

Let's recap where we are.

- Isotropic/homogenous universe leads to Robertson-Walker metric
- Using this metric in Einstein's General Relativity leads to the Friedmann Equations.
- The Friedmann equations determine the cosmic scale evolution of the universe.
- So are we done with this part of cosmology?
- Do question (1) on the worksheet and **STOP**

$$(1b) \quad 1 - \Omega(t) = \frac{H_o^2 (1 - \Omega_o)}{H(t)^2 a(t)^2}$$

$$(1d) \quad |1 - \Omega(t)| \leq 2 \times 10^{-6}$$

$$(1c) \quad 1 - \Omega(t) = \frac{(1 - \Omega_o) a^2}{\Omega_{r,o} + a \Omega_{m,o}}$$

$$(1e) \quad |1 - \Omega(t)| \leq 7 \times 10^{-16}$$

$$(1f) \quad |1 - \Omega(t)| \leq 2 \times 10^{-62}$$

This result is referred to as the *Flatness problem*. The problem simply stated is that to see the amount of *flatness* we see today requires  $\Omega$  to be within  $2 \times 10^{-62}$  of unity. That's an incredible number.

We keep saying that the universe is isotropic and homogeneous, and that's a good thing we've said. But consider

- At the time of the last scattering, the distance to the horizon is **0.251 Mpc**
- Today that distance translates to an angular separation of  **$1.1^\circ$**
- The largest temperatures fluctuations in the CMB are on the order of  **$30 \mu\text{K}$**
- There are about **40,000 patches** in the sky of angular separation  **$1.1^\circ$**
- Do question (2) on the worksheet and **STOP**

This problem is referred to as the **horizon** problem.

The monopole problem



But Maxwell's equation allow for  n (and/or s)

These **monopoles** are not seen even though (at least in most modern theories) they should be as prevalent as **dipoles**.

Inflation to the rescue. Do question (3 a) on the worksheet and **STOP**. Finish question (3) on the worksheet and **STOP**

$$(3b) \quad P < -\frac{\epsilon}{3}$$

$$(3c) \quad P = w \cdot \epsilon \Rightarrow w < -\frac{1}{3}$$

$$(3d) \quad \frac{\ddot{a}}{a} = \frac{\Lambda}{3}$$

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{\Lambda}{3}$$

solving for  $a$

$$a \propto e^{H_i t} \left( H_i = \sqrt{\frac{\Lambda}{3}} \right)$$

How does this work to solve problems?

Suppose universe had the following behavior

(pay attention to subscripts on  $t$ )

$$a(t) = \begin{cases} \sqrt{t/t_i} & t < t_i \\ ae^{H_i(t-t_i)} & t_i < t < t_f \\ ae^{H_i(t_f-t_i)} \sqrt{t/t_f} & t > t_f \end{cases}$$

Do question (4 a) and **STOP**

- This scenario describes a universe that up until time  $t_i$  is growing “normally”,
- then between  $t_i$  and  $t_f$  undergoes exponential expansion
- After  $t_f$ , it resumes “normal growth”

It is usual in studying inflation to compare the how much the scale factor changed between  $t_i$  and  $t_f$  by forming the ratio:

$$\frac{a(t_f)}{a(t_i)} = e^N; \text{ where } N \equiv H_i(t_f - t_i); H_i \equiv \sqrt{\frac{\Lambda}{3}}$$

$N$  is called the number of **e-foldings**

Finish problem 4

$$(4 \text{ b}) \quad \epsilon_{\Lambda} = \frac{3c^2}{8\pi G} H_i^2 \approx 10^{105} \text{ TeV m}^{-3} \quad (4 \text{ c}) \quad |1 - \Omega(t_f)| = 2^{-2N} |1 - \Omega(t_i)| \quad (4 \text{ d}) \quad |1 - \Omega(t_f)| = 2^{-2N}$$

The horizon problem. Recall that the horizon distance is given by the relation

$$d_{\text{hor}}(t) = a(t)c \int_0^t \frac{dt'}{a(t')}$$

Do question (5a) on the worksheet and **STOP**      Do question (5b) and (5c) and **STOP**

$$(5 \text{ b}) \quad d_{\text{hor}}(t_i) = a_i c \int_0^{t_i} \frac{dt'}{a_i \sqrt{t/t_i}} = 2ct_i \quad (5 \text{ c}) \quad d_{\text{hor}}(t_f) = a_i c e^N \int_0^{t_i} \frac{dt'}{a_i \sqrt{t/t_i}} + \int_{t_i}^{t_f} \frac{dt}{a_i \exp[H_i(t - t_i)]}$$

When  $N$  is large, then the integrals in (5c) yield  $d_{\text{hor}}(t_f) = e^N c (2t_i + H_i^{-1})$

Finish question (5) on the worksheet and **STOP**

$$d_{\text{hor}}(t_i) = 6 \times 10^{-28} \text{ m} \quad d_{\text{hor}}(t_f) = 15 \text{ m}$$