

# Galaxy Zoo Builder: Morphological Dependence of Spiral Galaxy pitch angle

Timothy Lingard<sup>1\*</sup>, Karen L. Masters<sup>2</sup>, Coleman Krawczyk<sup>1</sup>, Robert C. Nichol,<sup>1</sup>

<sup>1</sup>*Institute of Cosmology and Gravitation, University of Portsmouth, Dennis Sciama Building, Burnaby Road, Portsmouth, PO1 3FX, UK*

<sup>2</sup>*Haverford College, 370 Lancaster Ave., Haverford, PA 19041, USA*

Accepted XXX. Received YYY; in original form ZZZ

## ABSTRACT

Abstract

**Key words:** galaxies: evolution – galaxies: spiral – galaxies: photometry

## 1 INTRODUCTION

Spiral structure is present in a majority of massive galaxies (Buta 1989, Lintott et al. 2008) yet the formation mechanisms through which spiral structure originates are still hotly debated. Spirals are as diverse as the theories proposed to govern their evolution; from the quintessential pair of well-defined arcs of the grand design spiral, to the fragmented arm segments of the flocculent spiral, to the disjointed multi-armed spiral. These variations on structure account for 18%, 50% and 32% of the population respectively (Elmegreen et al. 2011, Buta et al. 2015). The Hubble classification scheme (Hubble 1926), and its revisions and expansions (Sandage 1961, de Vaucouleurs et al. 1991) contain detailed variations of different types of spiral galaxy, divided by the presence of a bar and ordered by how obvious spiral arm patterns are, how tightly they are wound and the prominence of a central bulge.

Arms of spiral galaxies are the source of the vast majority of star formation in the Universe, and spirals rearrange disc gas and can lead to the formation of disc-like bulges (e.g. Kormendy & Kennicutt 2004). Studies of spiral morphology have found interesting correlations between spiral morphology and other galactic properties, such as a correlation between spiral tightness and central mass concentration (Yu & Ho 2019, though Hart et al. 2017 found no such relation) and tightness and rotation curve shape (Seigar et al. 2005, with rising rotation curves creating more open spiral structure). These predictions and observations provide compelling reasons for investigating their underlying rules and dynamics, as doing so is essential for understanding the secular evolution of disc galaxies.

Our current understanding of the mechanisms which drive spiral growth and evolution suggest that each of the different forms of spiral galaxy may be triggered primarily by different processes. Grand Design spirals are thought to

have undergone a tidal interaction, be driven by a bar (as seen in gas simulations, Sanders & Huntley 1976, Rodriguez-Fernandez & Combes 2008, and suggested for stars by Manifold theory, Romero-Gómez et al. 2006, Athanassoula et al. 2009a, Athanassoula et al. 2009b), or be obeying (quasi-stationary) density wave theory (QSDW theory), in which spiral arms are slowly evolving, ever-present structures in the disc (Lin & Shu 1964). Flocculent spirals are thought to be formed through swing amplification (shearing of small gravitational instabilities in the disc), and be transient and recurrent in nature (Julian & Toomre 1966). It is recognised that methods of spiral formation are not mutually exclusive.

One of the fundamental assumptions of early work on spiral formation mechanisms (primarily QSDW) was that the disc of a galaxy, if unstable to spiral perturbations, would create a stable, static wave which would exist unchanging for many rotational periods (Lin & Shu 1964). The motivation for static waves with small numbers of arms (with a preference for  $m = 2$ ) was primarily observational; most galaxies show spiral structure, suggesting that spirals exist for a long time or are continually rebuilt.

Many simulations demonstrate that spirals do not maintain a constant pitch angle, and instead wind-up over time due to the differential rotation of the disc (Baba et al. 2013). Recent research suggests that spirals arms are transient in nature, and continually dissipate and re-form (Dobbs & Baba 2014). These spirals can be maintained through the same mechanisms that drive QSDW spirals (i.e. WASER, Mark 1976, swing amplification, Goldreich & Lynden-Bell 1965), but do not require the idealistic disc conditions required for the formation and maintenance of a stationary wave. The pitch angles of these transient spiral arms will decrease due to the differential rotation of the disk, with the density of the arm peaking at some critical pitch angle, before dissipating to be reformed.

In this dynamic picture of spiral arms, pitch angle monotonically decreases from a spiral arm's formation to its dissipation. Pringle & Dobbs (2019) proposes a simple

\* E-mail: tklingard@gmail.com

test of spiral arm winding, assuming the cotangent of the pitch angle of a spiral arm evolves linearly with time. They found that the distribution of pitch angles of their sample of 86 galaxies was consistent with this prediction, evidence against QSDW theory in favour of the dynamic spirals produced in many simulations.

Spiral evolution also appears to be influenced by the presence and strength of a bar; in barred grand-design spirals the arms often appear to start from the ends of the bar. Simulations of gas in barred galaxies often demonstrate that bars can drive long-term spiral evolution (Rodríguez-Fernández & Combes 2008), or boost transient spiral structure (Grand et al. 2012). Manifold theory is one attempt to determine the orbits of stars in bar-driven spiral arms: it proposes that stars in the vicinity of the unstable Lagrangian points at either end of the bar tend to escape along predictable orbits, governed by invariant manifolds. One of the primary factors influencing the shape of this invariant manifold is the relative strength of the non-axisymmetric forcing caused by the bar, with stronger bars resulting in spirals with larger pitch angles.

Many other systems contribute to spiral morphology, including potential ties to bulge fraction (Yoshizawa & Wakamatsu 1975, Savchenko & Reshetnikov 2013, Masters et al. 2019) and black hole mass (Seigar et al. 2008, Davis et al. 2017, Al-Baidhany et al. 2019). Stronger bulges and more massive central black holes have both been linked to more tightly wound spiral arms.

This paper makes use of the classification data and fitted photometric models obtained through the *Galaxy Builder* citizen science project for the 196 spiral galaxies present in [\[Lingard et al. \(2020\)\]](#). In this paper we focus on the use of measured spiral tightness (quantified using pitch angle, Binney & Tremaine 1987) as a probe into the dynamical mechanisms governing a spiral galaxy’s evolution. We make use of Bayesian hierarchical modelling to measure galaxy pitch angle from the spiral arm clusters produced by *Galaxy Builder*.

Section 3.3 investigates spiral arm winding using the test derived by Pringle & Dobbs (2019) (uniformity of galaxy pitch angle in  $\cot \phi$ ) and concludes using a marginalized Anderson-Darling test that we cannot unilaterally reject winding of this form at the 1% level. Section 3.2 examines the correlation between pitch angle and bulge size implied by the Hubble sequence, and pitch angle and bar strength implied by Manifold theory, and find no significant correlation.

## 2 METHOD

### 2.1 Measuring galaxy pitch angle

Many methodologies have been proposed and implemented to measure spiral arm properties, including visual inspection (Herrera-Endoqui et al. 2015), fourier analysis (i.e. 2DFFT, Davis et al. 2012), texture analysis (i.e. SpArcFiRe, Davis & Hayes 2014), and combinations of automated methods and human classifiers (Hart et al. 2017, Hewitt & Treuthardt 2020). One potentially underused method of obtaining measurements of spirals is through photometric fitting of spiral structure, as possible using tools such as GALFIT (Peng

et al. 2010) and *Galaxy Builder* ([\[Lingard et al. \(2020\)\]](#)). These methods attempt to localize light from an image of a galaxy into distinct subcomponents, such as a galaxy disc, bulge, bar and spiral arms, generally finding the optimum solution using computational optimization. This optimization process, however, is often not robust for complex, many-component models and requires significant supervision to converge to a physically meaningful result (Gao & Ho 2017). [\[Lingard et al. \(2020\)\]](#) propose a solution to this problem through the use of citizen science to provide priors on parameters used in computational fitting.

A common assumption when measuring galaxy pitch angle is that observed spiral arms have a constant pitch angle. Spirals of this kind are known as logarithmic spirals and are described by

$$r = A e^{\theta \tan \phi}, \quad (1)$$

where  $\phi$  is the arm’s pitch angle. One method used to obtain a pitch angle of a galaxy is to fit logarithmic spirals to individually identified arm segments and take the weighted mean of their pitch angles (which often vary by upwards of  $10^\circ$ , Davis & Hayes 2014). Weighting is determined by the length of the arc segment, with longer being assigned higher weights, i.e. for a galaxy where we have identified  $N$  arm segments, each with length  $L_i$  and pitch angle  $\phi_i$

$$\phi_{\text{gal}} = \left( \sum_{i=1}^N L_i \right)^{-1} \sum_{i=1}^N L_i \phi_i. \quad (2)$$

The most commonly used measurement of uncertainty of length-weighted pitch angles is the unweighted sample variance between the arm segments identified.

A notable drawback of length-weighted pitch angle is sensitivity to the number and quality of the spiral arm segments identified; Hart et al. (2017) found that only 15% of the arm segments identified by a leading algorithm (SPARCFIRE) were identified as “good” matches to real spiral arms by citizen science classifiers.

Fourier analysis in one- and two-dimensions (as performed by Díaz-García et al. 2019, Davis et al. 2012, Mutlu-Pakdil et al. 2018) is another widely used method of computationally obtaining galaxy pitch angles. Two-dimensional fourier methods generally decompose a deprojected image of a galaxy into a superpositions of logarithmic spirals between inner and outer annuli (Davis et al. 2012), and reports the pitch angle with the highest amplitude as the galaxy’s pitch angle. Hewitt & Treuthardt (2020) combined fourier analysis of spiral galaxies with human tracing of spiral arms, to great effect. It is uncertain how variation between pitch angles of individual arms impacts this measurement.

### 2.2 The Galaxy Sample

The galaxies analysed in this paper are the 198 galaxies from [\[Lingard et al. \(2020\)\]](#). These are a subset of the *stellar mass-complete sample* in Hart et al. (2017), a sample of low-redshift face-on spiral galaxies selected using data from the NASA-Sloan Atlas (Blanton et al. 2011) and Galaxy Zoo 2 (Willett et al. 2013).

Some galaxies in [\[Lingard et al. \(2020\)\]](#) were shown

to volunteers again in a repeat *validation subset* in order to create a second aggregate model used to test internal consistency. We combine the classifications of galaxies in this *validation subset* with the original classifications. Clustering of drawn spiral arms and cleaning of points was then performed as detailed in [[Lingard et al. (2020)]]. We remove any galaxies for which no spiral arms were identified, resulting in a hierarchical data structure of 139 galaxies, 261 spiral arms and 239,947 points.

Spiral arm points are deprojected to a face-on orientation using the disk inclination and position angle obtained through photometric model fitting performed in [[Lingard et al. (2020)]]. Arms are individually corrected to all have the same chirality (a pitch angle greater than or equal to zero) using the logarithmic spiral fit in [[Lingard et al. (2020)]].

### 2.3 Bayesian modelling of spiral arms in *Galaxy Builder*

We assume that a galaxy has a single value for pitch angle,  $\phi_{\text{gal}}$ , and that the pitch angles of spiral arms in the galaxy,  $\phi_{\text{arm}}$ , are constant (giving logarithmic spirals) and drawn from a normal distribution centred on  $\phi_{\text{gal}}$ , with some spread  $\sigma_{\text{gal}}$  common to all galaxies. We truncate this normal distribution between the physical limits of  $0^\circ$  (a ring) and  $90^\circ$  (a “spoke”), giving

$$\phi_{\text{arm}} \sim \text{TruncatedNormal}(\phi_{\text{gal}}, \sigma_{\text{gal}}, \min = 0, \max = 90). \quad (3)$$

Furthermore, we assume that the observed points in a *Galaxy Builder* spiral arm, once deprojected, follow a logarithmic spiral with gaussian radial error  $\sigma_r$ ,

$$\widetilde{r}_{\text{arm}} = \exp\left(\overrightarrow{\theta}_{\text{arm}} \tan \phi_{\text{arm}} + c_{\text{arm}}\right). \quad (4)$$

Where  $\widetilde{r}_{\text{arm}}$  is the model’s predictions for the radii of the deprojected points in a *Galaxy Builder* arm ( $\overrightarrow{r}_{\text{arm}}$ ), and  $\overrightarrow{\theta}_{\text{arm}}$  is the polar angles of the points.

We choose hyperpriors over  $\phi_{\text{gal}}$ ,  $\sigma_{\text{gal}}$ ,  $c_{\text{arm}}$  and  $\sigma_r$  of

$$\phi_{\text{gal}} \sim \text{Uniform}(\min = 0, \max = 90), \quad (5)$$

$$\sigma_{\text{gal}} \sim \text{InverseGamma}(\alpha = 2, \beta = 20), \quad (6)$$

$$c_{\text{arm}} \sim \text{Cauchy}(\alpha = 0, \beta = 10), \quad (7)$$

$$\sigma_r \sim \text{InverseGamma}(\alpha = 2, \beta = 0.5). \quad (8)$$

The inverse gamma distribution is used to aid the convergence of the Hamiltonian Monte Carlo (HMC) algorithm used (discussed later). The Cauchy distribution is equivalent to the Student’s t-distribution with one degree of freedom, and was chosen due to its fatter tails than the normal distribution. Our likelihood function for  $N$  arms, each of which with  $n_{\text{arm}}$  points is

$$\mathcal{L} = \prod_{\text{arm}=1}^N \left(2\pi\sigma_r^2\right)^{-n_{\text{arm}}/2} \exp\left(-\frac{\|\overrightarrow{r}_{\text{arm}} - \widetilde{r}_{\text{arm}}\|^2}{2\sigma_r^2}\right). \quad (9)$$

We assume that the radial error is Gaussian for simplicity of analysis, Shapiro-Wilk tests on the residuals of the logarithmic spirals fit in [[Lingard et al. (2020)]] suggest

that this is not a good assumption, and a more robust likelihood (such as the Student’s t-distribution) would possibly be more appropriate.

To perform inference, we make use of the No-U-Turn-Sampler (NUTS, Hoffman & Gelman 2011), implemented in PYMC3<sup>1</sup>, an open source probabilistic programming framework written in python (Salvatier et al. 2016). To aid the convergence of MC chains, we scale the radii of deprojected points to have unit variance.

## 3 RESULTS

### 3.1 Constraints on Galaxy Pitch angle

Our hierarchical model identifies the pitch angle of individual arms ( $\phi_{\text{arm}}$ ) with less than  $1.6^\circ$  of uncertainty for 95% of arms, assuming no error on disc inclination and position angle. The pitch angle of a galaxy as a whole ( $\phi_{\text{gal}}$ ), however, is not well constrained. This is primarily a result of only having pitch angles measurements for a small number of arms per galaxy, and reflects the difficulty in providing a single value for the pitch angle of a galaxy containing individual arms with very different pitch angles. For galaxies with two arms identified in *Galaxy Builder*, we have a mean uncertainty of ( $\sigma_{\phi_{\text{gal}}}$ ) of  $7.9^\circ$ , which decreases to  $6.8^\circ$  and  $6.0^\circ$  for galaxies with three and four arms respectively. This is roughly consistent with the standard error on the mean for a galaxy with  $N$  arms,

$$\sigma_{\phi_{\text{gal}}} = \frac{\sigma_{\text{gal}}}{\sqrt{N}}, \quad (10)$$

where  $\sigma_{\text{gal}}$  is our measure of inter-arm variability of pitch angle and has a posterior distribution of  $11.0^\circ \pm 0.9^\circ$ . This variability is similar to the finding of Davis & Hayes (2014) and emphasises the need for fitting algorithms to not assume all arms have the same pitch angle.

### 3.2 Dependence of pitch angle on Galaxy Morphology

In order to test the possible progenitor distribution of our estimated arm pitch angles, we repeatedly perform an Anderson-Darling test (Stephens 1974, implemented in SCIPY, Jones et al. 2001) over each draw present in the MC trace, resulting in a distribution of Anderson-Darling statistics. We will refer to this test as the *marginalized Anderson-Darling test*. We also make use of the two-sample Anderson-Darling (Scholz & Stephens 1987) test in a similar manner.

#### 3.2.1 Pitch angle vs. Bulge size

Morphological classification commonly links bulge size to spiral tightness, and such a link is implied by the Hubble Sequence (Sandage 2005, Gadotti 2009, Buta 2013). Some studies have indeed reported a link between measured spiral galaxy pitch angle and bulge size (i.e. Hart et al. 2017, Davis

<sup>1</sup> <https://docs.pymc.io/>

et al. 2019), while others have not found any significant correlation (Masters et al. 2019). We investigate this relationship here using a measure of bulge prominence from Galaxy Zoo 2, as Equation 3 in Masters et al. (2019):

$$B_{\text{avg}} = 0.2 \times p_{\text{just noticeable}} + 0.8 \times p_{\text{obvious}} + 1.0 \times p_{\text{dominant}}, \quad (11)$$

Where  $p_{\text{just noticeable}}$ ,  $p_{\text{obvious}}$  and  $p_{\text{dominant}}$  are the fractions of classifications indicating the galaxy’s bulge was “just noticeable”, “obvious” or “dominant” respectively.

We see no correlation between galaxy pitch angle derived from the hierarchical model and  $B_{\text{avg}}$ . The Pearson correlation coefficient between the expectation value of galaxy pitch angle ( $E[\phi_{\text{gal}}]$ ) and  $B_{\text{avg}}$  is 0.00.

We separate our sample into galaxies with weaker bulges ( $B_{\text{avg}} < 0.28$ , 83 galaxies) and those with stronger bulges ( $B_{\text{avg}} \geq 0.28$ , 54 galaxies), in order to test whether their pitch angles could be drawn from significantly different distributions. A marginalized two-sample Anderson-Darling test comparing the distributions of  $\phi_{\text{gal}}$  for the samples does not find evidence that galaxy pitch angles were drawn from different distributions: we reject the null hypothesis at the 1% level for only 1% of the samples. Similarly comparing arm pitch angles for galaxies in the different samples results in not rejecting the null hypothesis at the 1% level for any of the samples. The distributions of Anderson-Darling test statistic for  $\phi_{\text{gal}}$  and  $\phi_{\text{arm}}$  are shown in the upper panel of Figure 1 in blue and orange respectively.

These results may be caused by the galaxy sample not containing many galaxies with dominant bulges:  $B_{\text{avg}}$  only varied from 0.09 to 0.75 (the allowed maximum being 1.0), with only four galaxies having  $B_{\text{avg}} > 0.5$ . The split point of 0.28 was also chosen to produce evenly sized comparison samples rather than from some physical motivation. However, the lack of any form of correlation implies that there is no evidence in our data for the link between bulge size and pitch angle predicted by the Hubble sequence and observed in other studies.

### 3.2.2 Pitch angle vs. Bar Strength

One of the predictions of Manifold theory is that pitch angle increases with bar strength (Athanasoulas et al. 2009b). In order to investigate this relationship in our data, we make use of Galaxy Zoo 2’s bar fraction ( $p_{\text{bar}}$ ), which is widely viewed as a good measure of bar strength (Skibba et al. 2012, Masters et al. 2012) and therefore a measure of the torque applied on the disc gas.

We do not observe a correlation between  $p_{\text{bar}}$  and  $E[\phi_{\text{gal}}]$  (Pearson correlation coefficient of -0.05). Following Masters et al. (2012) and Skibba et al. (2012), we separate the sample into galaxies without a bar ( $p_{\text{bar}} < 0.2$ ), with a weak bar ( $0.2 \leq p_{\text{bar}} \leq 0.5$ ) and with a strong bar ( $p_{\text{bar}} > 0.5$ ). Performing marginalized three-sample Anderson-Darling tests does not find that pitch angles ( $\phi_{\text{gal}}$  or  $\phi_{\text{arm}}$ ) of galaxies with different bar strengths were drawn from different distributions; we do not reject the null hypothesis at the 1% level for any samples for the test of  $\phi_{\text{gal}}$ , and at the 10% level for the test of  $\phi_{\text{arm}}$ . The distributions of Anderson-Darling test statistic is shown in the lower panel of Figure 1.

We do not correct for variation in bulge size in this work,

however predictions from Manifold theory should not be affected as bulges do not provide a strong non-axisymmetric focusing. The fact that we do not find any link between bar strength and pitch angle suggests that the primary mechanism driving the evolution of the spirals in our sample is not Manifold theory. **[[This is possibly too strong a statement?]]**

### 3.3 Spiral Winding

For transient and recurrent spiral arms driven by self-gravity, Pringle & Dobbs (2019) suggest that spiral patterns form at some maximum pitch angle ( $\phi_{\text{max}}$ ), continually wind up over time and finally dissipate at some minimum pitch angle ( $\phi_{\text{min}}$ ). They propose that, under a set of very simple assumptions, the evolution of pitch angle would be governed by

$$\cot \phi = \left[ R \frac{d\Omega_p}{dR} \right] (t - t_0) + \cot \phi_{\text{max}}, \quad (12)$$

where  $\Omega_p$  is the radially dependant pattern speed of the spiral arm and  $t_0$  is the initial time at which it formed.

In QSDW theory, the pattern speed  $\Omega_p$  is a constant in  $R$ , as spiral arms obey rigid-body rotation. If  $\Omega_p$  instead varies with radius we would expect  $\cot \phi$  to be uniformly distributed between  $\cot \phi_{\text{max}}$  and  $\cot \phi_{\text{min}}$ .

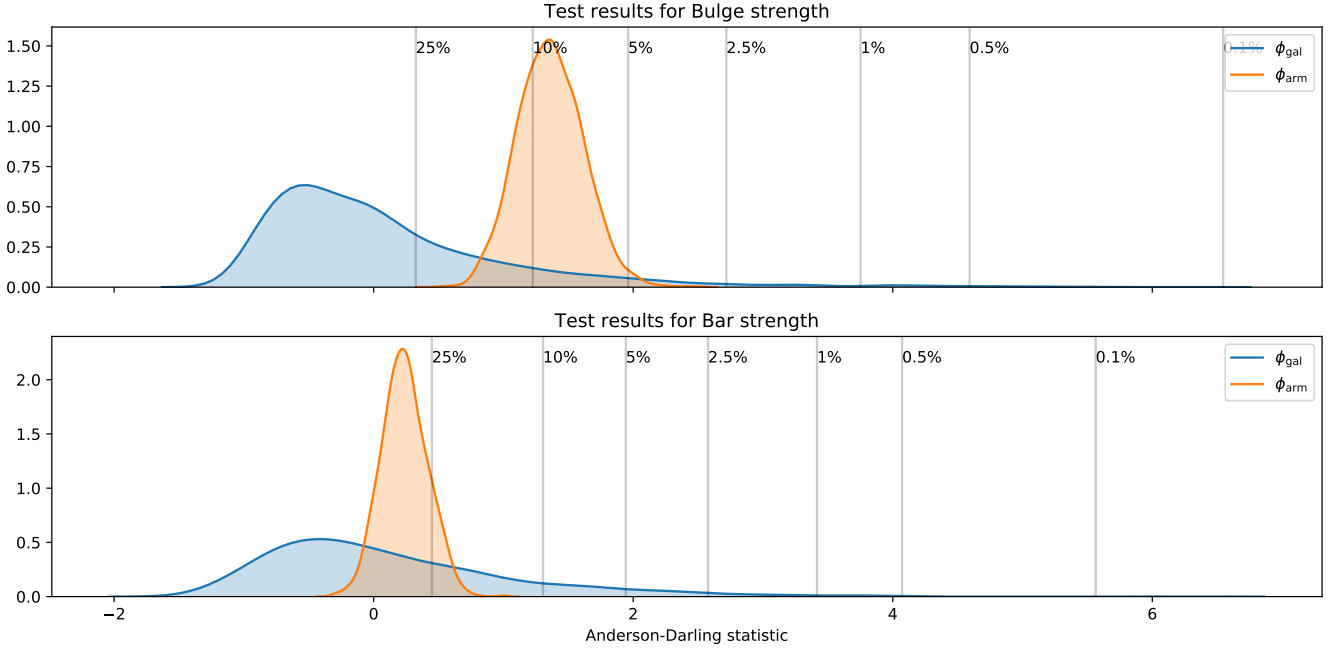
In order to test this theory, Pringle & Dobbs (2019) used a Kolmogorov-Smirnov test to examine the consistency of a sample of observed galaxy pitch angles with one uniform in  $\cot$ . Pitch angles were measured using discrete Fourier transformations in one- and two-dimensions, and as such do not account for inter-arm variations. They chose limits of  $\cot \phi \in [1.00, 4.75]$  (roughly  $11.9^\circ < \phi < 45.0^\circ$ ), motivated by examination of their data.

We aim to replicate their work here, using our sample and methods. We will make use of the marginalized Anderson-Darling test described above, and examine winding on a per-arm basis, as well as a per-galaxy basis. Observation of the distribution of arm pitch angles in our sample suggests limits of  $15^\circ < \phi < 50.0^\circ$ .

#### 3.3.1 Galaxy Pitch angle

Testing the uniformity of  $\cot \phi_{\text{gal}}$  between  $15^\circ$  and  $50^\circ$  using a marginalized Anderson-Darling test results in rejecting the null hypothesis at the 1% level for just 5% of samples, with a large spread in observed test values. The full distribution of Anderson-Darling statistics can be seen in the upper panel of Figure 2.

This result suggests that we cannot rule out a cot-uniform source distribution for galaxy pitch angle, but the large uncertainty in  $\phi_{\text{gal}}$  makes it difficult to make any conclusive statements. This result is also highly sensitive to the lower limit of  $\phi$ : decreasing it to  $10^\circ$  results in us rejecting the cot-uniform model at greater than the 0.1% level for 96% of the posterior samples. As we have no information available on the biases present in *Galaxy Builder* spiral arm classification, we choose to keep the less strict limit of  $15^\circ$ .



**Figure 1.** The results of marginalized two-sample Anderson-Darling tests examining whether pitch angles ( $\phi_{\text{gal}}$  in blue and  $\phi_{\text{arm}}$  in orange) for galaxies with  $B_{\text{avg}} < 0.28$  and  $B_{\text{avg}} \geq 0.28$  are drawn from the same distribution (top panel), and the results of marginalized three-sample Anderson-Darling tests for galaxies with no bar ( $p_{\text{bar}} < 0.2$ ), a weak bar ( $0.2 \leq p_{\text{bar}} \leq 0.5$ ) and a strong bar ( $p_{\text{bar}} > 0.5$ ) (bottom panel). Confidence intervals are shown, with moving rightwards indicating more confidence in rejecting the null hypothesis.

### 3.3.2 Arm Pitch angle

The inconclusive result for  $\phi_{\text{gal}}$  is perhaps unsurprising: were we to assume that spiral arms are transient and recurrent instabilities, there is little reason for all of the arms to be at precisely the same evolutionary stage at the same time. This is supported by the large observed spread in inter-arm pitch angles.

If we assume instead that spirals form and wind independently in a galaxy, and that their evolution over time can be described by Equation 12, the distribution of pitch angles of individual arms should be uniform in cot between our limits, rather than that of the galaxy’s pitch angle as a whole.

Using the marginalized Anderson-Darling test we cannot reject the null hypothesis at even the 5% level for any of the possible realizations of arm pitch angle. The resulting distribution of Anderson-Darling statistics is shown in the lower panel of Figure 2. This result is highly consistent with the model for spiral winding proposed by Pringle & Dobbs (2019), and can be seen as evidence that spirals are formed through local disc perturbations, and are primarily governed by local forces.

## 4 DISCUSSION

This paper presents a new Bayesian approach to estimate galaxy pitch angle, making use of citizen science results to measure spiral arms through photometric modelling. We introduce an adaptation of the Anderson-Darling test, which we name the *marginalized Anderson-Darling test*, to incorporate full Bayesian posterior probabilities and utilize this

test to investigate theories governing spiral formation and evolution.

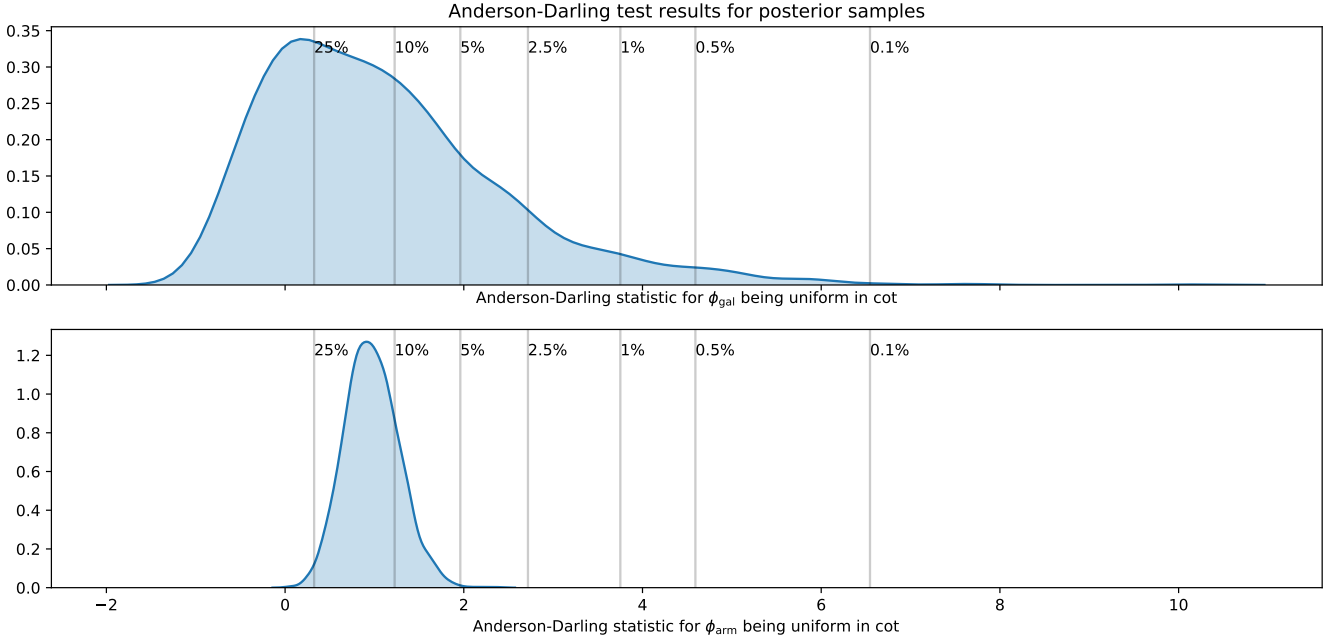
The statistical approach implemented in this paper allows a more thorough examination of pitch angle than length-weighted pitch angle calculation, and better accounts for the large variations observed in inter-arm pitch angle than fourier analysis, which assumes all arms in a given mode have the same pitch angle.

We do not see a relationship between bar strength and pitch angle in our data, as would have been predicted by Manifold theory, and do not find evidence for the relationship between bulge size and pitch angle predicted by the Hubble sequence. Graham & Driver (2007) propose that bulge Sérsic index is correlated with black hole mass, which has in turn been shown to correlate with spiral pitch angle, we do not investigate this here as *Galaxy Builder* could not accurately recover bulge Sérsic indices.

Our results are consistent with spiral winding of the form described by Pringle & Dobbs (2019), in which spiral arms are transient and recurrent, evolve through mechanisms such as swing-amplification (Goldreich & Lynden-Bell 1965) and which wind up over time. However, the assumptions of this model of spiral winding are highly simplistic, and it leaves many unanswered questions: what determines the limits on  $\phi$ ? Is the spiral arm equally apparent at all pitch angles, or is a selection effect present? This result is also not evidence against QSDW, as it is possible that our distribution of pitch angles is dictated by other factors such as disk shear and central mass concentration.

We do not account for observation effects, instead assuming that galaxy builder spiral arms are equally likely to be identified and recovered at all pitch angles within the lim-





**Figure 2.** The results of a marginalized Anderson-Darling test for uniformity in  $\cot$  for  $\phi_{\text{gal}}$  (top panel) and  $\phi_{\text{arm}}$  (bottom panel), with values corresponding to various confidence intervals shown. Moving rightwards on the x-axis implies greater confidence in rejecting the null hypothesis.

its specified above. We also assume that the galaxy sample is representative of the general spiral population, which is not necessarily the case. As with most analyses, the most impactful improvement it would be possible to make here would be to increase the cleanliness and volume of data analysed; due to time constraints the *Galaxy Builder* galaxies used are not guaranteed to be representative.

The methodology proposed here is a robust solution to the problems facing investigation of spiral morphology, namely that of reliably identifying spiral arms, and properly accounting for the spread in pitch angles of arms within a galaxy.

## 5 ACKNOWLEDGEMENTS

This publication made use of SDSS-I/II data. Funding for the SDSS and SDSS-II was provided by the Alfred P. Sloan Foundation, the Participating Institutions, the National Science Foundation, the U.S. Department of Energy, the National Aeronautics and Space Administration, the Japanese Monbukagakusho, the Max Planck Society, and the Higher Education Funding Council for England. The SDSS Web Site is <http://www.sdss.org/>.

This publication uses data generated via the Zooniverse.org platform, development of which is funded by generous support, including a Global Impact Award from Google, and by a grant from the Alfred P. Sloan Foundation. We would also like to thank the 2,340 volunteers who have submitted classifications to the *Galaxy Builder* project, especially user EliabethB, whose presence on the *Galaxy Builder* forum on top of a large number of galaxies modelled has been a huge help.

This project was partially funded by a Google Faculty Research Award to Karen Masters (<https://ai.google/research/outreach/faculty-research-awards/>), and Timothy Lingard acknowledges studentship funding from the Science and Technology Facilities Council (ST/N504245/1).

## References

- Al-Baidhany I. A., Chiad S. S., Jabbar W. A., Hussein R. A., Hussain F. F. K., Habubi N. F., 2019, in *Materials Science and Engineering Conference Series*. p. 012118, doi:10.1088/1757-899X/571/1/012118
- Athanassoula E., Romero-Gómez M., Masdemont J. J., 2009a, *MNRAS*, 394, 67
- Athanassoula E., Romero-Gómez M., Bosma A., Masdemont J. J., 2009b, *MNRAS*, 400, 1706
- Baba J., Saitoh T. R., Wada K., 2013, *ApJ*, 763, 46
- Binney J., Tremaine S., 1987, *Galactic dynamics*
- Blanton M. R., Kazin E., Muna D., Weaver B. A., Price-Whelan A., 2011, *AJ*, 142, 31
- Buta R., 1989, *Galaxy Morphology*. p. 151
- Buta R. J., 2013, *Galaxy Morphology*. p. 155
- Buta R. J., et al., 2015, *VizieR Online Data Catalog*, p. J/ApJS/217/32
- Davis D. R., Hayes W. B., 2014, *ApJ*, 790, 87
- Davis B. L., Berrier J. C., Shields D. W., Kenefick J., Kenefick D., Seigar M. S., Lacy C. H. S., Puerari I., 2012, *ApJS*, 199, 33
- Davis B. L., Graham A. W., Seigar M. S., 2017, *MNRAS*, 471, 2187
- Davis B. L., Graham A. W., Cameron E., 2019, *ApJ*, 873, 85
- Díaz-García S., Salo H., Knapen J. H., Herrera-Endoqui M., 2019, *arXiv e-prints*, p. arXiv:1908.04246
- Dobbs C., Baba J., 2014, *Publ. Astron. Soc. Australia*, 31, e035

- Elmegreen D. M., et al., 2011, *ApJ*, 737, 32
- Gadotti D. A., 2009, *MNRAS*, 393, 1531
- Gao H., Ho L. C., 2017, *ApJ*, 845, 114
- Goldreich P., Lynden-Bell D., 1965, *MNRAS*, 130, 125
- Graham A. W., Driver S. P., 2007, *ApJ*, 655, 77
- Grand R. J. J., Kawata D., Cropper M., 2012, *MNRAS*, 426, 167
- Hart R. E., et al., 2017, *MNRAS*, 472, 2263
- Herrera-Endoqui M., Díaz-García S., Laurikainen E., Salo H., 2015, *A&A*, 582, A86
- Hewitt I. B., Treuthardt P., 2020, *MNRAS*, 493, 3854
- Hoffman M. D., Gelman A., 2011, *arXiv e-prints*, p. arXiv:1111.4246
- Hubble E. P., 1926, *ApJ*, 64, 321
- Jones E., Oliphant T., Peterson P., et al., 2001, *SciPy: Open source scientific tools for Python*, <http://www.scipy.org/>
- Julian W. H., Toomre A., 1966, *ApJ*, 146, 810
- Kormendy J., Kennicutt Robert C. J., 2004, *ARA&A*, 42, 603
- Lin C. C., Shu F. H., 1964, *ApJ*, 140, 646
- Lintott C. J., et al., 2008, *MNRAS*, 389, 1179
- Mark J. W. K., 1976, *ApJ*, 205, 363
- Masters K. L., et al., 2012, *MNRAS*, 424, 2180
- Masters K. L., et al., 2019, *MNRAS*, 487, 1808
- Mutlu-Pakdil B., Seigar M. S., Hewitt I. B., Treuthardt P., Berrier J. C., Koval L. E., 2018, *MNRAS*, 474, 2594
- Peng C. Y., Ho L. C., Impey C. D., Rix H.-W., 2010, *AJ*, 139, 2097
- Pringle J. E., Dobbs C. L., 2019, *arXiv e-prints*, p. arXiv:1909.10291
- Rodriguez-Fernandez N. J., Combes F., 2008, *A&A*, 489, 115
- Romero-Gómez M., Masdemont J. J., Athanassoula E., García-Gómez C., 2006, *A&A*, 453, 39
- Salvatier J., Wiecki T. V., Fonnesbeck C., 2016, *PeerJ Computer Science*, 55
- Sandage A., 1961, *The Hubble Atlas of Galaxies*
- Sandage A., 2005, *ARA&A*, 43, 581
- Sanders R. H., Huntley J. M., 1976, *ApJ*, 209, 53
- Savchenko S. S., Reshetnikov V. P., 2013, *MNRAS*, 436, 1074
- Scholz F. W., Stephens M. A., 1987, *Journal of the American Statistical Association*, 82, 918
- Seigar M. S., Block D. L., Puerari I., Chorney N. E., James P. A., 2005, *MNRAS*, 359, 1065
- Seigar M. S., Kennefick D., Kennefick J., Lacy C. H. S., 2008, *ApJ*, 678, L93
- Skibba R. A., et al., 2012, *MNRAS*, 423, 1485
- Stephens M. A., 1974, *Journal of the American Statistical Association*, 69, 730
- Willett K. W., et al., 2013
- Yoshizawa M., Wakamatsu K., 1975, *A&A*, 44, 363
- Yu S.-Y., Ho L. C., 2019, *ApJ*, 871, 194
- de Vaucouleurs G., de Vaucouleurs A., Corwin Herold G. J., Buta R. J., Paturel G., Fouque P., 1991, *Third Reference Catalogue of Bright Galaxies*  
Appendix

This paper has been typeset from a  $\text{\LaTeX}$  file prepared by the author.