Note For Noether's Theorem

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1 Noether Theorem in Classical Mechanics

1.1 Lagrangian Formalism

In Lagrangian formalism, we've defined the action by lagrangian,

$$S = \int_{t_1}^{t_2} L[x(t), \dot{x}(t)] dt$$
 (1.1)

If action follows the *least action principle*, we would have the lagrangian be associated with the correct *equation of motion* of the system given by eq(1.1).

$$\frac{\partial L}{\partial x_i} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}_i} = 0 \tag{1.2}$$

Noether's theorem states that if given system has continuous symmetry, then there would be a conserved quantity associated with this symmetry.

What is continuous symmetry? One that leaves the action invariant even when the dynamical equations (equations of motion) are *not* satisfied, we call it **infinitesimal symmetry transformation**. We denote the infinitesimal symmetry transformation as δ_S to separate it from the variation δ .

We're interested in the case when the action doesn't change under a symmetry, and this implies there is special property behind the dynamics of the system.

If we would like to discuss the transformation on time, we should use another parameter,

$$t = t(\tau)$$

$$q(t(\tau)) = Q(\tau)$$
(1.3)

Generally, the transformation would be

$$t \to t' = \tau'(t')$$

$$q \to q'(t') = \mathcal{Q}(\tau')$$
(1.4)

and the variation of position and time¹,

$$\delta_{S}q = q'(t) - q(t) = \mathcal{Q}'(\tau) - \mathcal{Q}(\tau) = \delta_{S}\mathcal{Q}$$

$$\delta_{S}t = t' - t = \tau'(t) - \tau(t) = \delta_{S}\tau$$
(1.5)

Rewrite the original action,

$$S = \int_{t_1}^{t_2} dt \, L[q(t), \dot{q}(t)] = \int_{\tau_1}^{\tau_2} d\tau \, \left(\frac{dt}{d\tau}\right) L[\mathcal{Q}, \frac{d}{d\tau}\mathcal{Q}, t] \tag{1.6}$$

Redefine the new lagrangain,

$$\mathbb{L}\left[\mathcal{Q}, \frac{d}{d\tau}\mathcal{Q}, t, \frac{dt}{d\tau}\right] = \frac{dt}{d\tau}L\tag{1.7}$$

Now consider the infinitesimal symmetry transformation, and the new action would become,

$$S' = \int d\tau \, \mathbb{L} \left[\mathcal{Q}', \frac{d\mathcal{Q}'}{d\tau}, t', \frac{dt'}{d\tau} \right] + \int d\tau \frac{dK}{d\tau}$$
 (1.8)

Adding a term of total derivative for general consideration, and we know it wouldn't affect the equation of motion. Therefore we can do the variation of action and lagrangian,

$$\delta S = S' - S = 0 \tag{1.9}$$

The variation of action should vanish because we assume it is symmetry infinitesimal transformation. Let's consider the equation of motion corresponding to the new lagrangian $\mathbb{L}\left[\mathcal{Q},\frac{d\mathcal{Q}}{d\tau},t,\frac{dt}{d\tau}\right]$.

$$\frac{\partial \mathbb{L}}{\partial \mathcal{Q}} - \frac{d}{d\tau} \frac{\partial \mathbb{L}}{\partial d_{\tau} \mathcal{Q}} = 0$$

$$\frac{\partial \mathbb{L}}{\partial t} - \frac{d}{d\tau} \frac{\partial \mathbb{L}}{\partial d_{\tau} t} = 0$$
(1.10)

¹The parameter should be the same.