

# Magnetohydrodynamics

*origin, principles and applications in space physics*

# Topics/Schedule

Date	Lecture	Topic	Applications
	1	What is MHD?	
	2	Electrodynamics for MHD	ES/Inductive M-I Coupling
	3	Fluid Dynamics for MHD	LLBL
	4	MHD equations 1	Field-aligned currents
	5	MHD equations 2	Multi-fluid Gamera
	6	MHD waves 1	BBF
	7	MHD waves 2	Auroral physics
	8	MHD instability 1	Planetary KHI
	9	MHD instability 2	Interchange/ballooning
	10	MHD shocks	BS and MP
	11	MHD reconnection	GEM reconnection
	12	TBD	

# What is MHD?

Magnetohydrodynamics (MHD) is the study of the interaction between **magnetic fields** and moving, conducting **fluids**.

**MHD = Magnetic field + Hydrodynamics**

Maxwell's  
equations

Fluid  
Dynamics

## Who cares about MHD?

**MHD**

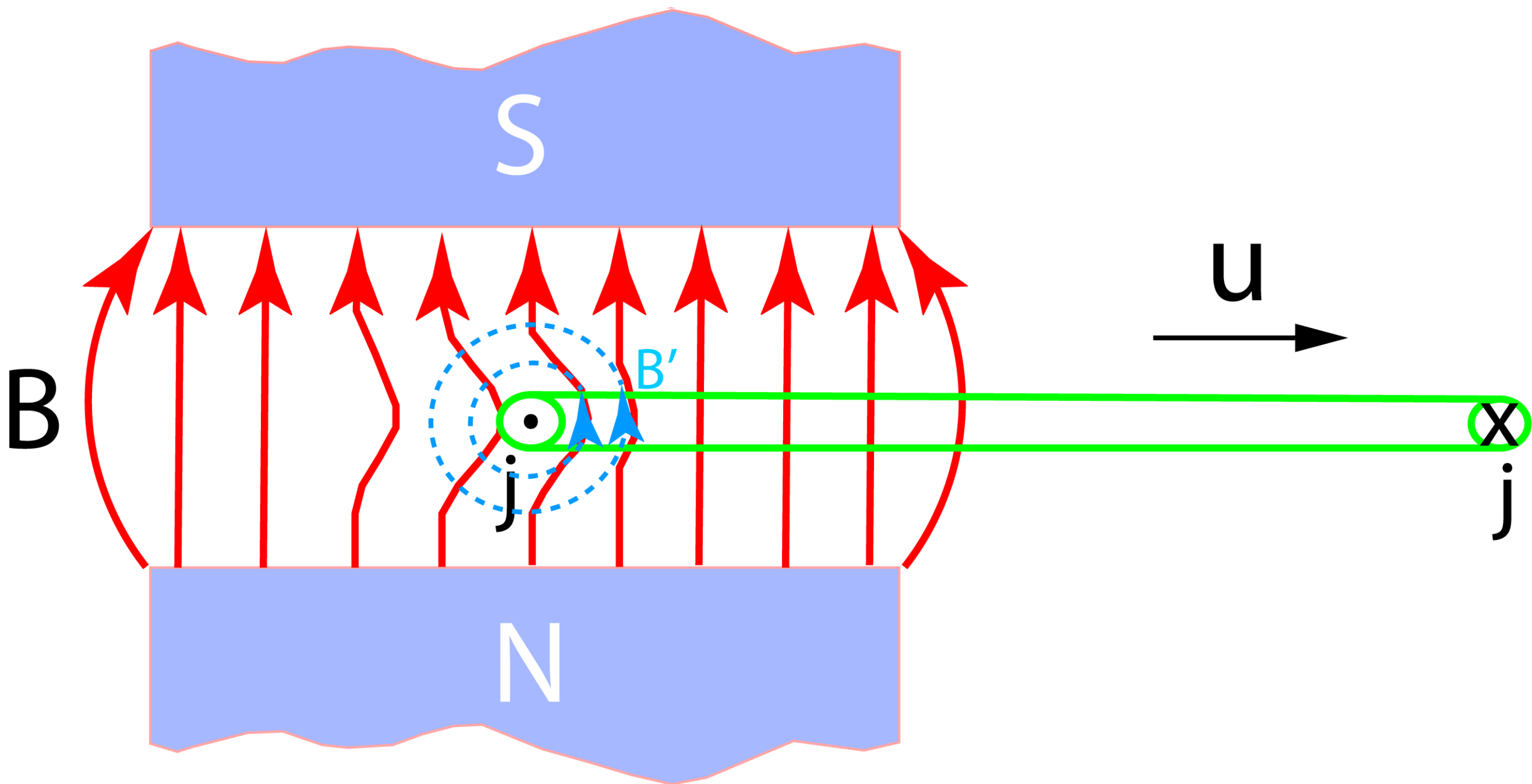
Engineering

Liquid Metal  
engineering pump

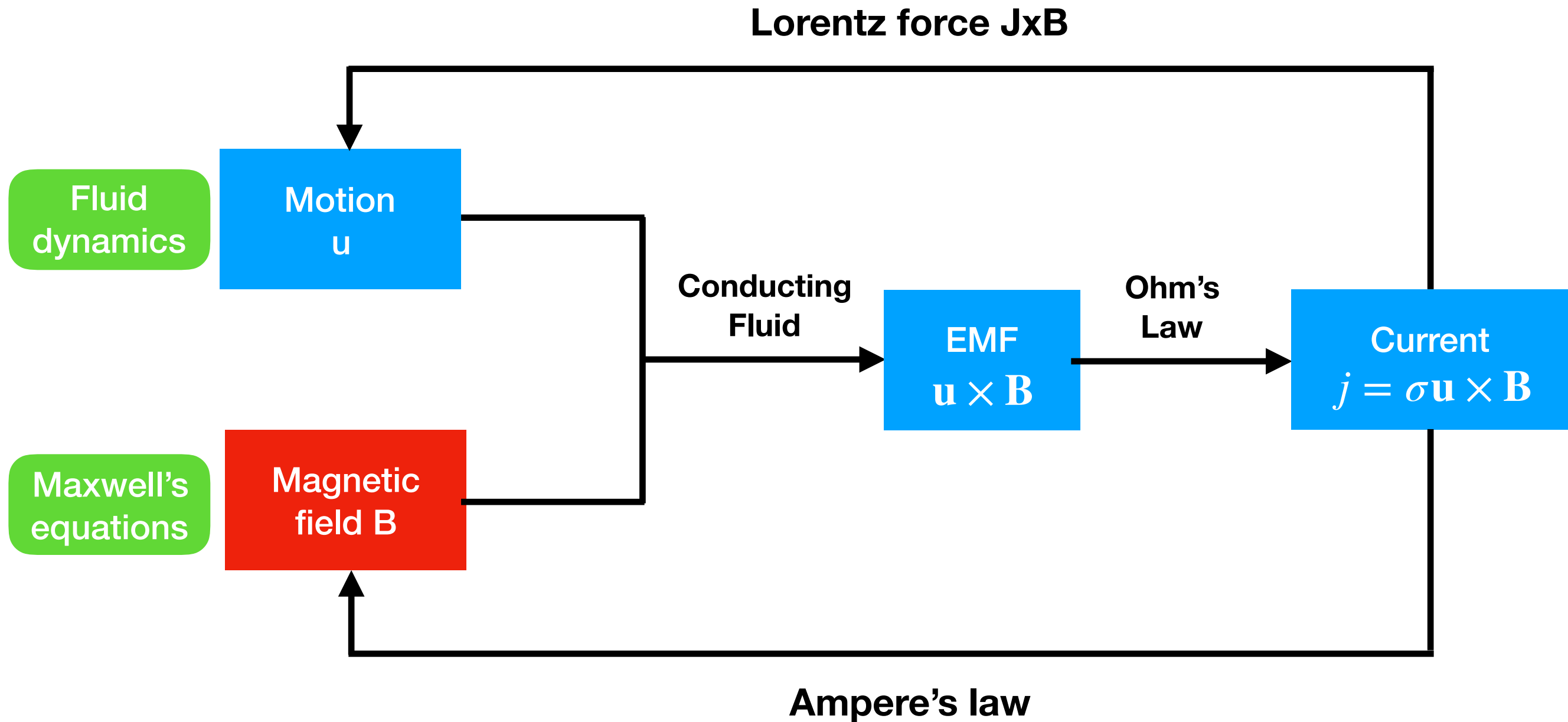
Science

Fusion plasma  
Space plasma  
Astrophysics

# What is MHD?



# What is MHD?



# Several Key Parameters to note

1.  $\mathbf{u}$ ,  $\mathbf{B}$ : fluid velocity and magnetic field
2.  $\sigma$ : conductivity determines the property of the **field-fluid** interaction
3.  $\frac{\Delta B}{B}$  ratio

Since  $\Delta \mathbf{B} \sim \oint \mathbf{E} \cdot d\mathbf{l} \sim \sigma |\mathbf{u} \times \mathbf{B}| l = \sigma u B l$

So  $\frac{\Delta B}{B} \sim \sigma u l$  (This is called the Magnetic Reynolds number  $R_m$ )

Two limits:  $R_m \gg 1$     Magnetospheric plasma

$R_m \ll 1$     Ionospheric plasma?

# From Electromagnetism to MHD

A few more important parameters

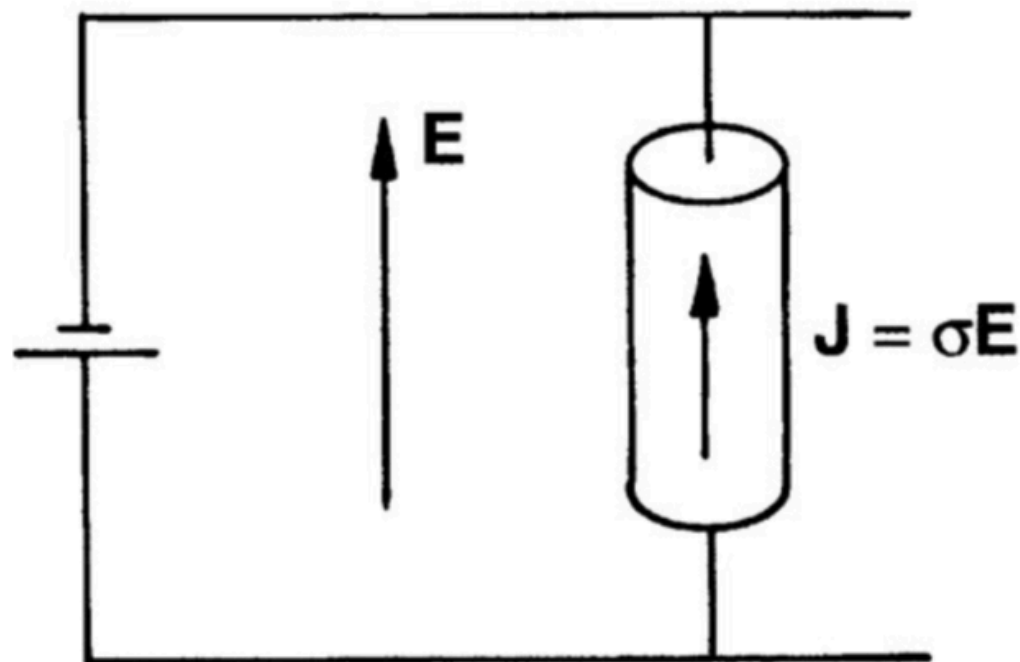
- Magnetic Reynolds number  $R_m = \mu \sigma u l$
- Magnetic damping time  $\tau = \left( \frac{\sigma B^2}{\rho} \right)^{-1}$
- Alfvén speed  $v_A = \frac{B}{\sqrt{\mu \rho}}$

	$R_m$	$v_A$	$\tau$
Solar Corona			
Solar Wind			
Magnetosphere			
Ionosphere			

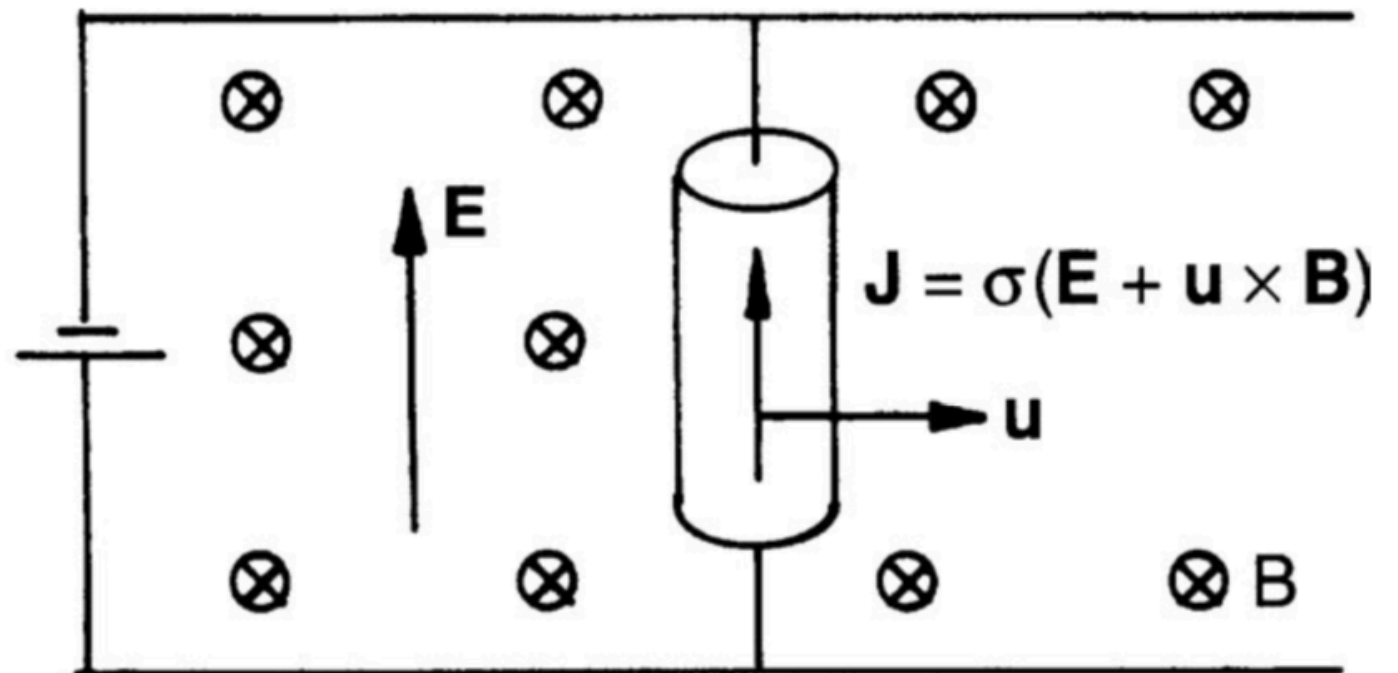
# From Electromagnetism to MHD

A few more important laws

## 1). Ohm's Law



Stationary conductor



Moving conductor

$$j = \sigma E$$

External electric field

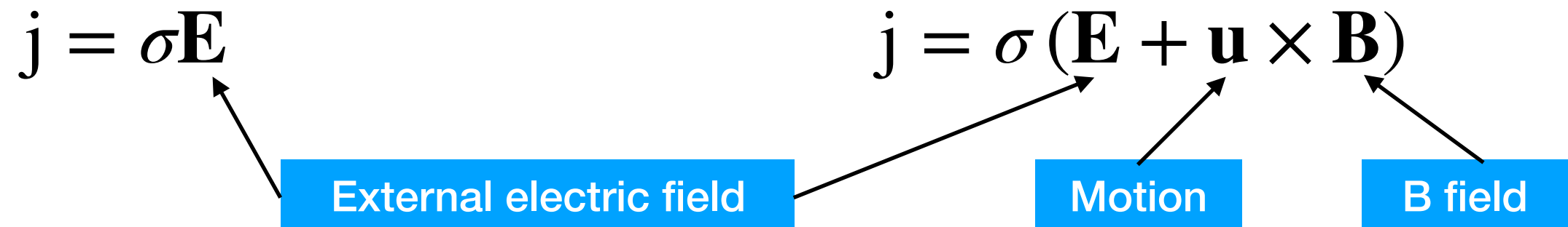
$$j = \sigma (E + u \times B)$$

Motion

B field



## Effective Electric field



The quantity  $\mathbf{E} + \mathbf{u} \times \mathbf{B}$ , which is the total electromagnetic force per unit charge, arises frequently in electrodynamics and it is convenient to give it a label  $\mathbf{E}_r$ :

$$\mathbf{j} = \sigma (\mathbf{E} + \mathbf{u} \times \mathbf{B}) = \sigma \mathbf{E}_r$$

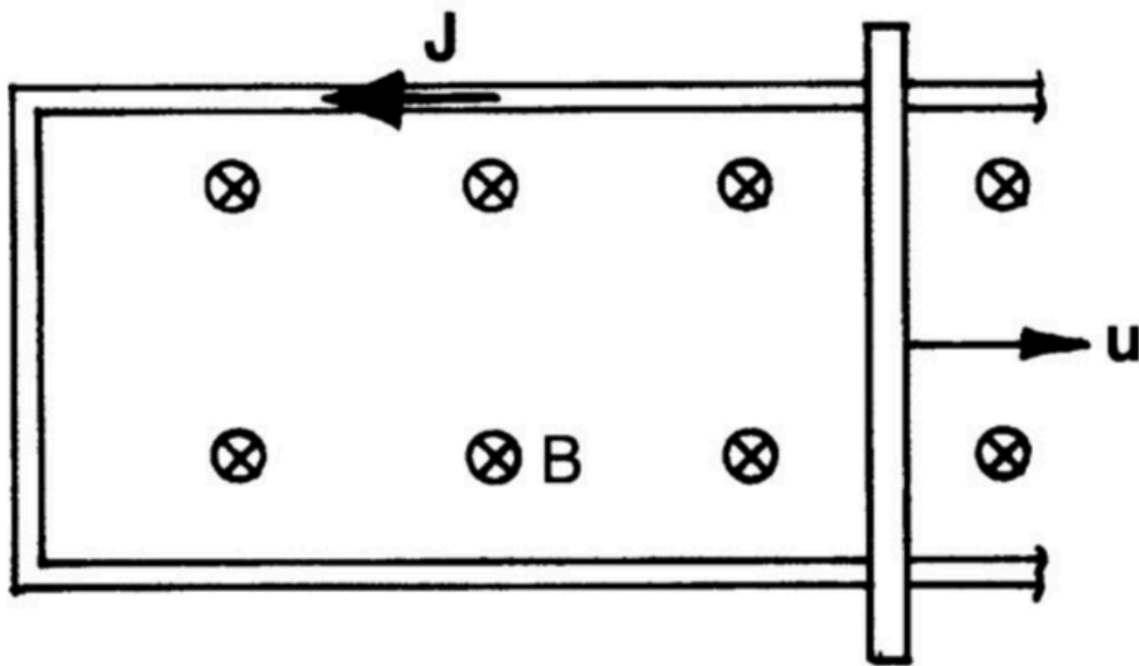
$\mathbf{E}_r$  is the electric field measured in a frame of reference moving with velocity  $\mathbf{u}$  relative to the laboratory frame

$$\mathbf{E}_r = \mathbf{E} + \mathbf{u} \times \mathbf{B}$$

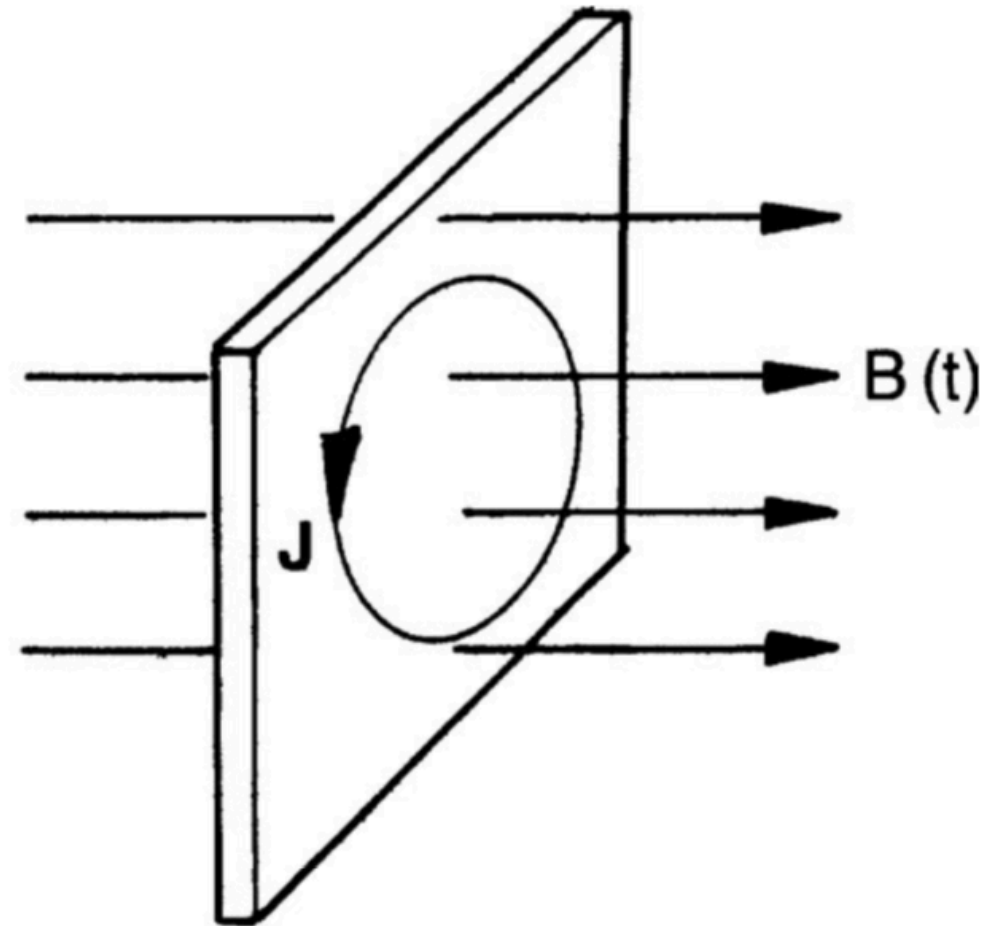
# From Electromagnetism to MHD

A few more important laws

## 2). Faraday's Law



EMF (electromotive force)

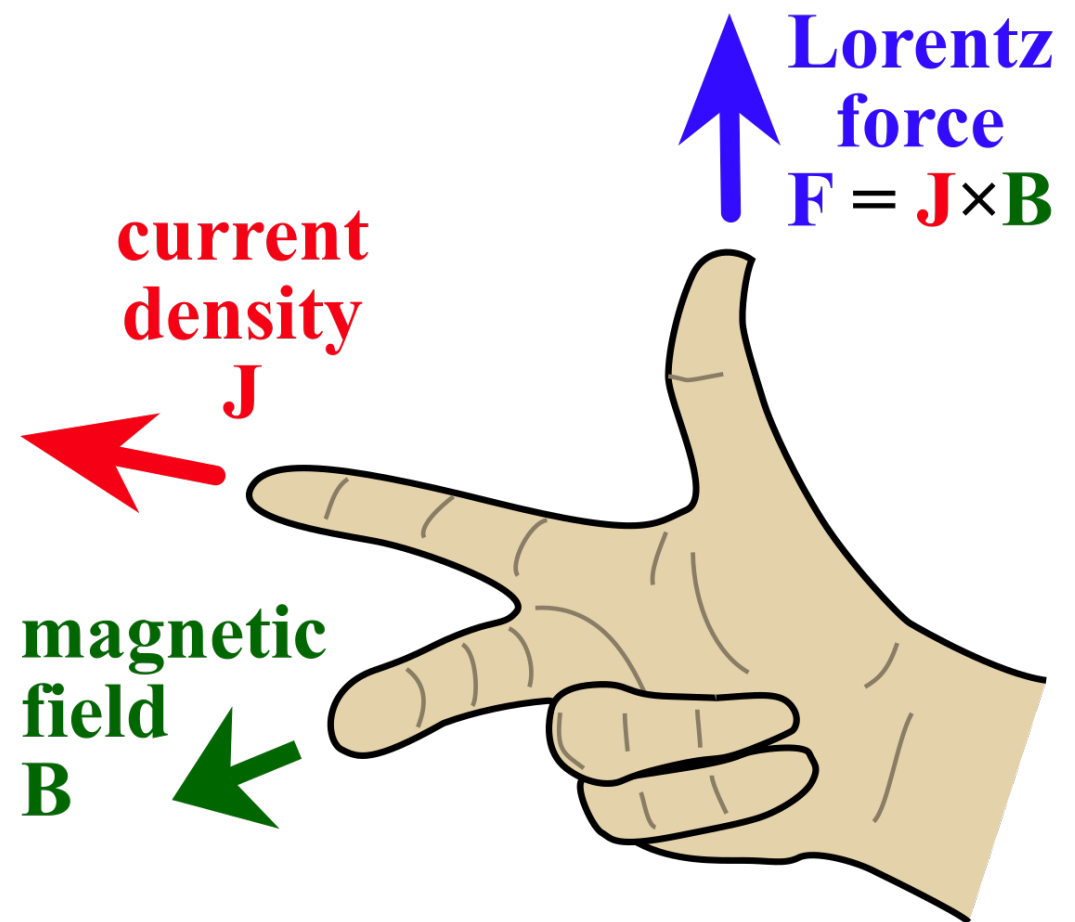
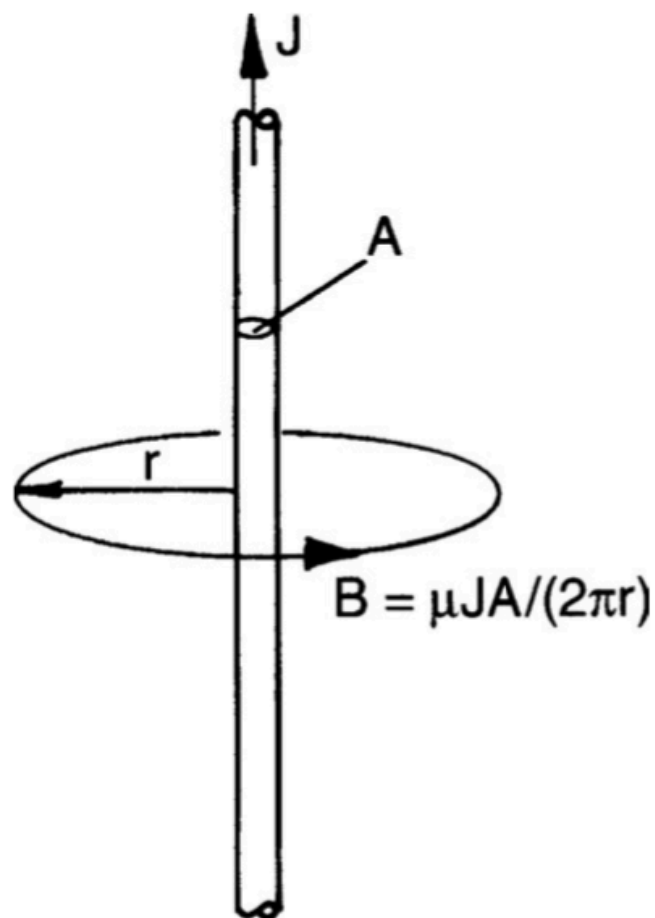


$$EMF = \oint_C \mathbf{E}_r \cdot d\mathbf{l} = - \frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S} = - \frac{d\Phi}{dt}$$

# From Electromagnetism to MHD

A few more important laws

## 2). Ampere's Law and Lorentz force

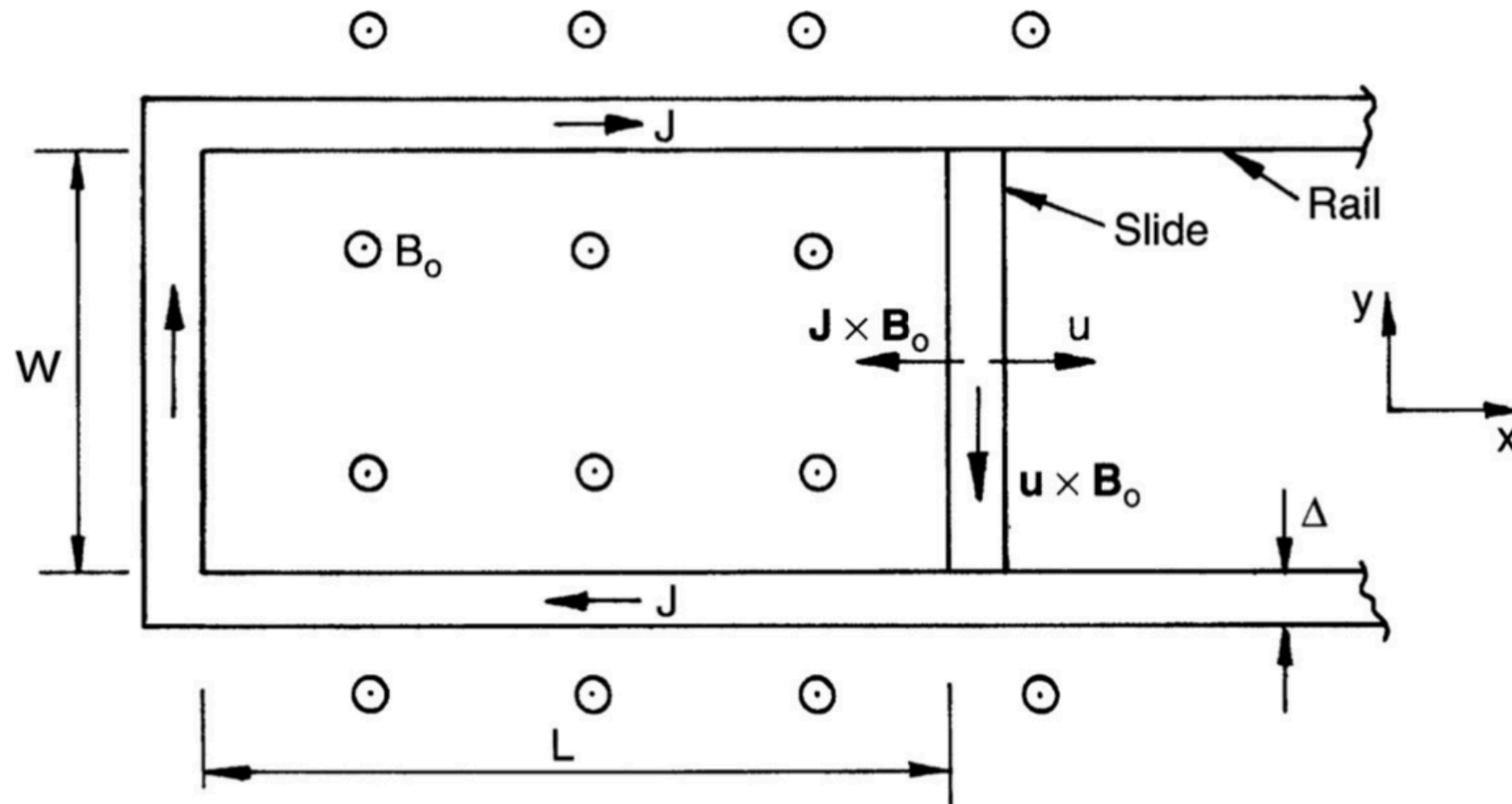


$$\oint_C \mathbf{B}_r \cdot d\mathbf{l} = \mu \int_S \mathbf{j} \cdot d\mathbf{S}$$

$$\mathbf{F} = \mathbf{j} \times \mathbf{B}$$

# From Electromagnetism to MHD

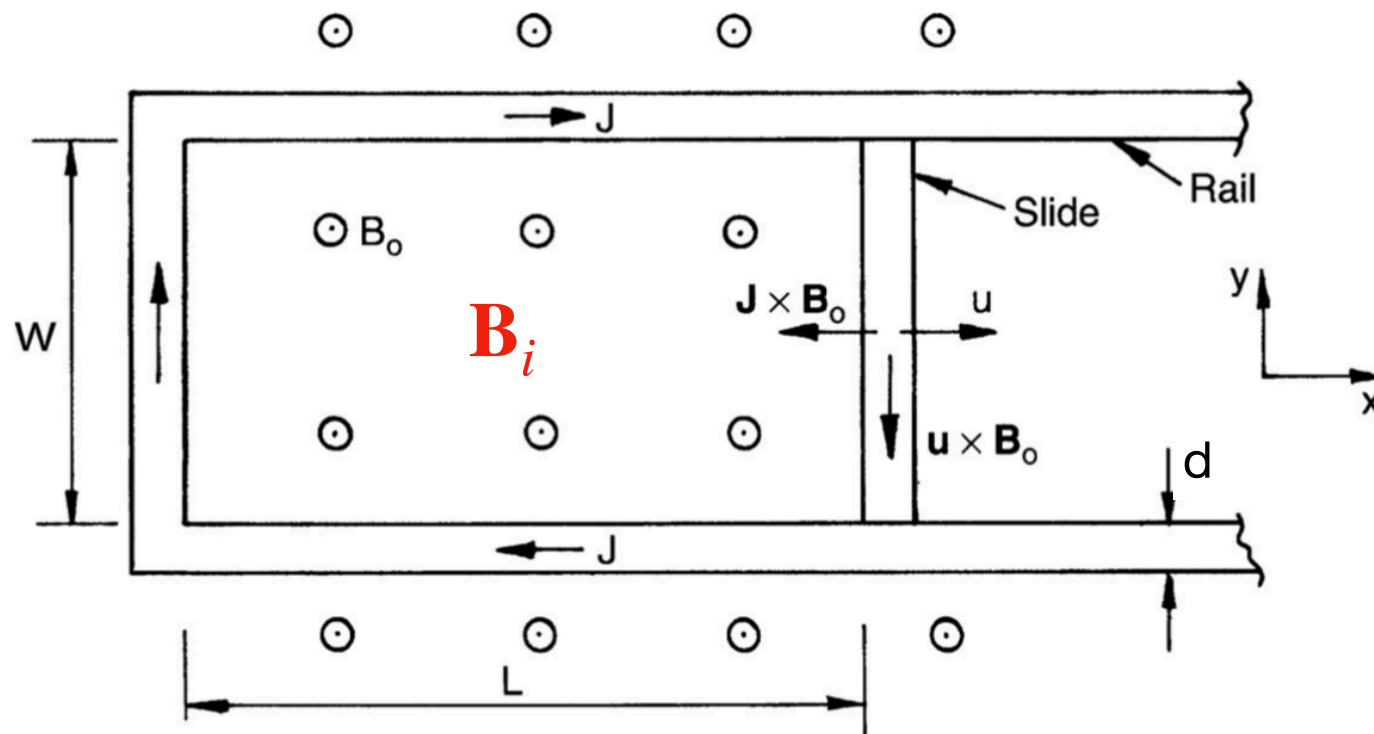
A familiar high-school experiment (高考题)



Width:  $W$   
Distance:  $L$   
Slide width:  $\Delta$   
Magnetic field:  $B_0$   
Conductivity:  $\sigma$

Question: when the motion of the slide is given a tap, what's the response?

# From Electromagnetism to MHD



Width:  $W$   
 Distance:  $L$   
 Slide width:  $d$   
 Magnetic field:  $B_0$   
 Conductivity:  $\sigma$

At  $t = 0$ , the slide is given a speed  $u$

**Ampere's law:**  $\mathbf{B}_i = -(\mu d J) \hat{\mathbf{z}}$  Note that the direction of  $B_i$  is such as to try to maintain a constant flux in the current loop

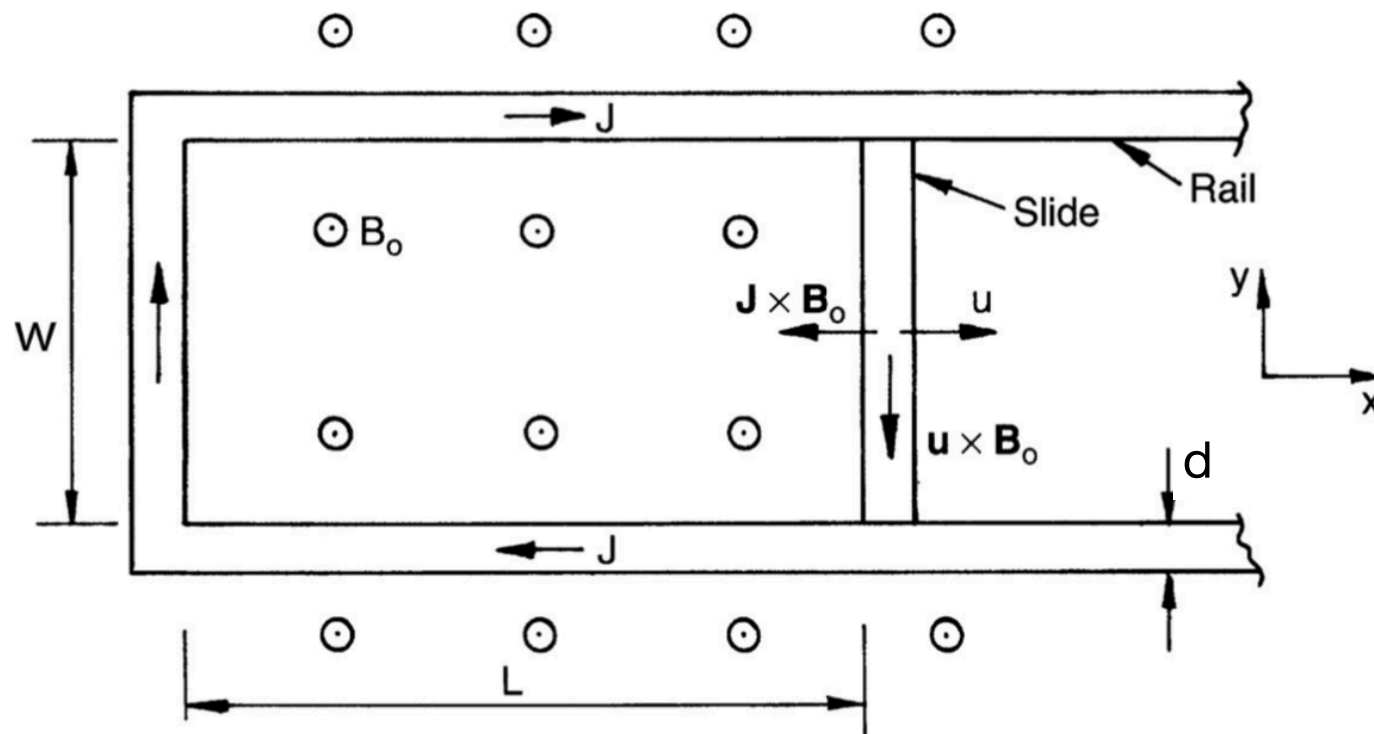
**Faraday's law+Ohm's law:**  $\frac{1}{\sigma} \oint_C \mathbf{j} \cdot d\mathbf{l} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S}$

**Use Ampere's law:**  $\longrightarrow \frac{d\Phi}{dt} = \frac{d}{dt} \left[ LW (B_0 - \mu d J) \right] = 2(L + W) \frac{J}{\sigma}$

**Lorentz force**  $\mathbf{F} = -J \left( B_0 - \frac{1}{2} \mu d J \right) dW \hat{\mathbf{x}}$

**Newton's law:**  $\rho \frac{d^2 L}{dt^2} = \rho \frac{du}{dt} = -J (B_0 - \mu d J / 2)$

# From Electromagnetism to MHD



**A couple of parameters:**

$$B_i = \mu d J \quad T = \mu \sigma d W$$

$$l = dW/L \quad R_m = \mu \sigma u l = uT/L$$

$$\frac{d\Phi}{dt} = \frac{d}{dt} \left[ LW (B_0 - \mu d J) \right] = 2(L + W) \frac{J}{\sigma} \longrightarrow \frac{d}{dt} \left[ L (B_0 - B_i) \right] = \frac{2(L + W) B_i}{T}$$

$$\rho \frac{d^2 L}{dt^2} = \rho \frac{du}{dt} = -J (B_0 - \mu d J / 2) \longrightarrow \rho d \frac{d^2 L}{dt^2} = \rho d \frac{du}{dt} = \frac{(B_0 - B_i)^2}{2\mu} - \frac{B_i^2}{2\mu}$$

**Now we got two non-linear PDEs, it's time to make assumptions based on  $R_m$**

$$R_m \gg 1$$

$$R_m \ll 1$$

# The perfect conducting case $R_m \gg 1$

In this case,  $R_m = \frac{u}{L} \gg T$ , or  $R_m = \mu \sigma u l \gg 1$ ,

$$\frac{d}{dt} \left[ L (B_0 - B_i) \right] = \frac{2(L+W)B_i}{T} \approx 0$$

the flux  $\Phi$  linking the current path is conserved during the motion:  $\Phi = LWB_0$

Let's expand the distance parameter:  $L = L_0 + \eta = \frac{\Phi}{B_0 W} + \eta$

substitute

$$\rho d \frac{d^2 L}{dt^2} = \frac{(B_0 - B_i)^2}{2\mu} - \frac{B_i^2}{2\mu} \xrightarrow{\text{Keep the leading order}} \boxed{\frac{d^2 \eta}{dt^2} + \frac{B_0^2}{\rho \mu d L_0} \eta = 0}$$

Wave equation

**Physical meaning:**

When the magnetic Reynolds number is high, the slide **oscillates** in an elastic manner  
With an angular frequency

$$\omega = \frac{v_A}{\sqrt{dL_0}}$$

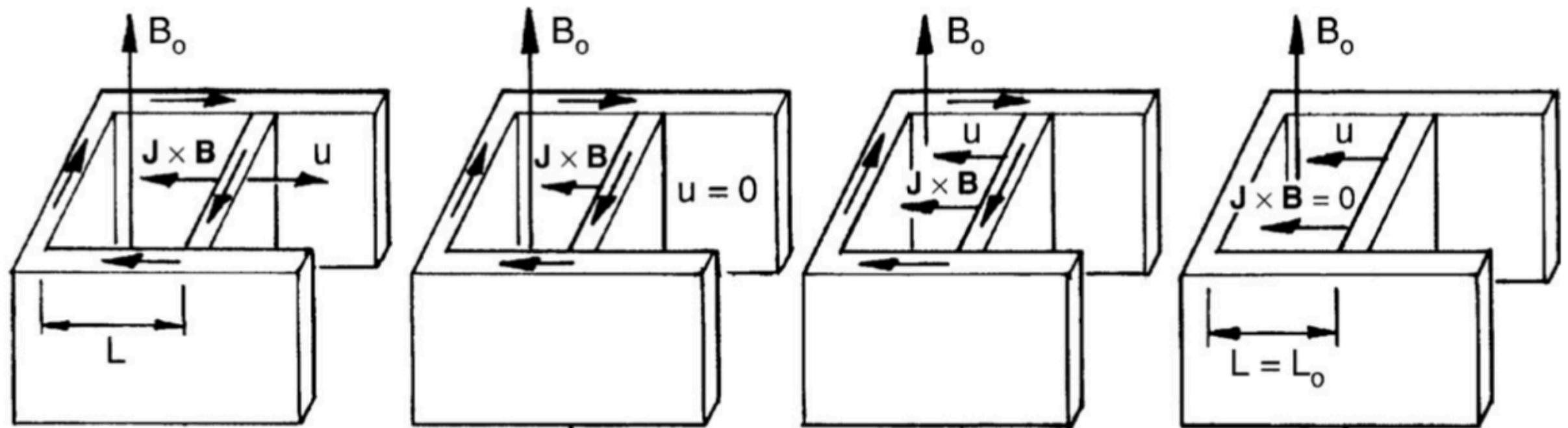
# The perfect conducting case $R_m \gg 1$

motion equation

$$d \frac{d^2 \eta}{dt^2} + \frac{B_0^2}{\rho \mu d L_0} \eta = 0$$

The key equation?

$$\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S} = - \frac{1}{\sigma} \oint_C \mathbf{j} \cdot d\mathbf{l}$$



Flux =  $(B_0 - B_i) LW$   
= constant

Slide reaches  
a halt

Slide  
reverses

$L$  returns to  $L_0$   
and  $J$  falls to zero



# The poor conducting case $R_m \ll 1$

In this case,  $R_m = \frac{u}{L} \ll \frac{1}{T}$ , or  $R_m = \mu\sigma ul \ll 1$ ,

The induction equation tells that  $B_i \ll B_0$

$$\frac{d}{dt} \left[ L (B_0 - \cancel{B_i}) \right] = \frac{2(L+W)B_i}{T} \approx \frac{dLB_0}{dt} = uB_0 \longrightarrow B_i = \frac{uT}{2(L+W)} B_0$$

substitute

$$\rho d \frac{du}{dt} = \frac{(B_0 - B_i)^2}{2\mu} - \frac{B_i^2}{2\mu} \longrightarrow \boxed{\frac{du}{dt} + \frac{W}{2(L+W)} \left( \frac{\sigma B_0^2}{\rho} \right) u = 0}$$

**Exponential decay equation**

## Physical meaning:

When the magnetic Reynolds number is low, the slide velocity **decays** exponentially on a time scale defined as

$$\tau = \left( \frac{\sigma B_0^2}{\rho} \right)^{-1}$$

# The poor conducting case $R_m \ll 1$

It is easy to show that the time rate of change in the kinetic energy is:

$$\frac{dE_{kin}}{dt} = - \int \frac{J^2}{\sigma} dV$$

The magnetic field now appears to play a dissipative role. Thus the mechanical energy of the fluid is lost to heat via **Ohmic (Joule)** dissipation.

In such a low conductivity (high resistivity) case, the induced magnetic field is small, then the Faraday's law becomes:

$$\nabla \times \mathbf{E} = - \frac{d\mathbf{B}}{dt} \approx 0 \quad \mathbf{E} \text{ is potential field}$$

Now the Ohm's law and the Lorentz force are written as

$$\mathbf{J} = \sigma (-\nabla V + \mathbf{u} \times \mathbf{B}_0)$$

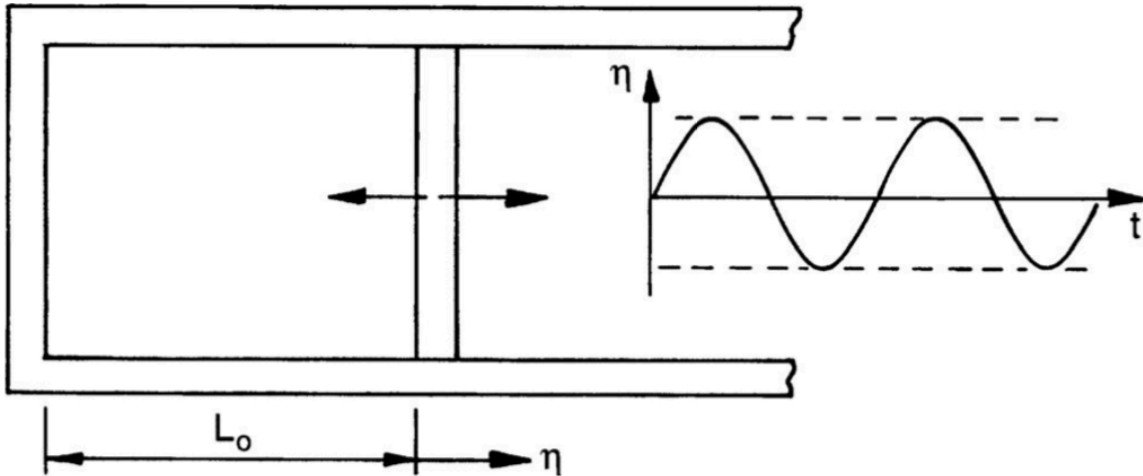
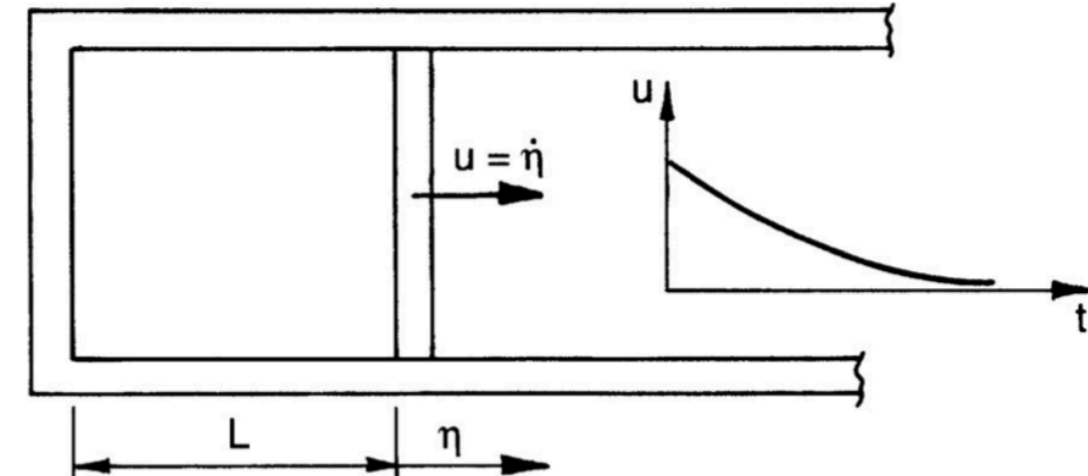
$$\mathbf{F} = \mathbf{J} \times \mathbf{B}_0$$

Electric  
potential

Constant  
Magnetic field

These two equations are the  
hallmark of low- $R_m$  MHD

# Summary of the Electromagnetism case

	Physical Picture	Interpretation
$R_m \gg 1$	 <p>The diagram shows a rectangular loop with a movable rod of length <math>L_0</math> that can slide along the rails with velocity <math>\eta</math>. To the right, a graph of <math>\eta</math> versus time <math>t</math> shows a sinusoidal wave oscillating between two constant values, indicating that the induced magnetic field is comparable to the external field <math>B_0</math>.</p>	<ul style="list-style-type: none"> <li>Induced B field comparable to <math>B_0</math></li> <li>Magnetic flux is conserved</li> <li><math>\mathbf{J} \times \mathbf{B}</math> force is non dissipative</li> </ul>
$R_m \ll 1$	 <p>The diagram shows a rectangular loop with a movable rod of length <math>L</math> that moves with velocity <math>u = \dot{\eta}</math>. To the right, a graph of <math>u</math> versus time <math>t</math> shows a decaying curve starting from a maximum value and approaching zero, indicating that the induced magnetic field is negligible.</p>	<ul style="list-style-type: none"> <li>Induced B field negligible</li> <li>Magnetic flux is not conserved</li> <li><math>\mathbf{J} \times \mathbf{B}</math> force is dissipative</li> </ul>

# Now When is MHD applicable?

validity of the MHD assumption

- **Slow, Non-relativistic:**  $u \ll c$
- **Fluid assumption**

Scale analysis:

$$\frac{\partial}{\partial t} f + A \cdot f = S \quad \xrightarrow[\frac{\partial}{\partial t} \sim \omega \sim \frac{1}{\tau_\omega}]{f \sim e^{i\omega t}} \quad \frac{1}{\tau_\omega} f + A \cdot f = S$$

$0, \text{ if } \frac{1}{\tau_\omega} \ll A$

$0, \text{ if } \frac{1}{\tau_\omega} \gg A$

**This is used extensively in the theory of MHD (kinetic as well)**

# Now When is MHD applicable?

## What does “Slow” mean?

- **Slow, Non-relativistic:  $u \ll c$**

In MHD “slow” means evolution on time scales longer than those on which individual particles are important, or on which the electrons and ions might evolve independently of one another.

$$\text{slow} \iff \frac{\partial}{\partial t} \ll \left\{ \begin{array}{ll} \omega_{pe} \equiv \sqrt{\frac{4\pi e^2 n_e}{m_e}} & \text{plasma frequency} \\ \Omega_i \equiv \frac{eB}{cm_p} & \text{Ion gyro-frequency} \\ \nu_{ei} \equiv \frac{\pi n_e e^4 \ln \Lambda}{\sqrt{m_e} (k_B T)^{3/2}} & \text{collision frequency} \end{array} \right.$$

				frequencies rad/sec			
	$n_e$ [cm <sup>-3</sup> ]	$T$ [K]	$B$ [G]	$\omega_{pe}$	$\Omega_i$	$\nu_{ei}$	MHD
magnetotail	1	10 <sup>7</sup>	10 <sup>-4</sup>	10 <sup>5</sup>	0.1	10 <sup>-9</sup>	$\Delta t \gg 10$ sec <sup>†</sup>
Solar corona (AR)	10 <sup>9</sup>	10 <sup>6</sup>	10 <sup>2</sup>	10 <sup>9</sup>	10 <sup>6</sup>	10 <sup>2</sup>	$\Delta t \gg 10$ ms
Solar interior (200 Mm)	10 <sup>23</sup>	10 <sup>6</sup>	10 <sup>5</sup>	10 <sup>16</sup>	10 <sup>9</sup>	10 <sup>16</sup>	$\Delta t \gg 10$ ns
Galaxy	1	10 <sup>4</sup>	10 <sup>-6</sup>	10 <sup>5</sup>	10 <sup>-2</sup>	10 <sup>-6</sup>	$\Delta t \gg 10$ days

此处忽略碰撞

## Ionosphere

# Now When is ideal MHD applicable?

## validity of the ideal MHD assumption

- **Assume the plasma behaves as a fluid:**
  - highly collisional
  - Distribution is Maxwellian - in thermal equilibrium
  - Only consider macroscopic behavior (low frequency, long wavelength)
- **Assume the ion gyro radius is small: (what about electrons?)**
- **Classic physics - ignore the most significant physics advances since 1860**
  - No displacement current
  - Non relativistic ( $u \ll c$ )
  - No quantum physics
- **Assume the plasma fully ionized:**
  - Limited applications in partially ionized plasmas like the photosphere, ionosphere
- **Ignore resistivity, viscosity, thermal conduction, radiation**

**Question: is the ideal MHD assumption valid for solar wind, magnetospheric, ionospheric plasmas?**