag{A.10}$$

now choose  $x^2$  tob e equal to 0:

$$\pi^2 = \frac{\pi^2}{3} + \sum_{n=1}^{\infty} \frac{4}{n^2} \tag{A.11}$$

subtract the  $a_0$  term and devide by 4:

$$\frac{\pi^2}{6} = \sum_{n=1}^{\infty} \frac{1}{n^2} \tag{A.12}$$

Complex Integral Way consider the following integral:

$$I = \int_C \frac{1}{z^2} f(z) dz \tag{A.13}$$

let f(z) be the Fermi function:

$$f(z) = \frac{1}{1 + e^{i\pi z}} \tag{A.14}$$

we have several first order poles at  $z = \pm (2n + 1) = p_n$  and a second order pole at z = 0.

The residue from the 1st order poles:

$$Res(f(z), p_n) = \lim_{z \to p_n} (z - p_n) f(z) = \frac{i}{\pi p_n^2}$$
 (A.15)

rom the 2nd order pole:

$$Res(f(z), 0) = \lim_{z \to 0} \frac{d}{dz}(z^2 f(z)) = \frac{-\pi i}{4}$$
 (A.16)

$$\int_{C} g(z)dz = 2\pi i \sum Res(g(z)) \tag{A.17}$$

$$\int_{C} f(z)dz = 2\pi i \left(\frac{-\pi i}{4}\right) + 2\pi i \sum_{n=0}^{\infty} \frac{2i}{\pi (2n+1)^{2}} = 0$$
 (A.18)

$$\frac{\pi i}{4} = \sum_{n=0}^{\infty} \frac{2i}{\pi (2n+1)^2} \tag{A.19}$$

$$\frac{\pi^2}{8} = \sum_{n=0}^{\infty} \frac{1}{(2n+1)^2} \tag{A.20}$$

$$= \frac{3}{4} \sum_{n=1}^{\infty} \frac{1}{n^2} \tag{A.21}$$

$$\frac{\pi^2}{6} = \sum_{1}^{\infty} \frac{1}{n^2} \tag{A.22}$$

- B Algorithms and Methods
- **B.1** Monte Carlo and Simulated Annealing
- B.2 Minimalization by Nelder-Mead algorithm