

Ex 3

Tubriums-Anmerkungen in blau (oder rose oder ss)

$$H_A = \hbar\omega_g |e\rangle\langle e| + \hbar\omega_g |s\rangle\langle s| + \frac{\hbar\Omega_1}{2} (e^{-i\omega_1 t} |e\rangle\langle g| + e^{i\omega_1 t} |g\rangle\langle e|) + \frac{\hbar\Omega_2}{2} (e^{-i\omega_2 t} |e\rangle\langle s| + e^{i\omega_2 t} |s\rangle\langle e|)$$

↓ $U(t)$

$$\tilde{H}_A = \hbar\Delta |e\rangle\langle e| - \hbar\Delta |s\rangle\langle s| + \frac{\hbar}{2}\Omega_1 (|e\rangle\langle g| + |g\rangle\langle e|) + \frac{\hbar\Omega_2}{2} (|e\rangle\langle s| + |s\rangle\langle e|)$$

$$\Delta = \omega_1 - \omega_g$$

$$\delta = \omega_1 - \omega_2 - \omega_g$$

Step 1 $U_1 = e^{-i\omega_1 t |e\rangle\langle e|} \rightarrow U H U^\dagger$ step by step.

$$\left[U H(t) U^\dagger + i \hbar \dot{U} U^\dagger \right] | \cdot \rangle$$

$$\tilde{H}(t)$$

→ $U|e\rangle\langle e|U^\dagger = |e\rangle\langle e|$, $U|e\rangle\langle g|U^\dagger = e^{i\omega_1 t}|e\rangle\langle g|$

→ $i\hbar\dot{U}U^\dagger = -\hbar\omega_1|e\rangle\langle e|$

$$U \hbar\omega_g |e\rangle\langle e| U^\dagger = |e\rangle\langle e| \hbar\omega_g$$

$$U \frac{\hbar\Omega_2}{2} e^{-i\omega_2 t} |e\rangle\langle s| U^\dagger = \frac{\hbar\Omega_2}{2} e^{-i\omega_2 t} e^{-i\omega_1 t} |e\rangle\langle s|$$

$$U \hbar\omega_g |s\rangle\langle s| U^\dagger = \hbar\omega_g |s\rangle\langle s|$$

$$U \frac{\hbar\Omega_1}{2} e^{i\omega_1 t} |e\rangle\langle g| U^\dagger = \frac{\hbar\Omega_1}{2} e^{i\omega_1 t} e^{i\omega_1 t} |e\rangle\langle g|$$

$$U \frac{\hbar\Omega_1}{2} e^{-i\omega_1 t} |g\rangle\langle e| U^\dagger = \frac{\hbar\Omega_1}{2} e^{-i\omega_1 t} e^{-i\omega_1 t} |g\rangle\langle e|$$

$$U \frac{\hbar\Omega_2}{2} e^{i\omega_2 t} |g\rangle\langle e| U^\dagger = \frac{\hbar\Omega_2}{2} e^{i\omega_2 t} e^{i\omega_1 t} |g\rangle\langle e|$$

$$U|e\rangle = e^{-i\omega_1 t |e\rangle\langle e|} |e\rangle = \begin{pmatrix} 1 & 0 \\ 0 & e^{-i\omega_1 t} \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ e^{-i\omega_1 t} \end{pmatrix} = e^{-i\omega_1 t} |e\rangle$$

$$U|g\rangle = e^{-i\omega_1 t |e\rangle\langle e|} |g\rangle = \begin{pmatrix} 1 & 0 \\ 0 & e^{-i\omega_1 t} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} |g\rangle$$

$$\langle e|U^\dagger = (U|e\rangle)^\dagger = (e^{-i\omega_1 t} |e\rangle)^\dagger = \langle e| e^{i\omega_1 t}$$

$$\langle g|U^\dagger = \langle g|$$

$$U|e\rangle\langle g|U^\dagger = e^{i\omega_1 t} |e\rangle\langle g|, U|g\rangle\langle e|U^\dagger = |g\rangle\langle e| e^{i\omega_1 t}$$

$$\Rightarrow U H(t) U^\dagger = \hbar\omega_g |e\rangle\langle e| + \hbar\omega_g |s\rangle\langle s| + \frac{\hbar\Omega_1}{2} (e^{-2i\omega_1 t} |e\rangle\langle g| + e^{2i\omega_1 t} |g\rangle\langle e|) + \frac{\hbar\Omega_2}{2} (e^{-i(\omega_1+\omega_2)t} |e\rangle\langle s| + e^{i(\omega_1+\omega_2)t} |s\rangle\langle e|)$$

$i\hbar\dot{U}U^\dagger$:

$$\dot{U} = i\omega_1 e^{i\omega_1 t |e\rangle\langle e|}$$

$$\Rightarrow \dot{U}U^\dagger = i\omega_1 e^{i\omega_1 t |e\rangle\langle e|} e^{-i\omega_1 t |e\rangle\langle e|} = i\omega_1 |e\rangle\langle e|$$

$$U = e^{i\omega_1 t |e\rangle\langle e|}$$

$$\Rightarrow i\hbar\dot{U}U^\dagger = -\hbar\omega_1 |e\rangle\langle e| \rightarrow \hbar\omega_g |e\rangle\langle e| - \hbar\omega_1 |e\rangle\langle e| = -\hbar\Delta |e\rangle\langle e|$$

$$\Rightarrow \tilde{H}_A(t) = -\hbar\Delta |e\rangle\langle e| + \hbar\omega_g |s\rangle\langle s| + \frac{\hbar\Omega_1}{2} (e^{-2i\omega_1 t} |e\rangle\langle g| + e^{2i\omega_1 t} |g\rangle\langle e|) + \frac{\hbar\Omega_2}{2} (e^{-i(\omega_1+\omega_2)t} |e\rangle\langle s| + e^{i(\omega_1+\omega_2)t} |s\rangle\langle e|)$$

Now step 2: $U_2 = e^{i(\omega_1 - \omega_2)t} |s\rangle\langle s|$

$$\Rightarrow -\hbar\omega_1 U |e\rangle\langle e| U^\dagger = -\hbar\omega_1 |e\rangle\langle e| \quad \left| \quad \frac{\hbar\Omega_1}{2} e^{-2i\omega_1 t} U |e\rangle\langle g| U^\dagger = \frac{\hbar\Omega_1}{2} e^{-2i\omega_1 t} |e\rangle\langle g| \right.$$

$$\hbar\omega_2 U |s\rangle\langle s| U^\dagger = \hbar\omega_2 |s\rangle\langle s| \quad \left| \quad \frac{\hbar\Omega_2}{2} e^{2i\omega_2 t} U |g\rangle\langle e| U^\dagger = \frac{\hbar\Omega_2}{2} e^{2i\omega_2 t} |g\rangle\langle e| \right.$$

$$\frac{\hbar\Omega_2}{2} e^{-i(\omega_1 + \omega_2)t} U |e\rangle\langle s| U^\dagger = \frac{\hbar\Omega_2}{2} e^{-i(\omega_1 + \omega_2)t} e^{-i(\omega_1 - \omega_2)t} |e\rangle\langle s| = \frac{\hbar\Omega_2}{2} e^{-2i\omega_1 t} |e\rangle\langle s|$$

$$\frac{\hbar\Omega_2}{2} e^{i(\omega_1 + \omega_2)t} U |s\rangle\langle e| U^\dagger = \frac{\hbar\Omega_2}{2} e^{i(\omega_1 + \omega_2)t} e^{i(\omega_1 - \omega_2)t} |s\rangle\langle e|$$

$$= \frac{\hbar\Omega_2}{2} e^{2i\omega_2 t} |s\rangle\langle e|$$

fast oscillating

fast osc.

$$\text{and } i\dot{U} = i(\omega_1 - \omega_2) e^{i(\omega_1 - \omega_2)t} |s\rangle\langle s| e^{-i(\omega_1 - \omega_2)t} |s\rangle\langle s| = -(\omega_1 - \omega_2) |s\rangle\langle s|$$

$$\Rightarrow \hat{H} = -\hbar\omega_1 |e\rangle\langle e| - \hbar\omega_2 |s\rangle\langle s| + \frac{\hbar\Omega_1}{2} (|e\rangle\langle g| + |g\rangle\langle e|) + \frac{\hbar\Omega_2}{2} (|e\rangle\langle s| + |s\rangle\langle e|)$$

wrong lol

$$\text{use } U_1 = e^{i\omega_1 t} |e\rangle\langle e|$$

$$U_2 = e^{i(\omega_1 - \omega_2)t} |s\rangle\langle s|$$

$$U_1 = e^{-i\omega_1 t} |g\rangle\langle g|$$

works well but then we have to energy shift

$$U_2 U_1 H U_1^\dagger U_2^\dagger = \hbar\omega_2 |e\rangle\langle e| + \hbar\omega_2 |s\rangle\langle s| + \Omega_1 (|e\rangle\langle g| + \text{h.c.}) + \Omega_2 (|e\rangle\langle s| + \text{h.c.}),$$

$$i\dot{U}_1 U_1^\dagger + i\dot{U}_2 U_2^\dagger = \omega_1 |g\rangle\langle g| + \omega_2 |s\rangle\langle s|$$

$$\hat{H} = \omega_2 |e\rangle\langle e| + \omega_1 |g\rangle\langle g| + \omega_2 |s\rangle\langle s|$$

$$H \rightarrow H - \omega_1 \mathbb{1} = H - \omega_1 (|e\rangle\langle e| + |g\rangle\langle g| + |s\rangle\langle s|)$$

$$\rightarrow \hat{H} - \omega_1 \mathbb{1} = |e\rangle\langle e| (\omega_2 - \omega_1) + |g\rangle\langle g| (\omega_1 - \omega_1) + |s\rangle\langle s|$$

→ what we did here

$$U_1 = e^{-i\omega_1 |g\rangle\langle g| t}$$

$$U_2 = e^{-i\omega_2 |S\rangle\langle S|}$$

$$H \rightarrow H - \omega_1 \mathbb{1}$$

$$U = e^{i\omega_1 \mathbb{1} t}$$

$$\xrightarrow{H \rightarrow i\hbar \frac{d}{dt}}$$

$$\begin{pmatrix} U H U^\dagger = H \\ -i\hbar U^\dagger = -\omega_1 \mathbb{1} \end{pmatrix}$$

$$U'_1 = e^{i\omega_1 t |e\rangle\langle e|}$$

$$U'_2 = e^{i(\omega_1 - \omega_2)t |S\rangle\langle S|}$$

$$\rightarrow U U_2 U_1 = e^{i\omega_1 t \mathbb{1}} e^{-i\omega_2 t |g\rangle\langle g|} e^{-i\omega_2 t |S\rangle\langle S|} = U'_2 U'_1$$

$$\hat{H} = \omega_{eg} |e\rangle\langle e| + \Omega (|e\rangle\langle g| + |g\rangle\langle e|)$$

$\Omega \ll \omega_{eg} \rightarrow$ use perturbation theory

$$\hat{H} = \omega_{eg} |e\rangle\langle e| + \Omega (e^{-i\omega_{eg} t} |e\rangle\langle g| + e^{i\omega_{eg} t} |g\rangle\langle e|)$$

can we still apply perturbation theory?

$$\Omega \ll \omega_{eg}$$

$$U_1 = e^{i\omega_{eg} t |e\rangle\langle e|}$$

$$H = \Omega (|e\rangle\langle g| + |g\rangle\langle e|) + (\omega_{eg} - \omega_{eg}) |e\rangle\langle e|$$

→ can't just do the same perturbation theory now since there's no ω_{eg} -term to compare it to

$$H = \Omega (|e\rangle\langle g| + |g\rangle\langle e|) + \Delta |e\rangle\langle e|$$

$\Delta \neq 0$

→ control Δ to adjust the magnitude of perturbation

best way to see if a term is big/negligible or not is to switch to a ref. frame where it is time-independent and compare it to other terms as well

the question is not only "is it fast oscillating" but rather "is it fast oscillating in a frequency that is away from the systems frequency?" (I think)

safer option to solve-ivp: often: just take $u = e^{-i\Omega t}$, especially if you're just interested in one specific t -point