

Laminar–Turbulent Sheet Conductivity Scaling

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Sheet-flow models

Table 1 lists the different sheet-flow options:

Table 1: Model descriptions	
Model	
Laminar	$\mathbf{q} = -k_s h^3 \nabla \phi$
Turbulent	$\mathbf{q} = -k_s h^{\alpha_s} \nabla \phi ^{-1/2} \nabla \phi$
Transition	$\mathbf{q} + \frac{\omega}{\nu} \left(\frac{h}{h_b} \right)^{3-2\alpha_s} \mathbf{q} \mathbf{q} = -k_s h^3 \nabla \phi$

Adjusting the sheet conductivity

Note that the units of k_s for the turbulent model depend on the value of α_s and are different from the units for the laminar and transition models. This means we need to find a relationship between the conductivity for the turbulent and laminar/transition models to make them comparable.

There are at least two ways to do this:

1. **Asymptotically match the behaviour.** The transition model asymptotically approaches the laminar model as $\text{Re} = \frac{q}{\nu} \rightarrow 0$ by definition. We could also tune the conductivity of the turbulent model so that the turbulent model matches the $\text{Re} \rightarrow \infty$ limit of the transition model.

Taking $\text{Re} \rightarrow \infty$, and taking the absolute value to simplify, the transition model becomes

$$\frac{\omega}{\nu} \left(\frac{h}{h_b} \right)^{3-2\alpha_s} q^2 = k_s h^{3/2} |\nabla \phi|^{1/2}. \quad (1)$$

This can be written to look like the turbulent model,

$$q = \left[k_s \frac{\nu}{\omega} (h_b)^{2\alpha_s-3} \right]^{1/2} h^{\alpha_s} |\nabla \phi|^{1/2}. \quad (2)$$

So, we've recovered the turbulent model as the $\text{Re} \rightarrow \infty$ limit of the transition model. We can then compute what k_s we should use in the transition/laminar model so that the entire coefficient in Eq. 2 matches the turbulent conductivity we have been using so far.

2. **Match laminar and turbulent models at the transition Re.** The idea is to have $\mathbf{q}_{\text{laminar}} = \mathbf{q}_{\text{turbulent}}$ at the transition between laminar and turbulent sheet flow. We do this by first computing the sheet thickness that gives $\omega \text{Re} = 1$. It's easiest to do this from the laminar model:

$$1 = \omega \text{Re} = \omega \frac{q}{\nu} = \frac{\omega}{\nu} k_s h_{\text{crit}}^3 |\nabla \phi|, \quad (3)$$

so that

$$h_{\text{crit}}^3 = \frac{\nu}{\omega k_s |\nabla \phi|}. \quad (4)$$

Now we substitute this expression for h_{crit}^3 into the turbulent model, setting $\omega \text{Re} = \omega \frac{g}{\nu} = 1$. For clarity, write the conductivity parameter for the turbulent model as k_t . Then,

$$1 = \frac{\omega}{\nu} k_t h_{\text{crit}}^{\alpha_s} |\nabla \phi|^{1/2}. \quad (5)$$

Simplifying,

$$k_s^\alpha = k_t^3 \left(\frac{\omega}{\nu} \right)^{3-\alpha_s} |\nabla \phi|^{(\frac{3}{2}-\alpha)}. \quad (6)$$

At this point you can go either direction to convert between k_s and k_t . If you want to continue with a conductivity that matches what you've been using for the turbulent model, use this for your k_t and solve for the laminar/transition conductivity k_s . Note that the $\nabla \phi$ dependence goes away if (and only if) the turbulent sheet conductivity $\alpha_s = 3/2$, so here you can use any value for $\nabla \phi$. For the turbulent model with $\alpha_s = 5/4$, you will need to assume a value for the hydraulic potential gradient and the result will depend on this value. The sensitivity is very low ($|\nabla \phi|^{1/12}$), so I've just used a taken the difference in potential (assuming $p_w = p_i$) between the terminus and highest point divided by the domain length.

For my simulations, it turned out that option (2) provided a more fair comparison between the models. For the same laminar/transition conductivity, option (1) results in a higher turbulent sheet conductivity, such that it produced fewer and smaller channels and extremely low winter water pressure.

In practice

For my simulations, parameters are in Table 2. I chose a laminar/transition sheet conductivity $k_s = 0.05 \text{ Pas}^{-1}$, and this gave a conductivity for the turbulent 5/4 model of 0.0071 (with appropriate units) using option (2) for scaling (Table 3).

Table 2: Constants (top group) and model parameters (bottom group) for GlaDS simulations.

Symbol	Description	Value	Units
ρ_w	Density of water	1000	kg m ⁻³
ρ_i	Density of ice	910	kg m ⁻³
g	Gravitational acceleration	9.81	m ³ s ⁻¹
c_w	Specific heat capacity of water	4.22×10^3	J kg ⁻¹
c_t	Pressure melting coefficient	-7.50×10^{-8}	K Pa ⁻¹
ν	Kinematic viscosity of water at 0°C	1.793×10^{-6}	m s ⁻²
k_s	Effective laminar sheet conductivity	0.1	Pa s ⁻¹
α_s	Sheet-flow exponent	$[\frac{5}{4}, \frac{3}{2}, 3]$	
β_s	Sheet-flow exponent	$[\frac{3}{2}, 2]$	
k_c	Channel conductivity	0.2	m ^{3/2} s ⁻¹
α_c	Channel-flow exponent	5/4	
β_c	Channel-flow exponent	3/2	
h_b	Bed bump height	0.5	m
l_b	Bed bump length	10	m
l_c	Width of sheet-flow contributing to channel	10	m
e_v	Englacial porosity	1×10^{-4}	
ω	Laminar–turbulent transition parameter	1/2000	
u_b	Basal velocity	30	m a ⁻¹
\tilde{A}^a	Rheological parameter for creep closure	1.78×10^{-25}	s ⁻¹ Pa ⁻³
\tilde{A}_s	Rheological parameter for creep when $N < 0$	0	s ⁻¹ Pa ⁻³
n	Ice-flow exponent	3	
\dot{m}_s	Basal melt rate	0.01	m w.e. a ⁻¹

^a \tilde{A} differs from the canonical rheology parameter A by a factor of $\frac{2}{27}$. The listed value for \tilde{A} corresponds to the recommended value $A = 2.4 \times 10^{-24}$ s⁻¹ Pa⁻³ for temperate ice.

Table 3: Comparable model conductivities. Caution, these values are conditioned on $\omega = 1/2000$ (Table 2) and a hydraulic potential gradient of $\nabla\phi = 135$ Pa m⁻¹

Model	Sheet conductivity	Units
Turbulent ($\alpha_s = 5/4$)	0.0071	m ^{7/4} kg ^{-1/2}
Turbulent ($\alpha_s = 3/2$)	0.0134	m kg ^{-1/2}
Laminar	0.05	Pa s ⁻¹
Transition 5/4	0.05	Pa s ⁻¹
Transition 3/2	0.05	Pa s ⁻¹