Unlocking New Dynamics in Paraxial Fluids of Light with an Optical Feedback Loop

Tiago D. Ferreira, Ariel Guerreiro and Nuno A. Silva

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What are paraxial fluids of light?

And why are paraxial fluids of light important in current research?

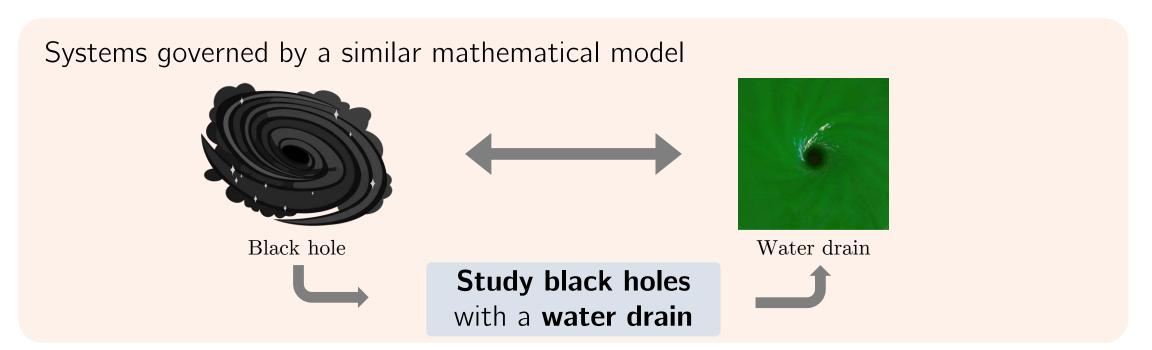
Analogue Simulations

Optical Analogues For Physical Simulations

Physics, nowadays, is interested in systems that hard or impossible to experiment with.

A **black hole** is perhaps the **best example** of this!

We can use **systems** that have a **similar mathematical description** for **emulating them**.



The application of analogue simulations

Analogue Simulations

Physical systems emulating other real phenomena or mathematical models!

Analogue simulations with light

Analogue simulations with **light** are perhaps the **most successful example** of this **branch**.

Light systems offer **great opportunities**:

- ✓ Parallelization
- ✓ High-speed capabilities
- ✓ **Low-power** requirements
- ✓ Precision control of the amplitude and phase profiles.
- ✓ Allow recovery of the experimental amplitude and phase profiles.

Analogue systems created with light for emulating Quantum Fluids.

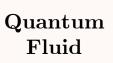
Quantum Fluids of Light

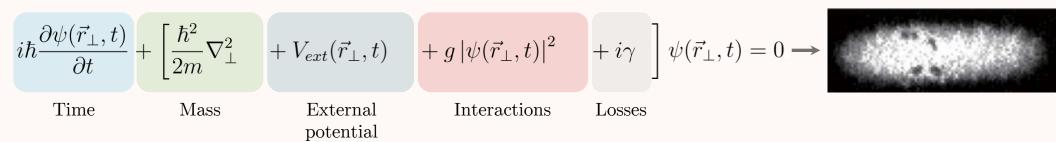
<u>Quantum fluids</u> are <u>hydrodynamical systems</u> that <u>exhibit quantum effects</u>, such as <u>superfluidity</u> or <u>quantum turbulence</u>.

Quantum fluids of light tries to emulate these systems with light.

Quantum Fluids of Light

Systems governed by a similar mathematical model





They support macroscopic quantum effects.

Vortex nucleation due to a defect

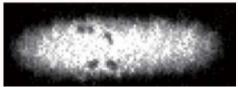
Quantum Fluids of Light

Systems governed by a similar mathematical model

Quantum Fluid

$$i\hbar \frac{\partial \psi(\vec{r}_{\perp}, t)}{\partial t} + \left[\frac{\hbar^2}{2m} \nabla_{\perp}^2 + V_{ext}(\vec{r}_{\perp}, t) + g |\psi(\vec{r}_{\perp}, t)|^2 + i\gamma\right] \psi(\vec{r}_{\perp}, t) = 0 \longrightarrow$$
Time Mass External Interactions Losses

Vortex nucleation due to a defect



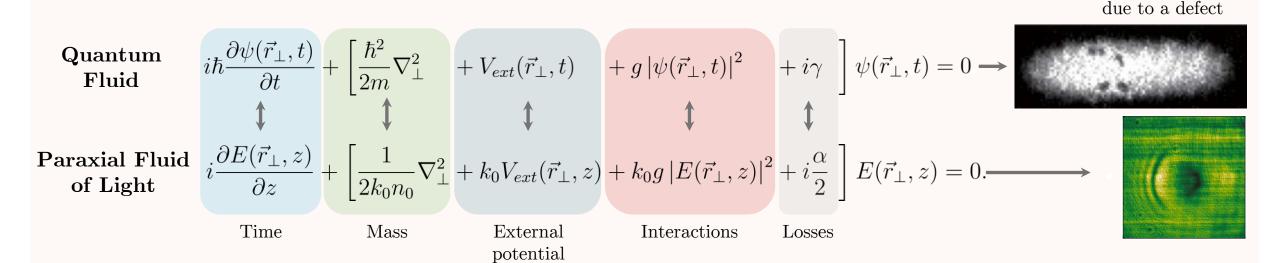
<u>However</u>, they require **low-temperatures** and **expensive experimental setups**.

Accessing all the fluid properties, such as the velocity field, is difficult.

potential

Quantum Fluids of Light

Systems governed by a similar mathematical model

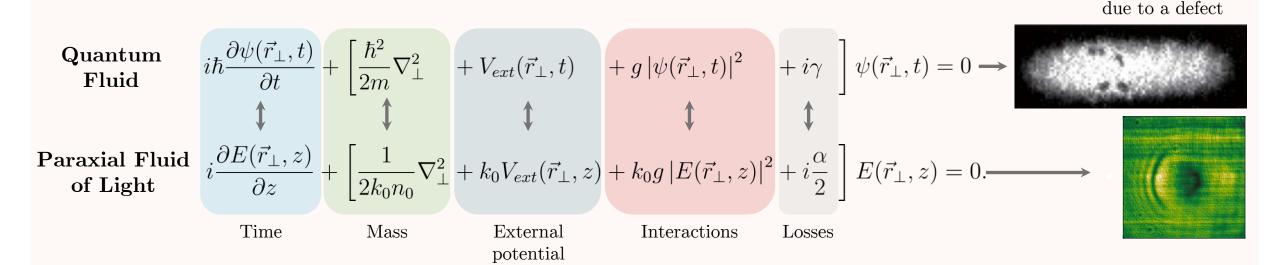


Paraxial fluids of light are suitable physical system capable of emulating quantum fluids.

Vortex nucleation

Quantum Fluids of Light

Systems governed by a similar mathematical model



More generally, Paraxial fluids of light are:

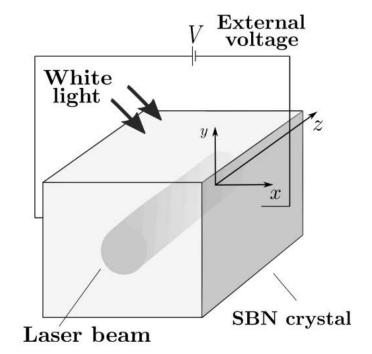
Physical Emulators of the Nonlinear Schrödinger equation.

Vortex nucleation

How can we use a photorefractive crystal to emulate quantum fluids?

A light beam propagating in a photorefractive crystal can be described by the following model:

$$i\frac{\partial E_f}{\partial z} + \frac{1}{2k_f n_e} \nabla_{\perp}^2 E_f + k_f \Delta n_{max} \frac{|E_f|^2}{|E_f|^2 + I_{sat}} E_f + i\frac{\alpha}{2} E_f = 0$$



$$\Delta n_{max} \propto V_{Ext}$$
 $I_{sat} \propto$ White light

Figure 1 — Photorefractive crystal. Image adapted from "Fluids of light in a nonlinear photorefractive medium", Omar Boughdad, 2020.

How can we use a photorefractive crystal to emulate quantum fluids?

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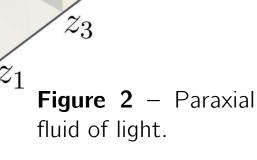
$$i\frac{\partial E_f}{\partial z} + \frac{1}{2k_fn_e}\nabla_\perp^2 E_f + k_f\Delta n_{max}\frac{|E_f|^2}{|E_f|^2 + I_{sat}}E_f + i\frac{\alpha}{2}E_f = 0$$
 Time

Laser beam

The **z** coordinate plays the role of an effective time.

The analogue fluid is two-dimensional.

Each slice along the propagation direction (z-axis) represents the fluid at different moments in time.



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Time Effective mass

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 Time Effective mass Interactions

The **nonlinearity mediates** the **photons interactions**.

These must be repulsive for the system to behave as a fluid.

Defocusing regime of the nonlinear material.

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Time Effective mass

Interactions

Losses

Experimental implementation



Figure 3 – Experimental setup

We start by **considering** a **collimated laser beam**.

Experimental implementation

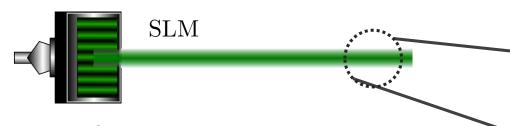
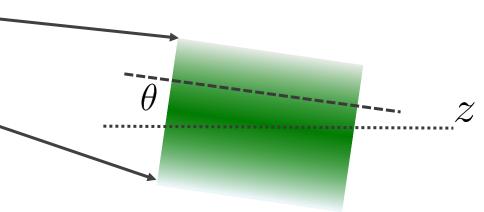


Figure 3 – Experimental setup.

Figure 4 – Angle between the beams at the input of the crystal.



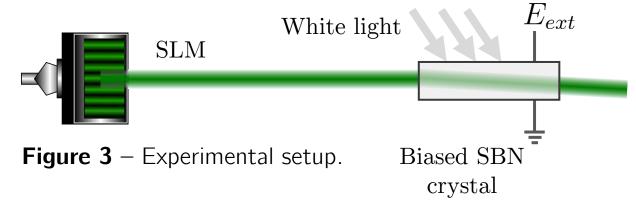
Analogue fluid velocity
$$\ \vec{v} = rac{1}{k_f n_e}
abla_\perp \phi$$

Uniform velocity
$$v pprox rac{ heta}{n_{t}}$$

The analogue **fluid velocity** corresponds to the angle between the beam and the propagation direction.

This **angle** is **controlled** with an **SLM**.

Experimental implementation



The **beam propagates inside** the **crystal**.

Experimental implementation

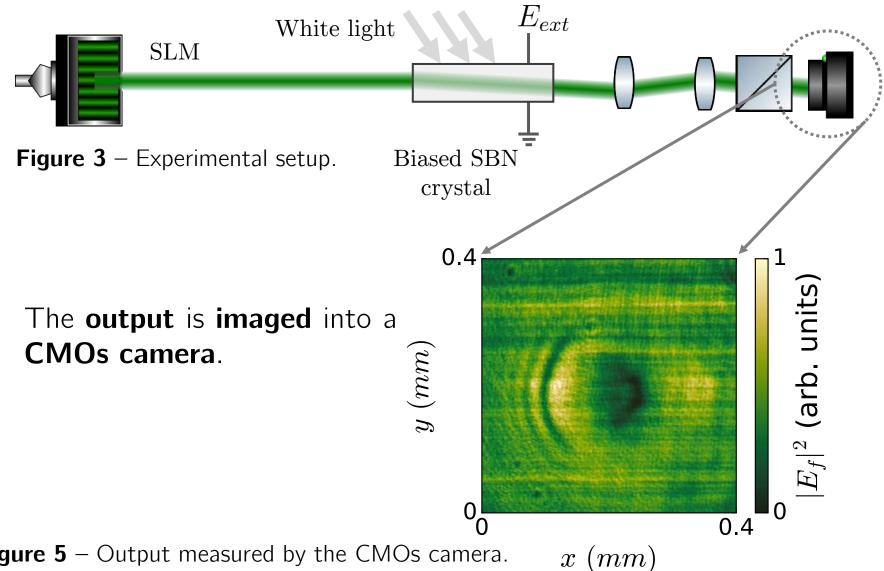


Figure 5 – Output measured by the CMOs camera.

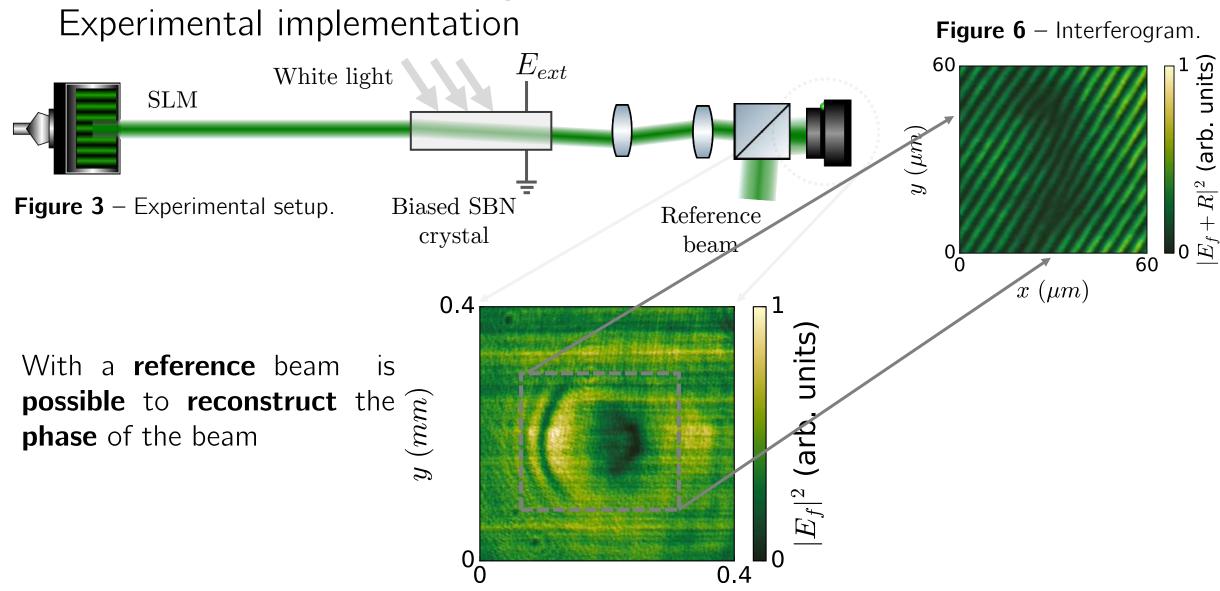


Figure 5 – Output measured by the CMOs camera.

Experimental implementation

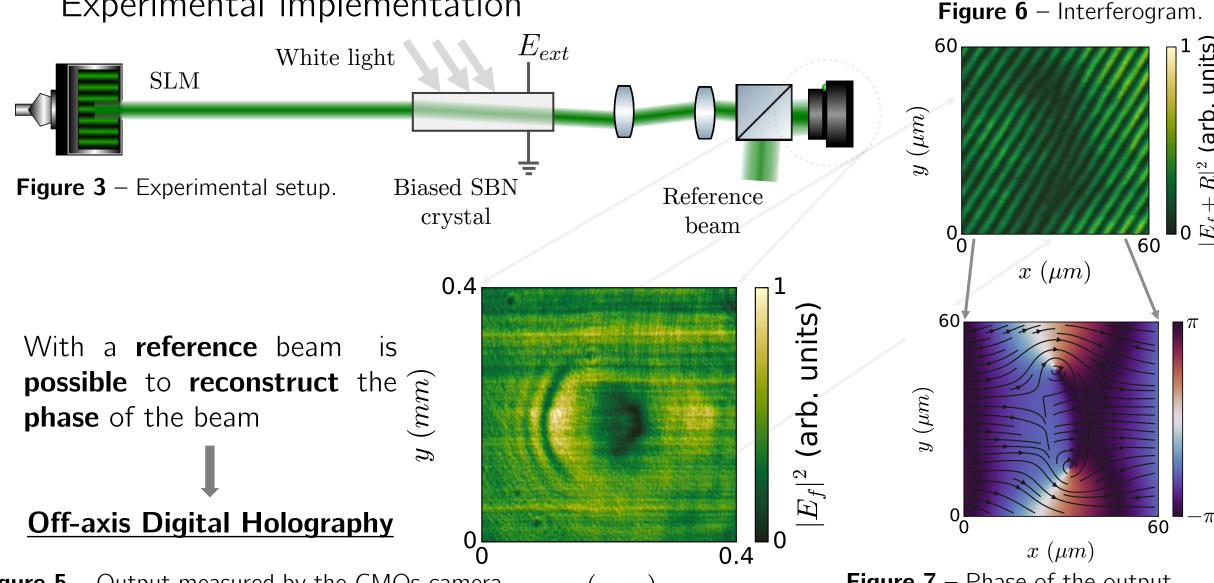


Figure 5 – Output measured by the CMOs camera.

x (mm)

Figure 7 – Phase of the output.

Optical Feedback Loop in Fluids of Light Motivation

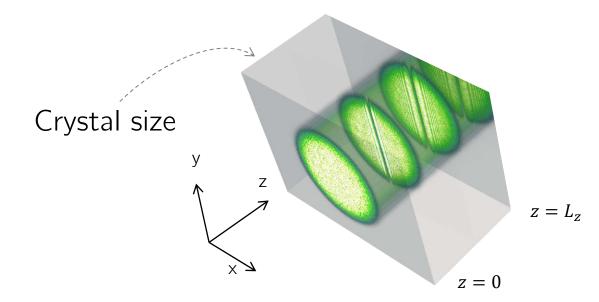
The size of the nonlinear optical medium <u>limits</u> the emulation time!

Remember:



Time is mapped along the z-axis!

Time

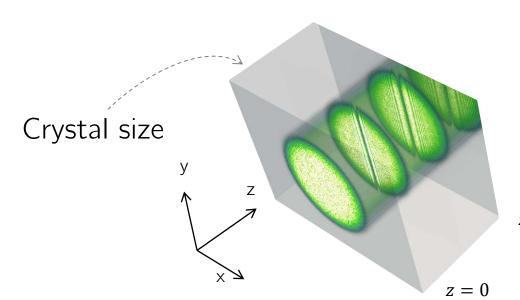


Motivation

The size of the nonlinear optical medium <u>limits</u> the emulation time!

Motivation: Extend the emulation beyond the size of the medium and probe the intermedia states.

Versatile Physical Emulators



 $z = 5L_z$

$$z = 4L_z$$

$$z = 3L_z$$

$$z = 2L_z$$

$$z = L_z$$

Experimental implementation

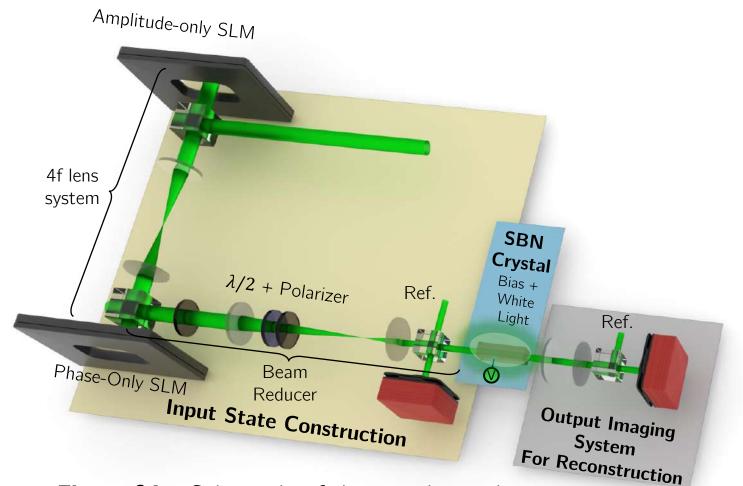


Figure 24 – Schematic of the experimental scheme used to implement the optical feedback loop.

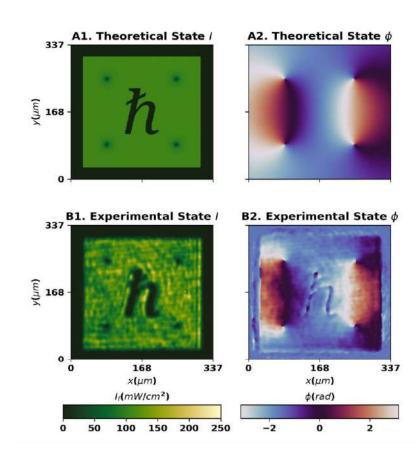
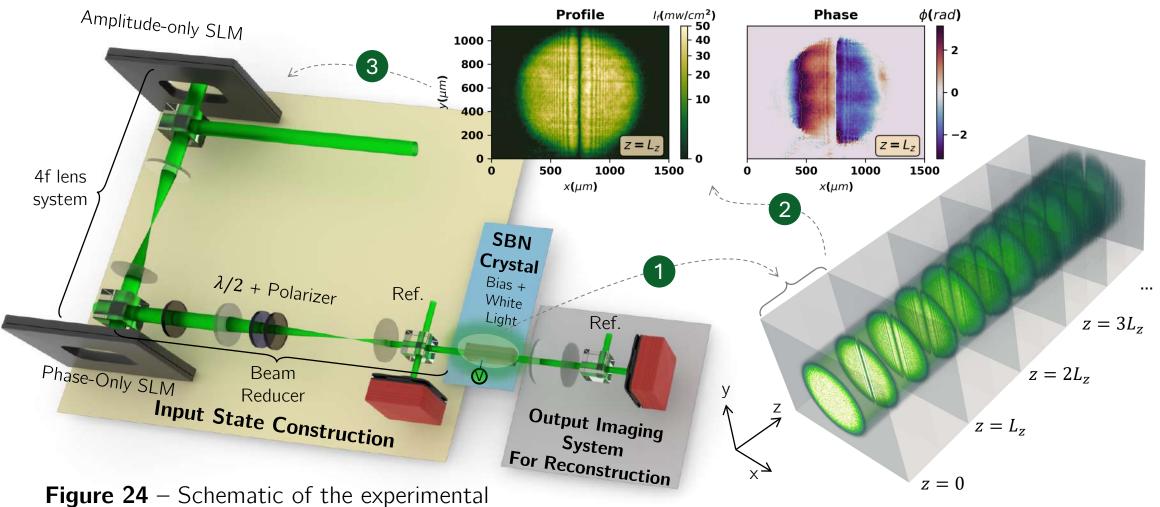


Figure 25 – Arbitrary state generation.

Experimental implementation



scheme used to implement the optical feedback loop.

Experimental results: Dark Soliton decay and shock waves expansion **Old method!**

Numerical / Experimental / Experimental ϕ 1125 A2. A1. y(µm) 562 $z = 1L_z$ $z = 1L_z$ $z = 1L_z$ 562 562 562 1125 0 1125 0 1125 $x(\mu m)$ *x*(μ*m*) $x(\mu m)$ $I_f(mW/cm^2)$ $\phi(rad)$ 25 30 35 40 -2.50.0 10 2.5

Figure 26 – Dark soliton decay, snake instability, and shock waves expansion.

Experimental results: Dark Soliton decay and shock waves expansion

Feedback Loop for 5 passages! Numerical / Experimental / Experimental ϕ 1125 A2. A1. y(µm) **562** $z = 1L_z$ $z = 1L_z$ $z = 1L_z$ 562 1125 0 562 1125 0 562 1125 $x(\mu m)$ $x(\mu m)$ $x(\mu m)$ $I_f(mW/cm^2)$ $\phi(rad)$ 25 30 35 40 -2.50.0 2.5

Figure 27 – Dark soliton decay, snake instability, and shock waves expansion.

Experimental results: Dark Soliton decay and shock waves expansion

The **shock waves expansion also follows** the **numerical simulations**.

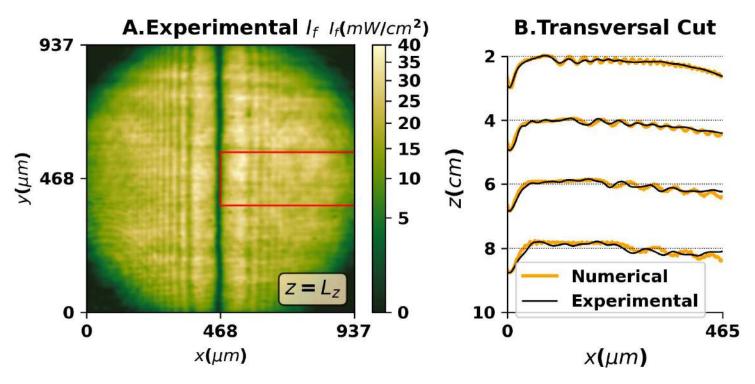


Figure 28 – Comparison between the experimental and numerical expansion of the shock waves.

Experimental results: Collision between three flat-top states

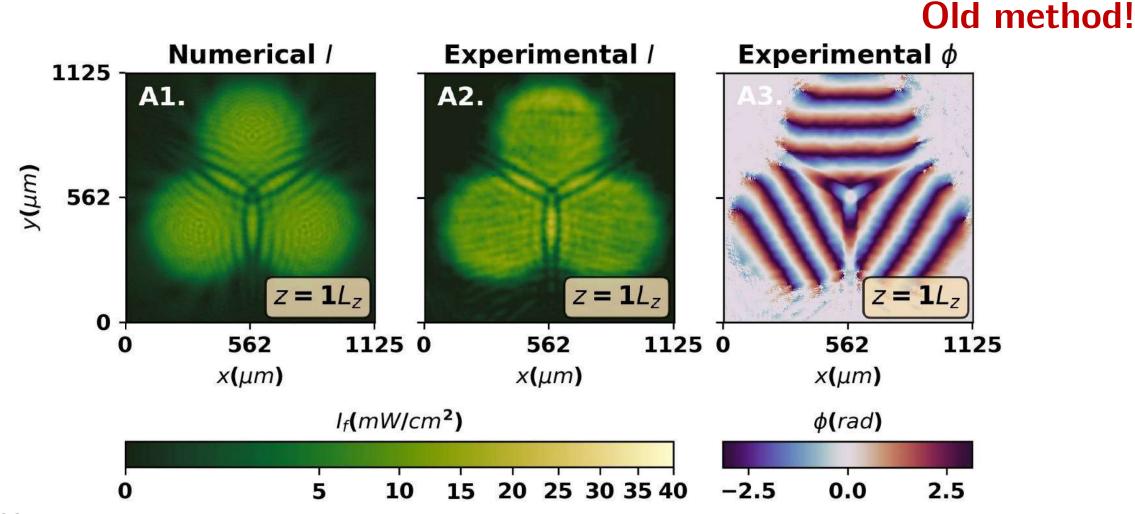


Figure 29 – Collision dynamics between three flat-top states.

Experimental results: Collision between three flat-top states

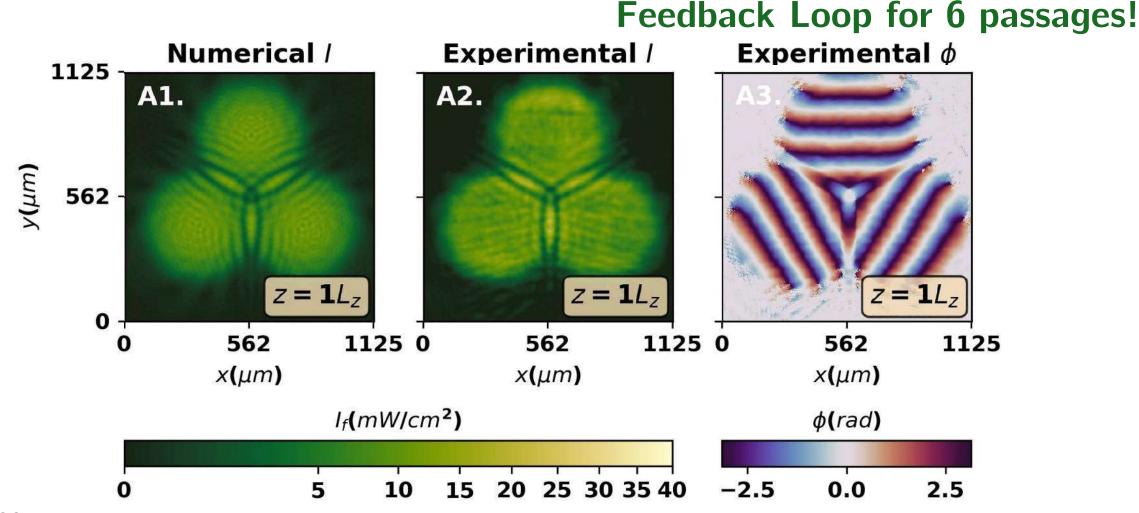


Figure 29 – Collision dynamics between three flat-top states.

Experimental results: Other preliminary results

Exploration of more complex and non-trivial configurations.

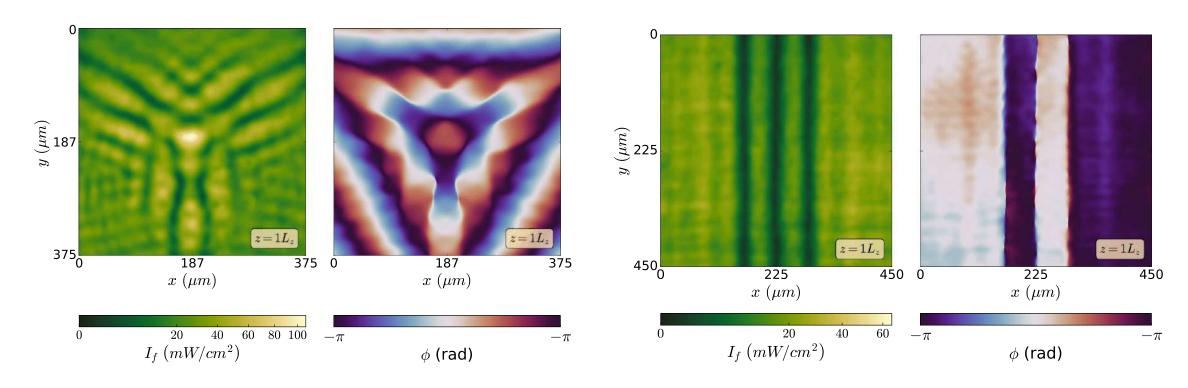


Figure 30 – Dynamics for the vortex formation and vortex collision with sound emission

Figure 32 – Dynamics for the decay of three dark solitons.

Important takeaways

The current version of the **Optical feedback loop** is **capable of**:

Create arbitrary states with arbitrary amplitude and phase profiles.

Perform the **loop** for **up to <u>6</u> passages** with **qualitative results** that **agree** with the **numerical simulations**.

Observe dynamics that were previously impossible with our experimental setups.

We now have a more versatile physical simulator of nonlinear Schrödinger equation.

Acknowledgments







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Thank you for your attention