

Last time we started to explore some of the consequences of

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3c^2}\epsilon - \frac{kc^2}{R_o^2 a^2} \quad (1)$$

$$\dot{\epsilon} + 3\frac{\dot{a}}{a}(\epsilon + P) = 0 \quad (2)$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3c^2}(\epsilon + 3P) \quad (3) \quad \text{with } P = w \epsilon$$

To start to build some physical intuition of the effect of each energy density component, we'll now explore *toy* universes composed of single components. We'll then put it all together to give us a picture of the universe. What do we mean by single component? **Epochs:**

$$\begin{matrix} \Omega_r & \Omega_m \\ \Omega_\Lambda & \Omega_{\text{other}} \end{matrix}$$

Current status: $\frac{\Omega_{\Lambda,o}}{\Omega_{m,o}} = 2.3$ so universe is currently dominated by cosmological constant. But this was not always the case

Do question (1) on the worksheet and **STOP**

$$(1) \quad \frac{\epsilon_{\Lambda}(a)}{\epsilon_m(a)} = \frac{\epsilon_{\Lambda,o}}{\epsilon_{m,o}/a^3} = a^3 \frac{\epsilon_{\Lambda,o}}{\epsilon_{m,o}} \quad \text{When, } \epsilon_{\Lambda}(a) = \epsilon_m(a) \text{ then } \frac{\epsilon_{\Lambda}(a)}{\epsilon_m(a)} = 1 = \frac{\epsilon_{\Lambda,o}}{\epsilon_{m,o}/a^3} = a^3 \frac{\epsilon_{\Lambda,o}}{\epsilon_{m,o}} \Rightarrow a_{m=\Lambda} \approx 0.75$$

Doing the same analysis for the matter/radiation terms gives

$$\frac{\epsilon_m(a)}{\epsilon_r(a)} = 1 = \frac{\epsilon_{m,o}}{\epsilon_{r,o}/a} = a \frac{\epsilon_{m,o}}{\epsilon_{r,o}} \Rightarrow a_{m=r} \approx 2.8 \times 10^{-4}$$

Recall that $a = \frac{1}{1+z}$ So that we can solve for z and determine that radiation was dominant when $z \approx 3600$ and dark energy becomes dominant when $z \approx 1/3$

We'll now work our way through several **toy** universes. From each of these we'll get important insights as to how the universe evolves.

We'll begin with what appears to be a silly example, a universe with nothing in it.

Do question (2) on the worksheet and **STOP**

- (2a) When $k = 0$, then $\dot{a} = 0$ and in an empty universe a static, spatially flat universe is permitted.
- (2b) $k > 0$, positively curved, as this would require an imaginary scale factor, a .
- (2c) $\dot{a} = \pm \frac{c}{R_o} \Rightarrow a(t) = \pm \frac{c}{R_o} t$ so that an empty universe must either expand or contract. We'll look at the expanding case.
- (2d) We have $a(t) = \frac{t}{t_o}$ ($t_o \equiv R_o/c$) so that $1 + z = \frac{1}{a(t_e)} = \frac{t_o}{t_e}$

and

$$H_o = \frac{\dot{a}}{a} = \frac{1/t_o}{t_e/t_o} = \frac{1+z}{t_o}$$

- (2e) Actually already done in (2d). We get that $1 + z = \frac{1}{a(t_e)} = \frac{t_o}{t_e} \Rightarrow t_e = \frac{t_o}{1+z}$

$$\begin{aligned} (2f) \quad d_p(t_o) &= c \int_{t_e}^{t_o} \frac{dt}{a(t)} \\ &= c \int_{t_e}^{t_o} dt \frac{t_o}{t} \\ &= ct_o \ln \left(\frac{t_o}{t_e} \right) \text{ or equivalently} \\ d_p(t_o) &= \frac{c}{H_o} \ln(1+z) \end{aligned} \qquad (2g) \quad d_p(t_e) = \frac{c}{H_o} \frac{\ln(1+z)}{1+z}$$

Next we'll add *stuff* to the universe and explore the consequences.

However, we'll first look at universes with no curvature, i.e. $k = 0$

In this case, the Friedmann equation is:

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3c^2}\epsilon - \frac{kc^2}{R_o^2 a^2} \Rightarrow \boxed{\dot{a}^2 = \frac{8\pi G\epsilon_o}{3c^2}a^{-1+3w}} \text{ where we have used that } \epsilon_w = -\epsilon_{w,o}a^{-3(1+w)}$$

Tricky to solve.

Let $a(t) = t^q$. Substituting into the Friedmann equation then results in

$$\begin{aligned} \dot{a}^2 &= \frac{8\pi G\epsilon_o}{3c^2}a^{-(1+3w)} \\ (qt^{q-1})^2 &= Kt^{-q(1+3w)} \text{ Now note that} \\ (qt^{q-1})^2 &= q^2t^{2q-2} \propto t^{2q-2} \text{ and substituting above gives} \\ t^{2q-2} &\propto t^{-q(1+3w)} \text{ and this can hold only if exponents are equal so} \\ -2 + \frac{2}{q} &= 1 + 3w \\ q &= \frac{2}{3 + 3w} \\ \text{or} \\ a(t) &= \left(\frac{t}{t_o}\right)^{2/(3+3w)} \quad w \neq -1; \quad t_o = \frac{1}{1+w} \left(\frac{c^2}{6\pi G\epsilon_o}\right) \end{aligned}$$

Do question (3) on worksheet
and **STOP**

$$(3a) \quad H_o = \frac{1}{t_o} \frac{2}{3(1+3w)} \quad (3b) \quad t_e = \frac{t_o}{(1+z)^{3(1+3w)/2}} = \frac{2}{3(1+w)H_o} \frac{1}{(1+z)^{3(1+w)/2}}$$

$$(3c) \quad d_p(t_o) = c \int_{t_e}^{t_o} \frac{dt}{a(t)} = ct_o \frac{3(1+w)}{1+3w} \left[1 - (t_e/t_o)^{(1+3w)/(3+3w)} \right] \quad (3d) \quad d_{hor}(t_o) = ct_o \frac{3(1+w)}{1+3w} = \frac{c}{H_o} \frac{2}{(1+3w)}$$

To study the behavior of single component universe is now just a matter of using the appropriate w .

Do question (4) on the worksheet and **STOP**

(4) Matter, $w = 0$ $t_0 = \frac{2}{3H_o}$ $d_p(t_o) = \frac{2c}{H_o} \left[1 - \frac{1}{\sqrt{1+z}} \right]$

$d_p(t_e) = \frac{2c}{H_o(1+z)} \left[1 - \frac{1}{\sqrt{1+z}} \right]$ $d_{hor} = \frac{2c}{H_o}$

Radiation, $w = 1/3$ $t_0 = \frac{1}{2H_o}$ $d_p(t_o) = \frac{c}{H_o} \frac{z}{1+z}$

$d_p(t_e) = \frac{c}{H_o} \frac{z}{(1+z)^2}$ $d_{hor} = \frac{c}{H_o}$

$\Lambda, w = -1$ $d_p(t_o) = \frac{c}{H_o} z;$ $d_p(t_e) = \frac{c}{H_o} \frac{z}{1+z}$

Do question (5) on the worksheet