

Homework 4: Phys 5210 (Fall 2021)

Tien Vo

September 24, 2021

Problem 1: Two particles move about each other in circular orbits under the influence of gravitational forces (so that the potential energy is $U = -\alpha/r$), with a period τ .

Their motion is suddenly stopped at a given instant of time, and they are then released and allowed to fall into each other.

Prove that they collide after a time $\tau/(4\sqrt{2})$.

Solution.

From the general solution derived in class, a circular orbit has a radius

$$R = \frac{1}{x_+} = \frac{1}{x_-} = \frac{l^2}{\mu\alpha} \quad (1.1)$$

where $l = \mu R^2 \omega$ is the angular momentum. Then the period is

$$\tau = \frac{2\pi}{\omega} = 2\pi \frac{l^3}{\mu\alpha^2} \quad (1.2)$$

Now, the particle is set to free fall with an initial energy $E_0 = -\alpha/R$. By energy conservation,

$$\frac{1}{2}\mu\dot{r}^2 - \frac{\alpha}{r} = E_0 = -\frac{\alpha}{R} \quad (1.3)$$

Thus, we can write the differential equation

$$dt = \sqrt{\frac{\mu}{2\alpha}} \left(\frac{1}{r} - \frac{1}{R} \right)^{-1/2} dr \quad (1.4)$$

Integrating both sides, we can find the time for the particles collide as

$$\tau' = \sqrt{\frac{\mu}{2\alpha}} \int_0^R \left(\frac{1}{r} - \frac{1}{R} \right)^{-1/2} dr = \frac{\pi}{2} \sqrt{\frac{\mu}{2\alpha}} R^{3/2} = \frac{\pi}{2\sqrt{2}} \frac{l^3}{\mu\alpha^2} = \frac{\tau}{4\sqrt{2}} \quad (1.5)$$

□

Problem 2: A central force potential is given by

$$U(r) = \begin{cases} 0, & r > a \\ -U_0, & r < a \end{cases} \quad (2.1)$$

Here $U_0 > 0$. We would like to study the scattering in this potential.

(a) If a particle approaches the potential with the kinetic energy E and the impact parameter s , find the angle φ and $\tilde{\varphi}$ as shown in the Figure above. The figure represents the potential as a circle of radius a , with the trajectory of a particle shown in blue. Use conservation of energy and conservation of angular momentum to find $\tilde{\varphi}$.

(b) Determine the scattering angle θ in terms of φ and $\tilde{\varphi}$.

(c) Express s as a function of θ . This step is tricky as it is easy to find θ in terms of s , harder the other way around. To do that, take the ratio $\sin \frac{\tilde{\varphi}}{\sin \varphi}$. On the one hand, use the results of part (a) to figure out what it is equal to. On the other hand, express $\tilde{\varphi}$ in terms of φ and θ , and work to solve the resulting equation for φ . Finally, use the previous result relating φ to s .

(d) Finally, find the differential cross section $d\sigma/d\Omega$ where $d\Omega = 2\pi \sin(\theta)d\theta$ for a particle scattering in this potential.

Solution.

(a) The angular momentum is defined as $\mathbf{L} = \mathbf{r} \times \mathbf{p}$. At the point of impact (when the particle enters the potential field U), the radius is $r = a$. Also, let v_0 be the initial velocity. Then the initial angular momentum is

$$L = mav_0 \sin \varphi \quad (2.2)$$

Let v be the velocity briefly after it has entered U . The radius is still a , but the angle between \mathbf{r} and \mathbf{p} is $\tilde{\varphi}$ and the angular momentum is

$$L = mav \sin \varphi \quad (2.3)$$

From conservation of the angular momentum, we can write

$$v = v_0 \frac{\sin \varphi}{\sin \tilde{\varphi}} \quad (2.4)$$

From energy conservation and the above result, we can write

$$E_0 = \frac{1}{2}mv_0^2 = \frac{1}{2}mv^2 - U_0 = \frac{1}{2}mv_0^2 \frac{\sin^2 \varphi}{\sin^2 \tilde{\varphi}} - U_0 \quad (2.5)$$

Inverting, we can write

$$\sin \tilde{\varphi} = \left(1 + \frac{U_0}{E_0}\right)^{-1/2} \sin \varphi \quad (2.6)$$

From geometry, $\sin \varphi = s/a$, so we can write $\tilde{\varphi}$ as

$$\tilde{\varphi} = \sin^{-1} \left[\frac{s}{a} \left(1 + \frac{U_0}{E_0}\right)^{-1/2} \right] \quad (2.7)$$

(b) From geometry, we can also write $\theta = 2(\varphi - \tilde{\varphi}) \Rightarrow \tilde{\varphi} = \varphi - \theta/2$.

(c) From part (b), we can rewrite (2.6) as

$$\sin\left(\varphi - \frac{\theta}{2}\right) = \sin\varphi \cos\left(\frac{\theta}{2}\right) - \cos\varphi \sin\left(\frac{\theta}{2}\right) = \left(1 + \frac{U_0}{E_0}\right)^{-1/2} \sin\varphi \quad (2.8)$$

Solving for φ , we have

$$\tan\varphi = f(\theta) = \frac{\sin(\theta/2)}{\cos(\theta/2) - (1 + U_0/E_0)^{-1/2}} = \frac{s}{\sqrt{a^2 - s^2}} \quad (2.9)$$

Inverting, we can write s as a function of θ

$$s^2 = a^2 \frac{f^2(\theta)}{1 + f^2(\theta)} = a^2 \frac{\sin^2(\theta/2)}{1 + C^2 - 2C \cos(\theta/2)} \quad (2.10)$$

where $C = \sqrt{E_0/(E_0 + U)}$.

(d) Taking the derivative of (2.10) with Mathematica, we can write the differential cross section as

$$\sigma_{\text{diff}} = \frac{1}{2 \sin \theta} \frac{ds^2}{d\theta} = \frac{a^2}{4} \frac{1 + C^2 - C[\cos(\theta/2) + \sec(\theta/2)]}{[1 + C^2 - 2C \cos(\theta/2)]^2} \quad (2.11)$$

Let $n = 1/C$, this can be simplified into

$$\begin{aligned} \sigma_{\text{diff}} &= \frac{a^2}{4} \frac{1 + \frac{1}{n^2} - \frac{1}{n} [\cos(\theta/2) + \sec(\theta/2)]}{[1 + \frac{1}{n^2} - \frac{2}{n} \cos(\theta/2)]^2} \\ &= \frac{n^2 a^2}{4} \frac{n^2 + 1 - n \cos(\theta/2) - \frac{n}{\cos(\theta/2)}}{[n^2 + 1 - 2n \cos(\theta/2)]^2} \\ &= \frac{n^2 a^2}{4} \frac{(n \cos(\theta/2) - 1)(n - \cos(\theta/2))}{(n^2 + 1 - 2n \cos(\theta/2))^2} \end{aligned} \quad (2.12)$$

□