

IB Physics Topic E1 Atomic Structure; SL & HL

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1 Nuclear Notation and Structure

- An atom has 0 net charge; the number of protons equals the number of electrons.
- An anion is negatively charged (with an excess of electrons)
- A cation is positively charged (with a deficit of electrons)

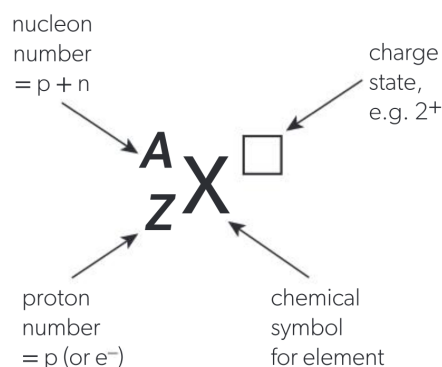


Figure 1: The atomic notation

- Let p = number of protons, n = number of neutrons, e = number of electrons.
- The nucleon number is the number of protons and neutrons, i.e. $A = p + n$
- The proton number is $Z = p$
- From this representation, one can work out the neutron number as $n = A - Z$
- In this notation,

– proton $\rightarrow {}^1_1p$

– neutron $\rightarrow {}^1_0n$

– electron $\rightarrow {}^0_{-1}e$

2 The Plum Pudding Model

This was first proposed by J.J. Thomson in 1897. It was based on the idea that the atom was a sphere of positive charge with electrons embedded in it, with an even distribution and uniform density.

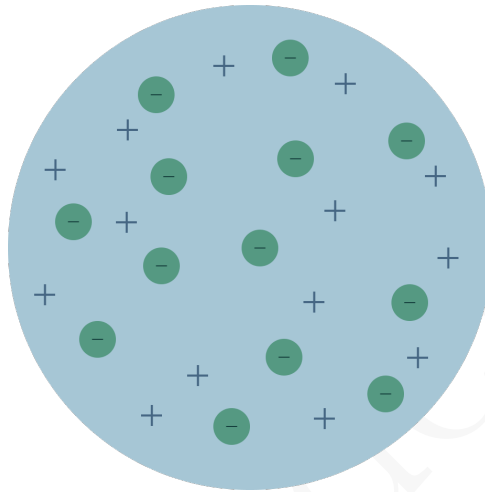


Figure 2: The plum pudding model

3 The Geiger-Marsden Experiment

An experiment carried out in 1911 to determine the structure of an atom, this disproved the plum pudding model.

The experiment is directed by Rutherford but carried out by Geiger and Marsden. The experiment involved firing alpha particles at a thin gold foil and observing the scattering pattern. The gold foil is made thin so that the alpha particles are not subject to the interference of many atoms — in fact, scientists believed that the thin foil was only two lines of atoms thick.

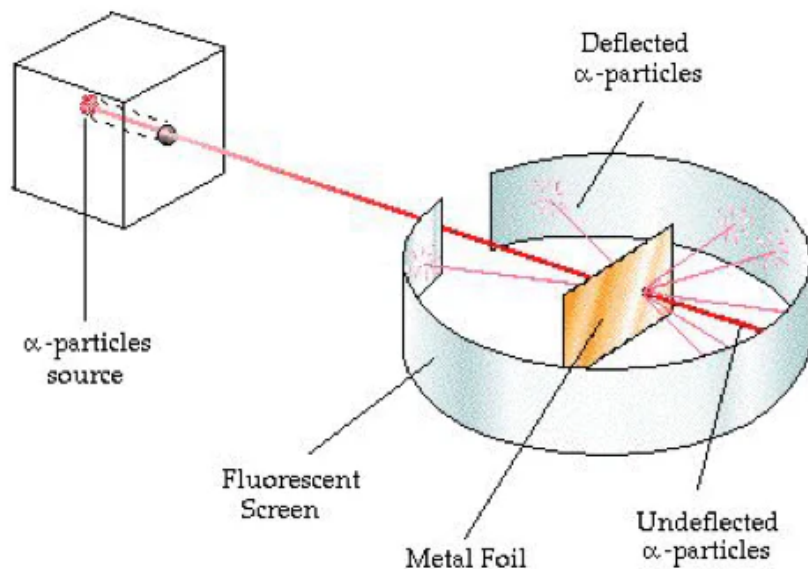


Figure 3: The Geiger-Marsden-Rutherford experiment

If the plum pudding model were correct, only one of the following would have happened:

- Case 1: if the gold atoms have **low density**, all the alpha particles would have effortlessly passed through the foil with minimal deflection.
- Case 2: if the gold atoms have **high density**, all the alpha particles would have been bounded back or deflected at large angles.

What actually happened was that most of the alpha particles passed through the foil with minimal deflection, but some were deflected at large angles, and a few were bounded back. This offers the following implications:

- most of the atom is a free space
- the atom contains small dense regions of positive electric charge

3.1 Deviations from the Scattering

Definition. 1

The distance of closest approach of an alpha particle is the minimum distance it gets to the center of the nucleus of an atom before being repelled by the electrostatic (Coulomb) force.

At high initial KE of the alpha particle E_α , the closest approach distance r_c becomes small.

For alpha particles with initial energy greater than 28 MeV, the scattering pattern observed by Geiger and Marsden was no longer true. The pattern assumed that the alpha particles only interact through electrostatic repulsion. Beyond that limit, the alpha particles will be close enough to the gold nucleus to also interact with the strong nuclear force. This then provides **evidence for the strong nuclear force**.

If $E_{\alpha,\max} \approx 28\text{MeV}$ is the maximum energy for which the scattering pattern still obeys the Rutherford model, then, the estimate gives the smallest r_c at which other nuclear forces to not operate — this is the effective size of a nucleus.

4 Distance of Closest Approach and Energy

1. As an alpha particle approaches the nucleus, WLOG, to the right, it will start to decelerate due to the repulsive electrostatic force and start to lose its kinetic energy.
2. At some point, the alpha particle will come to a stop; this is when all of its kinetic energy has been converted to electric potential energy. Therefore

$$E_{\alpha} = \frac{1}{2}m_{\alpha}v_{\alpha}^2 = k\frac{q_{\alpha}Q_{\text{gold}}}{r_c}$$

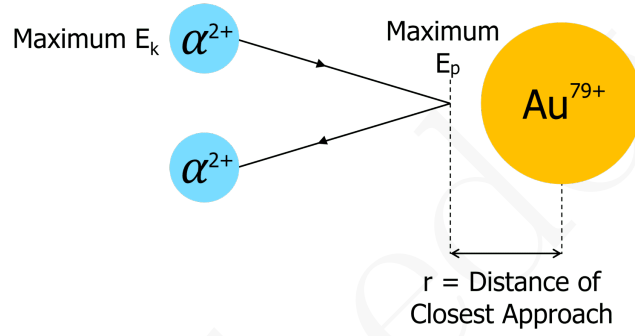


Figure 4: The distance of closest approach

- The charge of the gold nucleus is $Q_{\text{gold}} = Z \times e$, where Z is the atomic number of gold and e is the elementary charge.
- The charge of the alpha particle is $q_{\alpha} = 2e$.

3. Rearranging gives

$$\begin{aligned} r_c &= \frac{kq_{\alpha}Q_{\text{gold}}}{E_{\alpha}} \\ &= \frac{2kZe^2}{E_{\alpha}} \end{aligned}$$

5 Nuclear Density

The volume V of an atom is directly proportional to the nucleon number A

$$V \propto A$$

assuming the volume is a sphere, we then have the following, where R is the radius of the nuclide:

$$R^3 \propto A \iff R \propto A^{\frac{1}{3}}$$

The constant of proportionality is the Fermi radius $R_0 = 1.2 \times 10^{-15}\text{m}$, and thus

$$R = R_0 \sqrt[3]{A}$$

We can form an expression for the volume:

$$V_{\text{nucleus}} = \frac{4}{3}\pi AR_0^3$$

Assuming a proton and a neutron have roughly the same mass of $u = 1.66 \times 10^{-27}\text{kg}$ (this is the **unified atomic mass unit**), we can find the total mass of the nucleus:

$$M_{\text{nucleus}} = A \times u$$

The nuclear density is hence given by

$$\rho_{\text{nucleus}} = \frac{M_{\text{nucleus}}}{V_{\text{nucleus}}} = \frac{A \times u}{\frac{4}{3}\pi AR_0^3} = \frac{3u}{4\pi R_0^3}$$

A quick substitution gives the nuclear density as $\rho_{\text{nucleus}} = 2.4 \times 10^{17} \text{ kg m}^{-3}$. This is the same for all nuclides.

It also must be noted that the diameter of an atom is 10^5 times the diameter of the nucleus.

6 Emission and Absorption Spectra

6.1 Energy Levels of Electrons

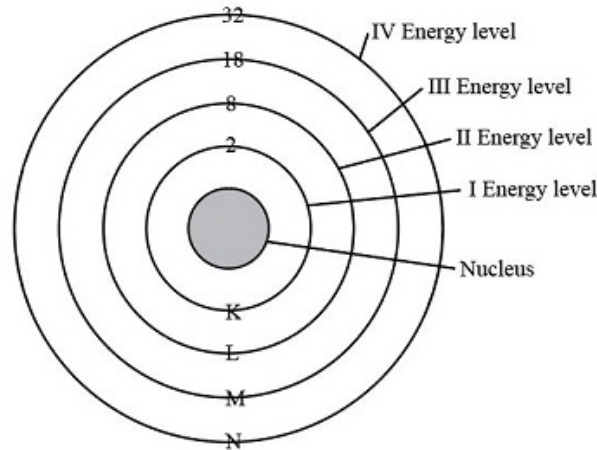


Figure 5: Energy levels of electrons

Electrons exist at **discrete energy levels** in an atom; this can be evidenced by the emission and absorption spectra of an atom.

- The ground state is the energy level at which the electron normally resides, this is the lowest energy level (max. negative).
- There is a limit of maximum energy level, beyond which, the electron is **ionized** and no longer part of the atom.
- When an electron exists at another energy level other than the ground state, it is said to be **excited**.
- An electron can be excited by:
 - absorbing a photon of energy
 - colliding with a nearby particle

Electrons do not remain excited for long, they will eventually relax and return to their ground state by emitting a photon of energy. The energy of the photon is equal to the difference in energy levels between the excited state and the ground state.

6.2 Emitting and Absorbing Photons

The energy of the photon emitted is given by

$$E = hf$$

where h is Planck's constant. The frequency of the photon is given by the equation

$$f = \frac{c}{\lambda}$$

where c is the speed of light and λ is the wavelength of the photon.

As previously mentioned, **the energy levels of electrons are discrete/quantized**, i.e. an electron cannot sit between two energy levels. This means that the energy of the photon emitted/absorbed is also discrete.

- If a photon whose energy does not match any of the energy differences between the discrete levels passes through the atom, **it will not be absorbed**.
- If we plot the wavelength of the emitted photons, we get an emission spectrum with a **black background and colored lines**.

6.2.1 Photons

Discrete packets of energy. The energy E is given by

$$E = hf$$

where h is Planck's constant. I.e., the energy in a photon is proportional to its frequency.

6.3 The Spectra

There is a one-to-one correspondence between every element and a spectrum — each element has a unique spectrum and each spectrum can uniquely determine an element. In other words, the spectrum of an element is its fingerprint that provides information about the chemical composition of an element.

6.3.1 The Absorption Spectrum

White light contains all wavelengths of light, and so it will occupy the entire spectrum.

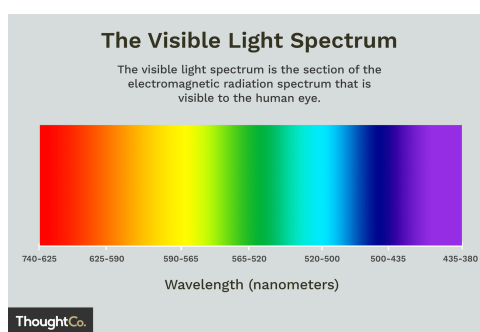


Figure 6: The visible light spectrum

If we shine the light through a gas of the element that we are studying, the gas will absorb photons of specific energy levels and hence wavelengths. Since they are absorbed, they will be missing from the observed spectrum.

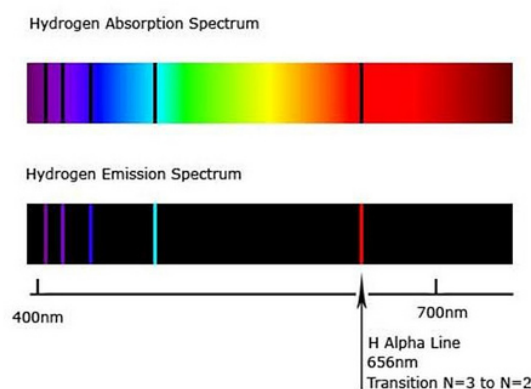


Figure 7: The emission and absorption spectra

6.4 Caveat with Energy Level Changes

Each energy change only occurs between a unique pair of levels in an atom, and doesn't occur twice in the same atom. This means that each line on an absorption or emission spectrum is associated with a unique pair of energy levels of the electron.

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7 The Bohr Model

7.1 Problem with the Rutherford Model & Classical Physics

- Classical electromagnetic theory predicted that accelerating charged particles, like electrons in orbit, should lose energy by emitting electromagnetic radiation.
- If this were true, the electrons would lose energy, **spiral into the nucleus**, and atoms would collapse, which contradicted experimental observations, whereby atoms are stable and do not collapse.

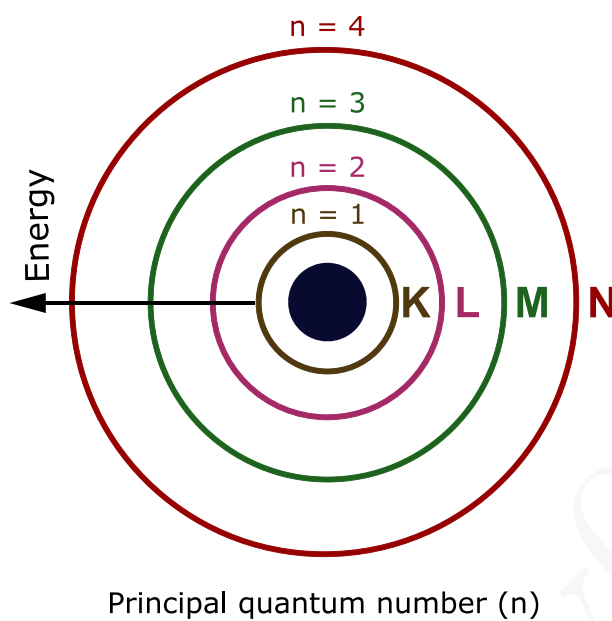
7.2 Bohr's Solution

Bohr's atomic model has the property where electrons orbit the nucleus in **circular orbits** at fixed radii. These orbits are **quantized**, i.e. the electron can only exist in certain orbits. These orbits are called **stationary states**, where the electron does not emit radiation, hence does not lose energy and spiral into the nucleus.

One of the limitations and reasons why this model does not hold is that it is **only applicable to hydrogen-like atoms**, i.e. atoms with only one electron. It fails to predict the atomic behavior for atoms with more than one electron in orbit.

7.3 Energy Levels

The energy levels of the electron in the Bohr model are quantized, i.e. the electron can only exist in certain energy levels. These levels are labeled by the principal quantum number $n \in \mathbb{Z}^+$, where $n = 1$ is the ground state, and $n = 2, 3, 4, \dots$ are the excited states.



The total (kinetic and electric potential) energy in an electron at the n th energy level is given by

$$E = -\frac{13.6}{n^2} \text{ eV}$$

Suppose energy possessed by an electron at the n th level is given by $-\frac{E_0}{n^2}$, then, for a transition between level $n = i$ and $n = j$ where $i < j$:

- Energy absorbed/emitted by the electron is $\Delta E = \frac{E_0}{i^2} - \frac{E_0}{j^2}$
- The frequency of the photon absorbed/emitted is given by $f = \frac{\Delta E}{h}$

7.4 Quantized Angular Momentum

The **de Broglie wavelength** of a charged particle is given by

$$\lambda = \frac{h}{p} = \frac{h}{mv} \quad (1)$$

where h is Planck's constant and p is the momentum of the particle.

One of the assumptions of Bohr that the angular momentum of an electron in a stationary state is **quantized** led to the implication that each orbital circumference is an integer multiple of the de Broglie wavelength of the electron "joined up". This means that, at the n th level, the orbital circumference is

$$n\lambda = 2\pi r \quad (2)$$

Combining (1) and (2), we get the following expression for the angular momentum

$$L = mvr = \frac{nh}{2\pi} \quad (3)$$

where m is the mass of the electron, v is the velocity of the electron, and r is the radius of the orbit. This quantization has the following implications

- The orbitals must have the nature of standing waves.
- Confirms the quantization of energy.

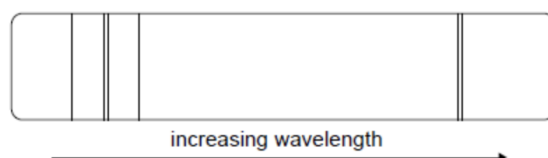
Also, if we rearrange [Equation \(2\)](#), we get the following expression for the orbital speed:

$$v = \frac{nh}{2\pi mr} \quad (4)$$

8 Exam Questions

8.1 MCQ – Energy Levels #1

The diagram shows the emission spectrum of an atom.



Which of the following atomic energy level models can produce this spectrum?

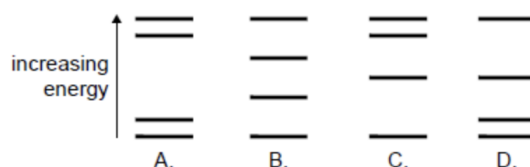


Figure 8: Exam question

Firstly, recall that

$$E = \frac{hc}{\lambda}$$

so the energy is inversely proportional to the wavelength. This means that, the larger transitions will be reflected in the smaller wavelengths that are on the left side. It is also important to recall how a line on the emission spectrum is produced: When an electron undergoes a transition from one energy level to another, it emits a photon of energy equal to the difference in energy levels. Thus, the number of possible transitions is equal to the number of lines on the emission spectrum. In general, with n energy levels, we have $\binom{n}{2} = \frac{n(n-1)}{2}$ possible transitions. If we picture all the possible transitions, for example, for A.:



Figure 9: Transitions

All that is left to do is to compare the lengths of the segments to the positions of the lines on the spectrum. There are many ways you could reason an answer:

- Consider the second and third longest transitions, they appear to be almost identical — this matches the pattern of the two extremely close lines on the left side of the spectrum.
- Now consider the two shortest transitions, they are also almost identical, which matches the two extremely close lines on the right side of the spectrum.

The two observations alone would suffice to deduce that A is indeed the answer.

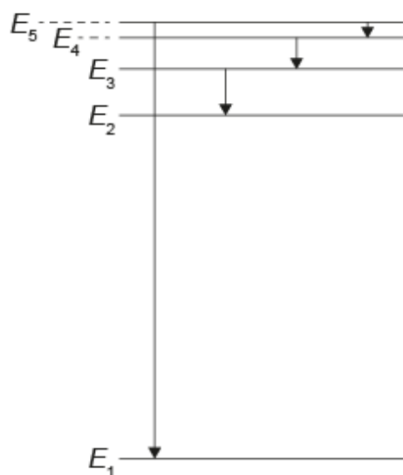
The rule of thumb for looking at energy transitions is that the longer the transition on an energy level diagram

- the shorter the wavelength
- the higher the energy
- the higher the frequency

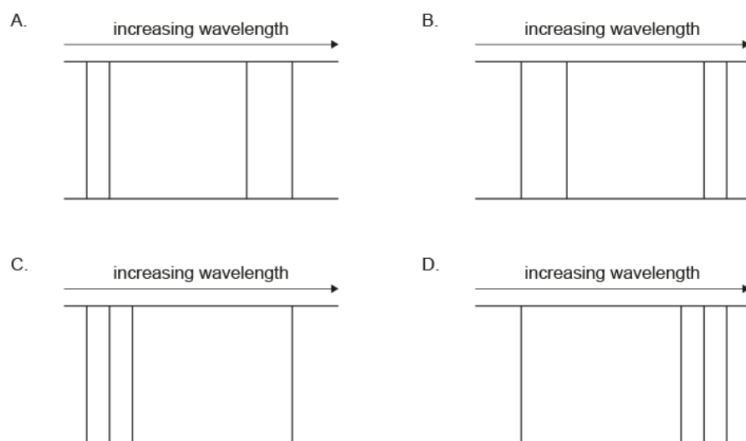
8.2 MCQ – Energy Levels #2

You should get this question within 10 seconds!

The energy levels E of an atom are shown.



Which emission spectrum represents the transitions?



[1]

Figure 10: Exam question

- There are three transitions with quite low energy releases (thus long wavelengths), and the small differences between them suggest that they should appear as three close lines on the spectrum on the right.
- Already that suffices to deduce D as the answer.

8.3 Angular Momentum – MCQ

In the Bohr model for hydrogen, the radius of the electron orbit in the $n = 2$ state is four times that of the radius in the $n = 1$ state.

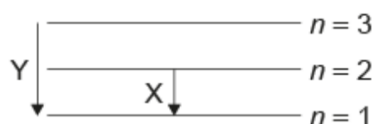
What is the ratio of the speed of the electron in the $n = 2$ state to that in the $n = 1$ state?

- We invoke $v_1 = \frac{nh}{2\pi mr}$
- For the second electron, $n \rightarrow 2n$, $r \rightarrow 4r$, and $v_1 \rightarrow (2)(\frac{1}{4})v_1 = \frac{1}{2}v_1$

8.4 Energy Levels – MCQ #3

Some energy levels for a hydrogen atom are shown.

diagram not to scale



What is the $\frac{\text{wavelength emitted in transition X}}{\text{wavelength emitted in transition Y}}$?

PDF

- A. $\frac{1}{2}$
- B. $\frac{27}{32}$
- C. $\frac{32}{27}$
- D. 2

[1]

Figure 11: Exam question

- Recall that the energy released is related to the change in energy levels by

$$\Delta E = \frac{E_0}{i^2} - \frac{E_0}{j^2}$$

- Then

$$\Delta E_X \equiv \frac{1}{4} - \frac{1}{1} = -\frac{3}{4}$$

- And for ΔE_Y :

$$\Delta E_Y \equiv \frac{1}{9} - \frac{1}{1} = -\frac{8}{9}$$

- Because we know that the energy released is inversely proportional to the wavelength, we now invert the fractions to get

$$\lambda_X = \frac{4}{3}$$

$$\lambda_Y = \frac{9}{8}$$

- The desired fraction is then

$$\frac{\lambda_X}{\lambda_Y} = \frac{4}{3} \div \frac{9}{8} = \frac{32}{27}$$