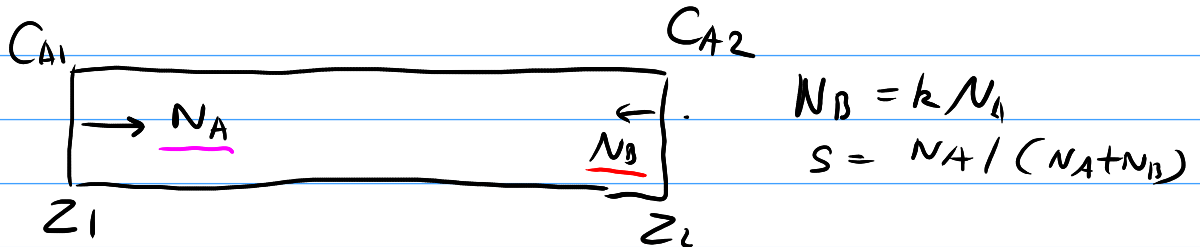


# CHE 318 L04 Theories of $D_{AB}$ (gas)

Jan -12 2026

1 Recap

General solution

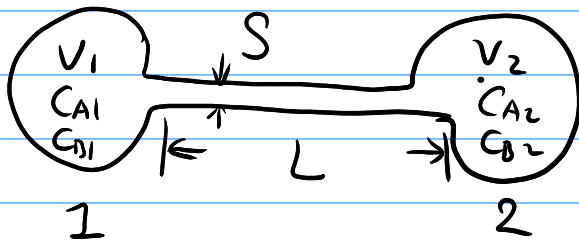


$$N_A = \frac{C_T D_{AB}}{z_2 - z_1} S \ln \left( \frac{S - x_{A2}}{S - x_{A1}} \right)$$

$N_A = \text{const}$

$x_A = \text{function of } z$

Measuring  $D_{AB}$  = two bulb Fitch



$t=0$  pure A in 1 Open  
pure B in 2

$t=t$  mixed Close

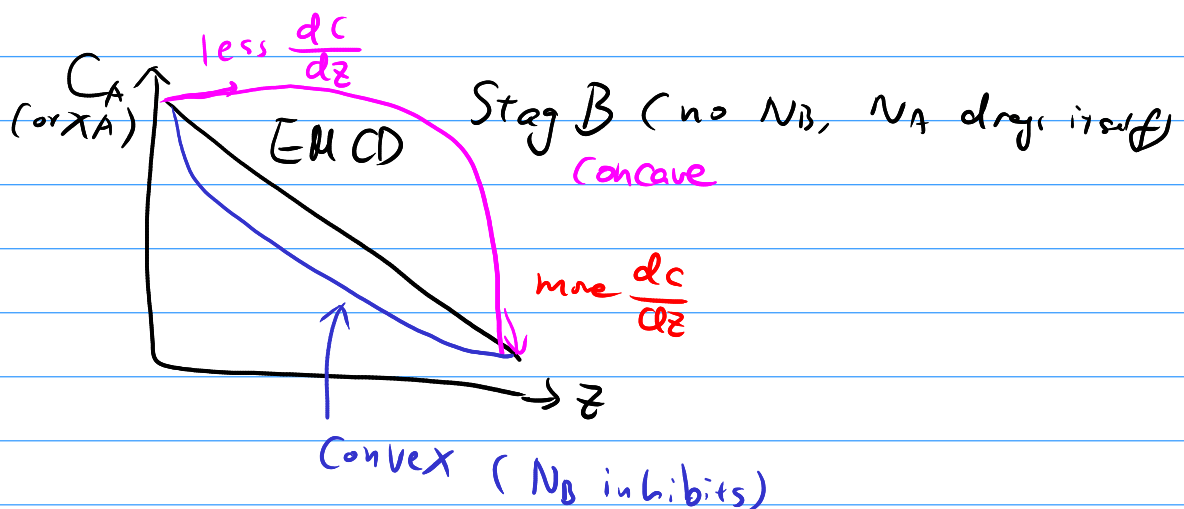
$$\ln \left[ \frac{C_{A,av} - C_{A2}^{t=t_e}}{C_{A,av} - C_{A1}^{t=t_0}} \right] = - D_{AB} \cdot \frac{V_T S}{V_1 V_2 L} \cdot t$$

$\ln()$

$D_{AB}$  from slope

For gas  $D_{AB} \approx 10^{-5} \text{ m}^2/\text{s}$

2. Demo ! See how flux change



All depends on  $S$  value

Stag B  $S=1$   $N_{\text{stag}} = N_{\text{EMCD}} \frac{1}{x_{Bm}} (x_{A1} - x_{A2})$

$0 < \log \text{mean} \uparrow < 1$

3. Today how do we get  $D_{AB}$ ?

without experiment

①  $D_{AB} = f(T, p)$

②  $D_{AB}(z) = \text{const}$

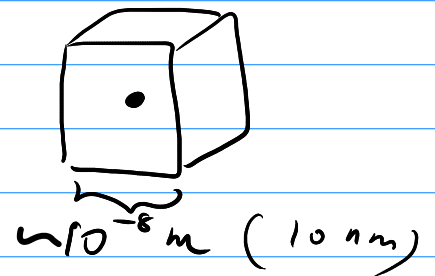
③  $D_{AB} = D_{BA}$  (see assignment 1 Q1)

Kinetic theory of gas

Diluted gas

$$C = \frac{n}{V} = \frac{P}{RT}$$

at 1 atm  $C = 10^{25} \text{ m}^{-3}$



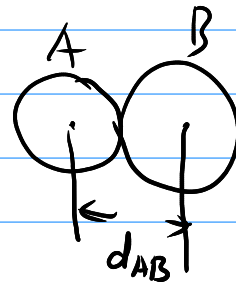
Assumptions

- ① A, B rigid spheres
- ② Far away only 2 can collide
- ③ Momentum conservation
- ④ Molecules travel very far ( $\lambda_{AB}$ ) before colliding

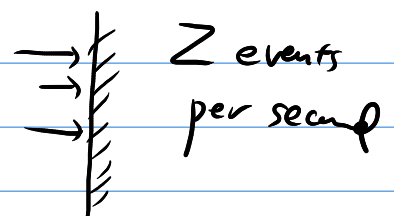
$$u = \sqrt{\frac{8k_B T}{\pi \bar{m}_{AB}}}$$

$$\bar{m}_{AB} = \left( \frac{1}{m_A} + \frac{1}{m_B} \right)^{-1}$$

$$\lambda_{AB} = \frac{1}{\sqrt{2} \pi d_{AB}^2 C_T}$$



$$Z = \frac{1}{4} C_A \bar{u}$$



$$a = \frac{2}{3} \lambda_{AB} \quad \begin{array}{l} \text{avg. path} \\ \text{collision} \longleftrightarrow \text{plane} \end{array}$$

$$\begin{array}{ccc} Z_1 & & Z_2 \\ | & & | \\ z = -a & z = 0 & z = a \\ \frac{1}{4} C_T x_{A1} \bar{u} & & \frac{1}{4} C_T x_{A2} \bar{u} \end{array} \quad \begin{array}{l} \text{In - Out} \\ x_{A2} - x_{A1} \\ = -2a \frac{dx_A}{dz} \end{array}$$

Molecular:  $a = \frac{2}{3} \lambda_{AB}$

$$J = - \frac{1}{4} C_T \cdot 2a \cdot \bar{u} \frac{dx_A}{dz}$$

Macro

$$J = - D_{AB} C_T \frac{dx_A}{dz}$$

$$\Rightarrow D_{AB} = \frac{1}{3} \bar{u} \lambda_{AB}$$

$$\propto \sqrt{\frac{T}{m_{AB}}} \frac{1}{d_{AB}^2 C_T} \quad C_T \propto \frac{P_T}{T}$$

$$\propto \sqrt{T^3} \left( \frac{1}{m_{AB}} \right)^{\frac{1}{2}} \cdot \frac{1}{d_{AB}^2} \cdot \frac{1}{P_T}$$

$$\propto T^{1.5} \left( \frac{1}{m_{AB}} \right)^{\frac{1}{2}} \frac{1}{d_{AB}^2} \frac{1}{P_T}$$

Note the powers!

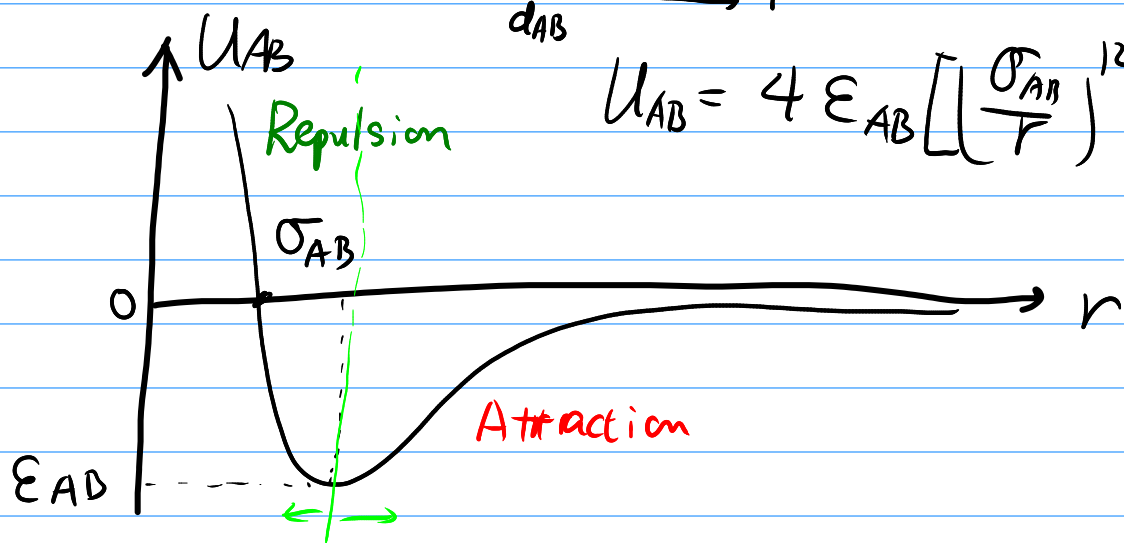
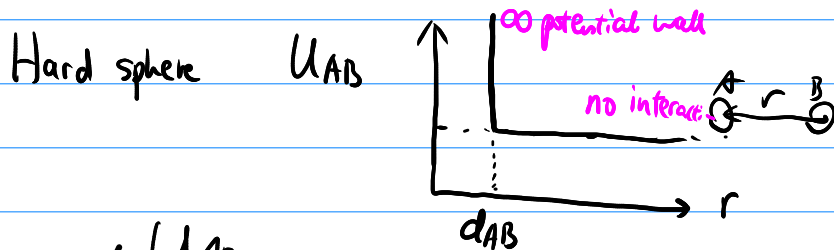
# Chapman-Enskog Theory

$$D_{AB} = \frac{1.8583 \times 10^{-7} T^{1.5}}{P_T \sigma_{AB}^2 \Omega_{D,AB}} \left( \frac{1}{m_A} + \frac{1}{m_B} \right)^{\frac{1}{2}}$$

↓
↓
↓
↓

In  $m^2/s$ 
In atm
In Å
kg/(kg mol)

(Chapman-Enskog theory is just hard-sphere kinetic theory on steroid!)

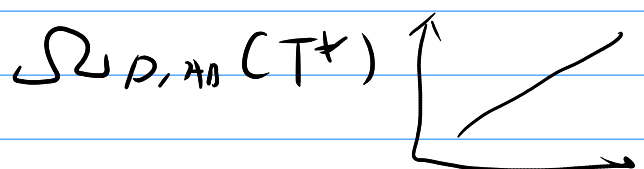


$$U_{AB} = 4 \epsilon_{AB} \left[ \left( \frac{\sigma_{AB}}{r} \right)^{12} - \left( \frac{\sigma_{AB}}{r} \right)^6 \right]$$

$$\sigma_{AB} = \frac{\sigma_A + \sigma_B}{2}$$

$$\epsilon_{AB} = \sqrt{\epsilon_A \epsilon_B}$$

$$T^* = \frac{k_B T}{\epsilon_{AB}}$$

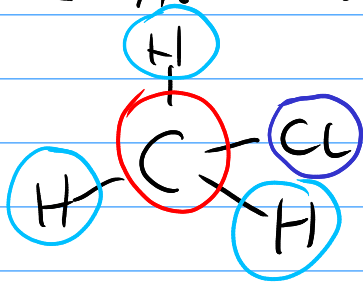


# Fuller Method

$$D_{AB} = \frac{1.0 \times 10^{-7} T^{1.75}}{P_T \left( \left( \sum V_{Ai} \right)^{\frac{1}{3}} + \left( \sum V_{Bi} \right)^{\frac{1}{3}} \right)^2}$$

$\sum V_{Ai}$  has meaning of volume  $\Rightarrow [\text{Length}]^3$   $\left( \frac{1}{m_A} + \frac{1}{m_B} \right)^{\frac{1}{2}}$   
 $\left[ \sum V_{Ai}^{\frac{1}{3}} + \sum V_{Bi}^{\frac{1}{3}} \right]^2 \Rightarrow$  has same meaning / unit  
as  $d_{AB}$

$\sum V_{Ai} =$  sum of volum contrib



$$\sum V_{CH_3Cl} = V_C + 3 \cdot V_H + 1 \cdot V_{Cl}$$