

- **Proper distance** roughly corresponds to where a distant object would be at a specific moment of cosmological time, which can change over time due to the expansion of the universe.
- **Comoving distance** between fundamental observers does not change with time, as comoving distance accounts for the expansion of the universe.
- **Luminosity distance** is the distance associated with the amount of flux one measures from a distant object
- **Transverse comoving distance**
- **Angular diameter distance**
- **Light Travel distance**

$$d_p \approx \frac{c}{H_o} z \left[ 1 - \frac{1 + q_o}{2} z \right] \quad q_o = -\frac{\ddot{a}(t_o)}{a(t_o)} \frac{1}{H_o^2} = -\frac{a(t_o) \ddot{a}(t_o)}{\dot{a}(t_o)} \quad f = \frac{L}{4\pi d_L^2}; \text{ or } d_L \equiv \left( \frac{L}{4\pi f} \right)^{1/2} \quad d_L = S_k(r)(1 + z)$$

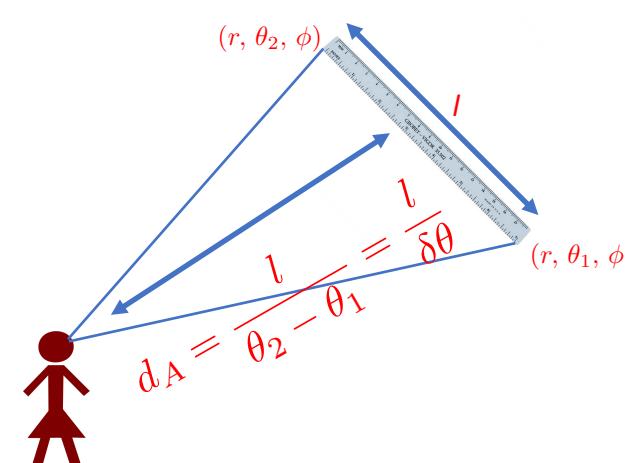
But remember that the universe is expanding.

This means that the light you see coming from the ends of the ruler has undergone expansion. So

$$ds = l = a(t_e) S_k(r) \delta\theta = \frac{S_k(r) \delta\theta}{1+z}$$

And so

$$d_A = \frac{l}{\delta\theta} = \frac{S_k(r)}{1+z}$$



Galaxies and clusters of galaxies are both large enough to serve as standard yardsticks at large distances.

Do question (1) on the worksheet and **S T O P**

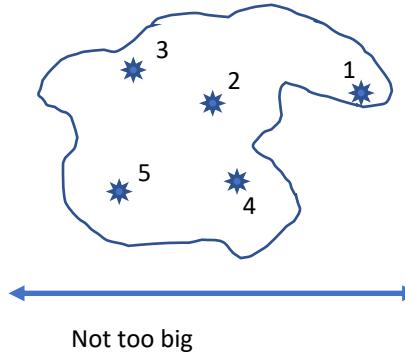
(1) Galaxies are not solid objects, cannot find the edges to find *l*.

We are back to this  $f = \frac{L}{4\pi d_L^2}$ ; or  $d_L \equiv \left(\frac{L}{4\pi f}\right)^{1/2}$ . But we need to find a standard candle that we can use to find  $L$

Recall we want a way to at least measure  $H_0$ , and to do so, we need a distance.

- Find a population of **standard candles**. This gives us  $L$ . They need to be bright enough that we can see them from far away.
- Measure the flux,  $f$ , and redshift,  $z$ , for each standard candle
- Compute,  $d_L = \sqrt{L/4\pi f}$  for each standard candle
- Plot  $cz$  versus  $d_L$
- Measure the slope of  $cz$  versus  $d_L$  when  $z \ll 1$ , this will give  $H_0$

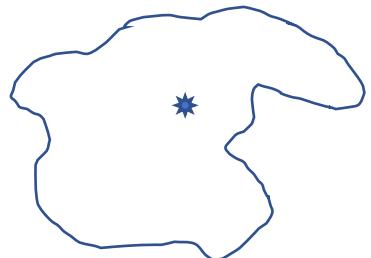
Cepheids, the work of Henrietta Leavitt



Star	Period (days)	Mean flux*
1	10	20
2	25	45
3	15	30
4	12	21
5	27	47

Difference in mean flux is due to their mean luminosity, not their distance.

So there is a period—luminosity relationship



Galaxy far away, star has a period 10 days, but a flux 100 times less than above.

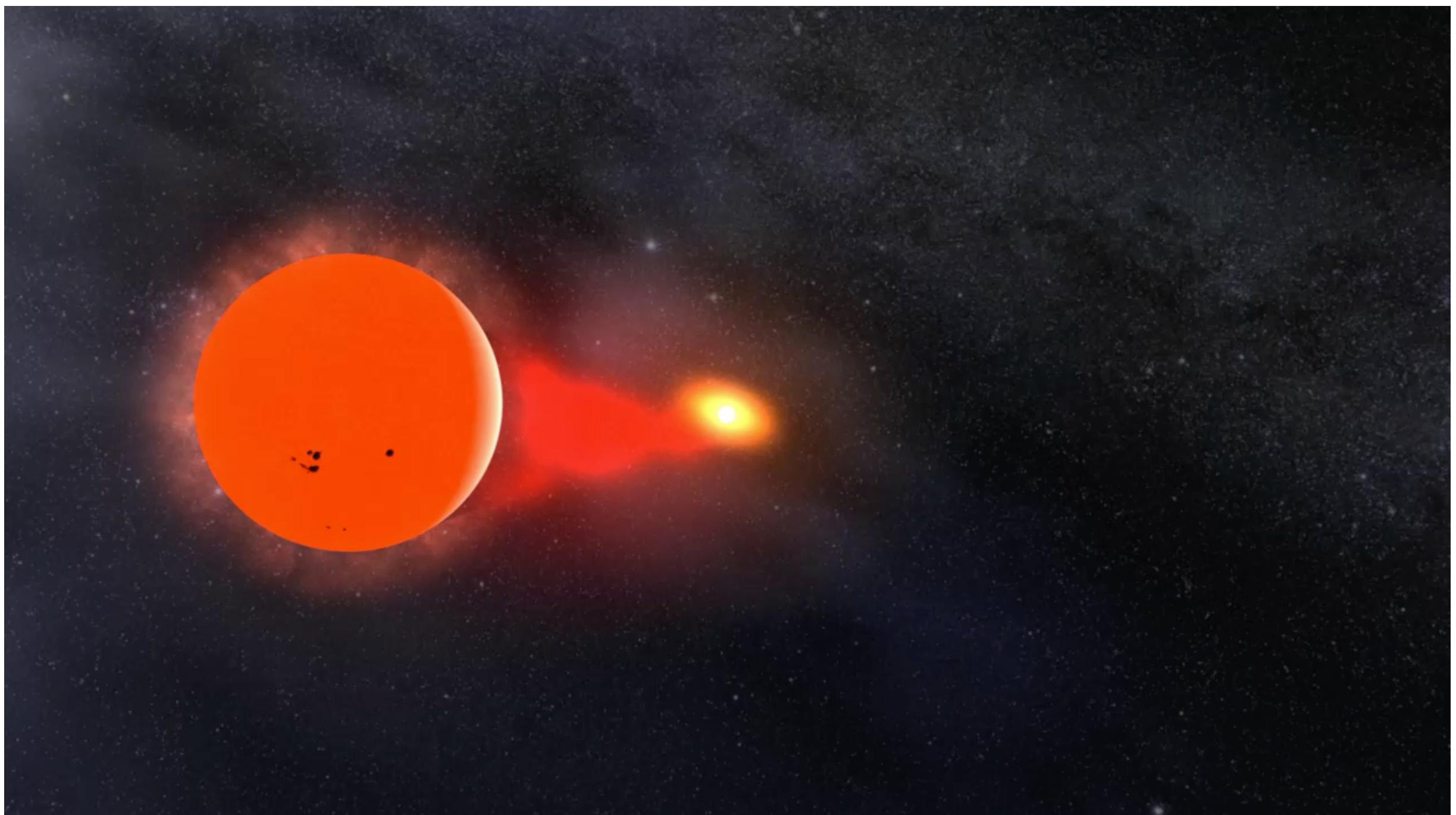
$$d_L = \left( \frac{L}{4\pi f} \right)^{1/2}$$

But from period, we know luminosity is the same as above, so

$$\frac{d_{L,1}}{d_{L,2}} = \left( \frac{\frac{L}{4\pi f_1}}{\frac{L}{4\pi f_2}} \right)^{1/2} = \left( \frac{f_2}{f_1} \right)^{1/2} = \sqrt{100} = 10 \text{ times further}$$

Unfortunately, this only gets us to about 30 Mpc, and the universe is definitely not isotropic/homogeneous on these scales. We can't use Cepheids as standard candles for cosmology

Do question (2) on  
the worksheet and  
**STOP**



## Type IA supernova

- Extremely luminous. A type IA supernova will outshine all the stars in the host galaxy for a few days
- We have a very good idea of their inherent luminosity, so we have a good standard candle.

Do question (3) on the worksheet and **S T O P**

(3)

- Find nearby Cepheid's whose distance can be found by accurate methods (parallax, etc) and determine their luminosity
- Find Cepheid's in nearby galaxies and measure the distances to those galaxies.
- Hope to find a Type 1A supernova in those galaxies and get their luminosity since we know the distance to the galaxy
- Use Type IA supernova to measure distances to many farther away galaxies, and determine  $H_o$

## Cosmic Microwave Background (CMB) Radiation

- We've pushed the *distance* base tests of cosmology about as far as we can. Time to change tactics
- We observe the universe on cosmological scales is
  - Isotropic and homogeneous
  - Obeys the Robertson-Walker Metric and Friedmann equations
  - Is close flat with  $\Omega_m = 0.3\text{-ish}$  and  $\Omega_\Lambda = 0.7\text{-ish}$
  - Filled with a background radiation

The universe is extremely bright in the microwave range. Some characteristics of the radiation include

- $T_o = 2.725 \text{ K}$
- $\epsilon_{\gamma,o} = 0.261 \frac{\text{MeV}}{\text{m}^3}$
- $n_{\gamma,o} = 4.11 \times 10^8 \text{ m}^{-3}$
- $n_{\text{baryons},o}/n_{\gamma,o} \approx 5 \times 10^{-10}$

The CMB is observed across the entire sky and there are **2 key important results** of these observations that will need explaining

1. The CMB is very close to a blackbody spectrum at  $T = 2.73$ . Recall that a blackbody spectrum is

$$I_\nu = \frac{2h\nu^3}{c^2} \left[ e^{h\nu/T} - 1 \right]^{-1}$$

2. After subtracting the effects of the motion of our galaxy (due to the local environment) we have that

$$\langle T \rangle = \frac{1}{4\pi} \int T(\theta, \phi) \sin \theta d\theta d\phi \approx 2.73$$

but there are small fluctuations in the temperature distribution,

$$\frac{\delta T}{T}(\theta, \phi) \equiv \frac{T(\theta, \phi) - \langle T \rangle}{\langle T \rangle} \approx 1.1 \times 10^{-5}$$

Do question (4) on the worksheet and **S T O P**

Wednesday—Week06

The physics. Three terms you need to be aware of: *recombination, decoupling, last scattering surface*.

- **Recombination**—the epoch during which previously ionized baryonic matter becomes neutral
- **Decoupling**—the time at which photons stop interacting with baryonic matter
- **Last scattering**— the time at which a CMB photon undergoes its last scattering from an electron

Some preliminaries

- $H$  is neutral hydrogen
- $p$  is ionized hydrogen (i.e. a proton)
- $e$  is electrons

Is all baryonic matter ionized. No. The fraction that is given by  $X \equiv \frac{n_p}{n_p + n_H} = \frac{n_p}{n_{\text{bary}}} = \frac{n_e}{n_{\text{bary}}}$

Hydrogen is ionized when  $E = 13.6 \text{ eV}$  so a photon with much energy can cause the reaction



An  $X$  depends on the details of these two reactions

Now let's play with the physics a bit. We'll go back to a moment,  $a = 10^{-5}$ ,  $z = 10^5$  and the universe is about 70 years old (using the benchmark model).

At that moment,  $T = 3 \times 10^5$  K;  $E_\gamma = 60$  eV and since there are so many photons than baryons, very little chance that neutral hydrogen is present, and  $X = 1$ .

With the universe fully ionized, the dominant reaction is  $\gamma + e \rightarrow \gamma + e$

Three facts:  $\sigma_e = 6.65 \times 10^{-29} \text{ m}^2$ ;  $\lambda = \frac{1}{n_e \sigma_e}$ ;  $n_e = n_{\text{baryons}} = \frac{n_{\text{baryons},0}}{a^3} = \frac{0.25 \text{ m}^{-3}}{a^3}$

Do question (5) on the worksheet

$$(5a) \quad \Gamma = \frac{c}{\lambda} = n_e \sigma_e c$$

$$(5b) \quad \Gamma = n_e \sigma_e c = 5.0 \times 10^{-6} \text{ s}^{-1}$$

$$(5c) \quad \frac{H^2}{H_o^2} = \frac{\Omega_{r,o}}{a^4} \Rightarrow H = \frac{2.1 \times 10^{-20} \text{ s}^{-1}}{a^2}$$

$$(5d) \quad H = 2.1 \times 10^{-10} \text{ s}^{-1}$$

The expansion is much slower than the collision rate, so photons are well-coupled at this time.